Zurich Phenomenology Workshop 2023

EFT-Diagrammatic Approach to Compact Binary Dynamics

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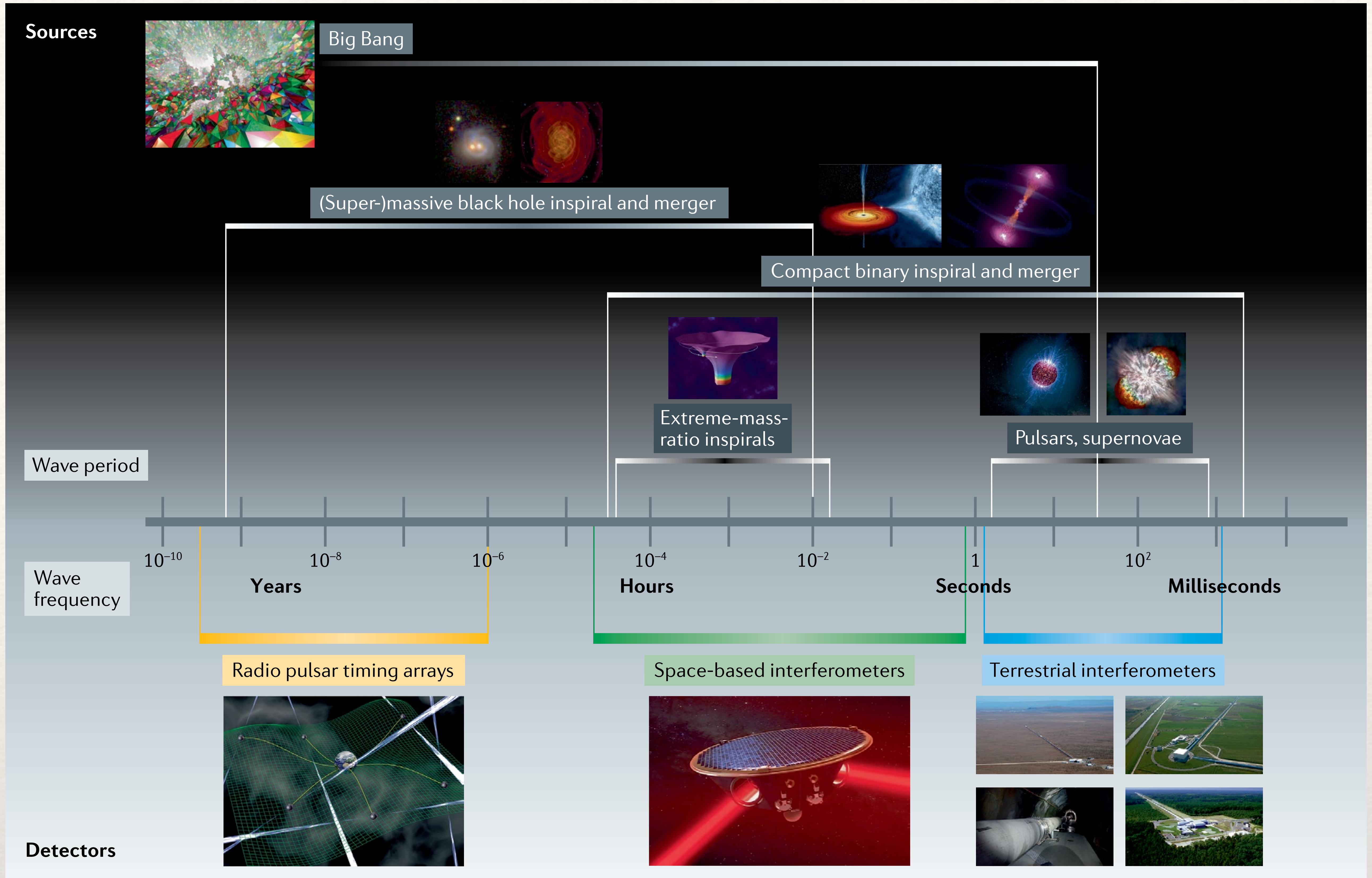
January 11, 2023

Motivation

- Gravitational Waves a new window on the Universe
- Two-body dynamics and radiative effects to exploring the most extreme conditions of spacetime and matter
- Next generation detectors, ground-based and in space, need of accurate waveform templates
- Precision Physics vs Precision Calculations: Multi-Loop Calculus, Scattering Amplitudes and General Relativity



[Bailes et al. 2021]

Motivation

- Gravitational Waves a new window on the Universe
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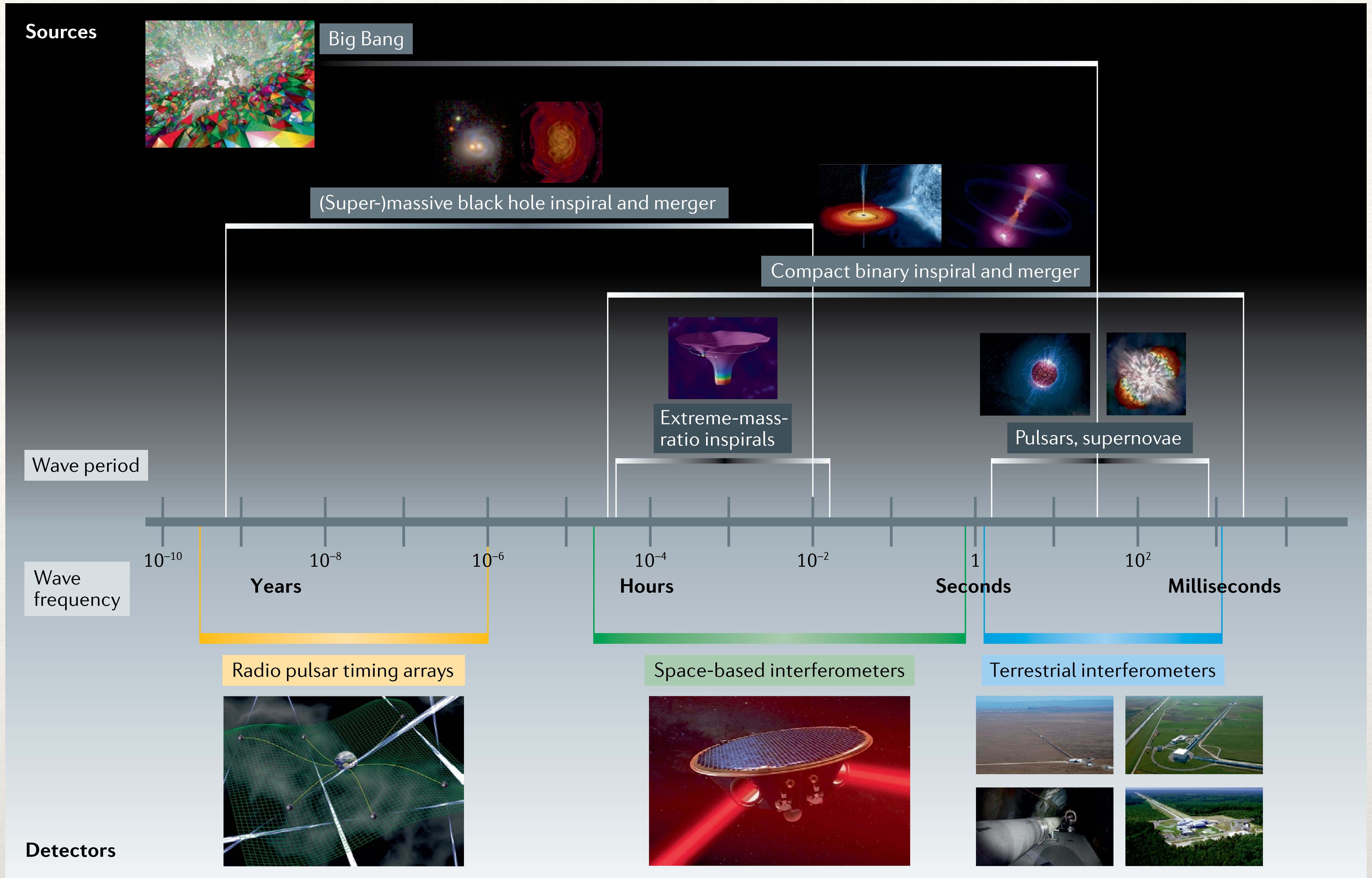


[Bailes et al. 2021]

Outline

1. Gravitational Waves Detection and Computational Techniques
2. Two-body problem in Classical GR and EFT Diagrammatic Approach
3. Conservative Effects from Near and Far Zone
4. Spin Effects

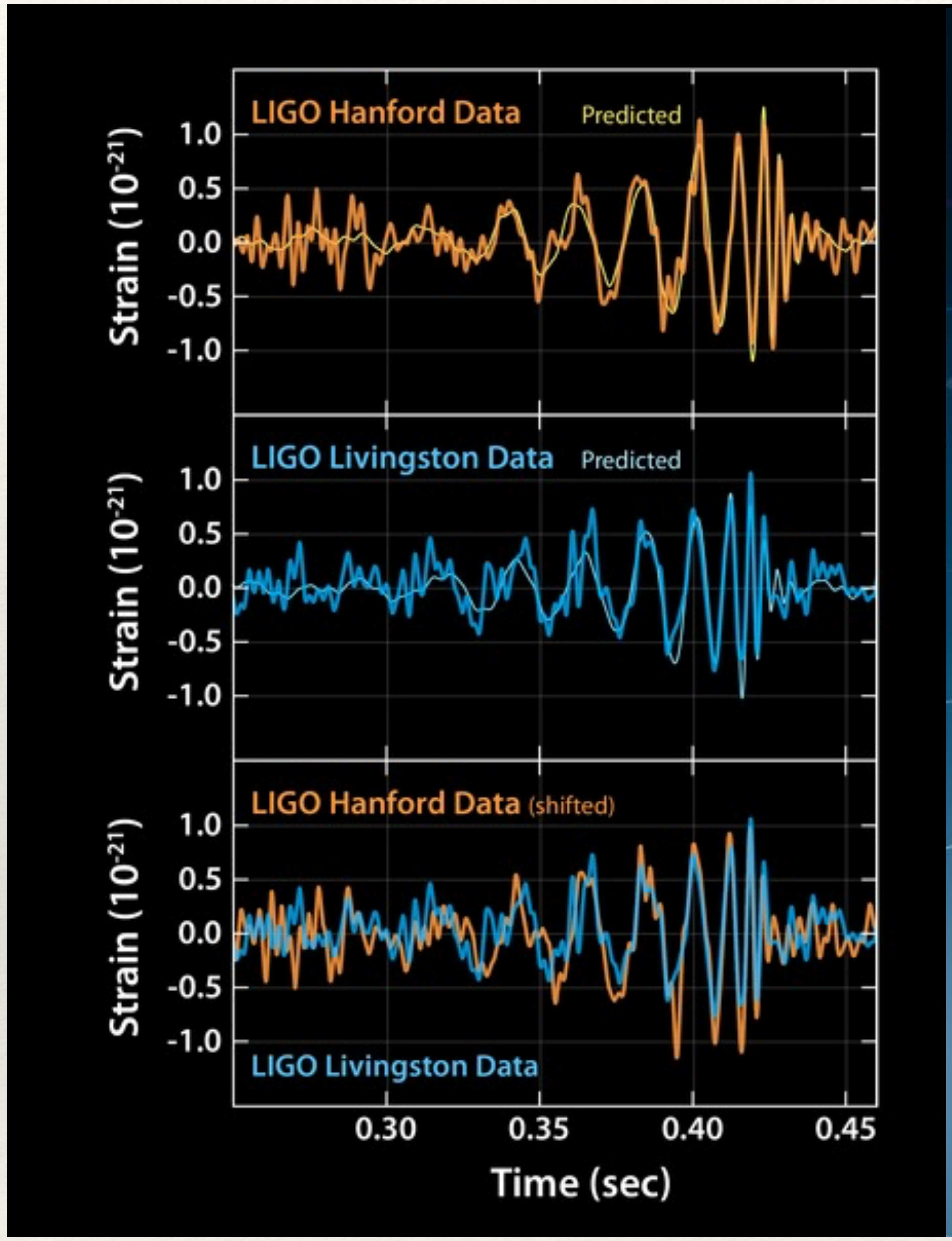
Based on collaborations with:

G. **Brunello**, J. Steinhoff, M.K. **Mandal**, R. **Patil**

D. Bini, T. Damour, A. **Geralico**, S. Laporta

S. Foffa, R. Sturani, C. Sturm, W.J. **Torres Bobadilla**

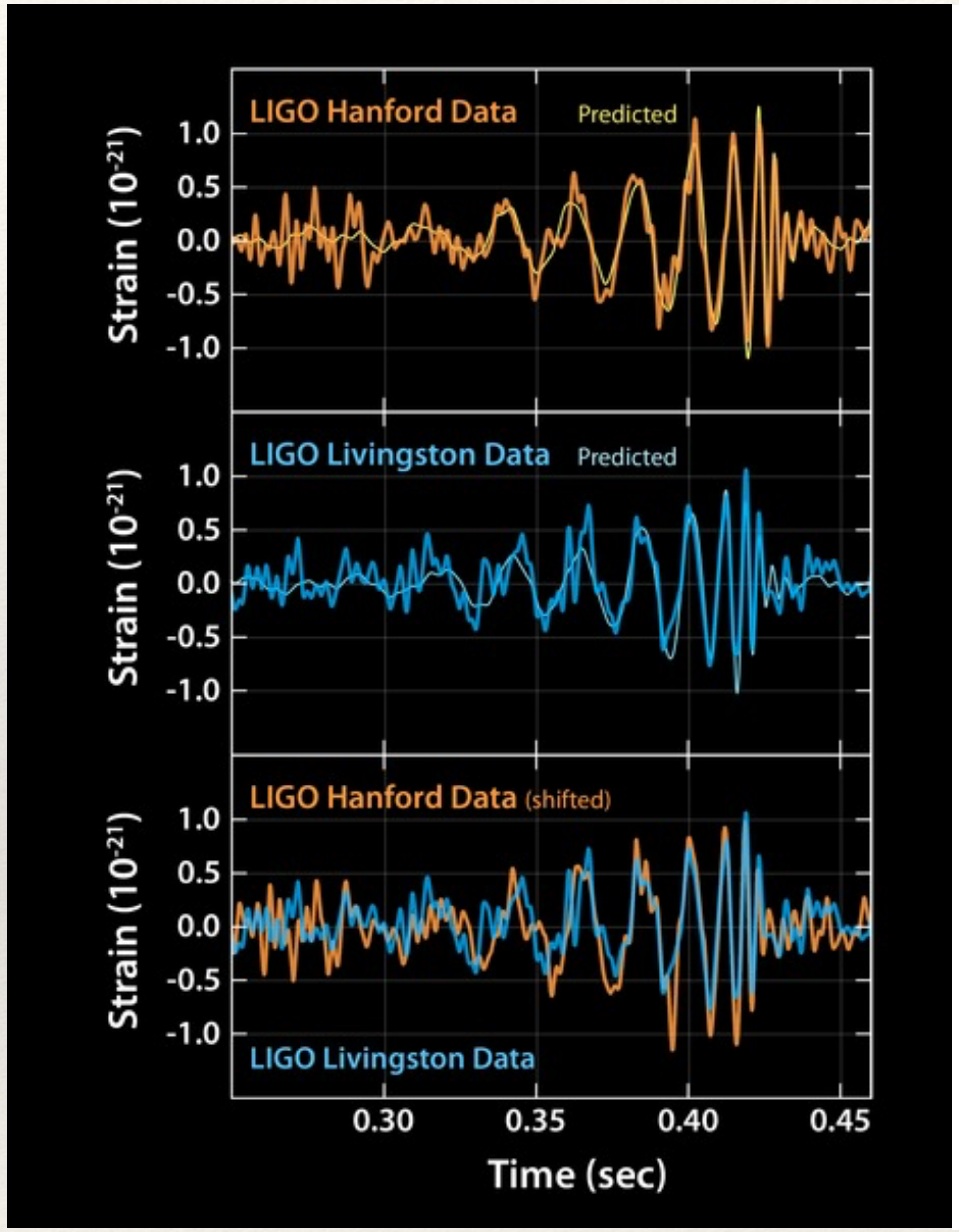
GW Detection

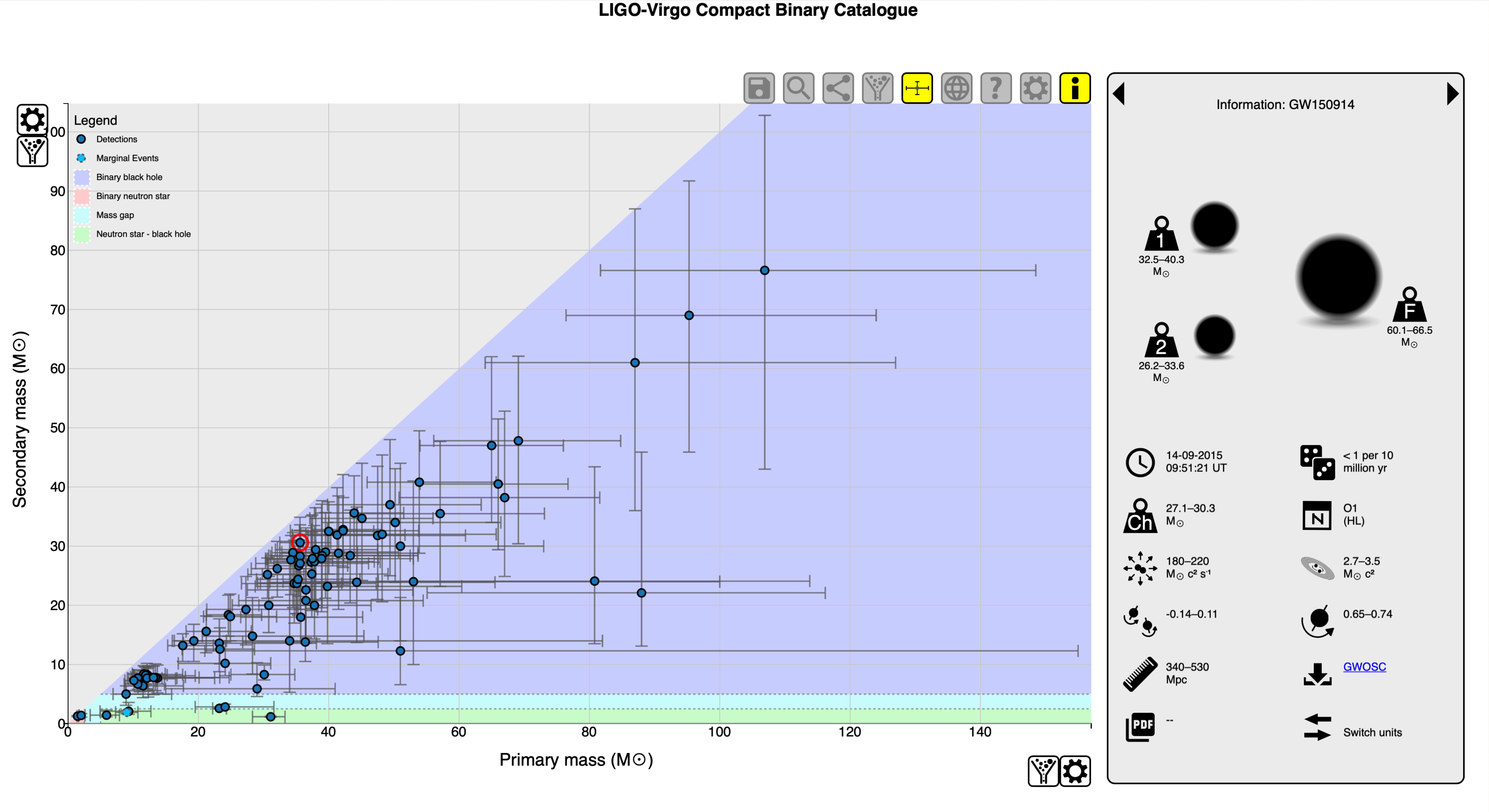


LIGO-Virgo Detection: **GW150914**

Zurich is here

GW Detection



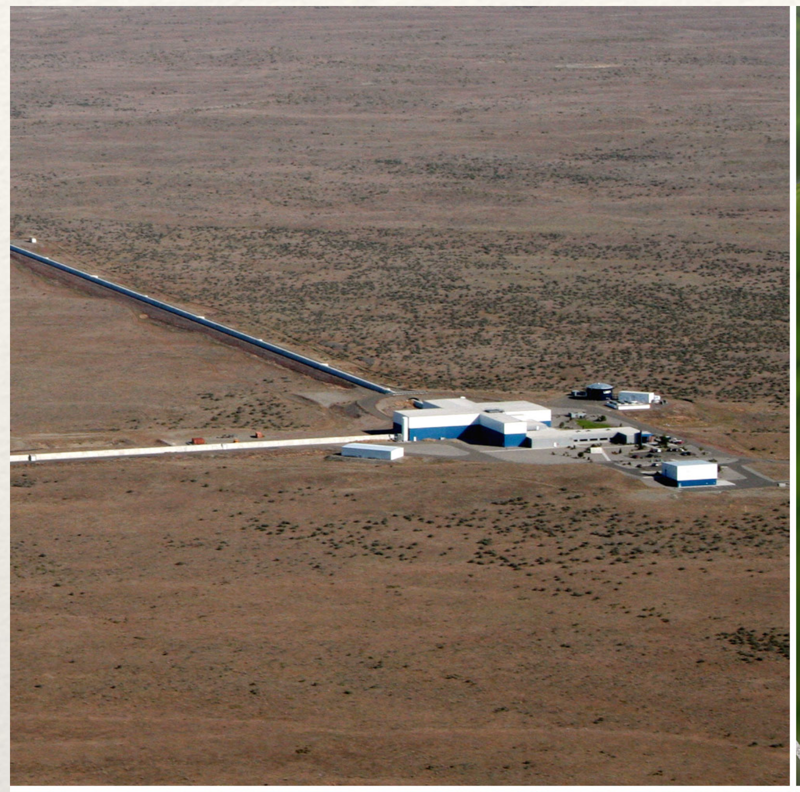


LIGO-Virgo Detection: **GW150914**

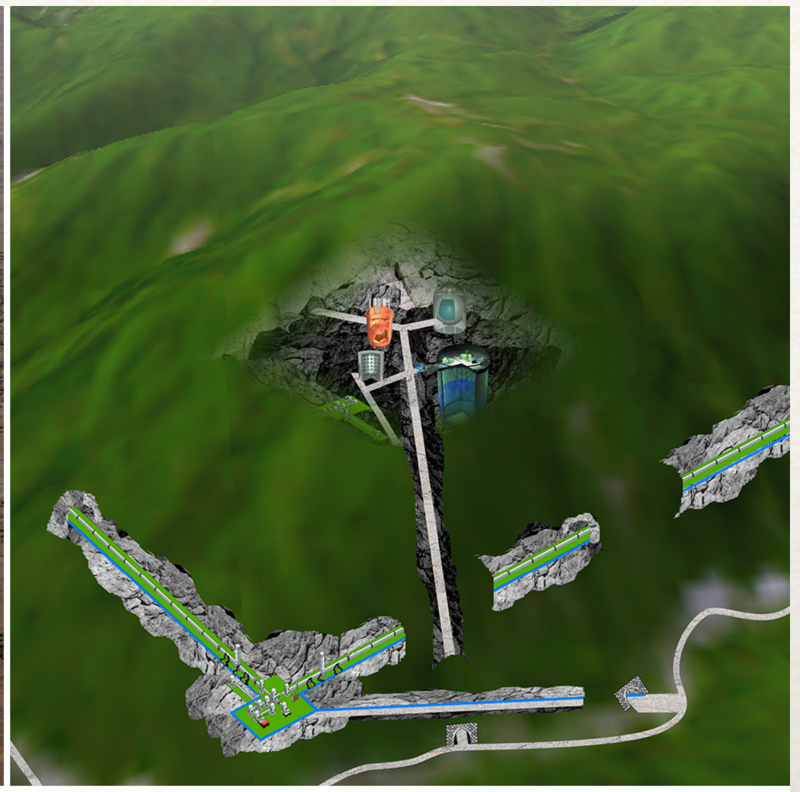
GW Detection

LIGO-Virgo-KAGRA Collaboration

Hanford



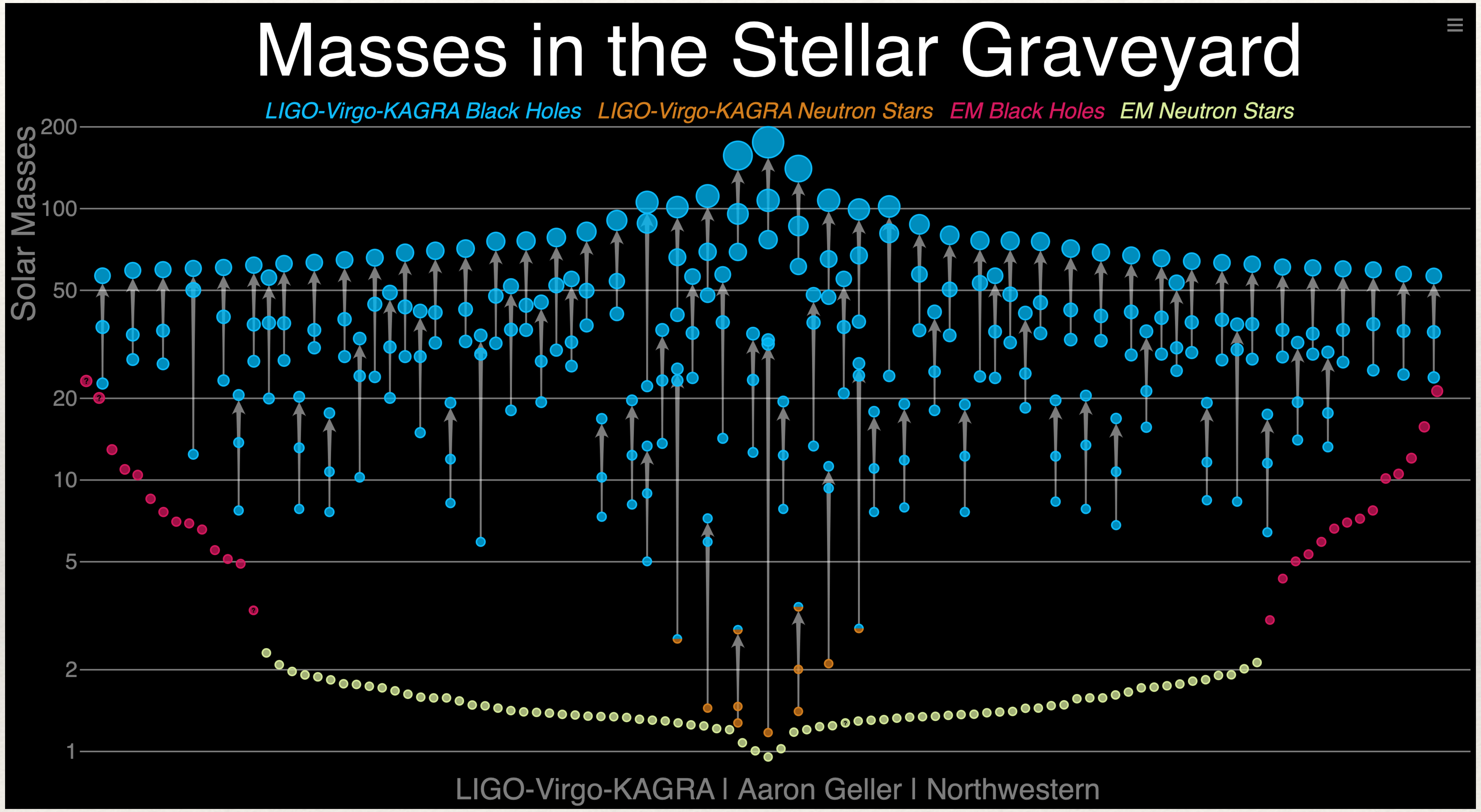
Kagra



Livingston

Virgo

O3b - Catalogue



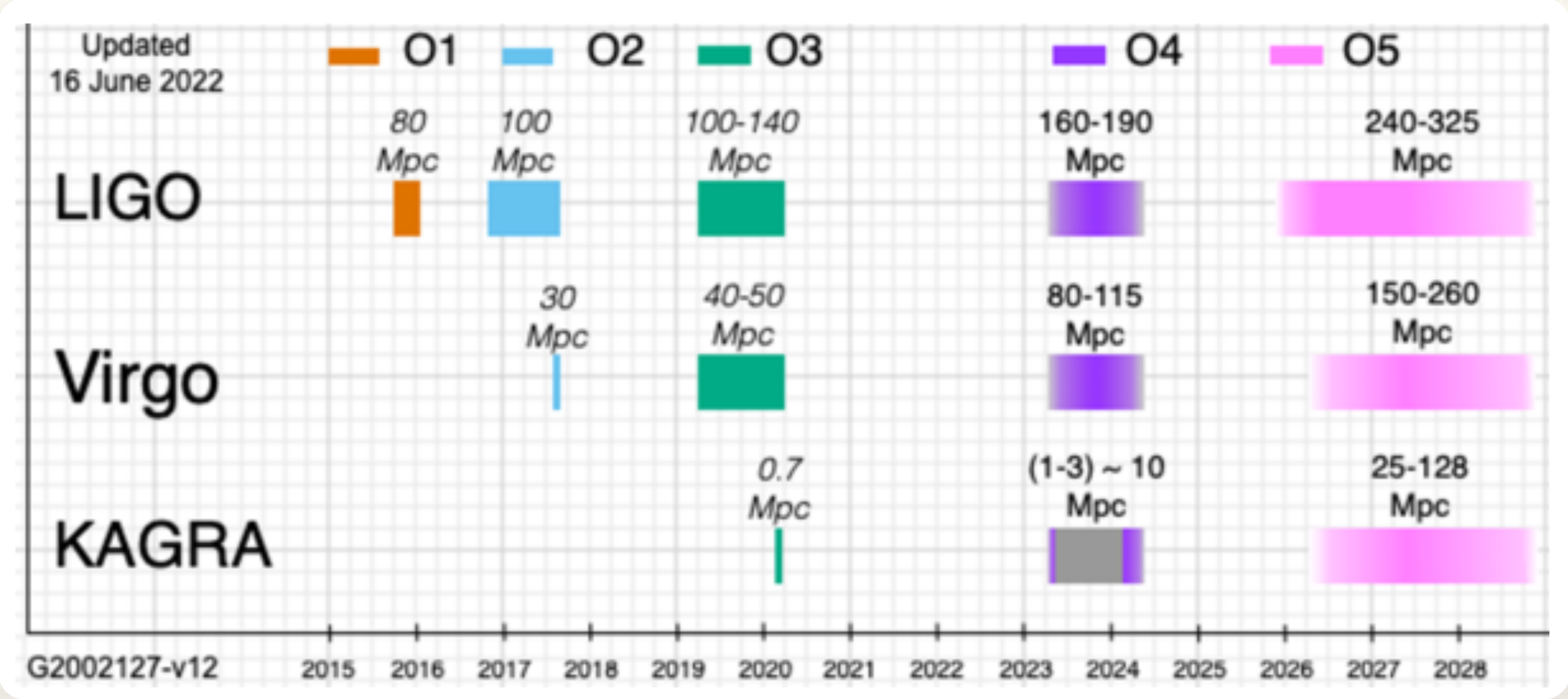
GW Detection

:: Current GW Detectors: advanced programs

Prospects for observations within advanced Programs

updated [Abbot et al. 2020]

O3 --> O5 <==> *O*(10) --> *O*(100) GW detections/year

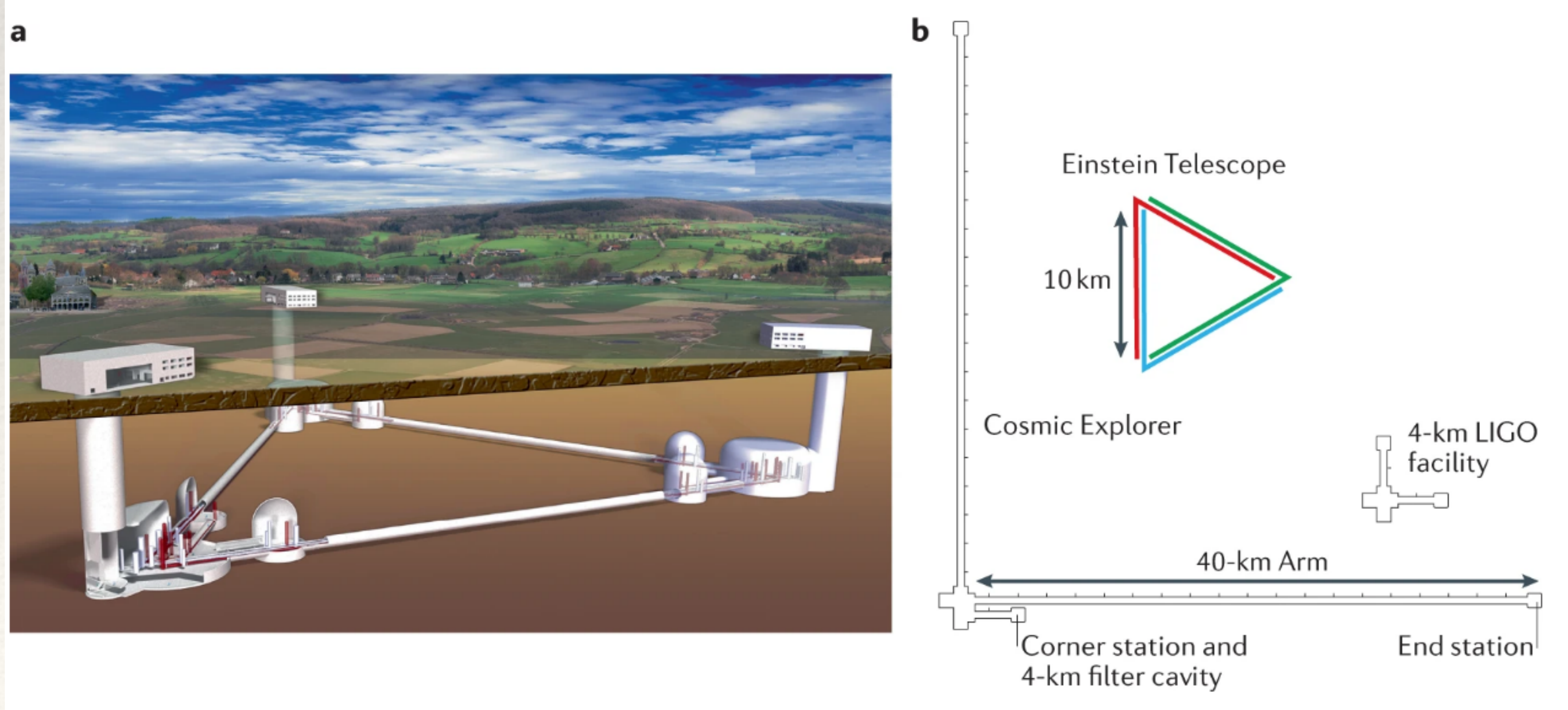


[credit: Ligo News]

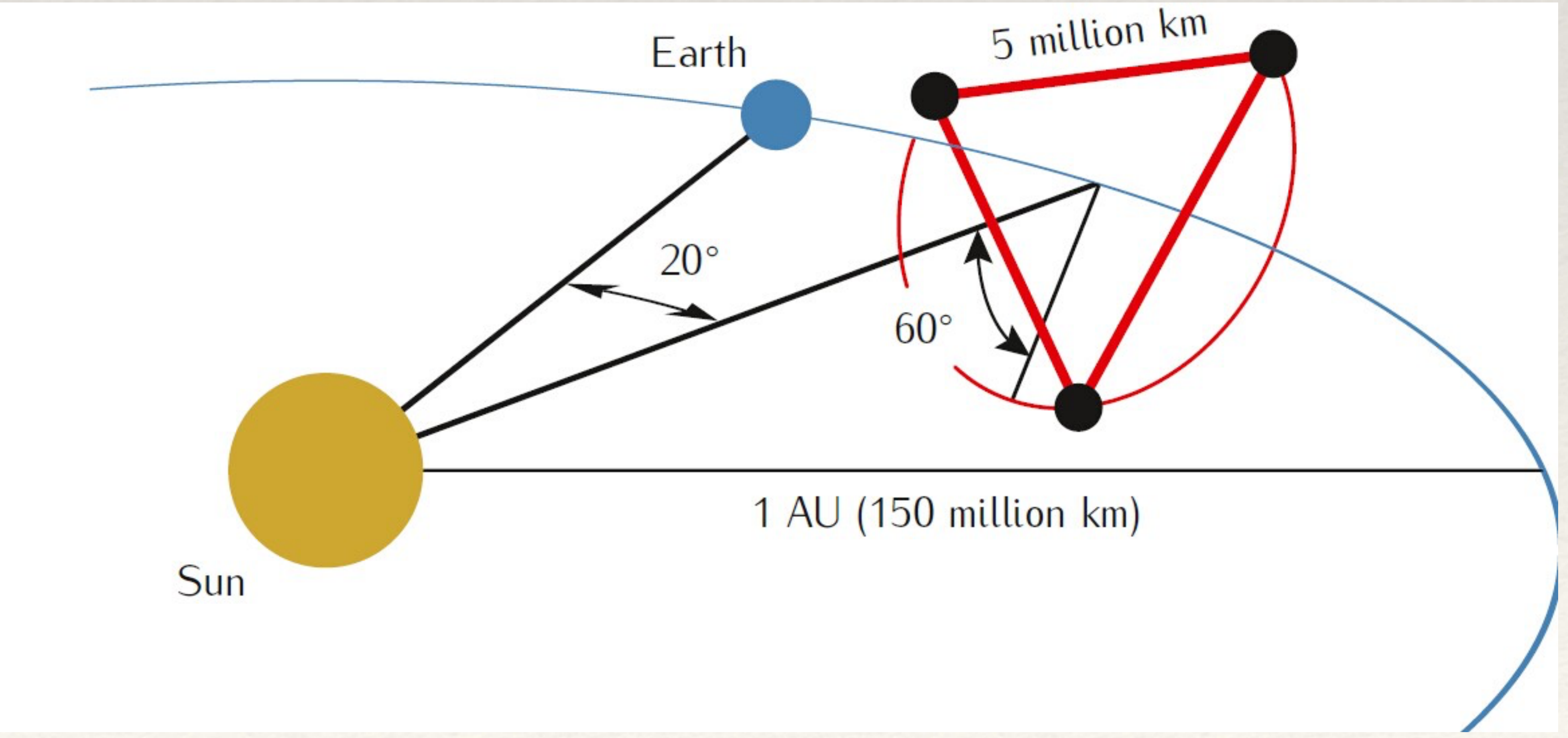
:: (some) Future GW Detectors

[Bailes et al. 2021]

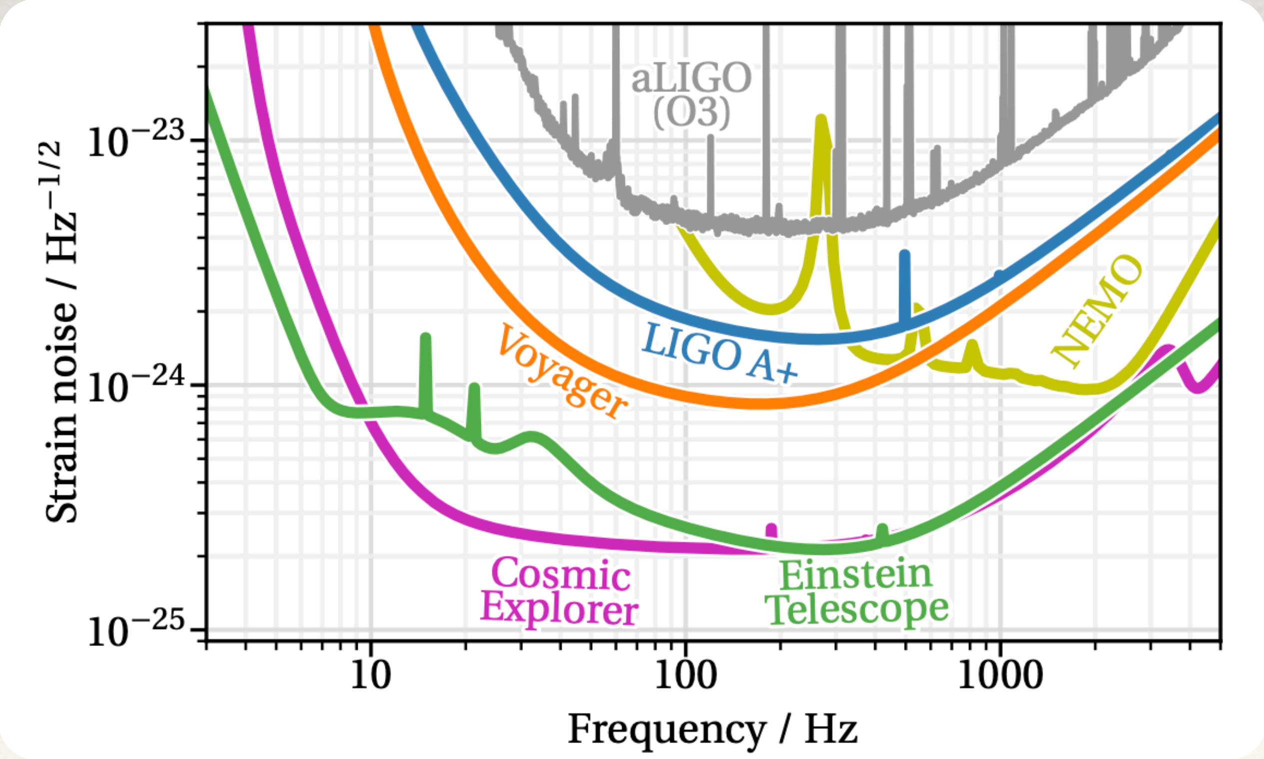
Einstein Telescope



Lisa Mission



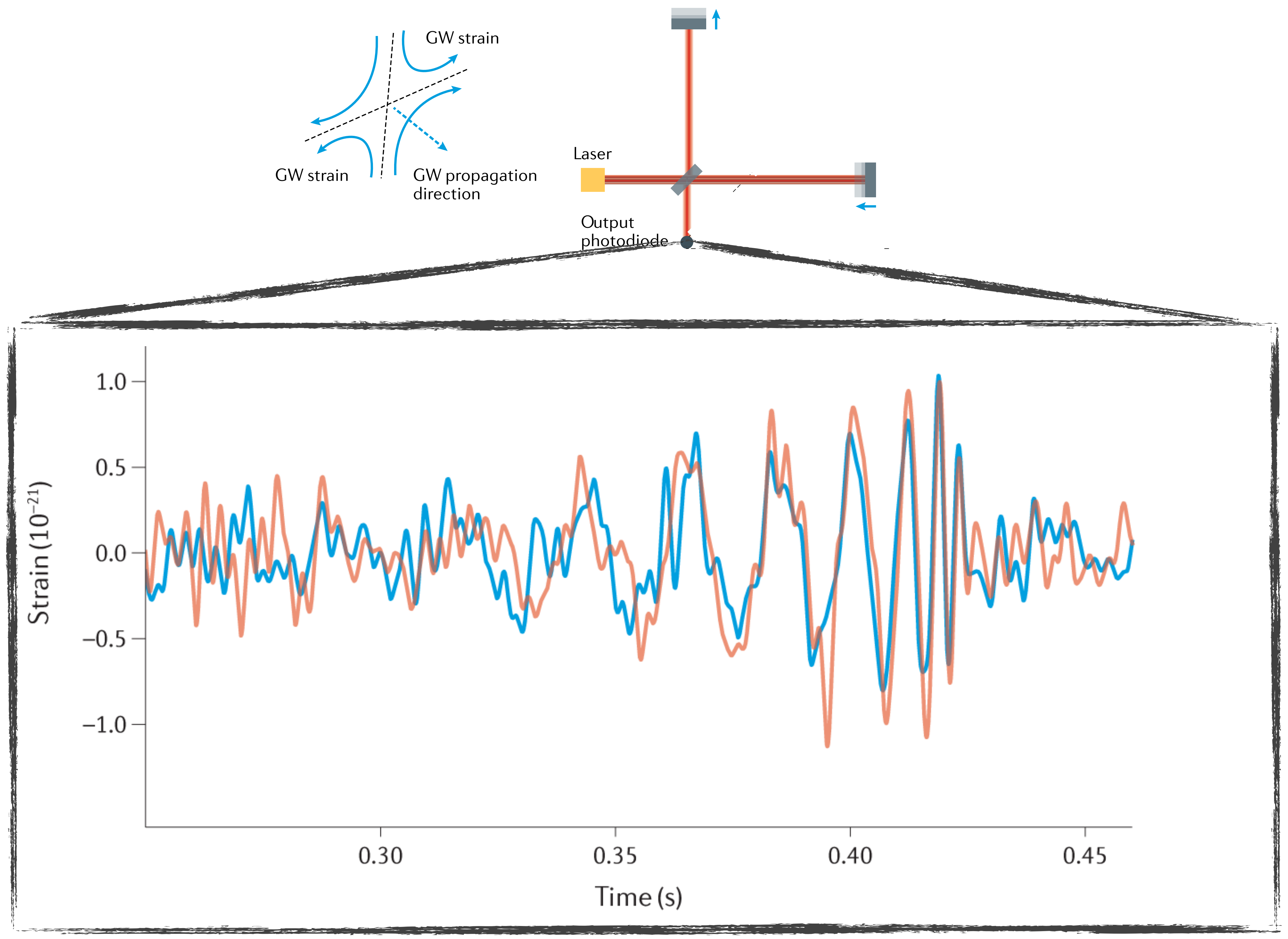
expected sensitivity



[credit: CE Consortium]

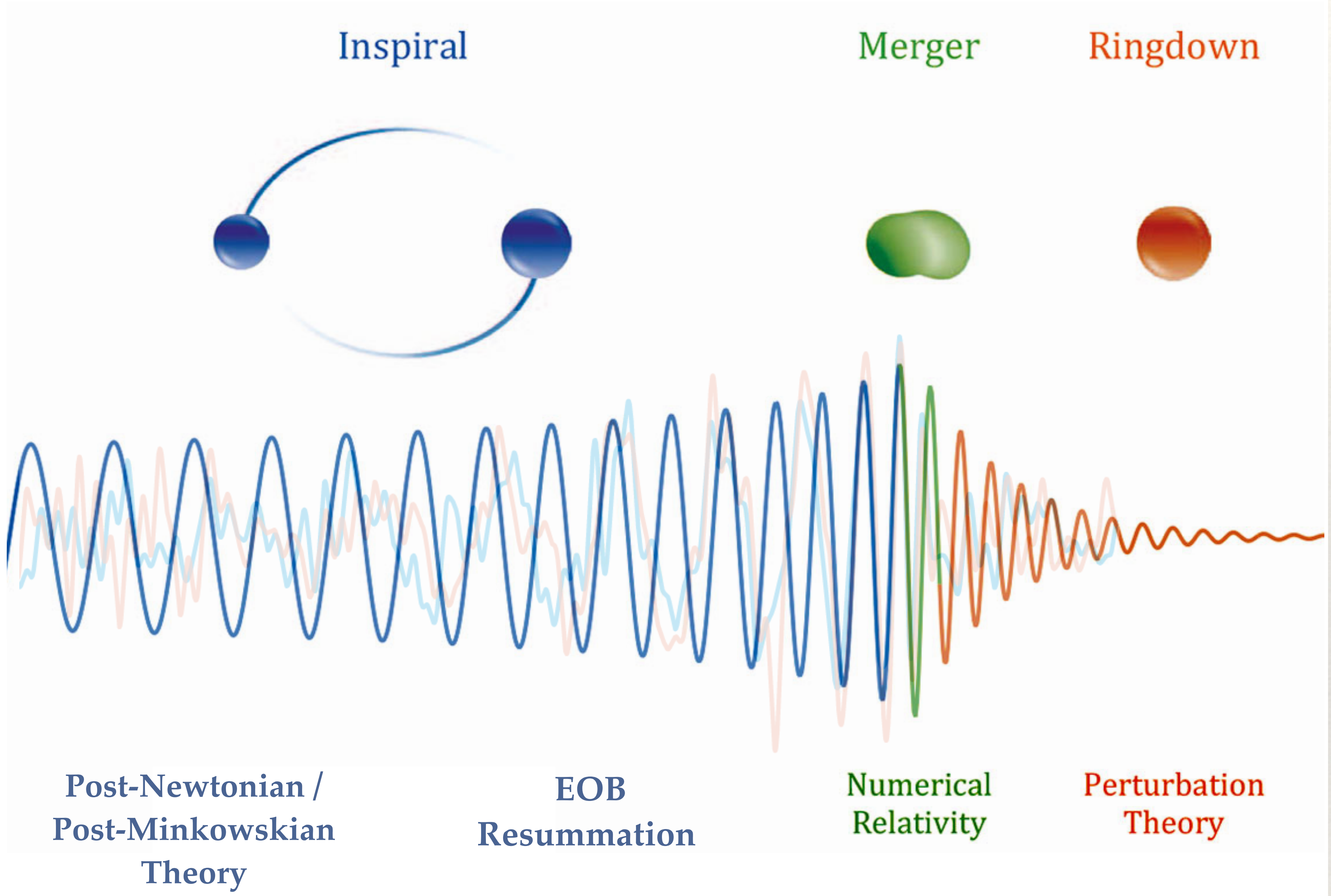
Two-Body Dynamics and GW Signal

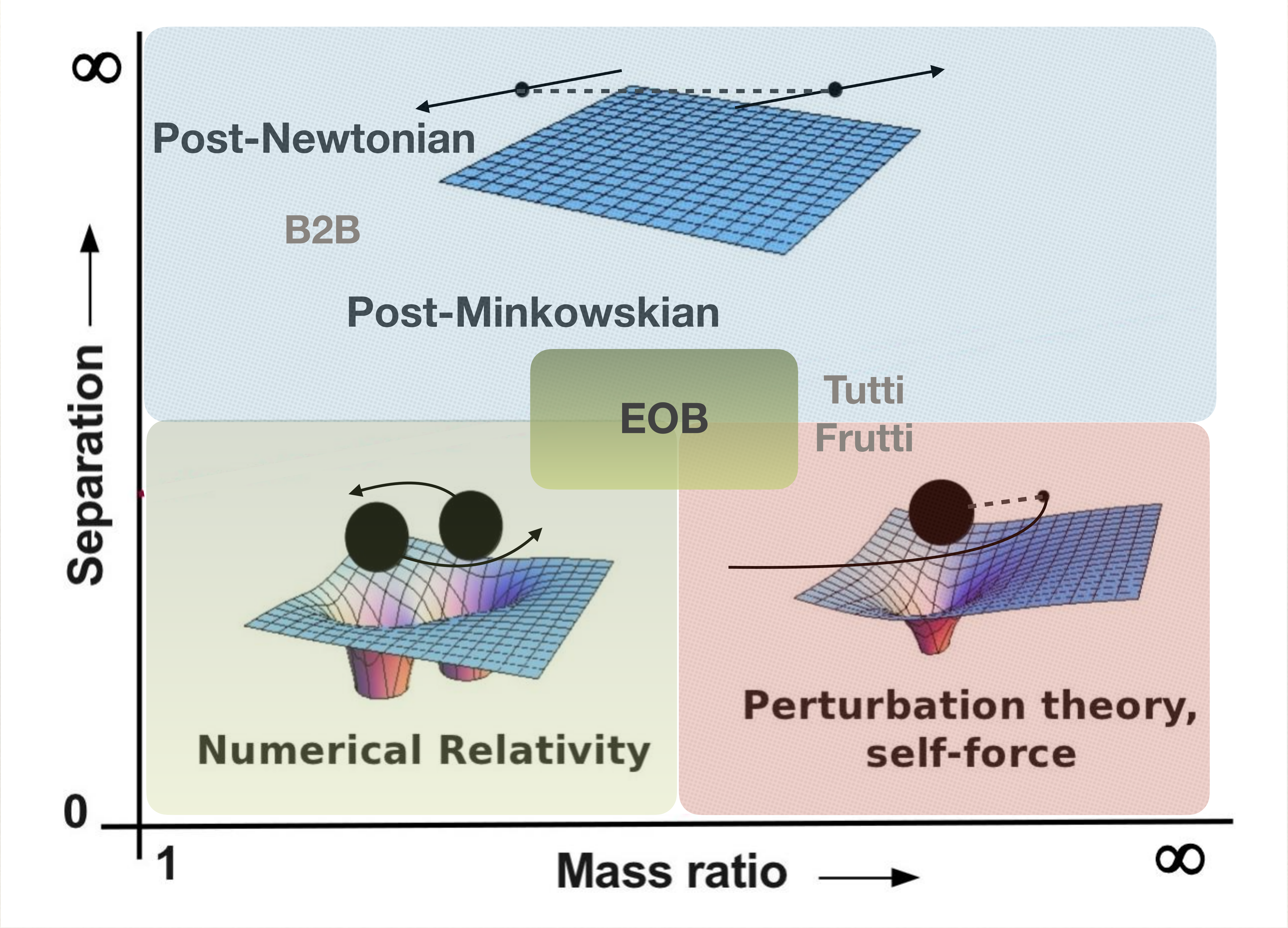
▸Real Event



Two-Body Dynamics and GW Signal

▸ **Waveform Model and Computing Techniques**

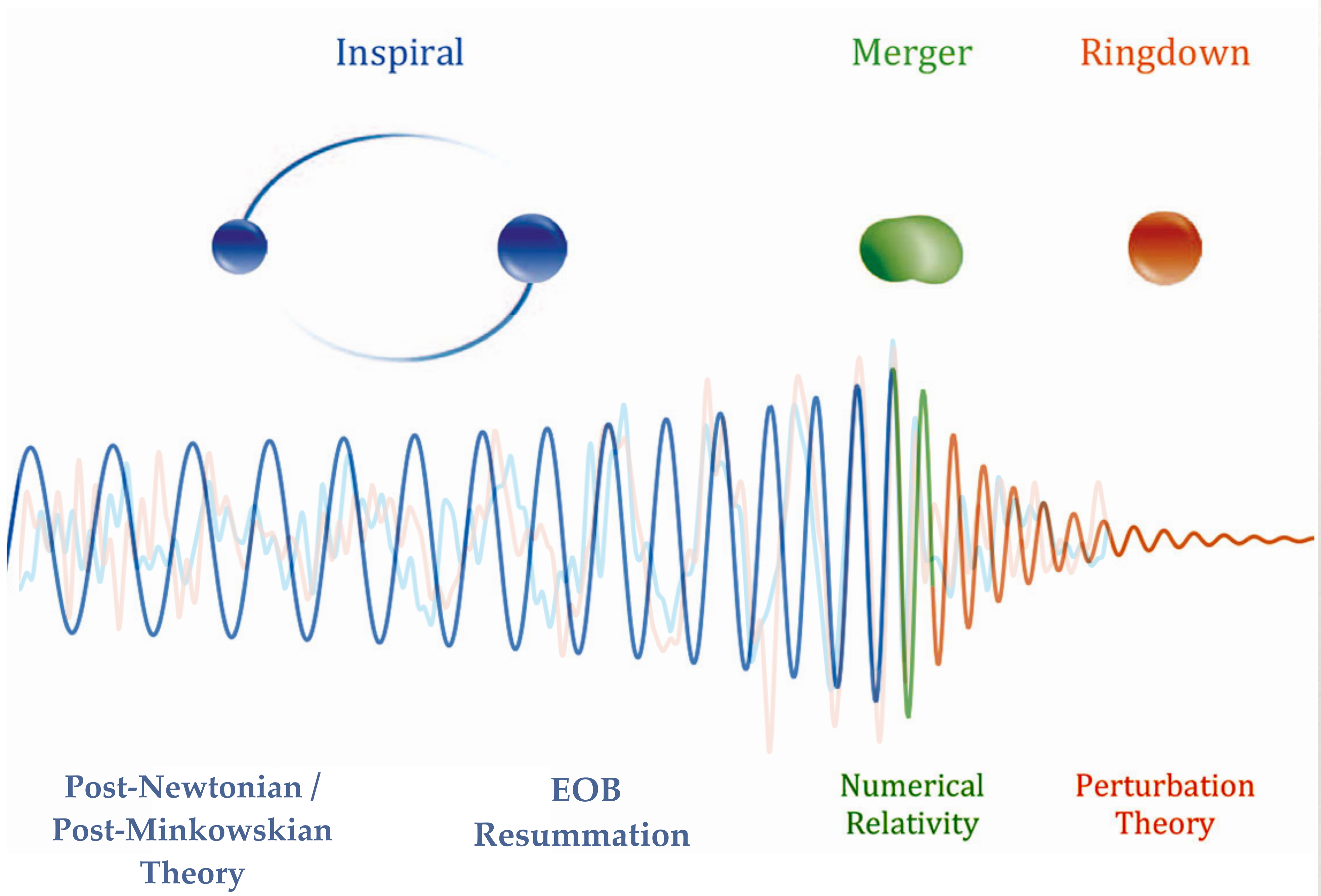




[adapted from: Barak]

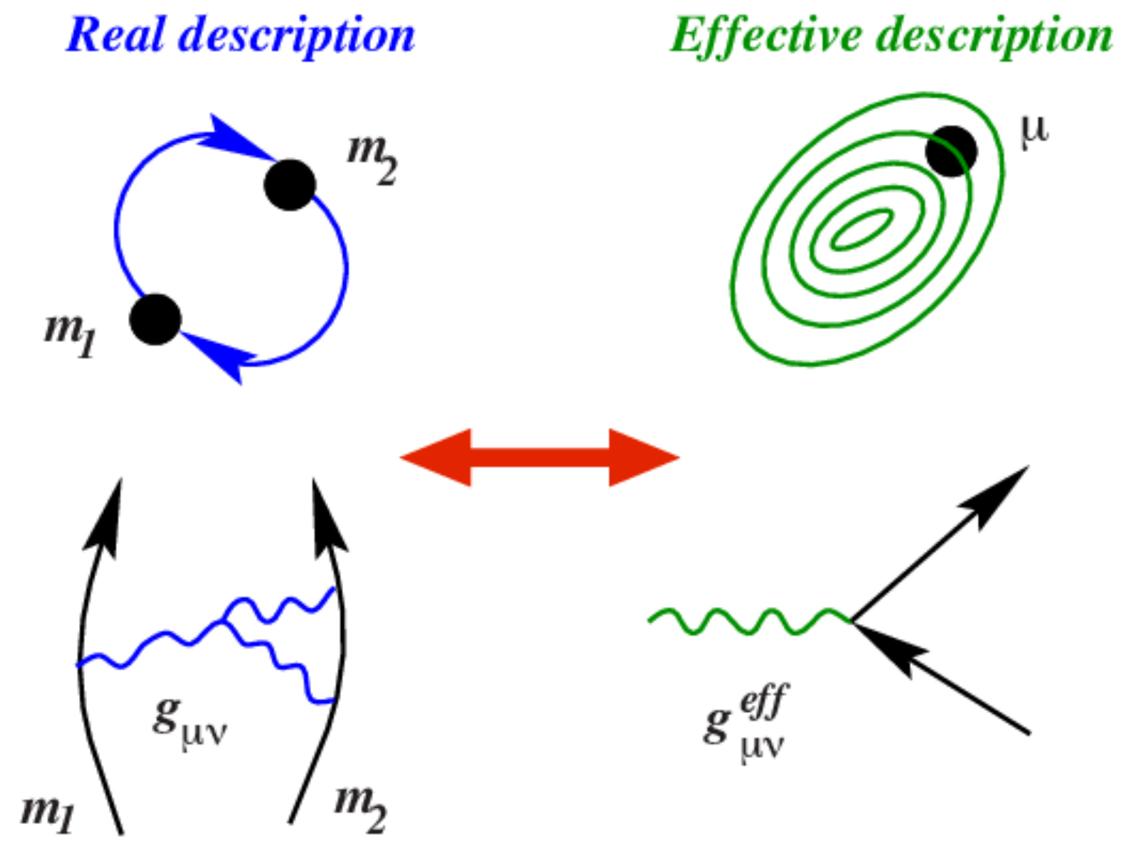
Two-Body Dynamics and GW Signal

Waveform Model and Computing Techniques



Effective One Body (EOB) Formalism [Buonanno Damour]

the contributions coming from different kinematic regions for combined and calibrated with Numerical Relativity



- Post-Minkowskian Expansion [relativistic scattering]

  $G_N \frac{m}{r} << v^2 \sim 1$

  Expansion in powers of G_N

- Post-Newtonian Expansion [non relativistic system]

  $G_N \frac{m}{r} \sim v^2 << 1$

  Expansion in powers of v/c

- BH perturbation theory / self force

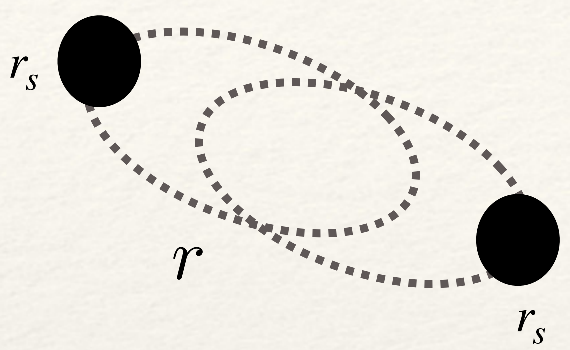
  Expansion for small metric deformation

  $\delta g_{\mu\nu} \sim \epsilon = m_2/m_1 << 1$

Effective Field Theory for General Relativity

Coalescing Binary System

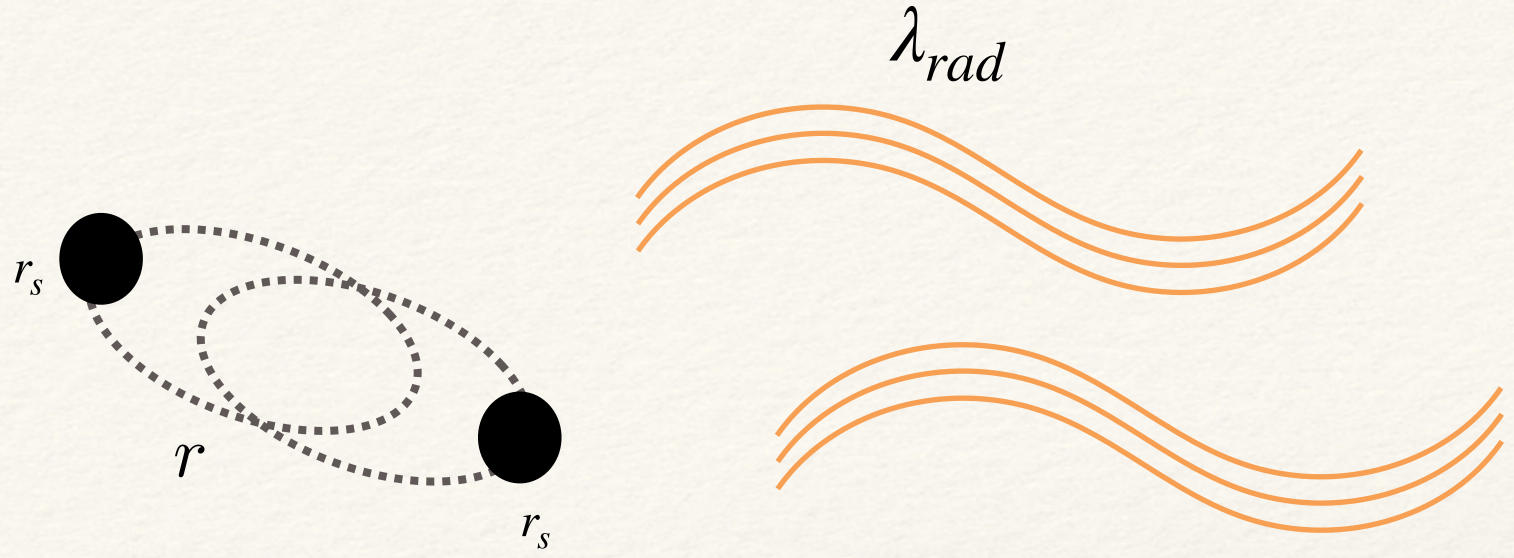
:: Double Hierarchy



Conservative system ::
~~GW emission~~

$$\boxed{r_s \ll r} \ll \lambda_{rad}$$

$$r_s \ll \boxed{r \ll \lambda_{rad}}$$

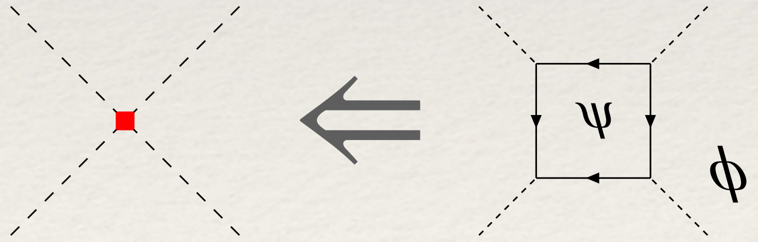


Dissipative system ::
GW emission

:: Effective Field Theory Approach

- **Fundamental [*complete*] theory** $S[\phi, \psi]$
  - Heavy fields ψ: Λ,
    short distance r_s
  - Light modes ϕ: $\omega \ll \Lambda$,
    large distance r

$$e^{\frac{iS_{eff}[\phi]}{\hbar}} = \int D\psi \, e^{\frac{iS[\phi,\psi]}{\hbar}}$$



- **Effective [*incomplete*] theory** $S_{eff}[\psi]$
  - Sensitive to the Lower-scale dynamics: $\omega \ll \Lambda$

GREFT / Action

$$S_{tot}[x_a, g] = S_{GR}[g] + S_{m_a}[x_a, g]$$

[Goldberger, Rothstein]

- **Einstein Hilbert + gauge fixing**

$$S_{GR}[g] = 2\Lambda^2 \int d^{d+1}x \, \sqrt{-g} \left(R - \frac{1}{2}\Gamma^\mu \Gamma_\mu \right)$$

$$\Lambda^{-1} = \sqrt{32\pi G_N}$$

- **Source/Worldline**

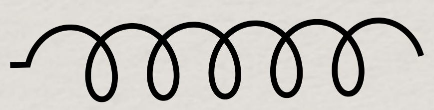
$$S_{m_a}[x_a, g] = S_{pp}[x_a, g] + \delta S_{m_a}[x_a, g]$$

$$= -m_a \int d\tau_a = -m_a \int dt \sqrt{-g_{\mu\nu}(x_a)\dot{x}_a^\mu \dot{x}_a^\nu}$$

▸ Non-relativistic approximation [method of regions]: [Beneke Smirnov]

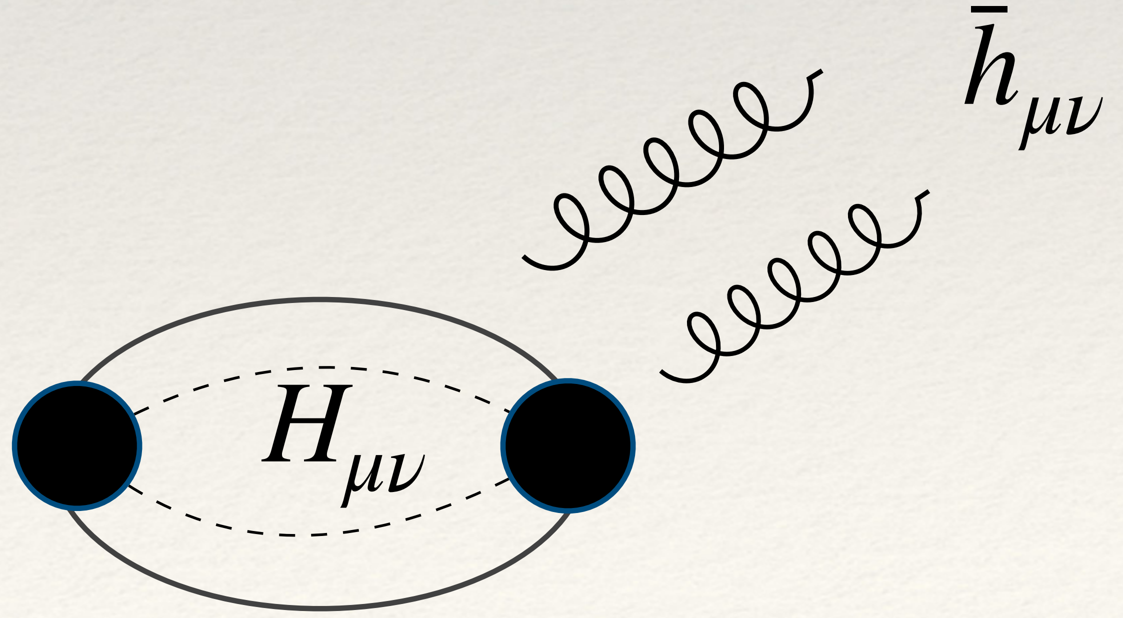
▸ Weak field expansion:

$$v \ll 1 \qquad\qquad g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} \qquad\qquad h_{\mu\nu} = H_{\mu\nu} + \bar{h}_{\mu\nu}$$

- **Potential gravitons** $H_{\mu\nu}$: $\quad (k_0, \mathbf{k}) \sim \left(\frac{v}{r}, \frac{1}{r} \right)$ - - - - - - - -
- **Radiation gravitons** $\bar{h}_{\mu\nu}$: $\quad (k_0, \mathbf{k}) \sim \left(\frac{v}{r}, \frac{v}{r} \right)$ 
- **Worldline/BH** x_a : ———————

▸ **Effective action** by **integrating out gravitons:**

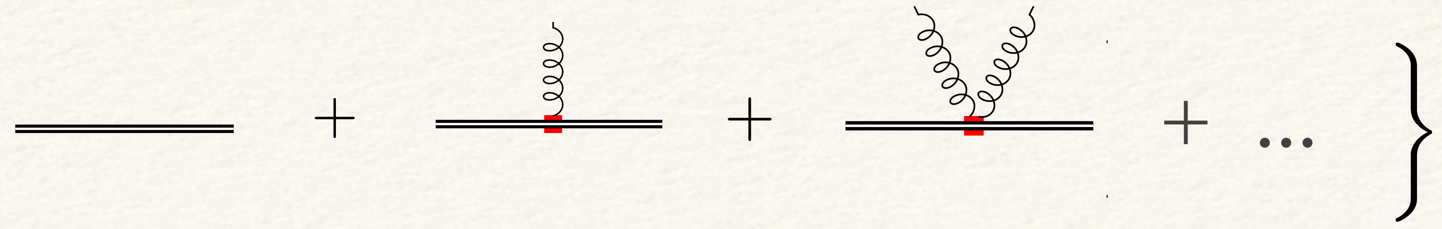
$$e^{iS_{eff}[x_a]} = \int D\bar{h} \int DH \, e^{iS_{tot}[x_a, H, \bar{h}]}$$

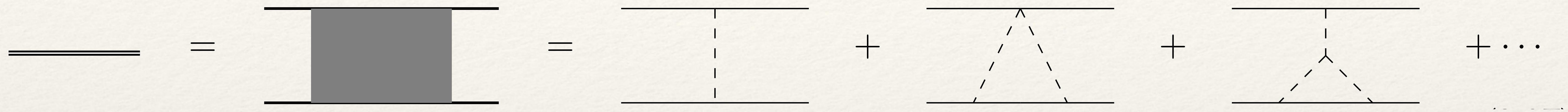


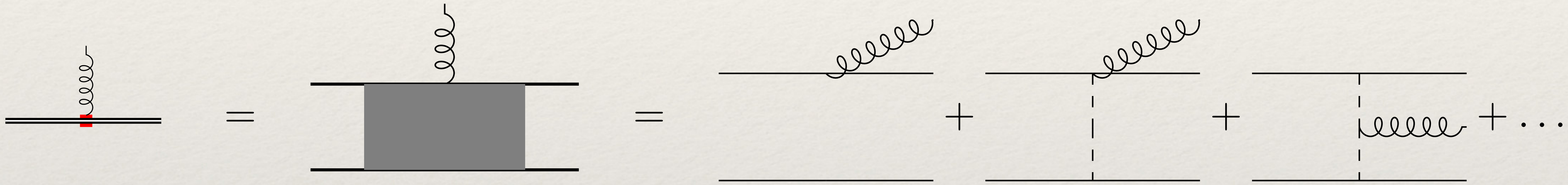
GREFT / Action / Near & Far Zone

[Goldberger, Rothstein]

$$e^{iS_{eff}[x_a]} = \int D\bar{h} \int DH\ e^{iS_{tot}[x_a,H,\bar{h}]} = \int D\bar{h}\ e^{\left\{ iS_{bulk}[\bar{h}] + \cdots \right\}}$$

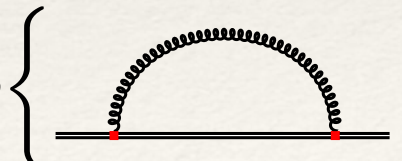






▸Near zone (r): $S_{pot}[x_a, g] = S_{GR}[g] + S_{m_a}[x_a, g]$ $\quad g_{\mu\nu} = \eta_{\mu\nu} + H_{\mu\nu}$ $\quad \int DH\ e^{iS_{tot}[x_a,H,\bar{h}=0]} = exp\{\ \}$

▸Far zone (λ_{rad}): $S_{rad}[g, \{Q_i\}] = S_{GR}[g] + S_{mult}[g, \{Q_i\}]$ $\quad g_{\mu\nu} = \eta_{\mu\nu} + \bar{h}_{\mu\nu}$ $\quad \int D\bar{h}\ e^{iS_{rad}[x_a,\bar{h}]} = exp\{\ + \ \ldots\}$



Conservative Dynamics :: **Near Zone** Spinless

Near Zone/EFT Diagrammatic Approach

[Goldberger, Rothstein]
[Gilmore, Ross]
[Foffa, Sturani]

$$S_{pot}[x_a, g] = S_{GR}[g] + S_{m_a}[x_a, g] \qquad\qquad S_{m_a}[x_a, g] = S_{pp}[x_a, g] + \cancel{\delta S_{m_a}[x_a, g]}$$

- **Kaluza-Klein parametrization:**

[Kol Smolkin] [Blanchet Damour]

$$g_{\mu\nu} = e^{2\phi/\Lambda}\begin{pmatrix} -1 & A_j/\Lambda \\ A_i/\Lambda & e^{-c_d\frac{\phi}{\Lambda}}\gamma_{ij} - A_iA_j/\Lambda^2 \end{pmatrix}$$

$$\gamma_{ij} = \delta_{ij} + \frac{\sigma_{ij}}{\Lambda} \qquad\qquad c_d = 2\frac{d-1}{d-2}$$

Graviton = Scalar + Vector + Sym. Tensor

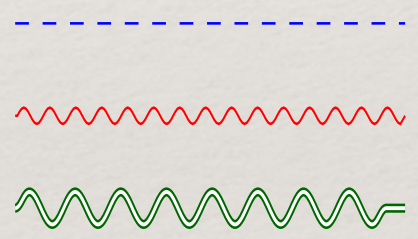
$10 \qquad 1 + 3 + 6$

$$g_{\mu\nu} \Rightarrow \quad \phi \quad A^i \quad \sigma^{ij}$$

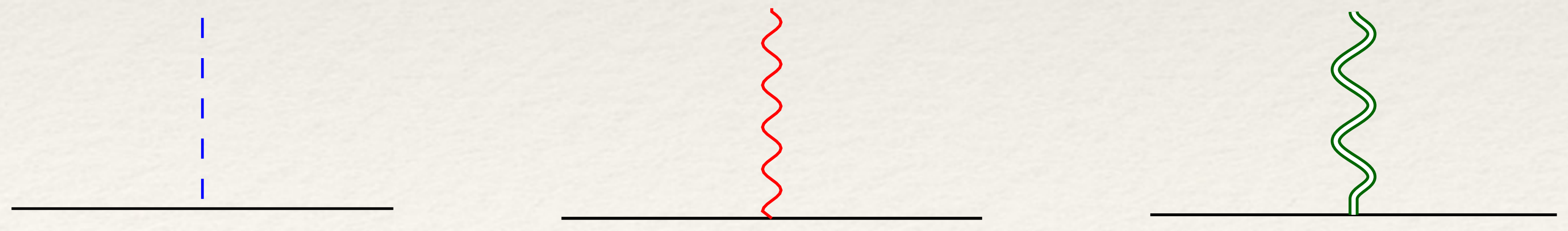
- **Feynman rules for:** $\phi \quad A^i \quad \sigma^{ij} \quad x_a$

Static / non-propagating source: x_a

Propagators: ϕ A^i σ^{ij}

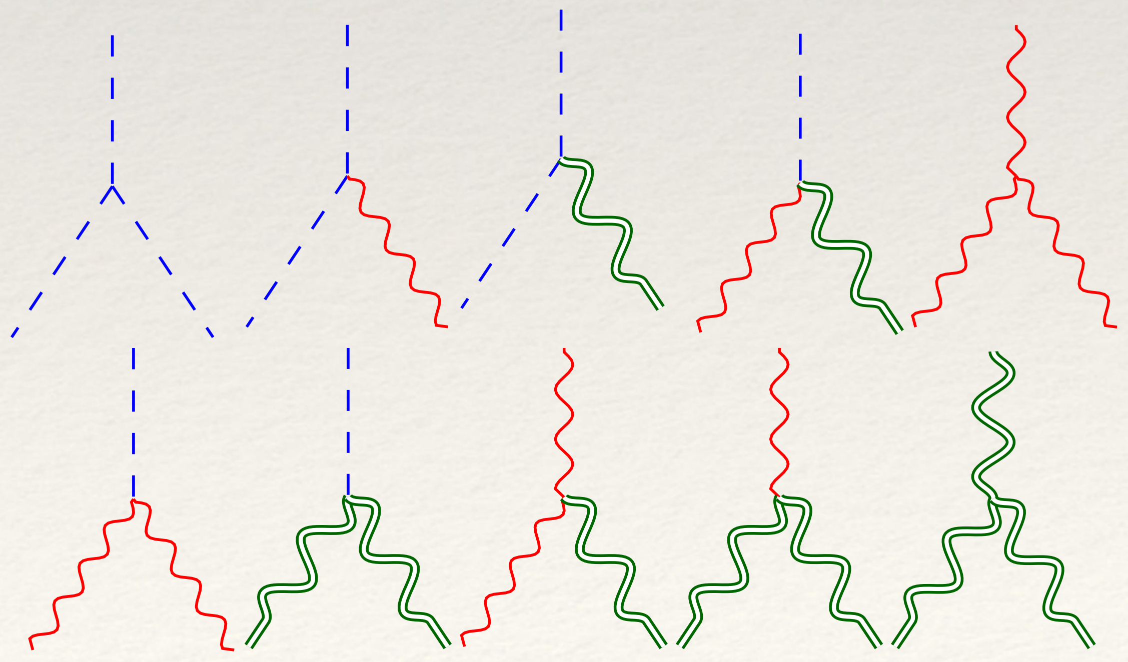


Source couplings:



$+ \ldots$

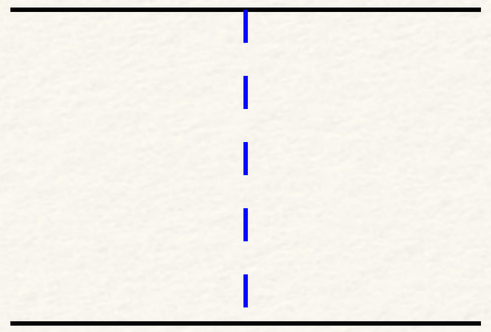
Self-interactions:

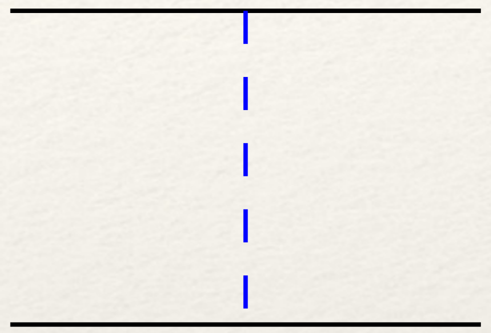


$+ \ldots$

Newton Potential

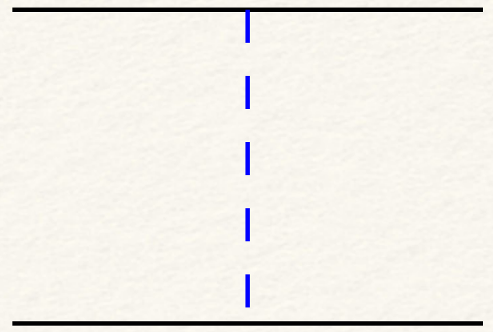
Diagrammatic approach

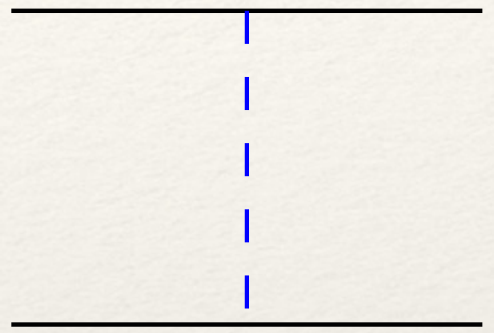
- Just 1 diagram: $\mathscr{M}_{0PN} =$  $= \dfrac{i m_1 m_2}{2 c_d \Lambda^2} \dfrac{1}{\mathbf{p}^2}$

- Fourier transform: from amplitude to the effective action: $\mathscr{L}_{0PN} = -i \lim\limits_{d \to 3} \int \dfrac{d^d \mathbf{p}}{(2\pi)^d} e^{i\mathbf{p}(x_1 - x_2)} \Bigg($  $\Bigg) = \dfrac{G_N m_1 m_2}{r}$

Newton Potential

Diagrammatic approach

- Just 1 diagram: $\mathscr{M}_{0PN} =$  $= \dfrac{i m_1 m_2}{2 c_d \Lambda^2} \dfrac{1}{\mathbf{p}^2}$

- Fourier transform: from amplitude to the effective action: $\mathscr{L}_{0PN} = -i \lim_{d \to 3} \int \dfrac{d^d \mathbf{p}}{(2\pi)^d} e^{i\mathbf{p}(x_1 - x_2)} \Bigg($  $\Bigg) = \dfrac{G_N m_1 m_2}{r}$

Corrections to the Newtonian potential:

- Non-relativistic velocities: $v^2 \ll 1$

- Dynamics in Post-Minkowskian perturbative scheme

- At nPM order: G_N^n

Astrophysicists/Cosmologists' whishlist

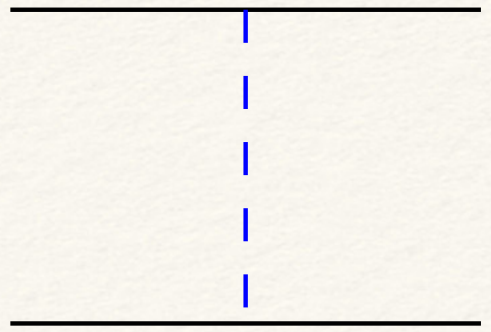
			0PN	1PN	2PN	3PN	4PN	5PN	6PN		
	G	(	1	$+\ v^2$	$+\ v^4$	$+\ v^6$	$+\ v^8$	$+\ v^{10}$	$+\ v^{12}$	$+\ \dots$)	1PM
1979-81 →	G^2	(	1	$+\ v^2$	$+\ v^4$	$+\ v^6$	$+\ v^8$	$+\ v^{10}$	$+\ v^{12}$	$+\ \dots$)	2PM
2019 →	G^3	(	1	$+\ v^2$	$+\ v^4$	$+\ v^6$	$+\ v^8$	$+\ v^{10}$	$+\ v^{12}$	$+\ \dots$)	3PM
2021 →	G^4	(	1	$+\ v^2$	$+\ v^4$	$+\ v^6$	$+\ v^8$	$+\ v^{10}$	$+\ v^{12}$	$+\ \dots$)	4PM
	G^5	(	1	$+\ v^2$	$+\ v^4$	$+\ v^6$	$+\ v^8$	$+\ v^{10}$	$+\ v^{12}$	$+\ \dots$)	5PM
	G^6	(	1	$+\ v^2$	$+\ v^4$	$+\ v^6$	$+\ v^8$	$+\ v^{10}$	$+\ v^{12}$	$+\ \dots$)	6PM
	G^7	(	1	$+\ v^2$	$+\ v^4$	$+\ v^6$	$+\ v^8$	$+\ v^{10}$	$+\ v^{12}$	$+\ \dots$)	7PM

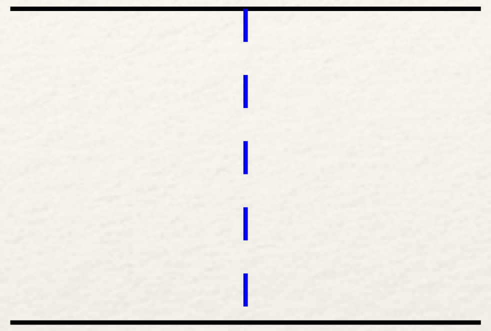
[credit: Bern et al.]

…Westphal, Damour, Cheung, Rothstein, Solon, Bern, Roiban, Shen, Zeng, Parra-Martinez, Ruf, Hermann, Buonanno, Porto, Dlapa, Kalin, Liu, Neef, Bjerrum-Bohr, Vanhove, Plante, Cristofoli, Damgaard, Guevara, Ochirov, Vines, Di Vecchia, Veneziano, Heisenberg, Russo, Plefka, Jakobsen, Mogull, Brandhuber, Travaglini, De Angelis, Accetulli-Huber, Luna, Kosmopoulos, and collaborators…

Newton Potential

Diagrammatic approach

- Just 1 diagram: $\mathscr{M}_{0PN} =$  $= \dfrac{i m_1 m_2}{2 c_d \Lambda^2} \dfrac{1}{\mathbf{p}^2}$

- Fourier transform: from amplitude to the effective action: $\mathscr{L}_{0PN} = -i \lim_{d \to 3} \int \dfrac{d^d \mathbf{p}}{(2\pi)^d} e^{i\mathbf{p}(x_1 - x_2)} \Big($  $\Big) = \dfrac{G_N m_1 m_2}{r}$

Corrections to the Newtonian potential:

- Non-relativistic velocities: $v^2 \ll 1$

- Virial theorem: $\dfrac{G_N m}{r} \approx v^2$

- Dynamics in Post-Newtonian perturbative scheme

- At nPN order: $G_N^{n-\ell} v^{2\ell}$

Astrophysicists/Cosmologists' whishlist

	0PN	1PN	2PN	3PN	4PN	5PN	6PN	7PN		
G	$(1$	$+\ v^2$	$+\ v^4$	$+\ v^6$	$+\ v^8$	$+\ v^{10}$	$+\ v^{12}$	$+\ v^{14}$	$+\ \cdots)$	1PM
G^2		$(1$	$+\ v^2$	$+\ v^4$	$+\ v^6$	$+\ v^8$	$+\ v^{10}$	$+\ v^{12}$	$+\ \cdots)$	2PM
G^3			$(1$	$+\ v^2$	$+\ v^4$	$+\ v^6$	$+\ v^8$	$+\ v^{10}$	$+\ \cdots)$	3PM
G^4				$(1$	$+\ v^2$	$+\ v^4$	$+\ v^6$	$+\ v^8$	$+\ \cdots)$	4PM
G^5					$(1$	$+\ v^2$	$+\ v^4$	$+\ v^6$	$+\ \cdots)$	5PM
G^6						$(1$	$+\ v^2$	$+\ v^4$	$+\ \cdots)$	6PM
G^7							$(1$	$+\ v^2$	$+\ \cdots)$	7PM

1687 1938 1980 2000 2014 2021

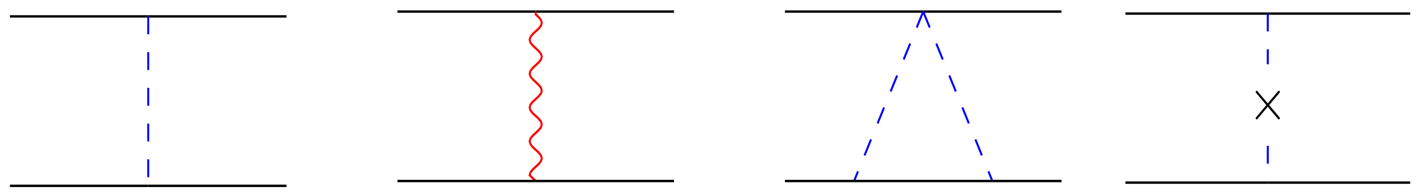
[credit: Bern et al.]

…Jaranowski, Schaefer, Damour, Blanchet, Faye, Porto, Rothstein, Goldberger, Foffa, Sturani, Bini, Buonanno, Geralico, Sturm, Torres Bobadilla, Bluemlein, Maier, Marquard, Levi, Steinhoff, Vines, Antonelli, Kavanagh, Khalil, Galley, von Hippel, McLeod, Edison, Kim, Morales, Yin, Mandal, Patil, Teng, P.M. …and collaborators ….

Post-Newtonian Corrections/EFT Potential

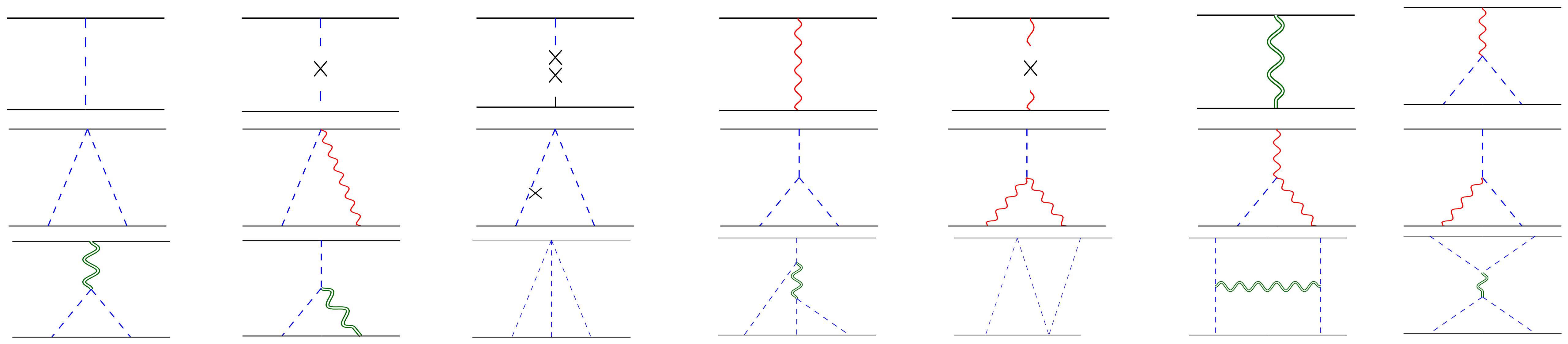
- **1PN corrections:**

  Einstein, Infeld, Hoffman (1938)



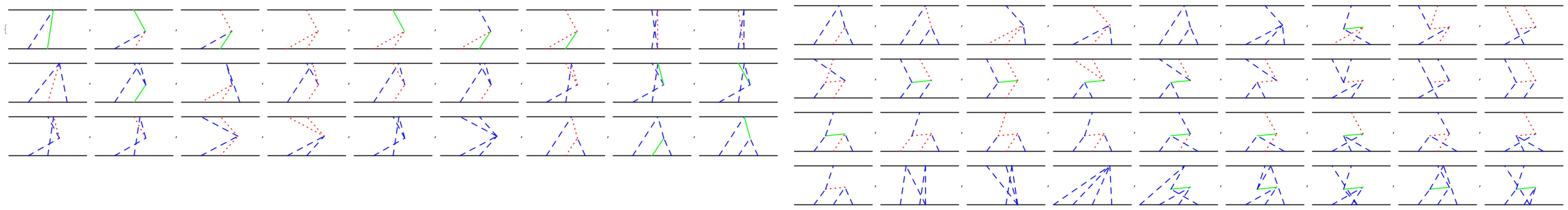
- **2PN corrections:**

  Ohta-Okamura-Kimura-Hiida (1974)
  Gilmore, Ross (2008)



- **3PN corrections:**

  Jaranowski,Schaefer (1997); Damour, Jaranowski, Schaefer (1997); Blachę, Faye (2000); Damour, Jaranowski Schaefer (2001); Foffa Sturani (2011)



- **4PN: corrections:**

  Damour, Jaranowski, Schaefer (2014);
  Bernard, Blanchet, Bohe, Faye, Marsa (2016);
  Foffa, Sturani, Sturm & P.M. (2016);
  Foffa, Porto, Rothstein, Sturani (2019)
  Blumlein, Maier, Marquard,Schaefer (2020)

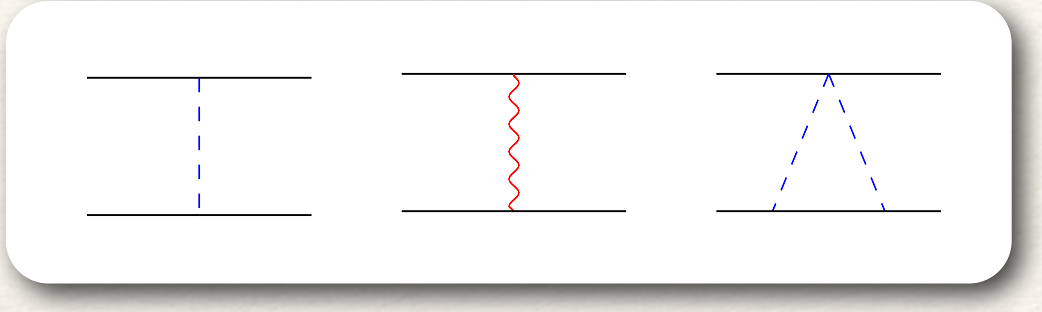
- **5PN: corrections:**

  Bini, Damour, Geralico (2019);
  Foffa, Sturani, Sturm, Torres Bobadilla & P.M. (2019);
  Blumlein, Maier, Marquard, Schaefer (2020,2021)

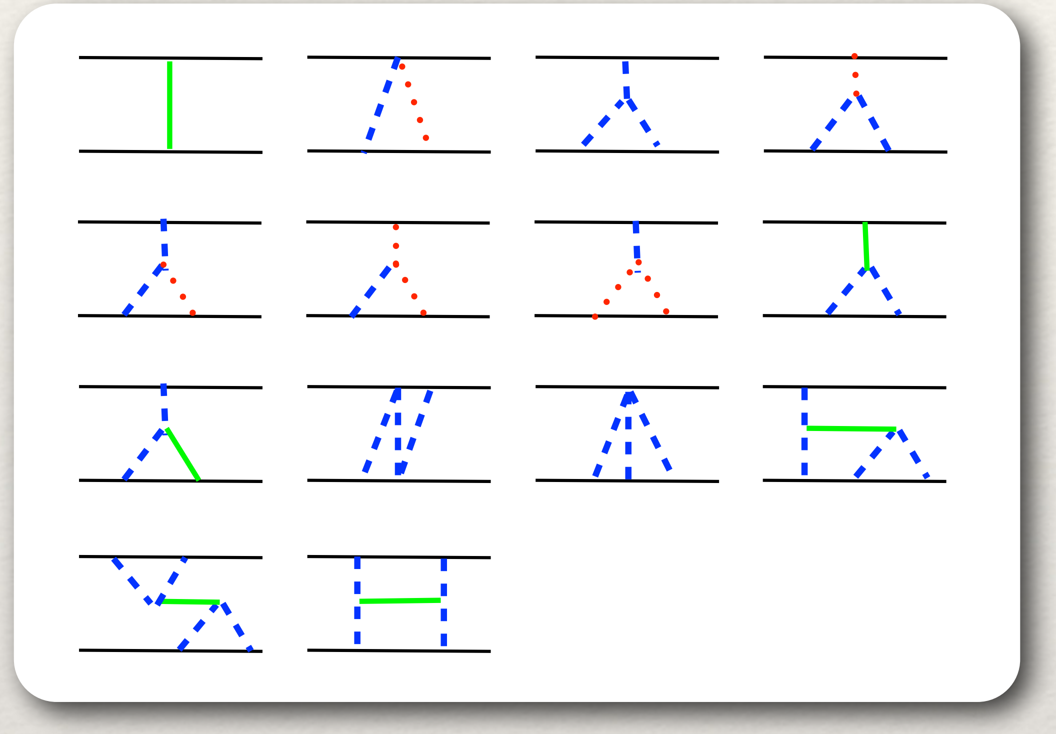
A closer look to 4PN anatomy

- Loop nr. $0 \leq \ell \leq n-1$

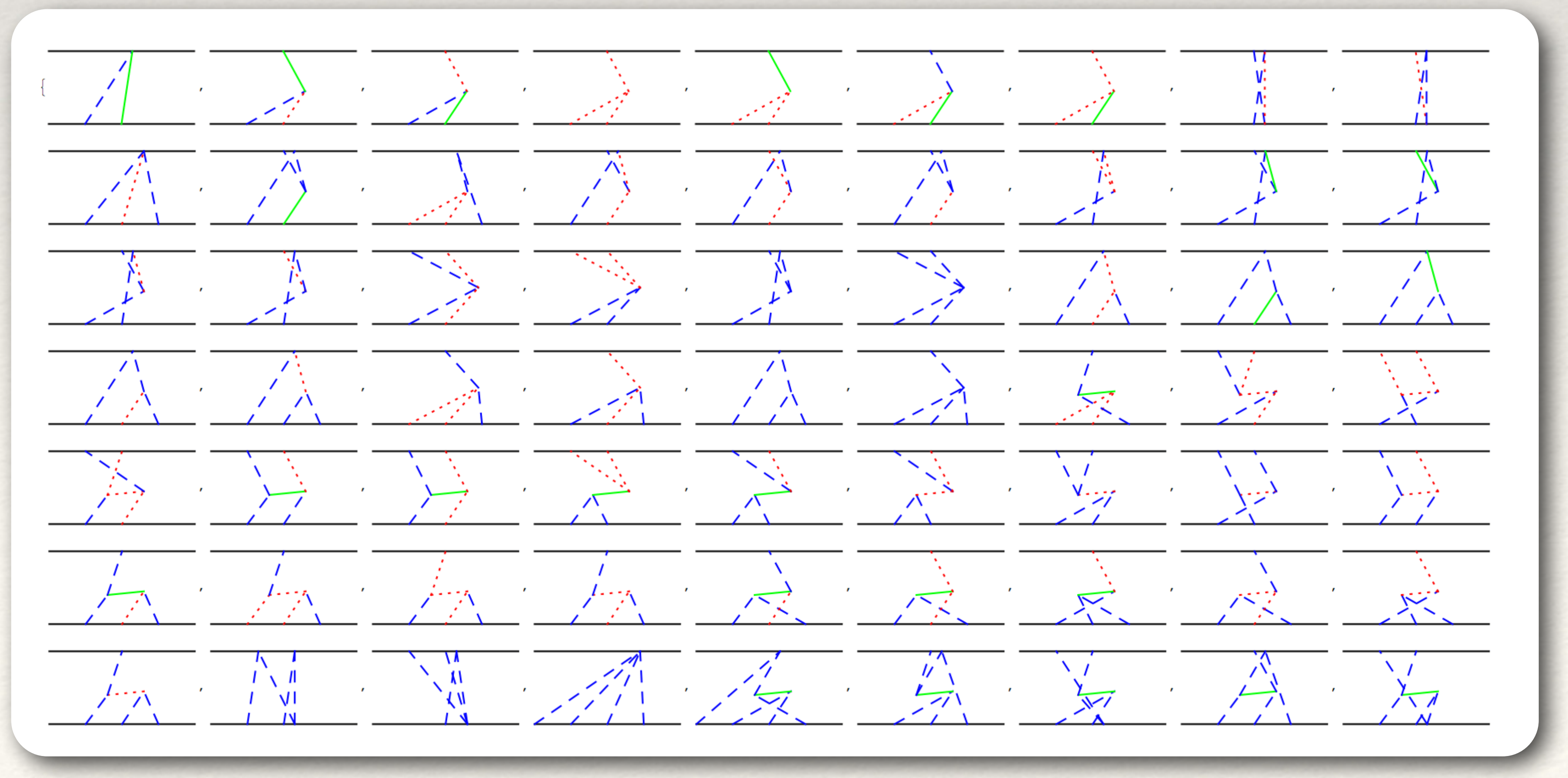
1PN



2PN

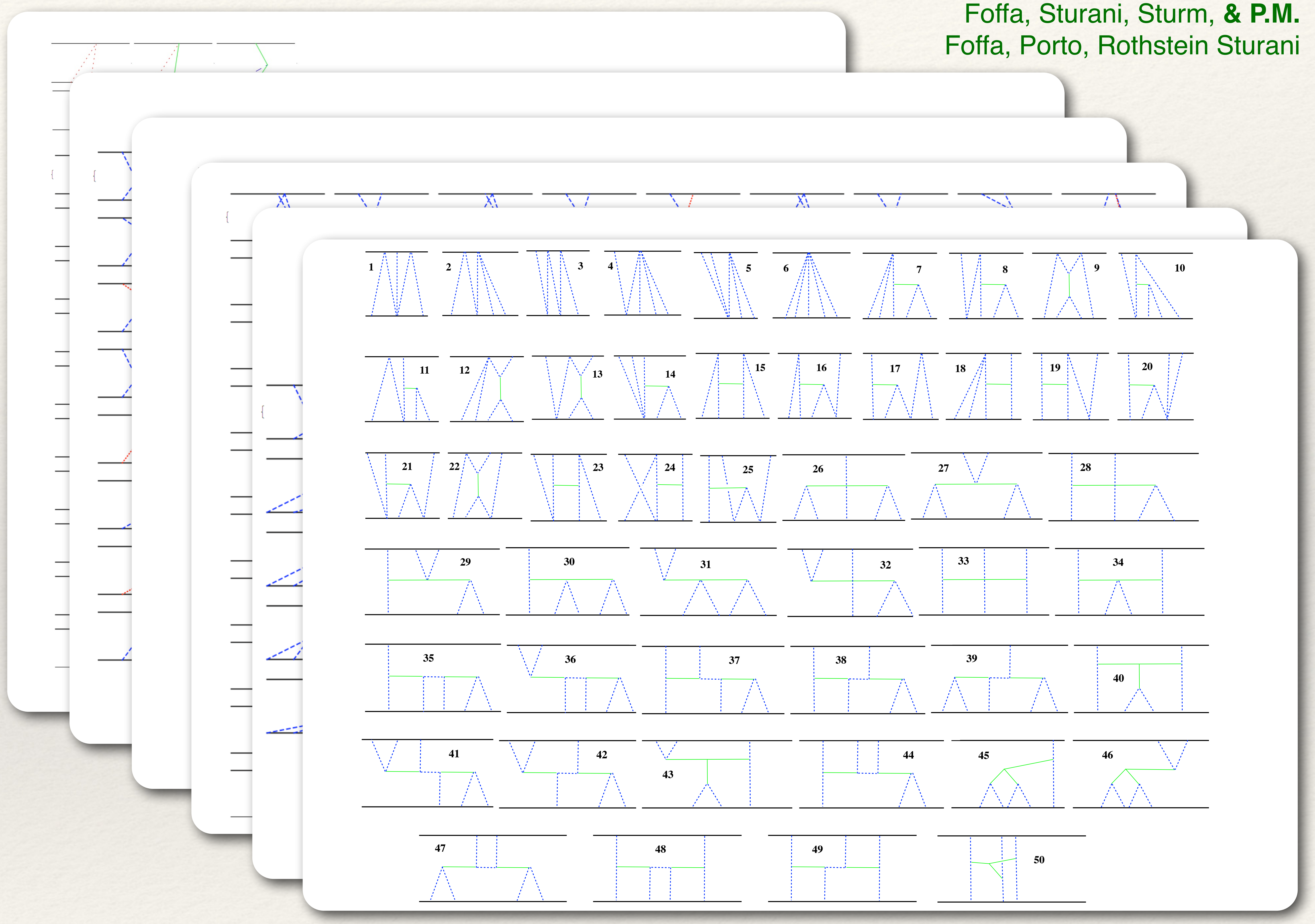


3PN



4PN : 605 GREFT diagrams (up-to 4-loops)

Foffa & Sturani
Foffa, Sturani, Sturm, **& P.M.**
Foffa, Porto, Rothstein Sturani

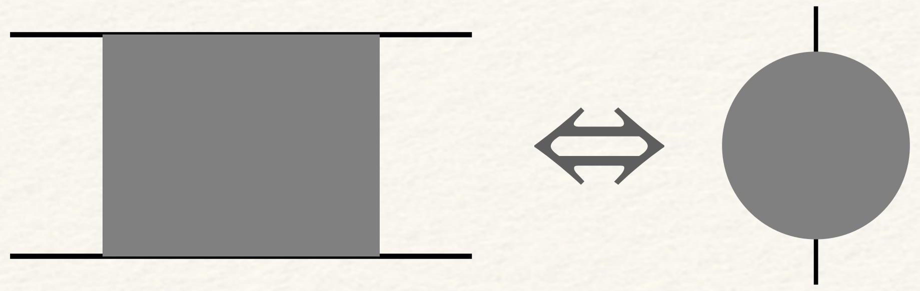


GREFT Diagrams & 2pt-QFT Integrals / a key observation

Foffa, Sturani, Sturm, & **P.M.** (2016)

Computational techniques:

- From Effective diagrams to QFT Amplitudes:
- World-lines are not propagating
- EFTGravity Amplitudes of order G_N^{ℓ} mapped into $(\ell - 1)$—loop 2-point functions with massless internal lines:
- Amplitudes evaluation with QFT multi-loop techniques



$$\mathscr{M} = \sum_i c_i \, I_i^{MI}$$

- Dimensional Regularization $d = 3 + \epsilon$
- Integration-by-parts (IBP) decomposition
- Master Integrals evaluation

- From QFT Amplitudes to Effective Lagrangians:

$$\mathscr{L}_{eff}[x_a] = -\, i \lim_{d\to 3} \int \frac{d^d\mathbf{p}}{(2\pi)^d} e^{i\mathbf{p}\cdot\mathbf{r}} \Bigg(\qquad \Bigg)$$

GREFT Diagrams & 2pt-QFT Integrals / a key observation

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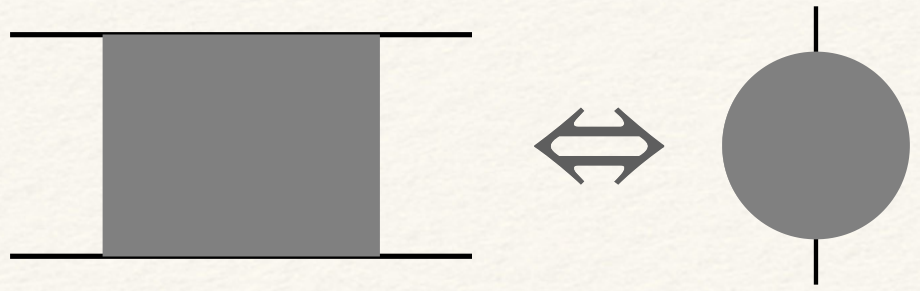

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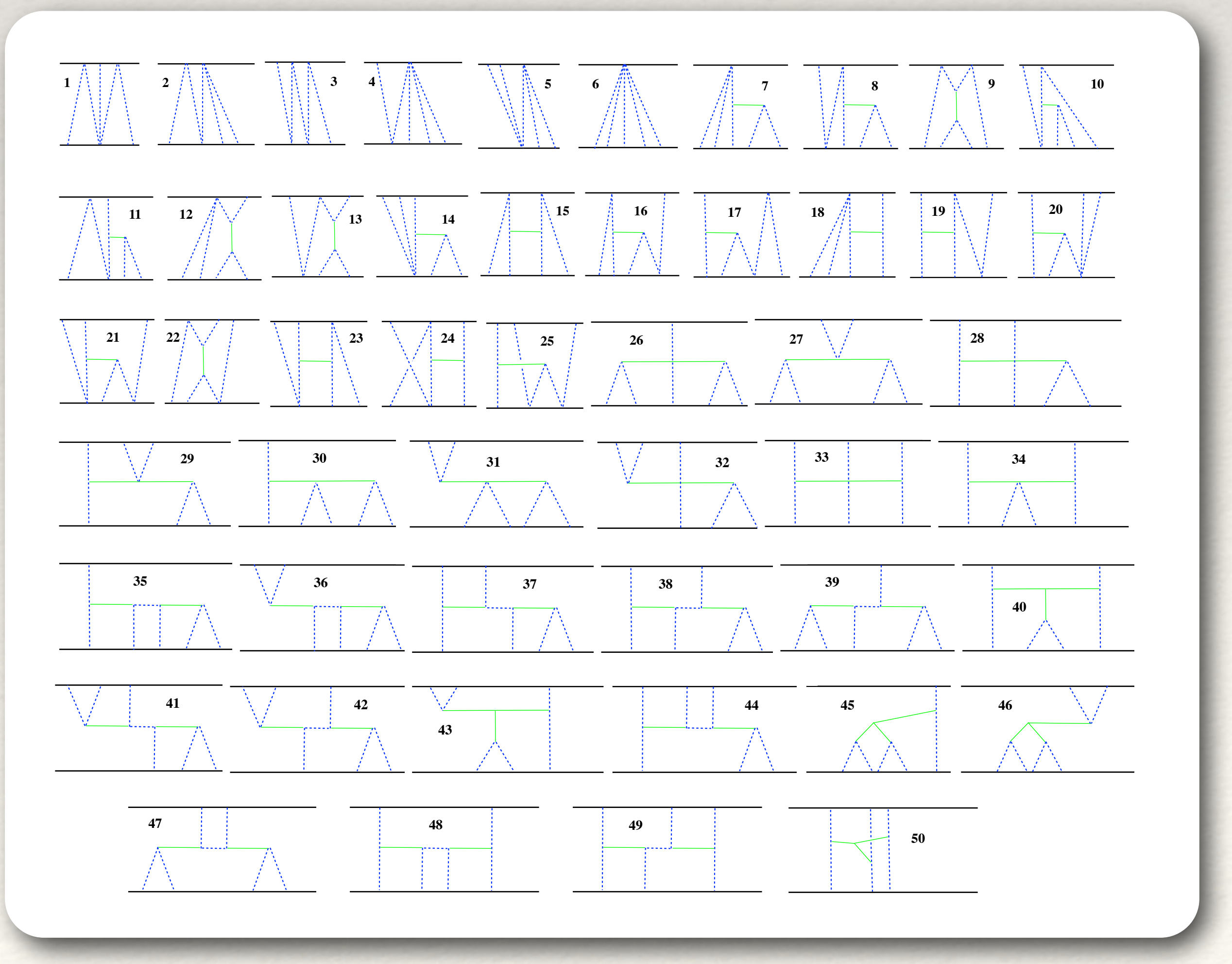
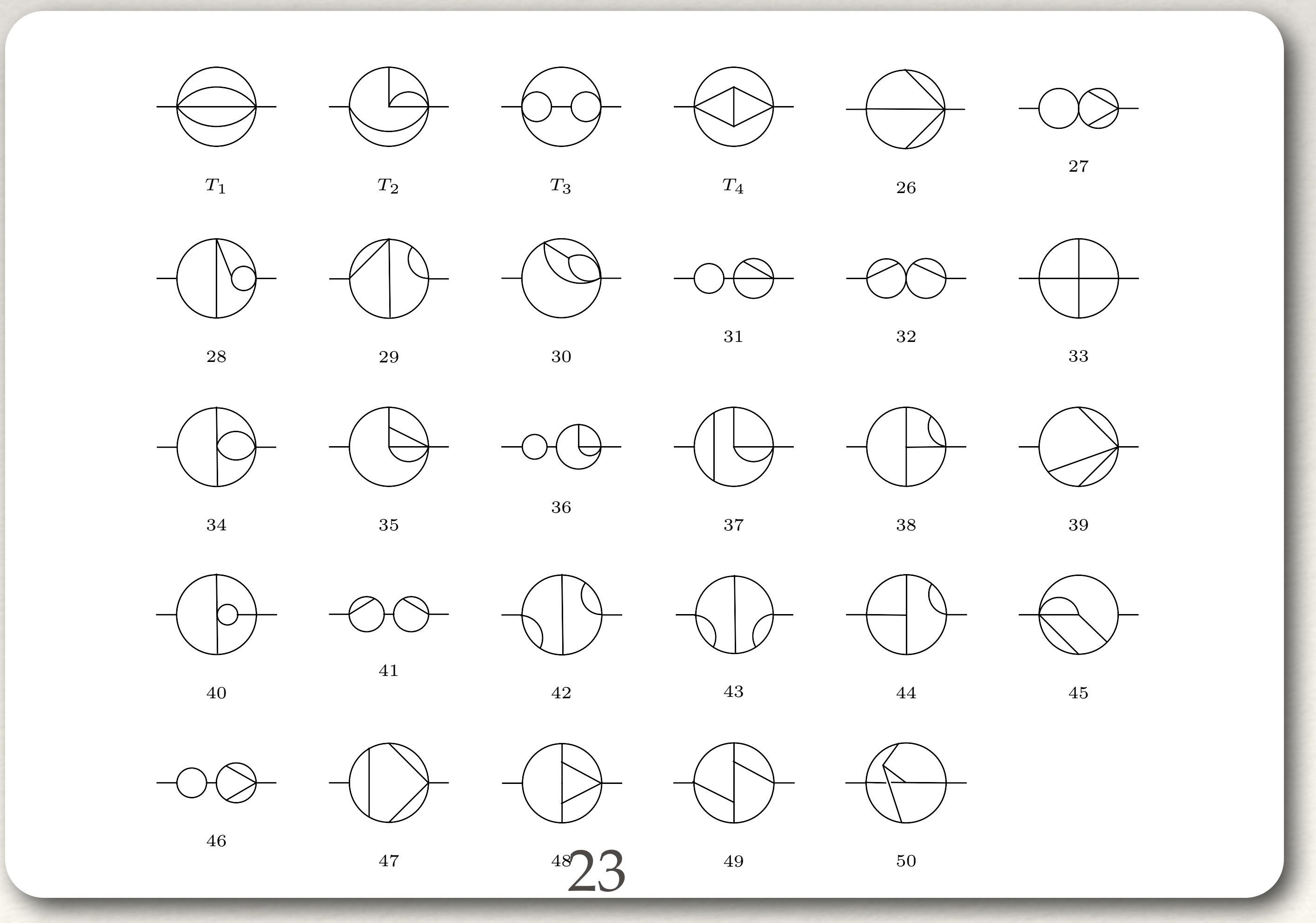
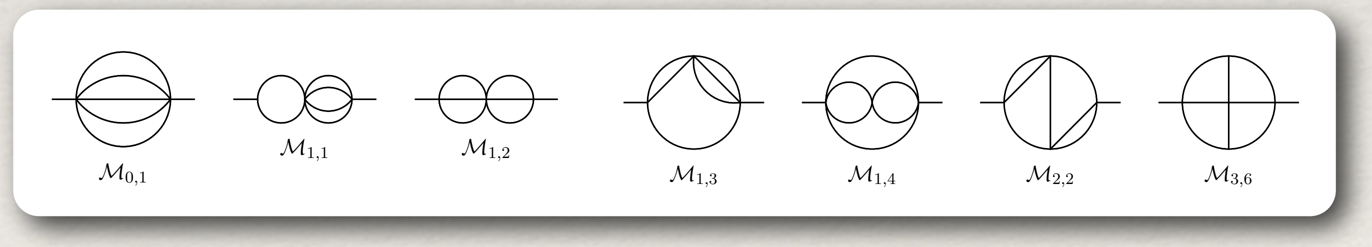
4PN static O(G^5): 50 4-loop GREFT diagrams



29 4-loop QFT diagrams

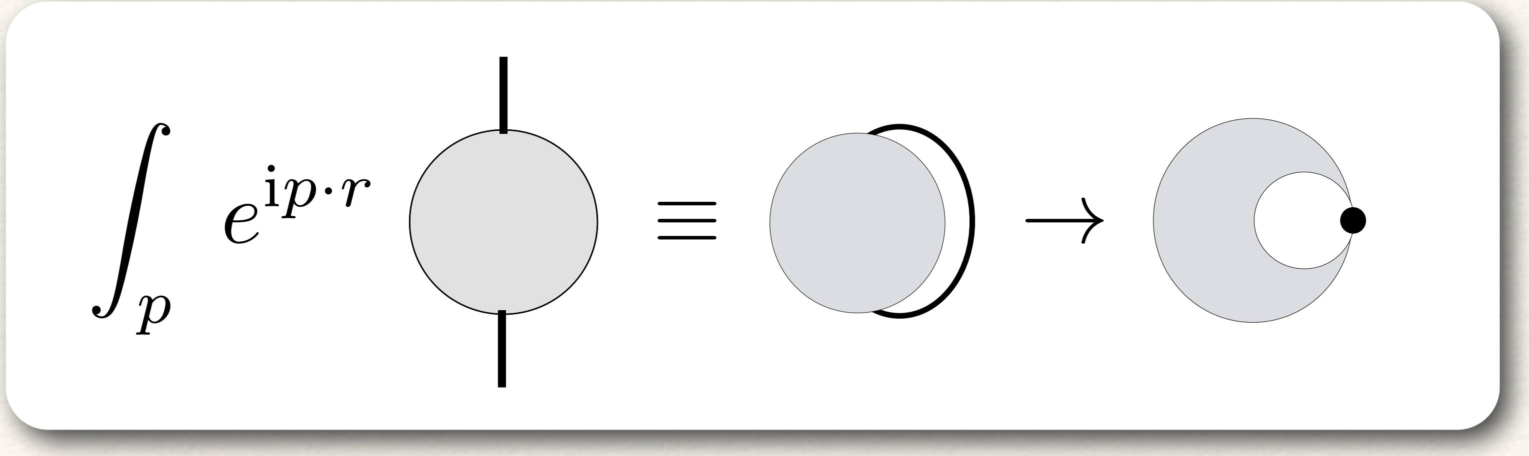


7 4-loop Master Integrals

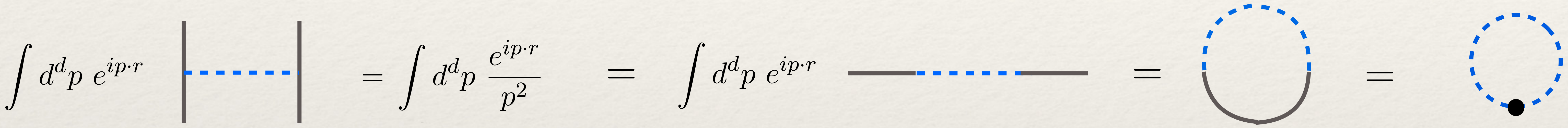


GREFT Diagrams & 2pt-QFT Integrals / Factorization Th'm

Foffa, Sturani, Sturm, Torres-Bobadilla & **P.M.** (2019)

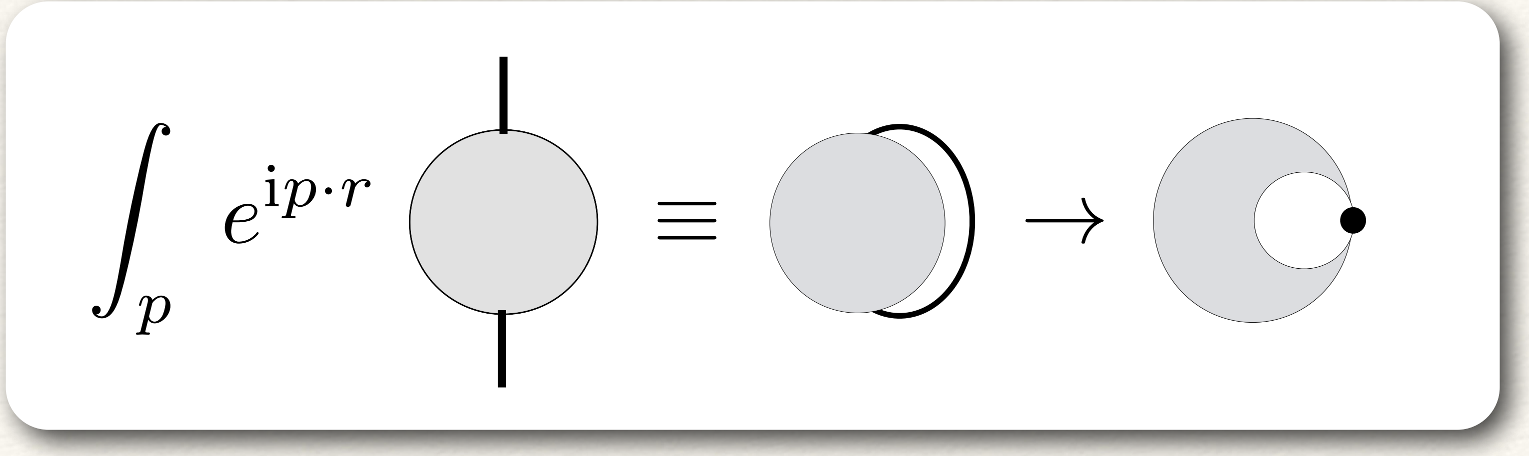


Newton Potential (reloaded):

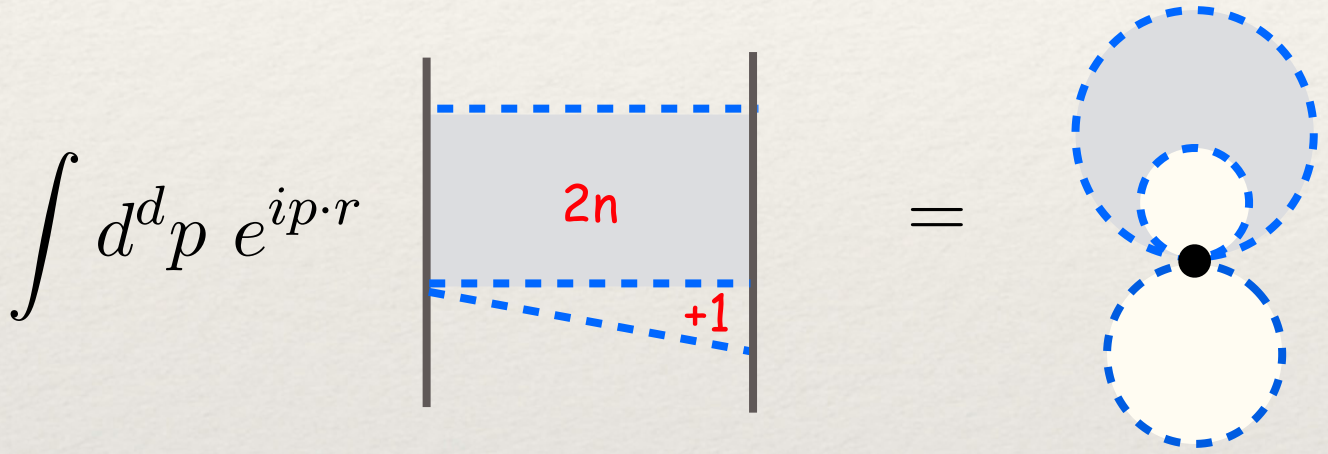


GREFT Diagrams & 2pt-QFT Integrals / Factorization Th'm

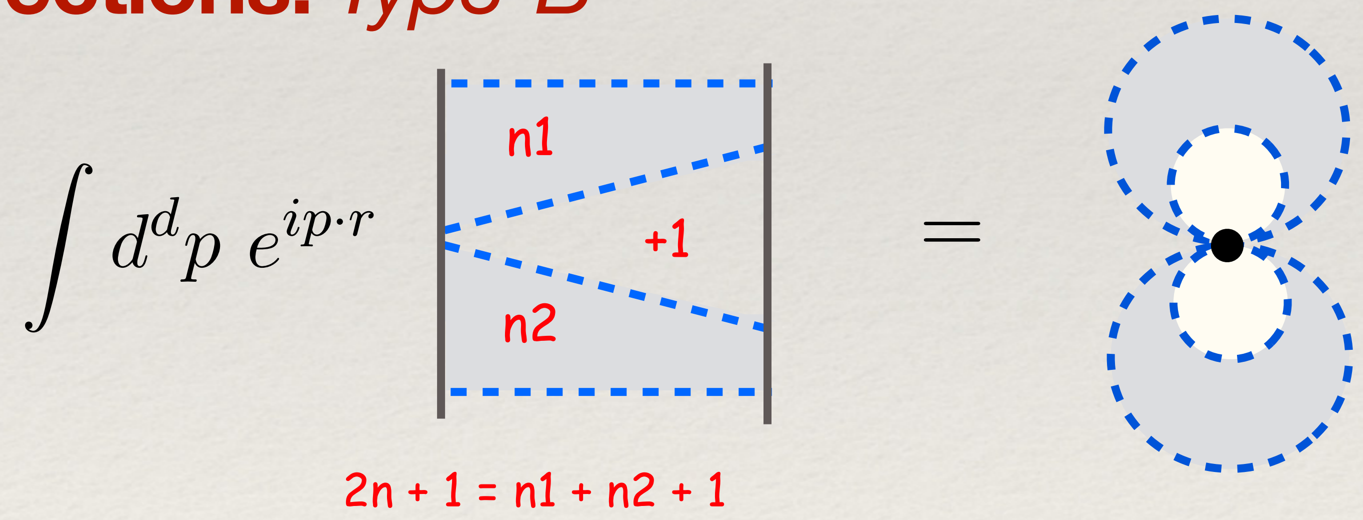
Foffa, Sturani, Sturm, Torres-Bobadilla & **P.M.** (2019)



(2n+1)-PN corrections: *Type-A*



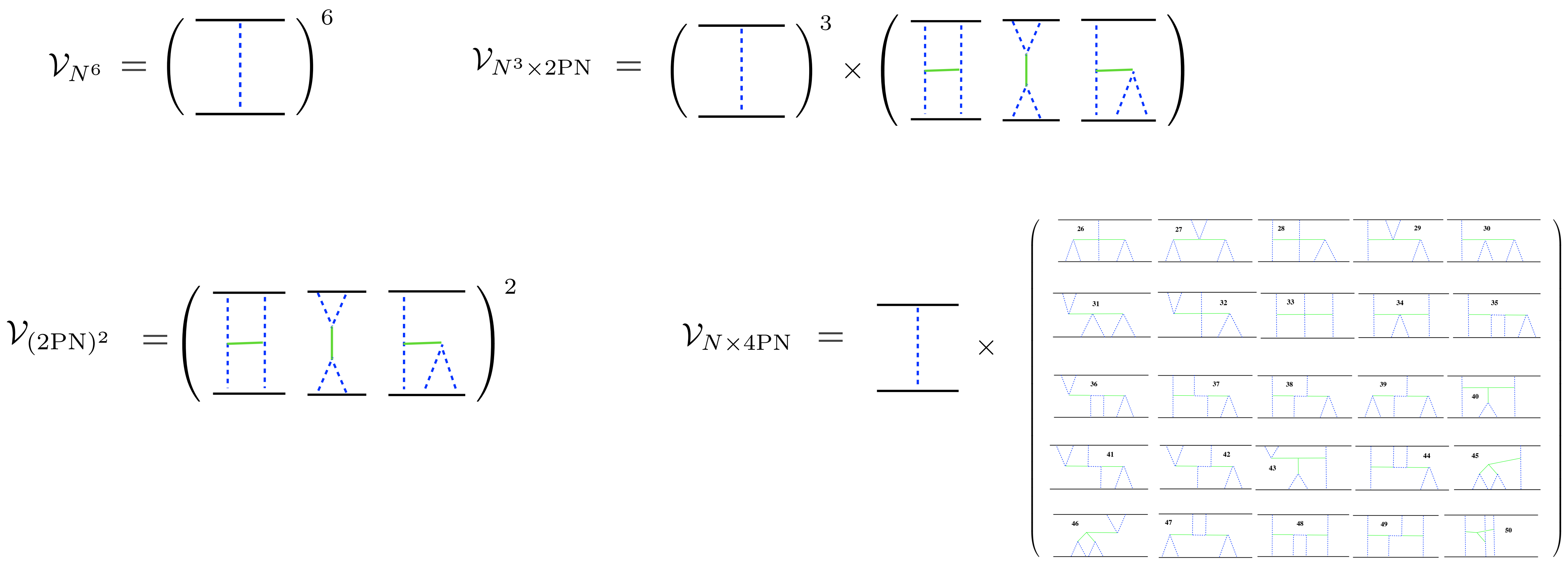
(2n+1)-PN corrections: *Type-B*



- static (2n+1)-PN Potential as product of lower-PN Potential terms

5PN static O(G^6): 154 5-loop GREFT diagrams

Foffa, Sturani, Sturm, Torres-Bobadilla & **P.M.**



$$\mathcal{V}^{(5\mathrm{PN})}_{\mathrm{static}} = \mathcal{V}_{N^6} + \mathcal{V}_{N^3\times 2\mathrm{PN}} + \mathcal{V}_{N\times 4\mathrm{PN}} + \mathcal{V}_{(2\mathrm{PN})^2}$$

5PN O(G^5 v^2): 1220 4-loop GREFT diagrams

Foffa, Sturani,, Torres-Bobadilla (2020)

- **Factorization Th'm: NO 5-loop diagram explicitly computed**

- Results confirmed and completed by explicit evaluation of 2pt-QFT 5-loop Integrals

Blümelein, Maier, Marquard, Schäfer (2019-21)

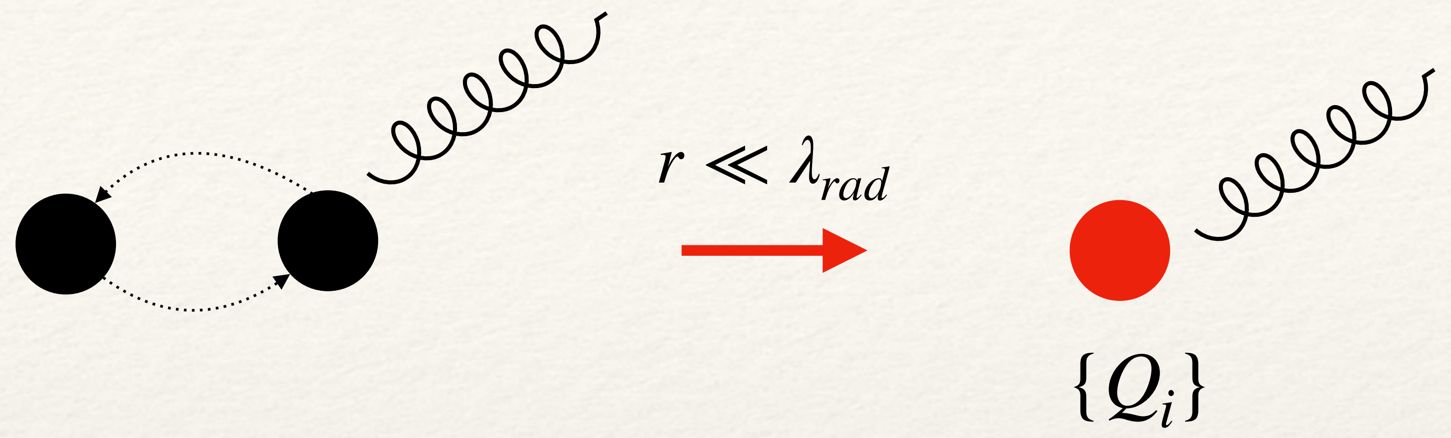
Conservative Dynamics :: **Far Zone** Spinless

Far Zone/EFT Diagrammatic Approach

Thorne (1980)
Goldberger, Rothstein (2005)
Goldberger, Ross (2009)
Galley, Tiglio (2009,2012)
Foffa, Sturani (2012); Ross (2012)
Galley, Leibovich, Porto, Ross (2015)
Leibovich, Maia, Rothstein, Yang (2019)
Blanchet et al.(2021)
....

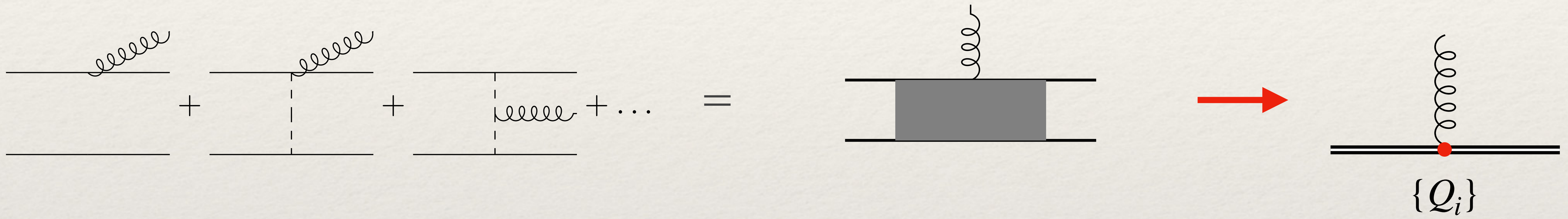
$$S_{rad}[g, \{Q_i\}] = S_{GR}[g] + S_{mult}[g, \{Q_i\}]$$

- **Far zone contributions** to the conservative dynamics are needed, starting at $4PN$ order



Multipole source emitting gravitons

- **Long-wavelength EFT:**

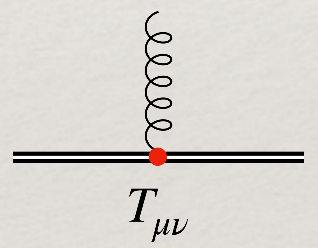


EFT matching

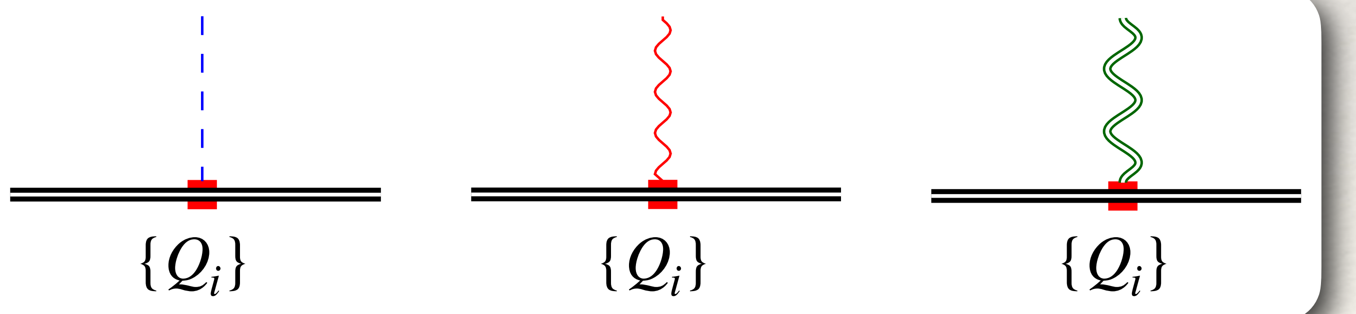
- **Multipole Action:**

Binary system as a linear source $T_{\mu\nu}$ of size r emitting $\bar{h}_{\mu\nu}$:

$$S_{mult} = -\frac{1}{2}\int d^4x T^{\mu\nu}\bar{h}_{\mu\nu}$$



$$S_{mult}[\bar{h}, \{Q_i\}] = \int dt \left[\frac{1}{2}E\bar{h}_{00} - \frac{1}{2}\epsilon_{ijk}L^i\bar{h}_{0j,k} - \frac{1}{2}Q^{ij}\mathscr{E}_{ij} - \frac{1}{6}O^{ijk}\mathscr{E}_{ij,k} - \frac{2}{3}J^{ij}B_{ij} + \dots\right]$$



- $\mathscr{E}_{ij}$, B_{ij} **are the electric and magnetic components of the Riemann tensor**

$$\mathscr{E}_{ij} = R_{0i0j} \approx -\frac{1}{2}\left(\bar{h}_{00,ij} + \ddot{\bar{h}}_{ij} - \dot{\bar{h}}_{0i,j} - \dot{\bar{h}}_{0j,i} + \mathcal{O}(\bar{h}^2)\right)$$

$$B_{ij} = \frac{1}{2}\epsilon_{ikl}R_{0jkl} \approx \frac{1}{4}\epsilon_{ikl}\left(\dot{\bar{h}}_{jk,l} - \dot{\bar{h}}_{jl,k} + \bar{h}_{0l,jk} - \bar{h}_{0k,jl} + \mathcal{O}(\bar{h}^2)\right)$$

- $\{Q_i\}$: **multipole moments** $E, L^i, Q^{ij}, O^{ijk}, J^{ij}$

Far Zone/EFT Diagrammatic Approach

Thorne (1980)
Goldberger, Rothstein (2005)
Goldberger, Ross (2009)
Galley, Tiglio (2009,2012)
Foffa, Sturani (2012); Ross (2012)
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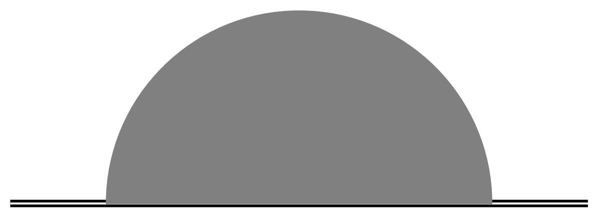
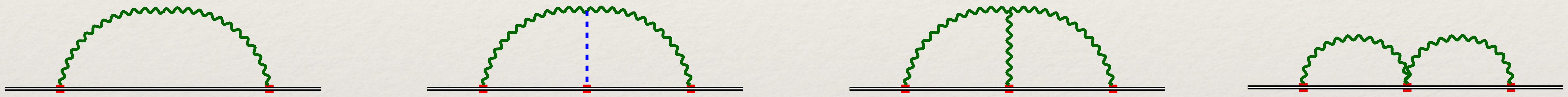
Leibovich, Maia, Rothstein, Yang (2019)
Blanchet et al.(2021)
Almeida, Foffa, Sturani (2021,2022)
Blumlein, Maier, Marquard, Schaefer (2021)
Edison, Levi (2022)
Brunello, Mandal, Patil & P.M. in progress

Hereditary Effects

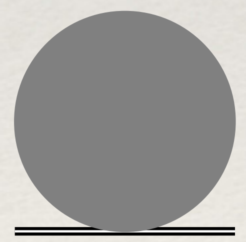
- Contributions to the conservative dynamics by integrating out radiation gravitons:

$$S_{eff}[\{Q_i\}] = -i \lim_{d\to 3}$$ 

- Hereditary Effects: GWs emitted by the source and then back-scattered into the system:



- EFTGravity Amplitude mapped into multi-loop 1-point functions with massive internal lines:

 $$= \int \prod_{i=1}^{n}\left[\frac{dk_0^i}{2\pi}\right]\Bigg($$  $$\Bigg)$$

Non propagating sources

Radiation gravitons propagator:

$$\frac{1}{\mathbf{k}^2 - k_0^2}$$

k_0 mass

$$\mathcal{M} = \sum_i c_i \, I_i^{MI}$$

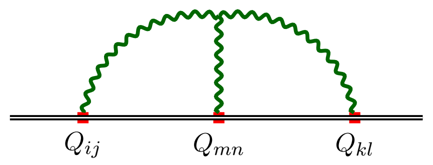
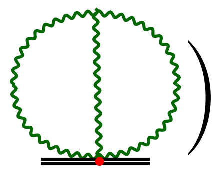
- Dimensional Regularization $d = 3 + \epsilon$
- Integration-by-parts (IBP) decomposition
- Master Integrals evaluation

Far Zone/EFT Diagrammatic Approach

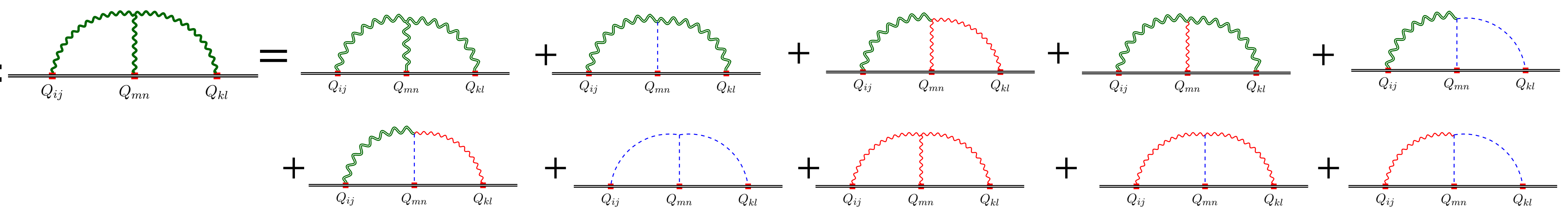
Foffa, Sturani
Blumlein, Maier, Marquard, Schaefer
Almeida, Foffa, Sturani
Brunello, Mandal, Patil & P.M. in progress

Memory Effect (5PN) within the *In-Out (causal)* formalism

Foffa, Sturani (2019)

- **Skeleton diagram:**  $= \int \frac{dk_0}{(2\pi)}\frac{dp_0}{(2\pi)}$ () $= \int \frac{dk_0}{(2\pi)}\frac{dp_0}{(2\pi)}\, \mathcal{M}^{Q^3}$

2 loop 1-point function

- **Diagram Generation:** 

(27 diagrams)

- **Tensor Decomposition:**

$$\mathcal{M}^{Q^3} = Tr[\ Q(k_0)Q(p_0)Q(q_0)\ ]\ \tilde{\mathcal{M}}^{Q^3}$$

- **IBP Decomposition:**

$$\tilde{\mathcal{M}}^{Q^3} = C_1\ (1)(3) + C_2\ (2)(3) + C_3\ (1)(2) + C_4\ \ominus$$

$C_4 \to = 0$

- **Total result:**

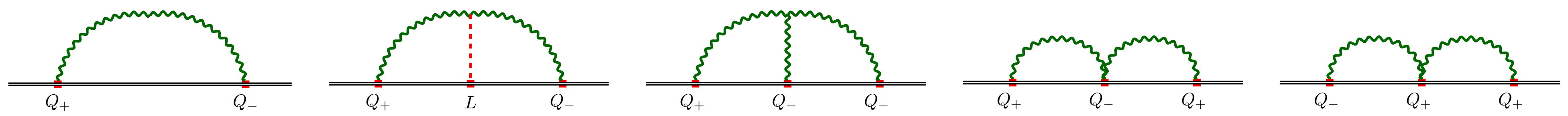
$$S^{Q^3}_{eff\ 5PN} = -i\int \frac{dk_0}{(2\pi)}\frac{dp_0}{(2\pi)}\,\tilde{\mathcal{M}}^{Q^3} = -\frac{G_N^2}{15}\int dt\ Tr\left[Q^{(4)}Q^{(4)}Q + \frac{4}{7}Q^{(3)}Q^{(3)}Q^{(2)}\right]$$

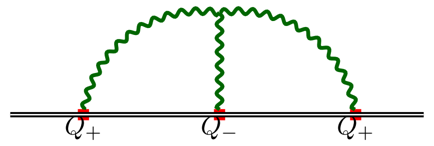
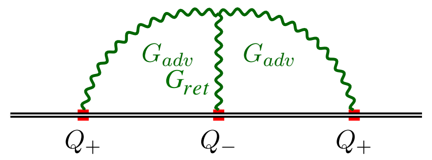
in agreement with
Foffa, Sturani (2019)

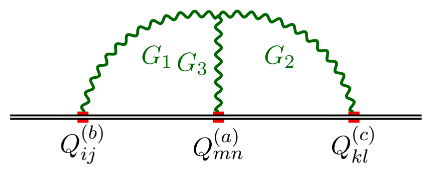
[credit: Brunello]

Memory Effect (5PN) within the *In-In* formalism

Blumlein, Maier, Marquard, Schaefer (2021)
Almeida, Foffa, Sturani (2022)



- **Sample diagram:**  $= \langle (V_{Qh})^{(+)}(V_{Qh})^{(-)}(V_{Qh})^{(+)}V_{h^3} \rangle = 3$ 

- **Tensor Decomposition:**  $= Tr[\ Q^{(a)}Q^{(b)}Q^{(c)}\ ]\ \tilde{\mathcal{M}}^{Q^3}(G_1, G_2, G_3)$

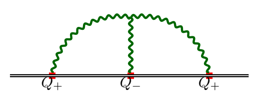
- **IBP Decomposition:**

$$\tilde{\mathcal{M}}^{Q^3}(G_1, G_2, G_3) = C_1\ (G_1)(G_2) + C_2\ (G_1)(G_3) + C_3\ (G_2)(G_3) + C_4\ (G_1, G_2, G_3)$$

IBPs are not sensitive to the prescription

$C_4 \to = 0$

- **Total result:**

$$S^{Q_+Q_+Q_-}_{eff\ 5PN} = -i\lim_{d\to 3}$$  $$= \frac{1}{35}\int dt\ Tr\left[8(\dddot{Q}_+)^2\ddot{Q}_- + 7(\ddot{\ddot{Q}}_+)^2 Q_- - 12\dddot{Q}_+\ddot{Q}_+\dddot{Q}_- - 14\ddot{\ddot{Q}}_+Q_+\dddot{Q}_-\right]$$

agreement with:
Almeida Foffa, Sturani (2022)
Bluemlein, Maier, Marquard, Schaefer (2022)

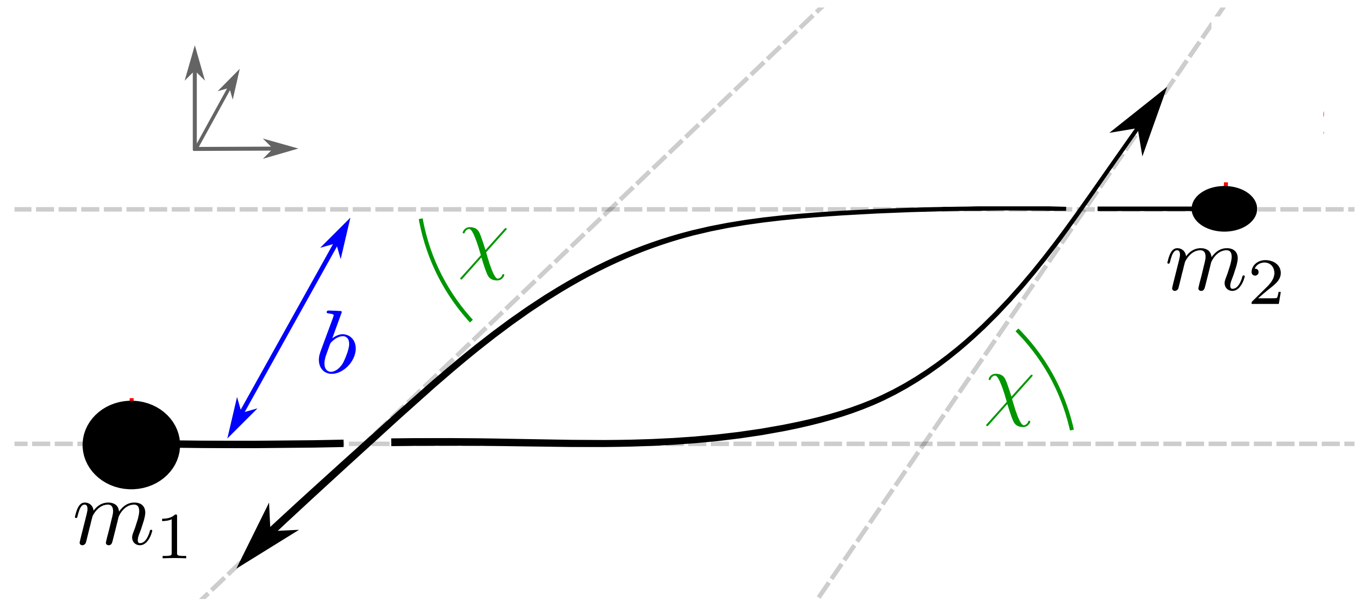
[credit: Brunello]

▸ known known FarZone-GREFT with causal propagators not adequate to describe Radiation/Hereditary effects

▸ known unknown: FarZone-GREFT within Keldysh-Schwinger "*in-in*" formalism under scrutiny

Scattering Angle

$$\chi = -2\int_{r_{\min}}^{\infty} dr \frac{\partial p_r}{\partial L} - \pi$$



[credit: Antornelli et al.]

$$M = m_1 + m_2 \qquad \mu = \frac{m_1 m_2}{m_1 + m_2} \qquad \nu = \frac{\mu}{M}$$

$$p_r = p_r(r, E, L, S_{(a)}) = p_r(r, v, b, S_{(a)})$$

$$\mathbf{p}^2 = p_r^2 + \frac{L^2}{r^2} = p_\infty^2 - V_{eff}\,, \qquad V_{eff}(r) = -\sum_{n\geq 1} f_n(E)\left(\frac{G_N}{r}\right)^n, \qquad p_r = \sqrt{p_\infty^2 - \frac{L^2}{r} - V_{eff}(r)}\,, \qquad V_{eff}(r \to \infty) \to 0\,.$$

$$H^{cons.} = H^{loc} + H^{nonloc.,cons.} \quad \Rightarrow \quad \chi = \chi^{loc} + \chi^{nonloc.}$$

- **PM-expansion:**

$$\frac{1}{2}\chi(b, E) = \sum_n \chi_b^{(n)}(E)\left(\frac{GM}{b}\right)^n = \sum_n \chi_j^{(n)}(E)\frac{1}{j^n}\,,$$

- **PN-expansion:**

$$\chi_b^{(n)} = \sum_{k\geq 0} \chi_b^{(n,k)}\left(\frac{v^2}{c^2}\right)^k$$

$$\chi_j^{(n)} = \hat{p}_\infty^n \chi_b^{(n)}\,, \qquad \hat{p}_\infty = p_\infty/\mu. \qquad j = \frac{L}{G_N M \mu}$$

$$E = M\,\Gamma \qquad \Gamma = \sqrt{1 + 2\nu(\gamma - 1)} \qquad \gamma = \frac{1}{\sqrt{1 - v_\infty^2}} \qquad p_\infty = \frac{m_1 m_2}{E}\sqrt{\gamma^2 - 1} = \mu^2\frac{\gamma^2 - 1}{\Gamma^2}$$

Scattering Angle / far zone (no spin) : 6PN & 7PN

Bini Damour Geralico (2020)

Bini Damour Geralico Laporta & P.M.(2020)

χ^{loc} analytically known

$$\chi^{nonloc.} = \frac{\partial}{\partial L} \int_{-\infty}^{\infty} dt \; H^{nonloc.}(t)$$

Multipole Radiation Formula

$$H^{nonloc.}(t) \propto \dddot{Q}_{ij}(t)\ \mathrm{PF}_T \int_{-\infty}^{\infty} \frac{d\tau}{|\tau|} \dddot{Q}_{ij}(t+\tau) \quad \text{+ higher-multipole terms}$$

$$\text{time scale } T \equiv 2r_{12}/c$$

Quadrupole moment

$$Q_{ij} \equiv \sum_a m_a \left(x_a^i x_a^j - \frac{1}{3}\delta^{ij}\mathbf{x}_a^2 \right) \text{+ PN corrections}$$

Hadamard Partie Finie

$$\mathrm{Pf}_T \int_0^{+\infty} \frac{dv}{v} g(v) \equiv \int_0^T \frac{dv}{v}[g(v) - g(0)] + \int_T^{+\infty} \frac{dv}{v} g(v)$$

(... similar to the plus-distribution formula)

▸ **PM-expansion:**

$$\frac{1}{2}\chi^{nonloc.} = \nu p_\infty^4 \left(\frac{A_0}{j^4} + \frac{A_1}{p_\infty j^5} + \frac{A_2}{p_\infty^2 j^6} + \ldots \right)$$

$$p_\infty \equiv \sqrt{\gamma^2 - 1}, \quad \text{and} \quad j = \frac{L}{G_N M \mu}$$

▸ **PN-expansion:**

$$A_m = \sum_{n \geq 0} \left(A_{mn} + A_{mn}^{\ln} \log(p_\infty/2) \right) p_\infty^n, \qquad A_{mn} = \sum_{k \geq 0} A_{mnk} \nu^k$$

$$A_{mnk} = \int_{-1}^{+1} \int_{-1}^{+1} \frac{dT dT'}{|T - T'|} a_{mnk}(T, T')$$

O(G^6) Coefficients

O(200) coefficients:
4 of them coefficients only numerically [Bini Damour Geralico]

▸ **Analytic evaluation**

A_{220}, A_{240}, A_{241}, A_{242} [Bini Damour Geralico Laporta & P.M.]

Extended to O(G^7) Coefficients

TABLE I: Numerical values of the Q_{nk} integrals with 200-digit accuracy.

Q_{20}	524.767292180212584342735955703101758476141999557369011937728711238498839830097712093907037158196060831706238995205677052067946783744966475134730111010455883184170170829347212071124106113165 8613485679
Q_{40}	544.493991570170677225845815854821570135584358333264830495936708341568294815861057428565302986287084252115923423339364981524722633807905033769432119691717874743144282677041484694939992691447 2804761699
Q_{41}	−1029.528875374038496846264209062889513113498910449676867454201338934155133394086571099160008096002700068331117924208151448401433450126671243392588753826600560395213100750620730514064621300 65024513617
Q_{42}	−802.885057050786642755886295069034459970736865058430654964178895902426423211047940727300850918742678716230784351105139654447667985252511824683509405317671631976450608758027815375931918602878433814664

TABLE II: PSLQ reconstruction of the various integrals.

Q_{20}	$\frac{25883}{1800} + \frac{22333}{140}\mathrm{K} - \frac{625463}{3360}\pi - \frac{361911}{560}\pi\ln 2 + \frac{99837}{160}\pi\zeta(3)$
d_{20}	$-\frac{32981}{112} - \frac{9216}{7}\ln 2 + \frac{99837}{64}\zeta(3)$
Q_{40}	$\frac{750674317}{1905120} + \frac{442237}{5040}\mathrm{K} - \frac{571787}{103680}\pi \frac{7207043}{6720}\pi\ln 2 - \frac{190489}{320}\pi\zeta(3)$
d_{40}	$\frac{725051}{1296} + \frac{19920}{7}\ln 2 - \frac{190489}{128}\zeta(3)$
Q_{41}	$-\frac{703435949}{1587600} - \frac{5747}{24}\mathrm{K} + \frac{1154149}{17280}\pi + \frac{1897771}{3360}\pi\ln 2 - \frac{306219}{640}\pi\zeta(3)$
d_{41}	$\frac{607867}{8064} + \frac{20224}{21}\ln 2 - \frac{306219}{256}\zeta(3)$
Q_{42}	$-\frac{59610947}{793800} - \frac{1499}{20}\mathrm{K} - \frac{402163}{2520}\pi + \frac{4497}{80}\pi\ln 2 - \frac{11871}{160}\pi\zeta(3)$
d_{42}	$-\frac{186743}{864} - \frac{11871}{64}\zeta(3)$

1. Numerical reconstruction w/200 digits

TABLE III: Independent sets of HPLs, at the point $x = i$, up to weight four.

$H_{-1}(i)$	$\frac{\ln 2}{2} + i\frac{\pi}{4}$
$H_0(i)$	$i\frac{\pi}{2}$
$H_1(i)$	$-\frac{\ln 2}{2} + i\frac{\pi}{4}$
$H_{0,-1}(i)$	$\frac{\pi^2}{48} + i\mathrm{K}$
$H_{0,1}(i)$	$-\frac{\pi^2}{48} + i\mathrm{K}$
$H_{-1,-1}(i)$	$-\frac{\pi^2}{32} + \frac{\ln^2 2}{8} + \frac{1}{8}i\pi\ln 2$
$H_{-1,1}(i)$	$-\frac{\pi^2}{32} - \frac{\ln^2 2}{8} - \frac{3}{8}i\pi\ln 2 + i\mathrm{K}$
$H_{1,-1}(i)$	$-\frac{\pi^2}{32} - \frac{\ln^2 2}{8} + \frac{3}{8}i\pi\ln 2 - i\mathrm{K}$
$H_{1,1}(i)$	$-\frac{\pi^2}{32} + \frac{\ln^2 2}{8} - \frac{1}{8}i\pi\ln 2$
$H_{0,-1,-1}(i)$	$\frac{29}{64}\zeta(3) - \frac{1}{4}\mathrm{K}\pi - i\mathrm{Q}_3$
$H_{0,-1,1}(i)$	$\frac{27}{64}\zeta(3) - \frac{1}{4}\mathrm{K}\pi + i\frac{\pi^3}{32} - 3i\mathrm{Q}_3 - 2i\mathrm{K}\ln 2$
$H_{0,1,-1}(i)$	$\frac{27}{64}\zeta(3) - \frac{1}{4}\mathrm{K}\pi - i\frac{\pi^3}{32} + 3i\mathrm{Q}_3 + 2i\mathrm{K}\ln 2$
$H_{0,1,1}(i)$	$\frac{29}{64}\zeta(3) - \frac{1}{4}\mathrm{K}\pi + i\mathrm{Q}_3$
$H_{0,-1,-1,-1}(i)$	$\frac{61}{15360}\pi^4 - \frac{35}{128}\zeta(3)\ln 2 + \frac{5}{384}\pi^2\ln^2 2 - \frac{5}{384}\ln^4 2 - \frac{5}{16}a_4 + \frac{\pi\mathrm{Q}_3}{4} + i\mathrm{Q}_4$
$H_{0,-1,-1,0}(i)$	$-\frac{\pi^4}{4608} + \frac{\mathrm{K}^2}{2} + \frac{\pi\mathrm{Q}_3}{2} + \frac{29}{128}i\pi\zeta(3) - \frac{7}{48}i\mathrm{K}\pi^2$
$H_{0,-1,1,-1}(i)$	$-\frac{97}{9216}\pi^4 + \frac{91}{128}\zeta(3)\ln 2 - \frac{13}{384}\pi^2\ln^2 2 + \frac{13}{16}a_4 - \frac{13}{384}\ln^4 2 + \frac{3}{4}\pi\mathrm{Q}_3 + \frac{\mathrm{K}^2}{2} - \frac{7}{8}i\pi\zeta(3) + \frac{1}{16}i\pi^3\ln 2 - 2i\mathrm{K}\ln^2 2$ $-\frac{5}{16}i\mathrm{K}\pi^2 + 5i\beta(4) - 6i\mathrm{Q}_3\ln 2 - 9i\mathrm{Q}_4$
$H_{0,-1,1,0}(i)$	$-\frac{71}{4608}\pi^4 + \frac{3}{2}\pi\mathrm{Q}_3 + \mathrm{K}\pi\ln 2 + \frac{\mathrm{K}^2}{2} + \frac{27}{128}i\pi\zeta(3) + \frac{1}{8}i\pi^3\ln 2 - \frac{5}{48}i\mathrm{K}\pi^2 - 3i\beta(4)$
$H_{0,1,-1,-1}(i)$	$\frac{169}{9216}\pi^4 - \frac{77}{128}\zeta(3)\ln 2 + \frac{9}{128}\pi^2\ln^2 2 - \frac{27}{16}a_4 - \frac{9}{128}\ln^4 2 - \frac{3}{4}\pi\mathrm{Q}_3 - \frac{1}{2}\mathrm{K}\pi\ln 2 - \frac{21}{128}i\pi\zeta(3) - \frac{1}{32}i\pi^3\ln 2$ $+i\mathrm{K}\ln^2 2 + i\beta(4) + 2i\mathrm{Q}_3\ln 2 + i\mathrm{Q}_4$
$H_{0,1,-1,0}(i)$	$\frac{73}{4608}\pi^4 + \frac{\mathrm{K}^2}{2} - \mathrm{K}\pi\ln 2 - \frac{3}{2}\pi\mathrm{Q}_3 + \frac{27}{128}i\pi\zeta(3) - \frac{1}{8}i\pi^3\ln 2 - \frac{7}{48}i\mathrm{K}\pi^2 + 3i\beta(4)$
$H_{0,1,1,-1}(i)$	$\frac{61}{9216}\pi^4 + \frac{21}{128}\zeta(3)\ln 2 + \frac{13}{384}\pi^2\ln^2 2 - \frac{13}{384}\ln^4 2 - \frac{13}{16}a_4 - \frac{1}{2}\mathrm{K}\pi\ln 2 + \frac{\mathrm{K}^2}{2} - \frac{1}{4}\pi\mathrm{Q}_3 + \frac{133}{128}i\pi\zeta(3) - \frac{1}{32}i\pi^3\ln 2$ $+i\mathrm{K}\ln^2 2 + \frac{3}{16}i\mathrm{K}\pi^2 - 5i\beta(4) + 4i\mathrm{Q}_3\ln 2 + 7i\mathrm{Q}_4$
$H_{0,1,1,0}(i)$	$-\frac{\pi^4}{4608} + \frac{\mathrm{K}^2}{2} - \frac{1}{2}\pi\mathrm{Q}_3 - \frac{5}{48}i\mathrm{K}\pi^2 + \frac{29}{128}i\pi\zeta(3)$
$H_{-1,-1,-1,0}(i)$	$-\frac{31}{15360}\pi^4 + \frac{1}{2}\zeta(3)\ln 2 - \frac{1}{32}\pi^2\ln^2 2 + \frac{5}{384}\ln^4 2 + \frac{5}{16}a_4 + \frac{29}{256}i\pi\zeta(3) + \frac{1}{96}i\pi\ln^3 2 - \frac{1}{96}i\pi^3\ln 2 - \frac{1}{32}i\mathrm{K}\pi^2$ $-\frac{1}{8}i\mathrm{K}\ln^2 2 - \frac{1}{2}i\mathrm{Q}_3\ln 2 - i\mathrm{Q}_4$
$H_{-1,-1,1,0}(i)$	$-\frac{115}{9216}\pi^4 + \frac{13}{16}\zeta(3)\ln 2 - \frac{1}{12}\pi^2\ln^2 2 + \frac{9}{128}\ln^4 2 + \frac{27}{16}a_4 + \frac{1}{4}\mathrm{K}\pi\ln 2 + \frac{1}{2}\pi\mathrm{Q}_3 + \frac{27}{256}i\pi\zeta(3) - \frac{1}{96}i\pi\ln^3 2$ $+\frac{1}{24}i\pi^3\ln 2 - \frac{1}{32}i\mathrm{K}\pi^2 - \frac{1}{8}i\mathrm{K}\ln^2 2 - i\beta(4) - \frac{1}{2}i\mathrm{Q}_3\ln 2 - i\mathrm{Q}_4$
$H_{-1,1,-1,0}(i)$	$\frac{91}{9216}\pi^4 - \frac{1}{2}\zeta(3)\ln 2 + \frac{17}{96}\pi^2\ln^2 2 - \frac{13}{384}\ln^4 2 - \frac{13}{16}a_4 - \pi\mathrm{Q}_3 + \frac{\mathrm{K}^2}{2} - \frac{3}{4}\mathrm{K}\pi\ln 2 + \frac{335}{256}i\pi\zeta(3) - \frac{1}{96}i\pi\ln^3 2$ $-\frac{3}{64}i\pi^3\ln 2 + \frac{13}{96}i\mathrm{K}\pi^2 + \frac{9}{8}i\mathrm{K}\ln^2 2 - 5i\beta(4) + \frac{9}{2}i\mathrm{Q}_3\ln 2 + 9i\mathrm{Q}_4$
$H_{-1,1,1,0}(i)$	$-\frac{79}{9216}\pi^4 + \frac{1}{16}\zeta(3)\ln 2 - \frac{7}{48}\pi^2\ln^2 2 + \frac{13}{384}\ln^4 2 + \frac{13}{16}a_4 + \frac{\mathrm{K}^2}{2} + \frac{1}{2}\pi\mathrm{Q}_3 + \frac{1}{2}\mathrm{K}\pi\ln 2 - \frac{195}{256}i\pi\zeta(3) + \frac{1}{96}i\pi\ln^3 2$ $+\frac{3}{64}i\pi^3\ln 2 - \frac{31}{96}i\mathrm{K}\pi^2 - \frac{7}{8}i\mathrm{K}\ln^2 2 + 5i\beta(4) - \frac{7}{2}i\mathrm{Q}_3\ln 2 - 7i\mathrm{Q}_4$
$H_{1,-1,-1,0}(i)$	$\frac{55}{9216}\pi^4 - \frac{7}{96}\pi^2\ln^2 2 - \frac{1}{16}\zeta(3)\ln 2 - \frac{13}{384}\ln^4 2 - \frac{13}{16}a_4 + \frac{1}{2}\mathrm{K}\pi\ln 2 + \frac{1}{2}\pi\mathrm{Q}_3 - \frac{\mathrm{K}^2}{2} - \frac{279}{256}i\pi\zeta(3) - \frac{1}{96}i\pi\ln^3 2$ $-\frac{7}{96}i\mathrm{K}\pi^2 - \frac{7}{8}i\mathrm{K}\ln^2 2 + 5i\beta(4) - \frac{7}{2}i\mathrm{Q}_3\ln 2 - 7i\mathrm{Q}_4$
$H_{1,-1,1,0}(i)$	$\frac{29}{9216}\pi^4 + \frac{1}{2}\zeta(3)\ln 2 + \frac{5}{48}\pi^2\ln^2 2 + \frac{13}{384}\ln^4 2 + \frac{13}{16}a_4 - \frac{3}{4}\mathrm{K}\pi\ln 2 - \frac{\mathrm{K}^2}{2} - \pi\mathrm{Q}_3 + \frac{167}{256}i\pi\zeta(3) + \frac{1}{96}i\pi\ln^3 2$ $-\frac{1}{32}i\pi^3\ln 2 + \frac{37}{96}i\mathrm{K}\pi^2 + \frac{9}{8}i\mathrm{K}\ln^2 2 - 5i\beta(4) + \frac{9}{2}i\mathrm{Q}_3\ln 2 + 9i\mathrm{Q}_4$
$H_{1,1,-1,0}(i)$	$\frac{91}{9216}\pi^4 - \frac{13}{16}\zeta(3)\ln 2 + \frac{5}{96}\pi^2\ln^2 2 - \frac{9}{128}\ln^4 2 - \frac{27}{16}a_4 + \frac{1}{4}\mathrm{K}\pi\ln 2 + \frac{1}{2}\pi\mathrm{Q}_3 + \frac{111}{256}i\pi\zeta(3) - \frac{1}{192}i\pi^3\ln 2$ $+\frac{1}{96}i\pi\ln^3 2 - \frac{1}{8}i\mathrm{K}\ln^2 2 - \frac{1}{32}i\mathrm{K}\pi^2 - i\beta(4) - \frac{1}{2}i\mathrm{Q}_3\ln 2 - i\mathrm{Q}_4$
$H_{1,1,1,0}(i)$	$\frac{71}{15360}\pi^4 - \frac{1}{2}\zeta(3)\ln 2 - \frac{5}{384}\ln^4 2 - \frac{5}{16}a_4 + \frac{29}{256}i\pi\zeta(3) + \frac{1}{192}i\pi^3\ln 2 - \frac{1}{32}i\mathrm{K}\pi^2 - \frac{1}{8}i\mathrm{K}\ln^2 2 - \frac{1}{96}i\pi\ln^3 2$ $-\frac{1}{2}i\mathrm{Q}_3\ln 2 - i\mathrm{Q}_4$
$\mathrm{Li}_4(1/2)$	a_4
$\mathrm{ImLi}_2(i)$	K
$\mathrm{ImLi}_4(i)$	$\beta(4)$
$\mathrm{Im}H_{0,1,1}(i)$	Q_3
$\mathrm{Im}H_{0,1,1,1}(i)$	Q_4

2. Analytic integration w/ **HPL's**

Far-Zone GR**EFT** / validation

- Mass polynomiality of the scattering angle: $\chi_4^{cons,tot} = \chi_4^{Schw} + \nu\chi_4^{\nu}$ $\qquad \nu = \frac{\mu}{M}$

- Compatible with “Tutti Frutti“ method and PM-Amplitudes-based calculations

  [Damour] [Bern et al.]

  [Damour, Bini, Geralico]

- GREFT calculations point at possible quadratic behaviour: $\chi_4^{cons,tot} = \chi_4^{Schw} + \nu\chi_4^{\nu} + \nu^2\chi_4^{\nu^2}\,, \qquad \chi_4^{\nu^2} \neq 0$

  [Bluemlein et al.]

  [Almeida et al.] [Porto et al.]

  [Brunello et al.]

- **known unknown: FarZone-GREFT is an challenging theoretical puzzle:**
  - Which effects do the GREFT diagrams contain?
  - Interplay between conservative and dissipative effects?
  - Double counting or missing contribution?
  - FarZone/Radiation and proper choice of Green-Functions

Conservative Dynamics :: **Near Zone** with Spin

Near Zone with Spin/PN Corrections

Porto (2013)
Levi, Steinhoff (2015)
......
Kim, Levi, Yin (2022)
Mandal, Patil, Steinhoff & P.M. (2022)
Levi, Morales, Yin (2022)
Levi, Yin (2022)

Action for Spinning compact object

$$S_{pot}[x_a, g] = S_{GR}[g] + S_{m_a}[x_a, g] \qquad S_{m_a}[x_a, g] = S_{pp}[x_a, g] + \delta S_{m_a}[x_a, g]$$

$$S_{m_a}[x_a, g] = \sum_{a=1,2} \int d\tau \left(-m_{(a)} c \sqrt{u^2_{(a)}} - \frac{1}{2} S_{(a)\mu\nu} \Omega^{\mu\nu}_{(a)} - \frac{S_{(a)\mu\nu} u^\nu_{(a)}}{u^2_{(a)}} \frac{du^\mu_{(a)}}{d\tau} + \mathcal{L}^{(R)}_{(a)} + \mathcal{L}^{(R^2)}_{(a)} + \ldots \right)$$

$$u^\mu_{(a)} \equiv \dot{x}^\mu_a$$

Wilson coefficients that describe the internal structure

$$\mathcal{L}^{(R)}_{(a)} = -\frac{1}{2m_{(a)}c} \left(C^{(0)}_{\mathrm{ES}^2} \right)_{(a)} \frac{E_{\mu\nu}}{u_{(a)}} \left[S^\mu_{(a)} S^\nu_{(a)} \right]_{\mathrm{STF}} + \ldots .$$

$$\mathcal{L}^{(R^2, S^0)}_{(a)} = \frac{1}{2} \left(C^{(2)}_{\mathrm{E}^2} \right)_{(a)} \frac{G^2_N m_{(a)}}{c^5} \frac{E_{\mu\nu} E^{\mu\nu}}{u^3_{(a)}} S^2_{(a)} + \ldots .$$

$$\mathcal{L}^{(R^2, S^2)}_{(a)} = \frac{1}{2} \left(C^{(0)}_{\mathrm{E}^2\mathrm{S}^2} \right)_{(a)} \frac{G^2_N m_{(a)}}{c^5} \frac{E_{\mu\alpha} E_\nu{}^\alpha}{u^3_{(a)}} \left[S^\mu_{(a)} S^\nu_{(a)} \right]_{\mathrm{STF}} + \ldots .$$

Electric and Magnetic components of Riemann tensor

$$E_{\mu\nu} \equiv R_{\mu\alpha\nu\beta} u^\alpha u^\beta$$

$$B_{\mu\nu} \equiv \frac{1}{2} \epsilon_{\alpha\beta\gamma\mu} R^{\alpha\beta}{}_{\delta\nu} u^\gamma u^\delta$$

STF = Symmetrized Trace-Free

	PN order		1.5		2.5		3.5		4.5		5.5		6.5
	0	1		2		3		4		5		6	
S^0	0PN	1PN		2PN		3PN		4PN		5PN		6PN	
S^1			LO		NLO		N2LO		N3LO		N4LO		
S^2				LO		NLO		N2LO		N3LO			
S^3							LO		NLO				
S^4								LO		NLO			
S^5											LO		NLO
S^6												LO	

(L+1)PM/loop order
- tree
- 1-loop
- 2-loop
- 3-loop
- 4-loop
- 5-loop
- 6-loop
- 7-loop

[credit: Vines/Roiban]

Near Zone with Spin/EFT Diagrammatic Approach

Kim, Levi, Yin (2022)
Mandal, Patil, Steinhoff & P.M. (2022)

$$S_{pot}[x_a, g] = S_{GR}[g] + S_{m_a}[x_a, g] \qquad S_{m_a}[x_a, g] = S_{pp}[x_a, g] + \boxed{\delta S_{m_a}[x_a, g]}$$

Spin dependence

- **Kaluza-Klein** parametrization: [Kol Smolkin]

$$g_{\mu\nu} = e^{2\phi/\Lambda}\begin{pmatrix} -1 & A_j/\Lambda \\ A_i/\Lambda & e^{-c_d\frac{\phi}{\Lambda}}\gamma_{ij} - A_iA_j/\Lambda^2 \end{pmatrix}$$

Graviton = Scalar + Vector + Sym. Tensor
10 = 1 + 3 + 6

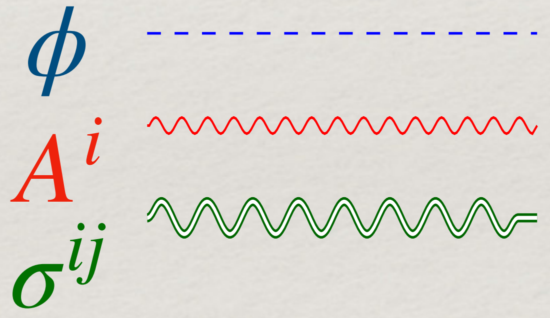
$$g_{\mu\nu} \Rightarrow \phi \quad A^i \quad \sigma^{ij}$$

$$\gamma_{ij} = \delta_{ij} + \frac{\sigma_{ij}}{\Lambda} \qquad c_d = 2\frac{d-1}{d-2}$$

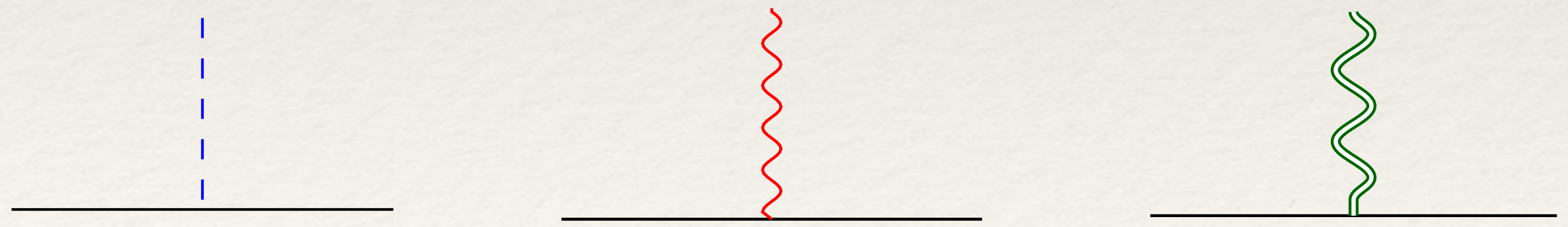
- **Feynman rules** for: $\phi \quad A^i \quad \sigma^{ij} \quad x_a$

Static / non-propagating source: x_a

Propagators:



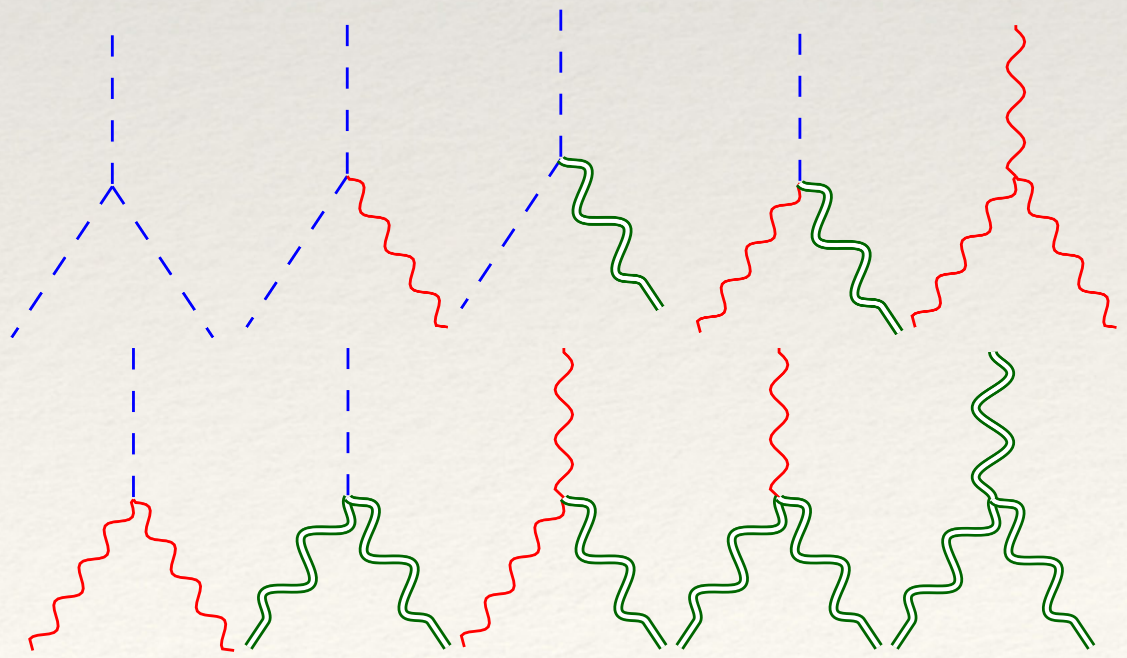
Source couplings:



$+ \ldots$

Spin dependence

Self-interactions:



$+ \ldots$

GREFT Diagrams & 2pt-QFT Integrals

Kim, Levi, Yin (2022)

Mandal, Patil, Steinhoff & P.M. (2022)

S^0

Order	Diagrams	Loops	Diagrams
0PN	1	0	1
1PN	4	1	1
		0	3
2PN	21	2	5
		1	10
		0	6
3PN	130	3	8
		2	75
		1	38
		0	9

(a) Non-spinning sector

S^1

Order	Diagrams	Loops	Diagrams
LO	2	0	2
NLO	13	1	8
		0	5
N^2LO	100	2	56
		1	36
		0	8
N^3LO	894	3	288
		2	495
		1	100
		0	11

(b) Spin-orbit sector

S^2

Order	Diagrams	Loops	Diagrams
LO	1	0	1
NLO	7	1	3
		0	4
N^2LO	58	2	27
		1	24
		0	7
N^3LO	553	3	125
		2	342
		1	76
		0	10

(a) Spin1-Spin2 and Spin1^2 (Spin2^2) sector

Order	Diagrams	Loops	Diagrams
LO	1	0	1
NLO	4	1	1
		0	3
N^2LO	25	2	7
		1	12
		0	6
N^3LO	168	3	15
		2	101
		1	43
		0	9

(b) ES^2 sector

Order	Loops	Diagrams
LO	1	1

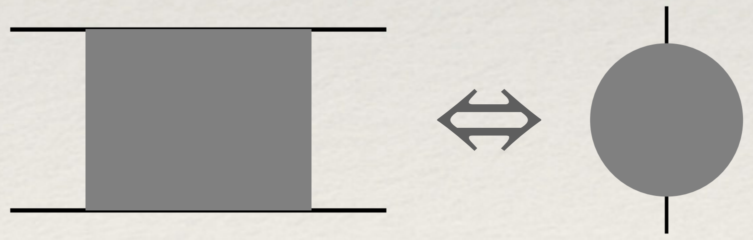
(c) E^2 sector

Order	Loops	Diagrams
LO	1	1

(d) E^2S^2 sector

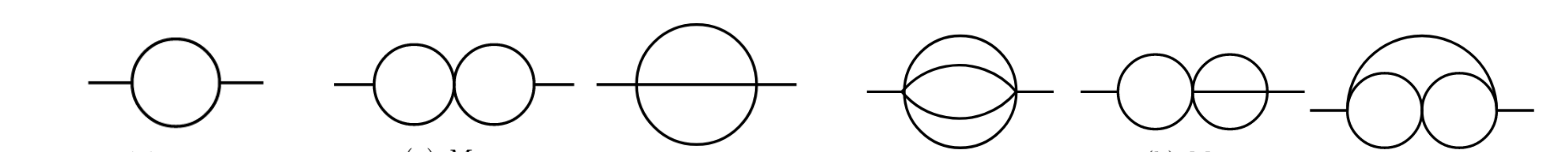
- **Mapping to 2-point Functions**

$$\mathscr{L}_{eff}[x_a, \dot{x}_a, \ddot{x}_a, \dots, S_a, \dot{S}_a, \dots] = -i \lim_{d\to 3} \int \frac{d^d\mathbf{p}}{(2\pi)^d} e^{i\mathbf{p}\cdot\mathbf{r}} \Big(\quad \Big)$$



$$\mathscr{M} = \sum_i c_i \, I_i^{MI}$$

- Dimensional Regularization $d = 3 + \epsilon$
- Integration-by-parts (IBP) decomposition
- Master Integrals evaluation



(a) $M_{1,1}$ (a) $M_{2,1}$ (b) $M_{2,2}$ (a) $M_{3,1}$ (b) $M_{3,2}$ (c) $M_{3,3}$

Conservative Dynamics :: **Near Zone** with Spin

Kim, Levi, Yin (2022)

Mandal, Patil, Steinhoff & P.M. (2022)

▸ Elimination of higher-order time derivatives / equation of motion

borrowed from: Damour, Schafer, Barker, O'Connell

$$\mathbf{x}_{(a)} \to \mathbf{x}_{(a)} + \delta\mathbf{x}_{(a)}$$

$$\delta\mathcal{L} = \left(\frac{\delta\mathcal{L}}{\delta\mathbf{x}^i_{(a)}}\right)\delta\mathbf{x}^i_{(a)} + \frac{1}{2}\left(\frac{\delta^2\mathcal{L}}{\delta\mathbf{x}^i_{(a)}\delta\mathbf{x}^j_{(a)}}\right)\delta\mathbf{x}^i_{(a)}\delta\mathbf{x}^j_{(a)} + \mathcal{O}\left(\delta\mathbf{x}^3_{(a)}\right)$$

$$\boldsymbol{\Lambda}^{ij}_{(a)} \to \boldsymbol{\Lambda}^{ij}_{(a)} + \delta\boldsymbol{\Lambda}^{ij}_{(a)} \qquad \boldsymbol{S}^{ij}_{(a)} \to \boldsymbol{S}^{ij}_{(a)} + \delta\boldsymbol{S}^{ij}_{(a)} \qquad \delta\boldsymbol{\Lambda}^{ij}_{(a)} = \boldsymbol{\Lambda}^{ik}_{(a)}\boldsymbol{\omega}^{kj}_{(a)} + \mathcal{O}\left(\boldsymbol{\omega}^2_{(a)}\right) \qquad \delta\boldsymbol{S}^{ij}_{(a)} = 2\boldsymbol{S}^{k[i}_{(a)}\boldsymbol{\omega}^{j]k}_{(a)} + \mathcal{O}\left(\boldsymbol{\omega}^2_{(a)}\right)$$

$$\delta\mathcal{L}' = -\left(\frac{1}{c}\right)\frac{1}{2}\dot{\boldsymbol{S}}^{ij}_{(a)}\boldsymbol{\omega}^{ij}_{(a)} - \left(\frac{1}{c}\right)\frac{1}{2}\boldsymbol{S}^{ij}_{(a)}\dot{\boldsymbol{\omega}}^{ik}_{(a)}\boldsymbol{\omega}^{kj}_{(a)} - \left(\frac{\delta V}{\delta\boldsymbol{S}^{ij}_{(a)}}\right)\delta\boldsymbol{S}^{ij}_{(a)} + \mathcal{O}\left(\boldsymbol{\omega}^3_{(a)}, \delta\boldsymbol{S}^2_{(a)}\right)$$

➡ $\mathcal{L}'' = \mathcal{L} + \delta\mathcal{L} + \delta\mathcal{L}'$ free of higher-order time derivatives

▸ Elimination of 1/(d-3) divergences and spurious Logarithmic terms / canonical transformations

$$\mathcal{H}(\mathbf{x},\mathbf{p},\mathbf{S}) = \sum_{a=1,2}\mathbf{p}^i_{(a)}\dot{\mathbf{x}}^i_{(a)} - \mathcal{L}''(\mathbf{x},\dot{\mathbf{x}},\mathbf{S})$$

may contain divergences and spurious logarithmic term

➡ $\mathcal{H}' = \mathcal{H} + \{\mathcal{H},\mathcal{G}\}$ — educated guess

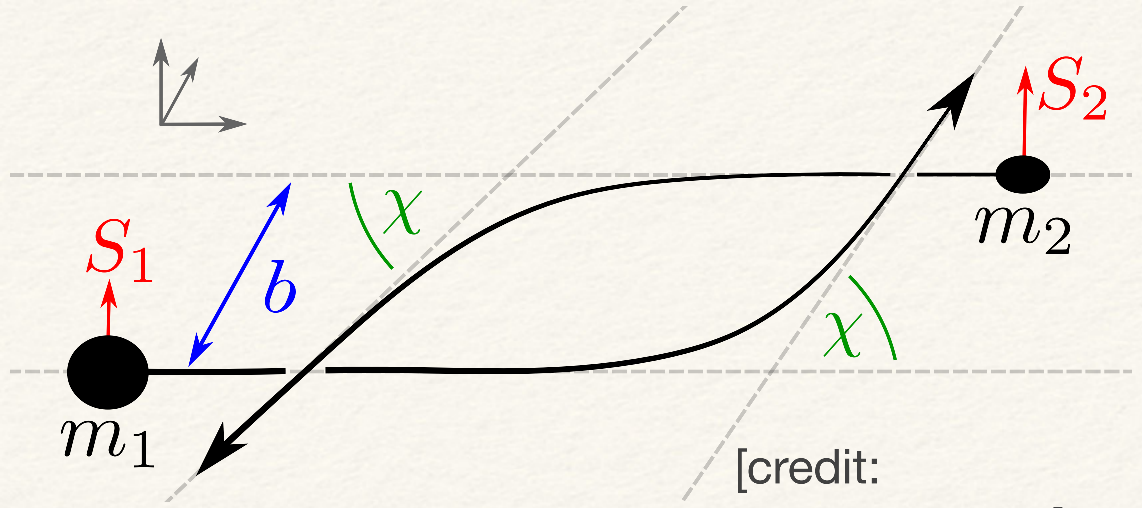
Effective Hamiltionian free of unphysical terms

Scattering Angle :: **Near Zone** with Spin

Mandal, Patil, Steinhoff & P.M. (2022)

▸Aligned spins

$$\chi = -2\int_{r_{\min}}^{\infty} dr\,\frac{\partial p_r}{\partial L} - \pi \qquad\qquad \chi = \chi^{loc}$$



[credit:
Antornelli et al.]

$$\chi(v, b, S_{(a)}) = \chi_{\text{pp}}(v, b) + \chi_{\text{SO}}(v, b, S_{(a)}) + \chi_{\text{SS}}(v, b, S_{(a)})$$

$$\begin{aligned}\chi_{\text{SS}}(v, b, S_{(a)}) &= \chi_{\text{S1S2}}(v, b, S_{(a)}) + \chi_{\text{S}^2}(v, b, S_{(a)}) + \chi_{\text{ES}^2}(v, b, S_{(a)})\\ &\quad + \chi_{\text{E}^2\text{S}^2}(v, b, S_{(a)}) + \chi_{\text{E}^2}(v, b, S_{(a)})\end{aligned}$$

▸**PM**-expansion:

$$\frac{1}{2}\chi(b, E) = \sum_n \chi_b^{(n)}(E)\left(\frac{GM}{b}\right)^n$$

▸**PN**-expansion:

$$\chi_b^{(n)} = \sum_{k\geq 0} \chi_b^{(n,k)}\left(\frac{\mathrm{v}^2}{c^2}\right)^k$$

$$\begin{aligned}\frac{\chi_{\text{pp}}}{\Gamma} &= \left(\frac{G_N M}{v^2 b}\right)\left\{2 + 2\left(\frac{v^2}{c^2}\right) + \mathcal{O}\left(\frac{v^8}{c^8}\right)\right\}\\
&\quad + \pi\left(\frac{G_N M}{v^2 b}\right)^2\left\{3\left(\frac{v^2}{c^2}\right) + \frac{3}{4}\left(\frac{v^4}{c^4}\right) + \mathcal{O}\left(\frac{v^8}{c^8}\right)\right\}\\
&\quad + \left(\frac{G_N M}{v^2 b}\right)^3\left\{-\frac{2}{3} + 2\frac{15-\nu}{3}\left(\frac{v^2}{c^2}\right) + \frac{60-13\nu}{2}\left(\frac{v^4}{c^4}\right) + \frac{40-277\nu}{12}\left(\frac{v^6}{c^6}\right) + \mathcal{O}\left(\frac{v^8}{c^8}\right)\right\}\\
&\quad + \pi\left(\frac{G_N M}{v^2 b}\right)^4\left\{15\frac{7-2\nu}{3}\left(\frac{v^4}{c^4}\right) + \left(\frac{105}{4} - \frac{437}{8}\nu + \frac{123}{128}\pi^2\nu\right)\left(\frac{v^6}{c^6}\right) + \mathcal{O}\left(\frac{v^8}{c^8}\right)\right\}\end{aligned}$$

$$\begin{aligned}\frac{\chi_{\text{SO}}}{\Gamma} = \frac{v}{bc}\begin{bmatrix}a_{(+)} & \delta a_{(-)}\end{bmatrix}\cdot\Bigg(&\left(\frac{G_N M}{v^2 b}\right)\left\{\begin{bmatrix}-4\\0\end{bmatrix} + \mathcal{O}\left(\frac{v^8}{c^8}\right)\right\}\\
&+ \pi\left(\frac{G_N M}{v^2 b}\right)^2\left\{-\frac{1}{2}\begin{bmatrix}7\\1\end{bmatrix} - \frac{3}{4}\begin{bmatrix}7\\1\end{bmatrix}\left(\frac{v^2}{c^2}\right) + \mathcal{O}\left(\frac{v^8}{c^8}\right)\right\}\\
&+ \left(\frac{G_N M}{v^2 b}\right)^3\left\{-2\begin{bmatrix}5\\1\end{bmatrix} - 20\begin{bmatrix}5-\nu/2\\1\end{bmatrix}\left(\frac{v^2}{c^2}\right) - 10\begin{bmatrix}5-77\nu/20\\1\end{bmatrix}\left(\frac{v^4}{c^4}\right)\right.\\
&\qquad\left. + \frac{1}{4}\begin{bmatrix}177\nu\\0\end{bmatrix}\left(\frac{v^6}{c^6}\right) + \mathcal{O}\left(\frac{v^8}{c^8}\right)\right\}\\
&+ \pi\left(\frac{G_N M}{v^2 b}\right)^4\left\{\frac{3}{4}\begin{bmatrix}-91+13\nu\\-21+\nu\end{bmatrix}\left(\frac{v^2}{c^2}\right) - \frac{1}{8}\begin{bmatrix}1365-777\nu\\315-45\nu\end{bmatrix}\left(\frac{v^4}{c^4}\right)\right.\\
&\qquad\left. - \frac{1}{32}\begin{bmatrix}1365-\left(\frac{23717}{3}-\frac{733}{8}\pi^2\right)\nu\\315-\left(\frac{257}{3}+\frac{251}{8}\pi^2\right)\nu\end{bmatrix}\left(\frac{v^6}{c^6}\right) + \mathcal{O}\left(\frac{v^8}{c^8}\right)\right\}\Bigg)\end{aligned}$$

$$\begin{aligned}\frac{\chi_{\text{S1S2}} + \chi_{\text{S}^2}}{\Gamma} = \frac{1}{b^2c^2}\begin{bmatrix}a_{(+)}^2 & \delta a_{(+)}a_{(-)} & a_{(-)}^2\end{bmatrix}\cdot\Bigg(&\left(\frac{G_N M}{v^2 b}\right)\left\{\begin{bmatrix}1\\0\\-1\end{bmatrix} + \begin{bmatrix}1\\0\\-1\end{bmatrix}\left(\frac{v^2}{c^2}\right)\right\}\\
&+ \pi\left(\frac{G_N M}{v^2 b}\right)^2\left\{\frac{3}{4}\begin{bmatrix}1\\0\\-1\end{bmatrix} + \frac{3}{16}\begin{bmatrix}41\\10\\-19\end{bmatrix}\left(\frac{v^2}{c^2}\right) + \frac{3}{128}\begin{bmatrix}55\\-10\\-41\end{bmatrix}\left(\frac{v^4}{c^4}\right)\right\}\\
&+ \left(\frac{G_N M}{v^2 b}\right)^3\left\{2\begin{bmatrix}1\\0\\-1\end{bmatrix} + \begin{bmatrix}102-2\nu\\32\\-38-6\nu\end{bmatrix}\left(\frac{v^2}{c^2}\right) + \frac{1}{14}\begin{bmatrix}2332-499\nu\\688\\-748-37\nu\end{bmatrix}\left(\frac{v^4}{c^4}\right)\right.\\
&\qquad\left. + \frac{1}{140}\begin{bmatrix}1704-9925\nu\\-288\\-1096+5377\nu\end{bmatrix}\left(\frac{v^6}{c^6}\right)\right\}\\
&+ \pi\left(\frac{G_N M}{v^2 b}\right)^4\left\{\frac{15}{16}\begin{bmatrix}63-2\nu\\22\\-21-6\nu\end{bmatrix}\left(\frac{v^2}{c^2}\right) - \frac{15}{448}\begin{bmatrix}-11063+2638\nu\\-4206+308\nu\\2657+790\nu\end{bmatrix}\left(\frac{v^4}{c^4}\right)\right.\\
&\qquad\left. - \frac{3}{229376}\begin{bmatrix}19102720+(-29293696+135555\pi^2)\\-256(-24640+11507\nu)\\-4730880+(8911744-139965\pi^2)\nu\end{bmatrix}\left(\frac{v^6}{c^6}\right)\right\}\Bigg)\end{aligned}$$

$$\begin{aligned}\frac{\chi_{\text{ES}^2}}{\Gamma} = \frac{1}{b^2c^2}\begin{bmatrix}a_{\text{ES}^2(+)}^2 & \delta a_{\text{ES}^2(-)}^2\end{bmatrix}\cdot\Bigg\{&\left(\frac{G_N M}{v^2 b}\right)\left\{\begin{bmatrix}2\\0\end{bmatrix} + \begin{bmatrix}2\\0\end{bmatrix}\left(\frac{v^2}{c^2}\right)\right\}\\
&+ \pi\left(\frac{G_N M}{v^2 b}\right)^2\left\{\frac{1}{2}\begin{bmatrix}3\\0\end{bmatrix} + \frac{1}{4}\begin{bmatrix}27\\3\end{bmatrix}\left(\frac{v^2}{c^2}\right) + \frac{1}{64}\begin{bmatrix}117\\57\end{bmatrix}\left(\frac{v^4}{c^4}\right)\right\}\\
&+ \left(\frac{G_N M}{v^2 b}\right)^3\left\{\begin{bmatrix}4\\0\end{bmatrix} + \begin{bmatrix}76-4\nu\\8\end{bmatrix}\left(\frac{v^2}{c^2}\right) + \frac{1}{7}\begin{bmatrix}748-131\nu\\216\end{bmatrix}\left(\frac{v^4}{c^4}\right) + \frac{1}{70}\begin{bmatrix}1096-5013\nu\\704\end{bmatrix}\left(\frac{v^6}{c^6}\right)\right\}\\
&+ \pi\left(\frac{G_N M}{v^2 b}\right)^4\left\{\frac{15}{4}\begin{bmatrix}11-\nu\\1\end{bmatrix}\left(\frac{v^2}{c^2}\right) + \frac{15}{224}\begin{bmatrix}2867-708\nu\\627-28\nu\end{bmatrix}\left(\frac{v^4}{c^4}\right)\right.\\
&\qquad\left. + \frac{15}{114688}\begin{bmatrix}1026816-(2184832+27111\pi^2)\nu\\128(3262-395\nu)\end{bmatrix}\left(\frac{v^6}{c^6}\right)\right\}\Bigg\}\end{aligned}$$

$$\frac{\chi_{\text{E}^2\text{S}^2}}{\Gamma} = \frac{1}{b^2c^2}\begin{bmatrix}a_{\text{E}^2\text{S}^2(+)}^2 & \delta a_{\text{E}^2\text{S}^2(-)}^2\end{bmatrix}\cdot\left\{\pi\left(\frac{G_N M}{v^2 b}\right)^4\frac{15}{32}\nu\begin{bmatrix}1\\1\end{bmatrix}\left(\frac{v^6}{c^6}\right)\right\} + \mathcal{O}\left(G_N^5, \frac{v^8}{c^8}\right)$$

$$\frac{\chi_{\text{E}^2}}{\Gamma} = \frac{1}{b^2c^2}\begin{bmatrix}a_{\text{E}^2(+)}^2 & \delta a_{\text{E}^2(-)}^2\end{bmatrix}\cdot\left\{\pi\left(\frac{G_N M}{v^2 b}\right)^4\frac{45}{16}\nu\begin{bmatrix}1\\1\end{bmatrix}\left(\frac{v^6}{c^6}\right)\right\} + \mathcal{O}\left(G_N^5, \frac{v^8}{c^8}\right)$$

In agreement with:
Antonelli et al. (2020)
Kim et al. (2022)

Binding Energy :: **Near Zone** with Spin

Mandal, Patil, Steinhoff & P.M. (2022)

- Circular Orbit and aligned spins

$$E(x,\widetilde{S}_{(a)}) = E_{\rm pp}(x) + E_{\rm SO}(x,\widetilde{S}_{(a)}) + E_{\rm SS}(x,\widetilde{S}_{(a)})$$

$$E_{\rm SS}(x,\widetilde{S}_{(a)}) = E_{\rm S1S2}(x,\widetilde{S}_{(a)}) + E_{\rm S^2}(x,\widetilde{S}_{(a)}) + E_{\rm ES^2}(x,\widetilde{S}_{(a)}) + E_{\rm E^2S^2}(x,\widetilde{S}_{(a)}) + E_{\rm E^2}(x,\widetilde{S}_{(a)})$$

$$\begin{aligned} E_{\rm pp}(x) = & -x\frac{1}{2} + x^2\left\{\frac{3}{8}+\frac{\nu}{24}\right\} + x^3\left\{\frac{27}{16}-\frac{19}{16}\nu+\frac{1}{48}\nu^2\right\} \\ & + x^4\left\{\frac{675}{128}+\left(-\frac{34445}{1152}+\frac{205\pi^2}{192}\right)\nu+\frac{155}{192}\nu^2+\frac{35}{10368}\nu^3\right\} \end{aligned}$$

$$\begin{aligned} E_{\rm SO}(x,\widetilde{S}) = & \, x^{5/2}\left\{S^*(-\nu)+S\left(-\frac{4}{3}\nu\right)\right\} \\ & + x^{7/2}\left\{S^*\left(-\frac{3}{2}\nu+\frac{5}{3}\nu^2\right)+S\left(-4\nu+\frac{31}{18}\nu^2\right)\right\} \\ & + x^{9/2}\left\{S^*\left(-\frac{27}{8}\nu+\frac{39}{2}\nu^2-\frac{5}{8}\nu^3\right)+S\left(-\frac{27}{2}\nu+\frac{211}{8}\nu^2-\frac{7}{12}\nu^3\right)\right\} \\ & + x^{11/2}\left\{S^*\left(-\frac{135}{16}\nu+\frac{565}{8}\nu^2-\frac{1109}{24}\nu^3-\frac{25}{324}\nu^4\right)\right. \\ & \left. + S\left(-45\nu+\left(\frac{19679}{144}+\frac{29\pi^2}{24}\right)\nu^2-\frac{1979}{36}\nu^3-\frac{265}{3888}\nu^4\right)\right\}, \end{aligned}$$

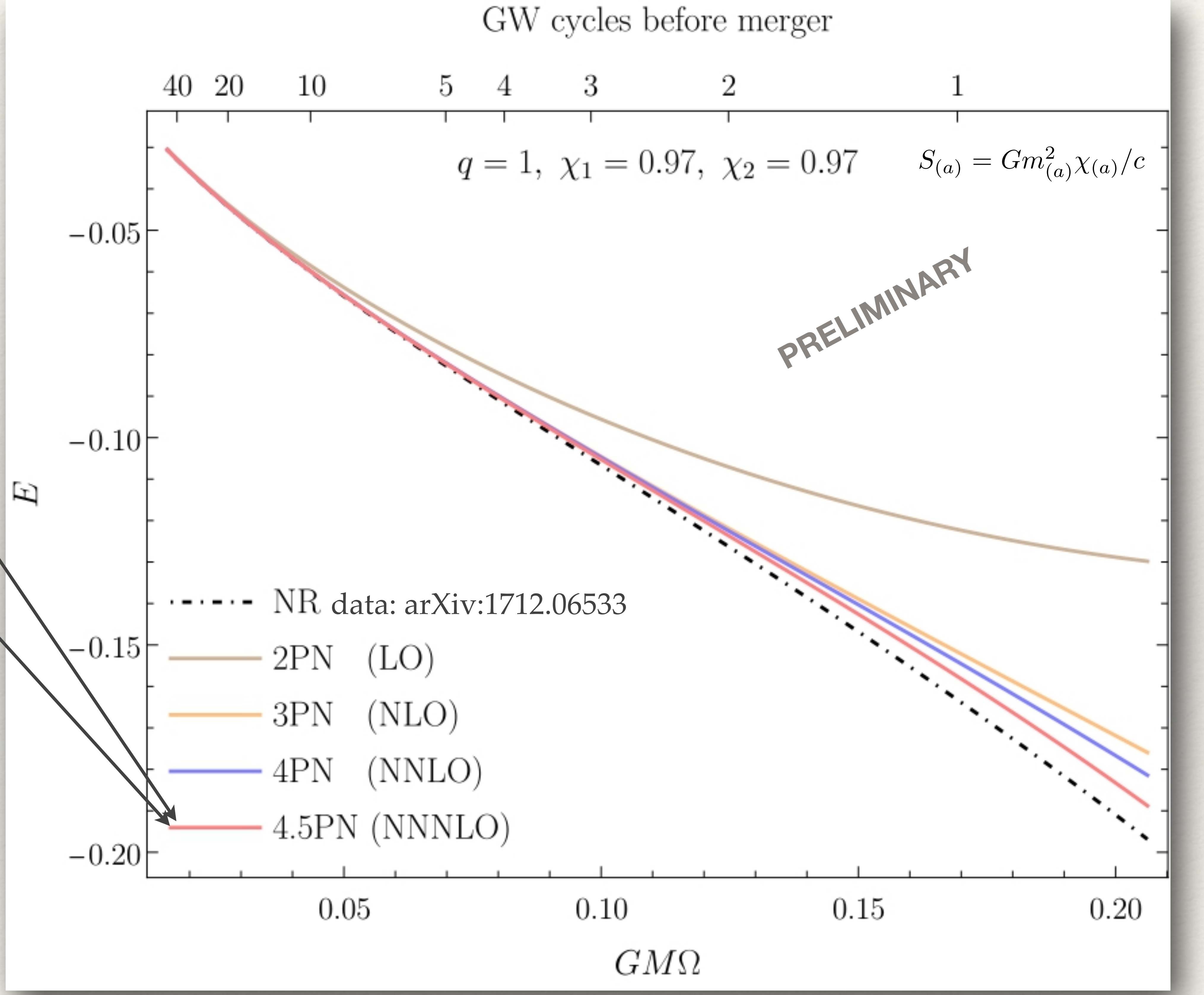
$$\begin{aligned} E_{\rm S1S2}(x,\widetilde{S}) = & \, \widetilde{S}_{(1)}\widetilde{S}_{(2)}\left\{x^3\left\{\nu\right\}+x^4\left\{\frac{5}{6}\nu+\frac{5}{18}\nu^2\right\}+x^5\left\{\frac{35}{8}\nu-\frac{1001}{72}\nu^2-\frac{371}{216}\nu^3\right\}\right. \\ & \left. + x^6\left\{\frac{243}{16}\nu-\left(\frac{2107}{16}-\frac{123}{32}\pi^2\right)\nu^2+\frac{147}{8}\nu^3+\frac{13}{16}\nu^4\right\}\right\}, \end{aligned}$$

$$\begin{aligned} E_{\rm S^2}(x,\widetilde{S}) = & \, \widetilde{S}^2_{(1)}\left\{x^4\left\{\frac{25}{18}\nu^2+\frac{1}{q}\left(-\frac{5}{2}\nu+\frac{5}{6}\nu^2\right)\right\}\right. \\ & + x^5\left\{\frac{10}{3}\nu^2-\frac{749}{108}\nu^3+\frac{1}{q}\left(-\frac{21}{4}\nu-\frac{7}{6}\nu^2-\frac{217}{36}\nu^3\right)\right\} \\ & + x^6\left\{\frac{1947}{112}\nu^2-\frac{48357}{560}\nu^3+\frac{159}{16}\nu^4\right. \\ & \left.\left. +\frac{1}{q}\left(-\frac{243}{16}\nu+\left(\frac{747}{16}-\frac{189\pi^2}{2048}\right)\nu^2-\frac{13731}{280}\nu^3+\frac{153}{16}\nu^4\right)\right\}\right\} \\ & + (1\leftrightarrow 2), \end{aligned}$$

$$\begin{aligned} E_{\rm ES^2}(x,\widetilde{S}) = & \left(C^{(0)}_{\rm ES^2}\right)_{(1)}\widetilde{S}^2_{(1)}\left\{x^3\left\{\frac{1}{q}\frac{\nu}{2}\right\}+x^4\left\{\frac{5}{3}\nu^2+\frac{1}{q}\left(\frac{5}{4}\nu+\frac{5}{4}\nu^2\right)\right\}\right. \\ & + x^5\left\{\frac{31}{4}\nu^2-\frac{35}{18}\nu^3+\frac{1}{q}\left(\frac{63}{16}\nu+\frac{77}{48}\nu^2-\frac{91}{48}\nu^3\right)\right\} \\ & + x^6\left\{\frac{789}{28}\nu^2-\frac{156}{7}\nu^3+\frac{5}{8}\nu^4\right. \\ & \left.\left. +\frac{1}{q}\left(\frac{405}{32}\nu+\left(\frac{3747\pi^2}{2048}-\frac{2389}{32}\right)\nu^2-\frac{555}{56}\nu^3+\frac{21}{32}\nu^4\right)\right\}\right\} \\ & + (1\leftrightarrow 2), \end{aligned}$$

$$E_{\rm E^2}(x,\widetilde{S}) = \left(C^{(2)}_{\rm E^2}\right)_{(1)}\widetilde{S}^2_{(1)}x^6\left\{9\nu^3\left(1+\frac{1}{q}\right)\right\}+(1\leftrightarrow 2),$$

$$E_{\rm E^2S^2}(x,\widetilde{S}) = \left(C^{(0)}_{\rm E^2S^2}\right)_{(1)}\widetilde{S}^2_{(1)}x^6\left\{\frac{3\nu^3}{2}\left(1+\frac{1}{q}\right)\right\}+(1\leftrightarrow 2).$$



[credit: Patil]

In agreement with:
Antonelli et al. (2020)
Kim et al. (2022)

Conclusion

- GW Astronomomy: a growing research field, where accuracy is not an option
- Compact objects evolution can benefit of the interplay between Cosmology, Astrophysics, and High-Energy Theoretical Physics
- Remarkable combination of traditional methods developed for the GR two-body problem and methods developed for elementary particle scattering to improve the GW waveforms modelling
- Under a diagrammatic viewpoint, Gravity is not so different from the other Fundamental Interactions

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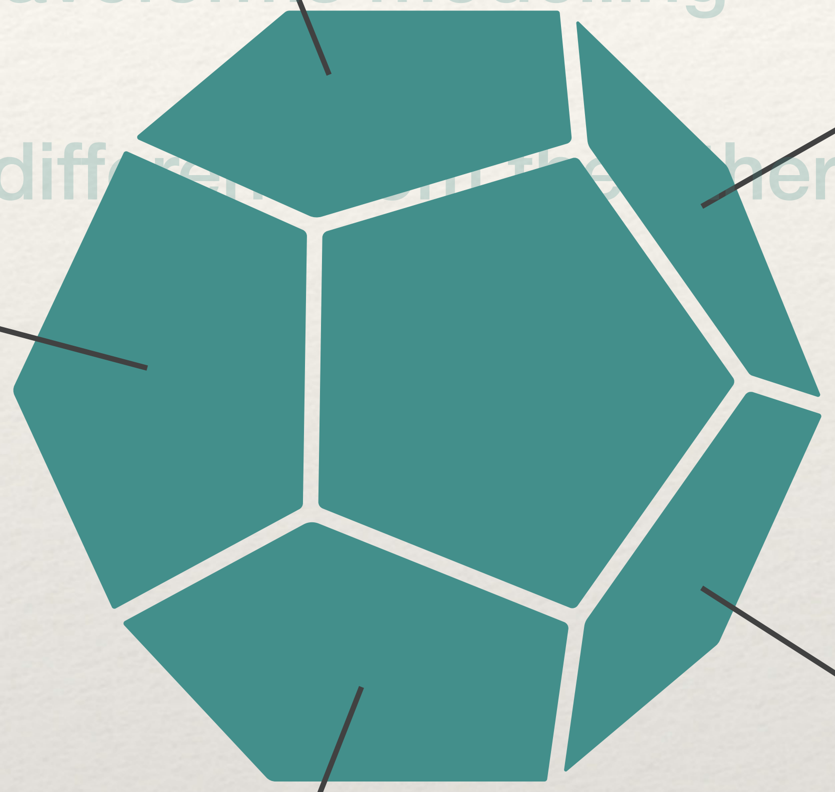
EFT - NRGR
- PN & PM EFT
- In-in formalism
- NRQFT
- HQEFT

GR-Techniques
- PN & PM
- Self-Force
- EOB
- Numerical Relativity
- Tutti Frutti

Amplitudes
- Unitarity-based methods
- Double-copy & BCJ relations
- Higher-spin
- Classical Scattering



Multiloop Techniques
- IBPs
- Difference & Differential Equations
- Theory of Special Functions
- High Precision arithmetics and Finite Fields
- Numerical Integration
- Asymptotic expansions

Observable-based methods
- Eikonal approach
- Inclusive & differential formalisms
- Radial action
- S-matrix

Definition. *Physics is a part of mathematics devoted to the calculation of integrals of the form $\int g(x)e^{f(x)}dx$. Different branches of physics are distinguished by the range of the variable x and by the names used for $f(x)$, $g(x)$ and for the integral.* [...]

Schwarz, Shapiro (2018)

Of course this is a joke, physics is not a part of mathematics. However, it is true that the main mathematical problem of physics is the calculation of integrals of the form

$$I(g) = \int g(x)e^{-f(x)}dx$$

[...] If f can be represented as $f_0 + \lambda V$ where f_0 is a negative quadratic form, then the integral $\int g(x)e^{f(x)}\, dx$ can be calculated in the framework of perturbation theory with respect to the formal parameter λ. We will fix f and consider the integral as a functional $I(g)$ taking values in $\mathbb{R}[[\lambda]]$. It is easy to derive from the relation

$$\int \partial_a(h(x)e^{f(x)})dx = 0$$

that the functional $I(g)$ vanishes in the case when g has the form

$$g = \partial_a h + (\partial_a f)h.$$

▸ Addressing a common math problem might be useful to make progress in different disciplines