

Constraining Lorentz Invariance Violation with Future Long-Baseline Experiments

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Introduction

Lorentz Invariance Violation (LIV) is a Planck-scale ($M_P \sim 10^{19}$ GeV) effect that can be studied at low energies using the so-called Standard Model Extension Lagrangian given by,

$$\mathcal{L}' \supseteq \frac{\lambda}{(M_P)^k} \langle T \rangle \bar{\psi} \Gamma(i\partial)^k \psi + h.c.$$

$M_P \rightarrow$ Planck-mass (10^{19} GeV)

$k = 0 \rightarrow$ CPT-conserving LIV parameters

$k = 1 \rightarrow$ CPT-violating LIV parameters

$$\mathcal{L}_{\text{LIV}} = -\frac{1}{2} [a_{\alpha\beta}^\mu \bar{\psi}_\alpha \gamma_\mu P_L \psi_\beta - i c_{\alpha\beta}^{\mu\nu} \bar{\psi}_\alpha \gamma_\mu \partial_\nu P_L \psi_\beta] + h.c.$$

We probe the time-like components ($\mu, \nu = 0$) of both the CPT-violating and CPT-conserving LIV parameters with standalone DUNE, Hyper-K, and their combination.

$$(a_L^0)_{\alpha\beta} \equiv a_{\alpha\beta}$$

$$(c_L^{00})_{\alpha\beta} \equiv c_{\alpha\beta}$$

Neutrino propagation with LIV

$$\mathcal{H}_{\text{eff}} = \underbrace{\frac{1}{2E} U \begin{pmatrix} 0 & 0 & 0 \\ 0 & \Delta m_{21}^2 & 0 \\ 0 & 0 & \Delta m_{31}^2 \end{pmatrix} U^\dagger}_{H_{\text{vac}}} + \underbrace{\sqrt{2} G_F N_e \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}}_{H_{\text{SI}}}$$

$$+ \underbrace{\begin{pmatrix} a_{ee} & a_{e\mu} & a_{e\tau} \\ a_{e\mu}^* & a_{\mu\mu} & a_{\mu\tau} \\ a_{e\tau}^* & a_{\mu\tau}^* & a_{\tau\tau} \end{pmatrix}}_{\text{CPT-violating}} - \underbrace{\frac{4}{3} E \begin{pmatrix} c_{ee} & c_{e\mu} & c_{e\tau} \\ c_{e\mu}^* & c_{\mu\mu} & c_{\mu\tau} \\ c_{e\tau}^* & c_{\mu\tau}^* & c_{\tau\tau} \end{pmatrix}}_{\text{CPT-conserving}}$$

$$H_{\text{LIV}}$$

$$P_{\mu e} \simeq P_{\mu e}(\text{SI}) + P_{\mu e}(a_{e\beta}/c_{e\beta}) \quad \& \quad P_{\mu\mu} \simeq P_{\mu\mu}(\text{SI}) + P_{\mu\mu}(a_{\mu\tau}/c_{\mu\tau})$$

$P_{\mu e}(\text{SI}), P_{\mu\mu}(\text{SI}) \rightarrow$ probabilities in presence of the standard interactions only.

$$P_{\mu e}(a_{e\beta}) \simeq 2|a_{e\beta}| L \sin \theta_{13} \sin 2\theta_{23} \sin \Delta [\mathbb{Z}_{e\beta} \sin(\delta_{\text{CP}} + \phi_{e\beta}) + \mathbb{W}_{e\beta} \cos(\delta_{\text{CP}} + \phi_{e\beta})]$$

$$P_{\mu e}(c_{e\beta}) \simeq \frac{-8}{3} |c_{e\beta}| L \sin \theta_{13} \sin 2\theta_{23} \sin \Delta [\mathbb{Z}_{e\beta} \sin(\delta_{\text{CP}} + \phi_{e\beta}) + \mathbb{W}_{e\beta} \cos(\delta_{\text{CP}} + \phi_{e\beta})]$$

$$P_{\mu\mu}(a_{\mu\tau}/c_{\mu\tau}) = \frac{\sin^2 2\theta_{23}}{2} [2 \sin^2 \theta_{13} \Delta - \mathbb{S}] \sin 2\Delta$$

Notations

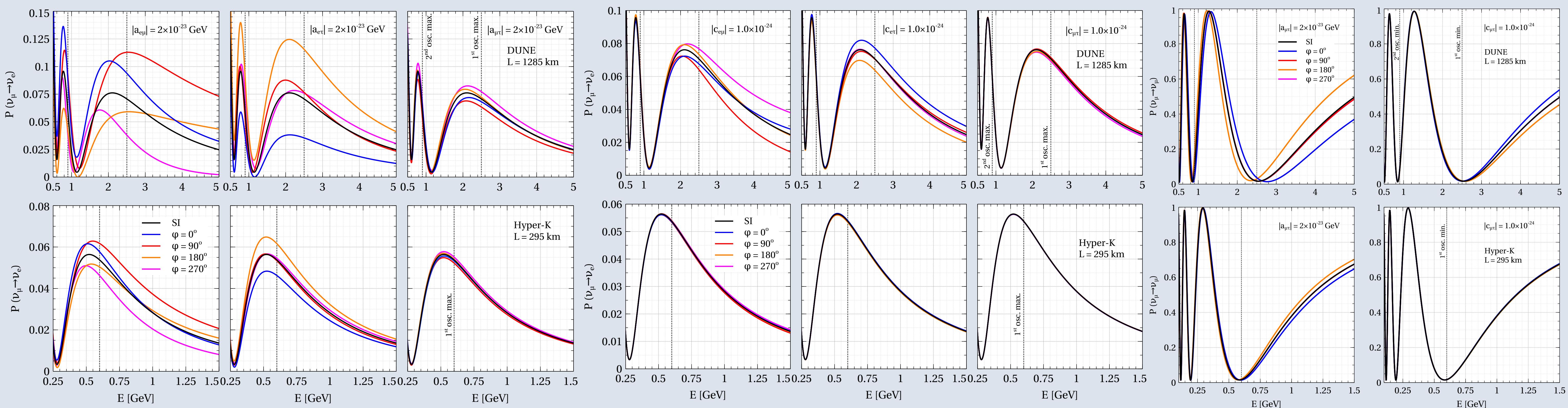
$$\mathbb{Z}_{e\beta} = \begin{cases} -c_{23} \sin \Delta, & \text{if } \beta = \mu. \\ s_{23} \sin \Delta, & \text{if } \beta = \tau. \end{cases}$$

$$\mathbb{W}_{e\beta} = \begin{cases} c_{23} \left(\frac{s_{23}^2 \sin \Delta}{c_{23}^2 \Delta} + \cos \Delta \right), & \text{if } \beta = \mu. \\ s_{23} \left(\frac{\sin \Delta}{\Delta} - \cos \Delta \right), & \text{if } \beta = \tau. \end{cases}$$

$$\mathbb{S}(a_{\mu\tau}) = 2 L \sin 2\theta_{23} |a_{\mu\tau}| \times \cos \phi_{\mu\tau}$$

$$\mathbb{S}(c_{\mu\tau}) = -\frac{8}{3} L \sin 2\theta_{23} |c_{\mu\tau}| \times \cos \phi_{\mu\tau}.$$

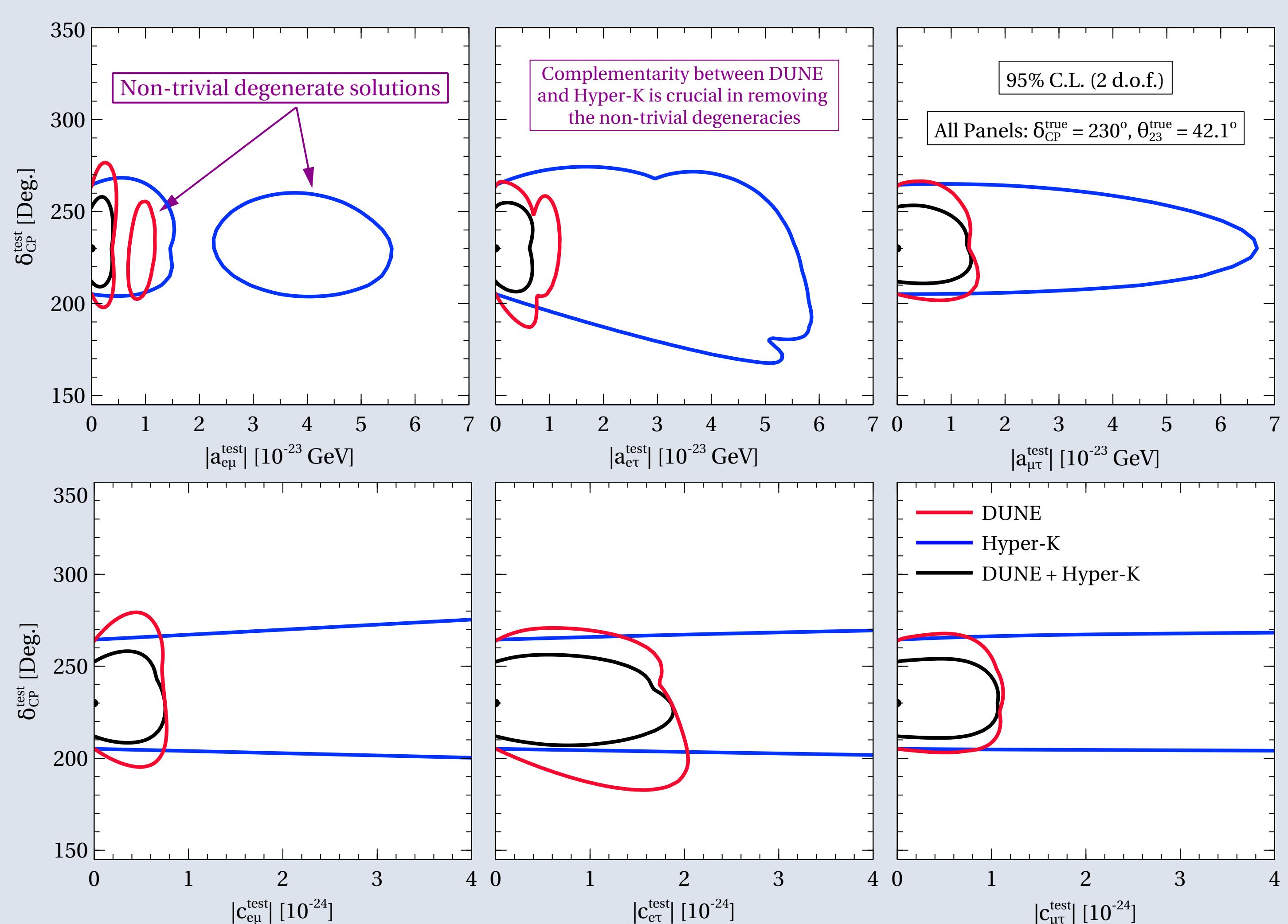
Neutrino oscillation probability in the presence of LIV



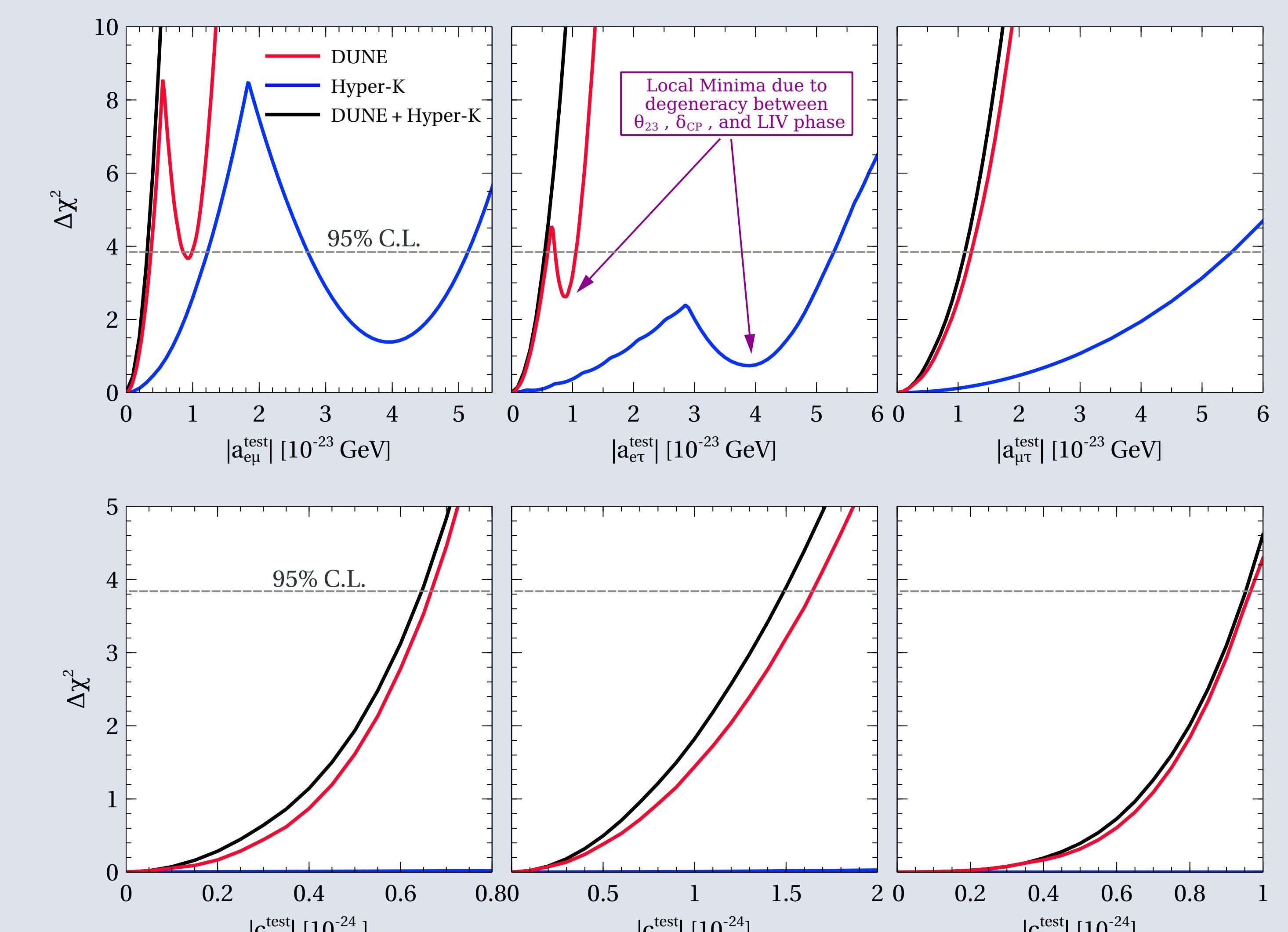
- Appearance probability is affected maximally by $a_{e\mu}$ ($c_{e\mu}$) followed by $a_{e\tau}$ ($c_{e\tau}$)
- $a_{\mu\tau}$ ($c_{\mu\tau}$) impacts the appearance probability marginally compared to $a_{e\mu}$ ($c_{e\mu}$) and $a_{e\tau}$ ($c_{e\tau}$)
- The effects of LIV phases are significant in DUNE compared to Hyper-K

- $a_{\mu\tau}$ ($c_{\mu\tau}$) mostly affects the disappearance probability
- The effects of $c_{\alpha\beta}$ are opposite in comparison to those of $a_{\alpha\beta}$
- Hyper-K is almost blind to the CPT-conserving LIV parameters

Correlations in $(\delta_{\text{CP}} - |a_{\alpha\beta}|/|c_{\alpha\beta}|)$ plane



Constraints on the LIV parameters



95% C.L. bounds

| | DUNE | Hyper-K | DUNE+Hyper-K | T2K+NOνA |
|--|--------|---------|--------------|----------|
| $ a_{e\mu} [10^{-23} \text{ GeV}]$ | < 1.0 | < 5.15 | < 0.32 | < 6.1 |
| $ a_{e\tau} [10^{-23} \text{ GeV}]$ | < 1.05 | < 5.3 | < 0.55 | < 7.0 |
| $ a_{\mu\tau} [10^{-23} \text{ GeV}]$ | < 1.26 | < 5.5 | < 1.1 | < 8.3 |
| $ c_{e\mu} [10^{-24}]$ | < 0.66 | < 17.1 | < 0.64 | < 11.0 |
| $ c_{e\tau} [10^{-24}]$ | < 1.65 | < 71.1 | < 1.49 | < 37.5 |
| $ c_{\mu\tau} [10^{-24}]$ | < 0.97 | < 42.4 | < 0.95 | < 29.0 |

Conclusions

- Due to longer baseline and access to multi-GeV energies, DUNE has a better reach in probing both CPT-violating and CPT-conserving LIV parameters
- Hyper-K, which mostly deals with sub-GeV neutrinos, is almost insensitive to the CPT-conserving LIV parameters
- The degeneracy between θ_{23} , δ_{CP} , and the LIV phases lead to the deterioration of the bounds for the individual setups; however, when we add the data from both, these degeneracies disappear

References

Agarwalla et al., Constraining Lorentz invariance violation with next-generation long-baseline experiments, arXiv: 2302.12005, JHEP 07 (2023) 216

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