LHC sensitivity to BSM triple Higgs couplings in the 2HDM

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DESY.

2HDM: The model

[R. Santos: arXiv:hep-ph/9701257 '97]

CP conserving 2HDM with two complex doublets: $\Phi_1 = \begin{pmatrix} \phi_1^+ \\ \frac{v_1+\rho_1+i\eta_1}{\sqrt{2}} \end{pmatrix}, \Phi_2 = \begin{pmatrix} \phi_2^+ \\ \frac{v_2+\rho_2+i\eta_2}{\sqrt{2}} \end{pmatrix}$

Softly broken $\mathbb{Z}_2$ symmetry ($\Phi_1 \to \Phi_1$; $\Phi_2 \to \Phi_2$) entails 4 Yukawa types (Types I and II were analyzed).

h (m_h = 125 GeV), **H** - CP even, **A** - CP odd, **H^+**, **H^-**

Pontental:

$$V_{2HDM} = m_{11}^2(\Phi_1^\dagger\Phi_1) + m_{22}^2(\Phi_2^\dagger\Phi_2) - m_{12}^2(\Phi_1^\dagger\Phi_2 + \Phi_2^\dagger\Phi_1) + \frac{\lambda_1}{2}(\Phi_1^\dagger\Phi_1)^2 + \frac{\lambda_2}{2}(\Phi_2^\dagger\Phi_2)^2 +$$
$$\lambda_3(\Phi_1^\dagger\Phi_1)(\Phi_2^\dagger\Phi_2) + \lambda_4(\Phi_1^\dagger\Phi_2)(\Phi_2^\dagger\Phi_1) + \frac{\lambda_5}{2}((\Phi_1^\dagger\Phi_2)^2 + (\Phi_2^\dagger\Phi_1)^2)$$

Free parameters: $m_h,\ m_A,\ m_H,\ m_{H^\pm},\ m_{12}^2, v,\ \cos(\beta-\alpha),\ \tan\beta$

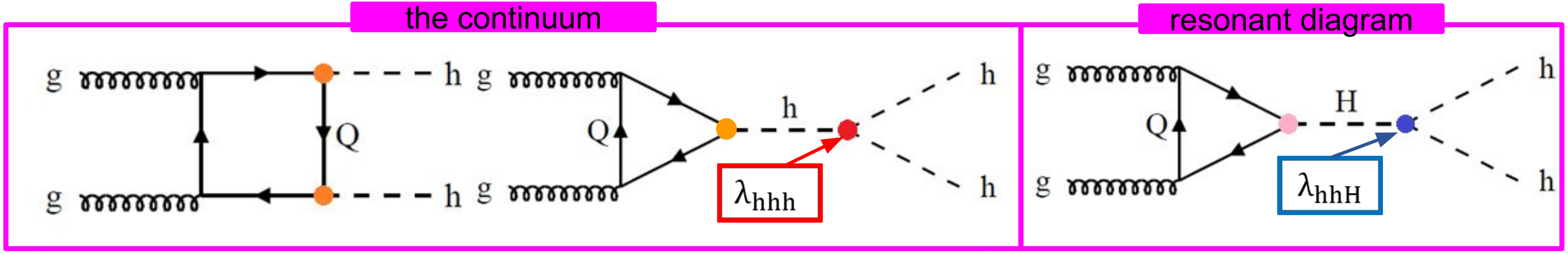
$\tan\beta = v_2/v_1$
$v^2 = v_1^2 + v_2^2 \sim (246\ \text{GeV})^2$

→ couplings to fermions and gauge bosons:

These lead to different phenomenology w.r.t the SM but also the contribution of the heavy Higgses in the loops.

Di-Higgs production ($gg \to hh$)

Dominant process at the LHC → Gluon Fusion



Diagrams that exist in the SM: They have a negative interference

Diagrams that are sensitive to triple Higgs couplings

$\sigma_{SM} \sim$ 38 fb at NLO

[Abouabid et al. arXiv:2112.12515 '21]

To obtain the cross section prediction for this process we use a modified version of the code **HPAIR** that contains the **2HDM** model.

→ We will study the **invariant mass distribution** of two 125 GeV in the final state.

Triple Higgs Couplings

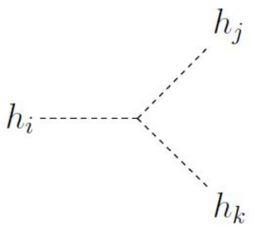
Can have **large deviations** from SM predictions in BSM:

$$\kappa_\lambda = \lambda_{hhh}/\,\lambda_{hhh}^{SM}$$

$$\lambda_{hhh}^{SM} = \frac{m_h^2}{2v^2} \cong 0.129$$

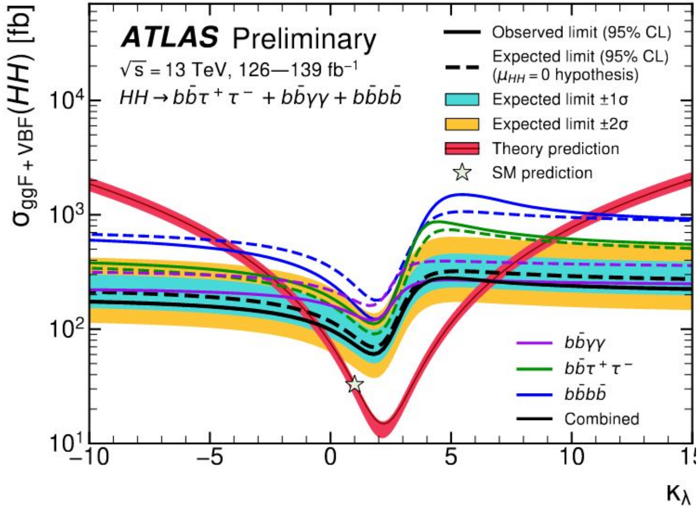
[$-0.4 < \kappa_\lambda < 6.3$] (95% CL at LHC Run II)

Notation:



$$= -i\; v\; n!\; \lambda_{h_i h_j h_k}$$

n = number of identical Higgses



[ATLAS-CONF-2022- 050]

Large λ's were found: far from the alignment limit and/or for large masses.

Type I

κ_λ = [-0.5, 1.3]

λ_{hhH} = [-1.7, 1.6]

[F. Arco et al. arXiv2003.12684 '22]

Invariant mass distribution: effects of deviations in κ_λ

BP: Type I, $\cos(\beta-\alpha)=0.1$, $\tan\beta=10$, $m_{12}^2=m_H^2\cos^2\alpha/\tan\beta$, $m_H=m_A=m_{H^\pm}$

Prediction for BSM couplings:

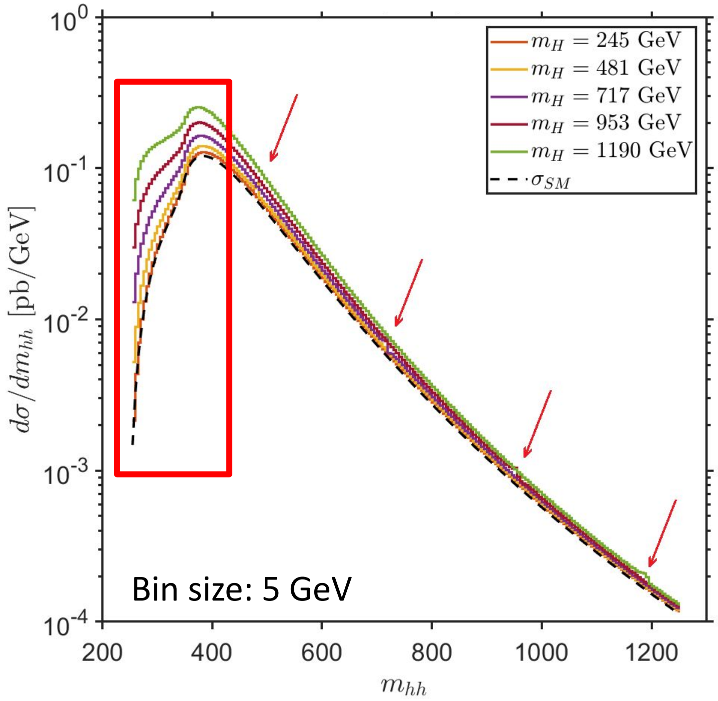
$\kappa_\lambda = 0.97,\ \lambda_{hhH} = 0.05$

$\kappa_\lambda = 0.85,\ \lambda_{hhH} = 0.19$

$\kappa_\lambda = 0.67,\ \lambda_{hhH} = 0.42$

$\kappa_\lambda = 0.41,\ \lambda_{hhH} = 0.74$

$\kappa_\lambda = 0.08,\ \lambda_{hhH} = 1.15$



- Larger sensitivity to κ_λ in the low mhh region.

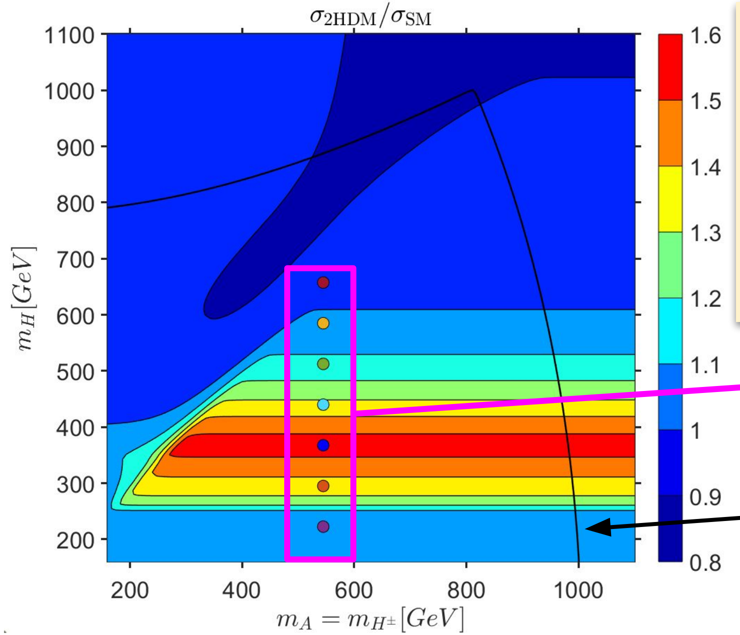
- Resonant contribution very suppressed due to very small top Yukawa $\xi^t_H \sim 10^{-4}$.

$\xi^t_H = \cos(\beta-\alpha) - \sin(\beta-\alpha)/\tan(\beta)$

Effect of the mass of the heavy Higgs

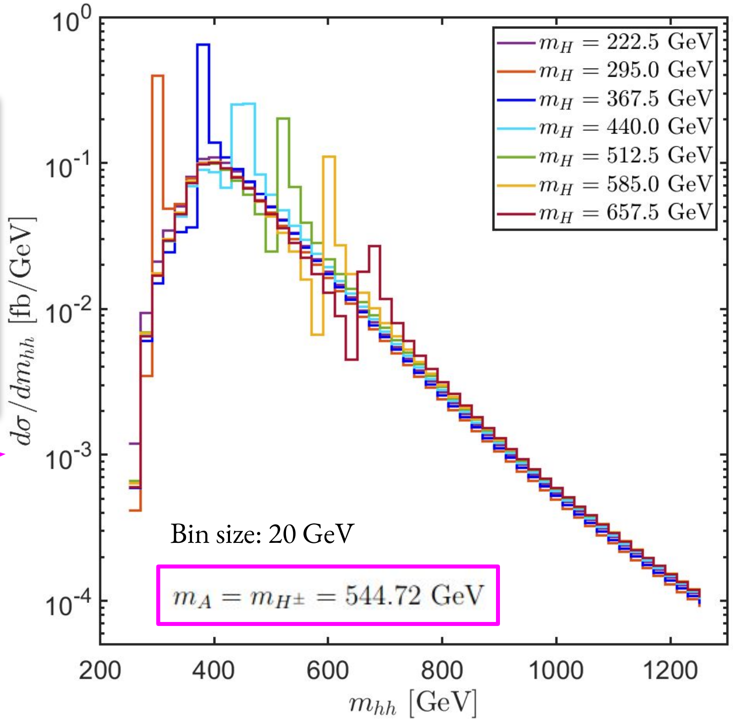
- We vary the mass of the heavy Higgs boson leaving the rest of the parameters of the model fixed.

BP: Type I, $\cos(\beta - \alpha) = 0.2$, $\tan\beta = 10$, $m_{12}^2 = m_H^2 \cos^2\alpha / \tan\beta$



Enhancement in the total cross section is resonance dominated. **Location** of the resonance is related to the mass of **H.**

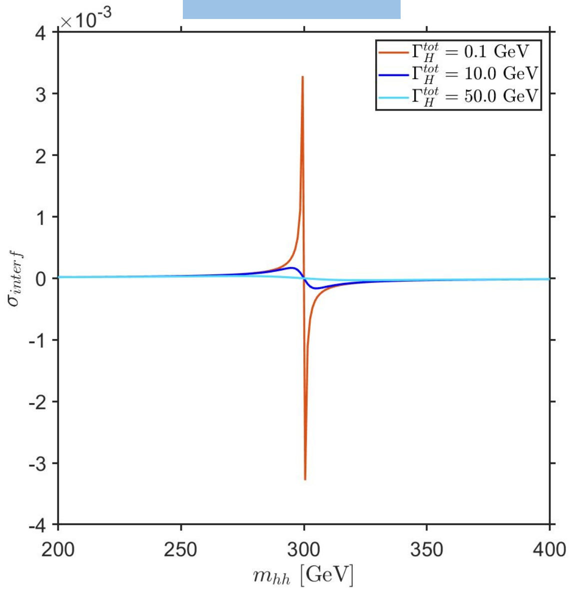
Allowed region inside the black contour.



Effect of the total decay width

$$\frac{1}{Q^2 - M^2_{h/H} + i\Gamma_{h/H} M_{h/H}}$$

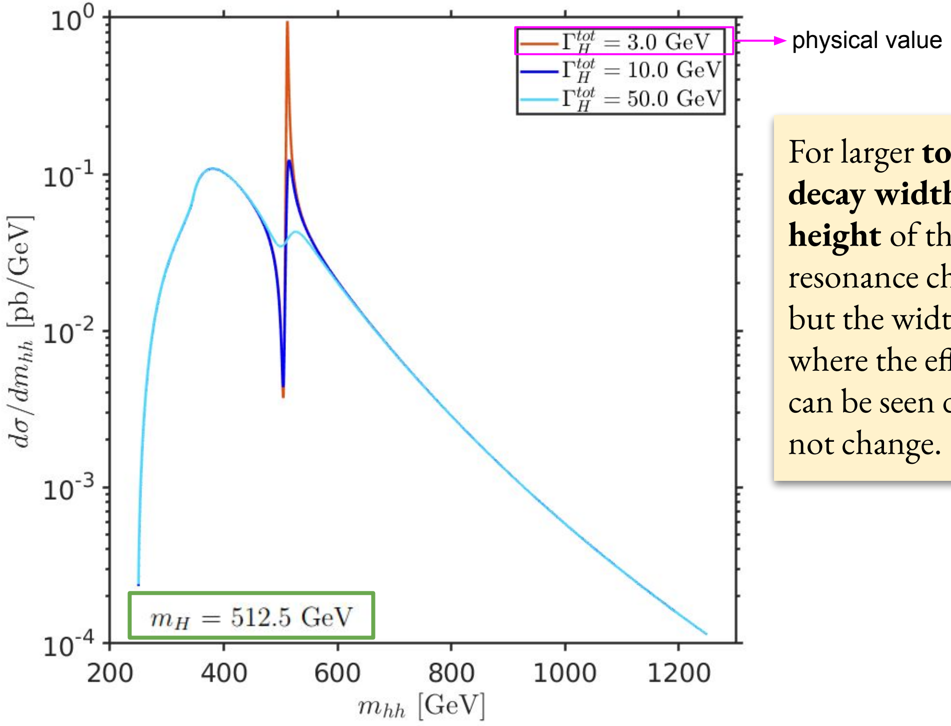
Toy model



$$\sigma_{\text{interf}} \propto \frac{Q^2 - m_H^2}{(Q^2 - m_H^2)^2 + m_H^2 \Gamma_H^2}$$

- For the green point of the previous benchmark plane

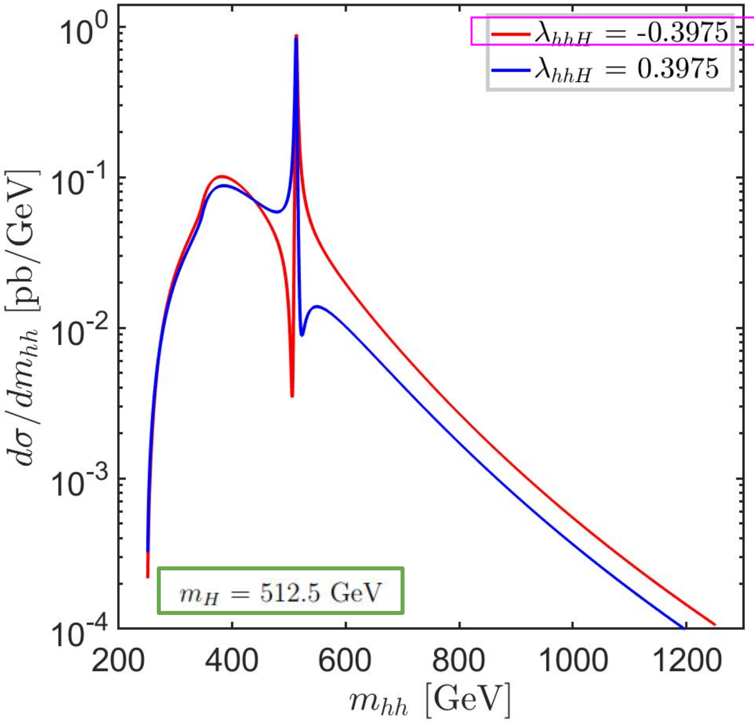
we artificially change the total decay width of the heavy Higgs H:



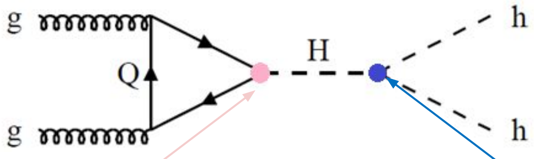
For larger **total decay widths** the **height** of the resonance changes but the width where the effect can be seen does not change.

Effect of the couplings

- What is the effect of the couplings involved in the resonant diagram on the invariant mass distributions ?



physical value



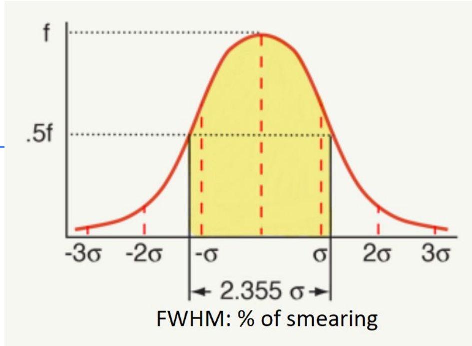
$\xi_H^t = \cos(\beta-\alpha) - \sin(\beta-\alpha)/\tan(\beta) = 0.104$

λ_{hhH}

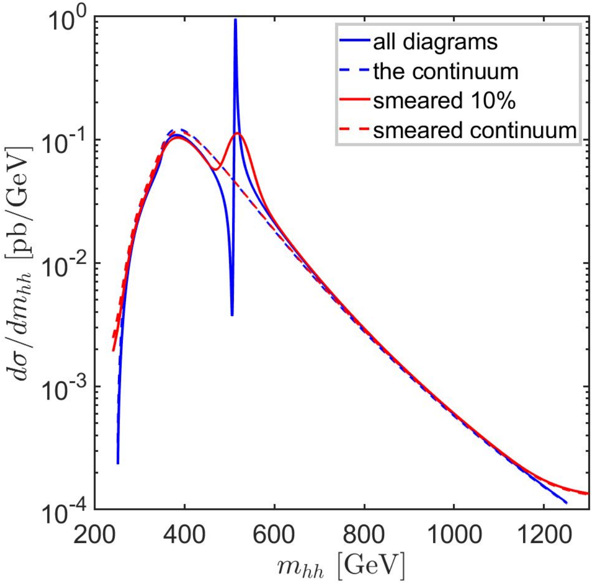
The **relative sign** of the top Yukawa and the BSM coupling to the heavy Higgs gives a **structure** to the resonance:

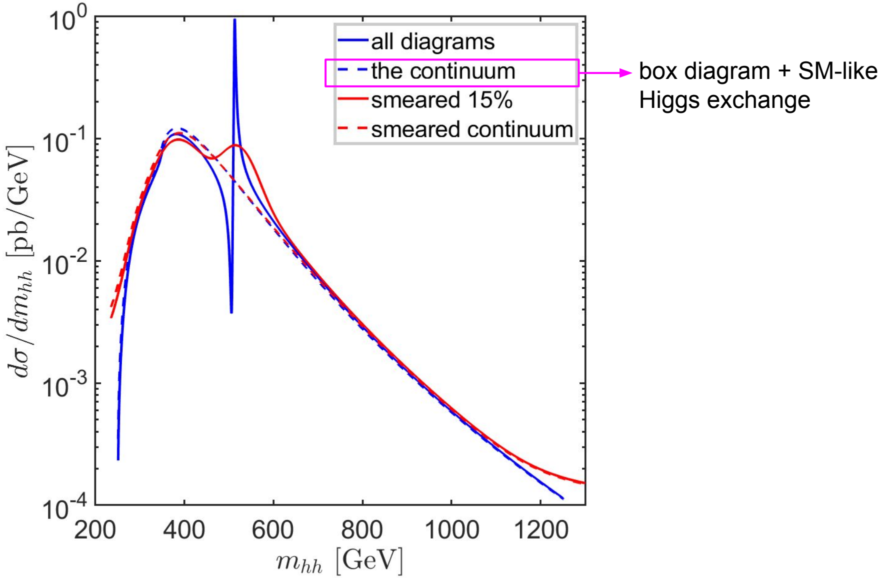
sign $(\lambda_{hhH} \cdot \xi_H^t)$	structure
+	peak-dip
-	dip-peak

Experimental challenges: smearing

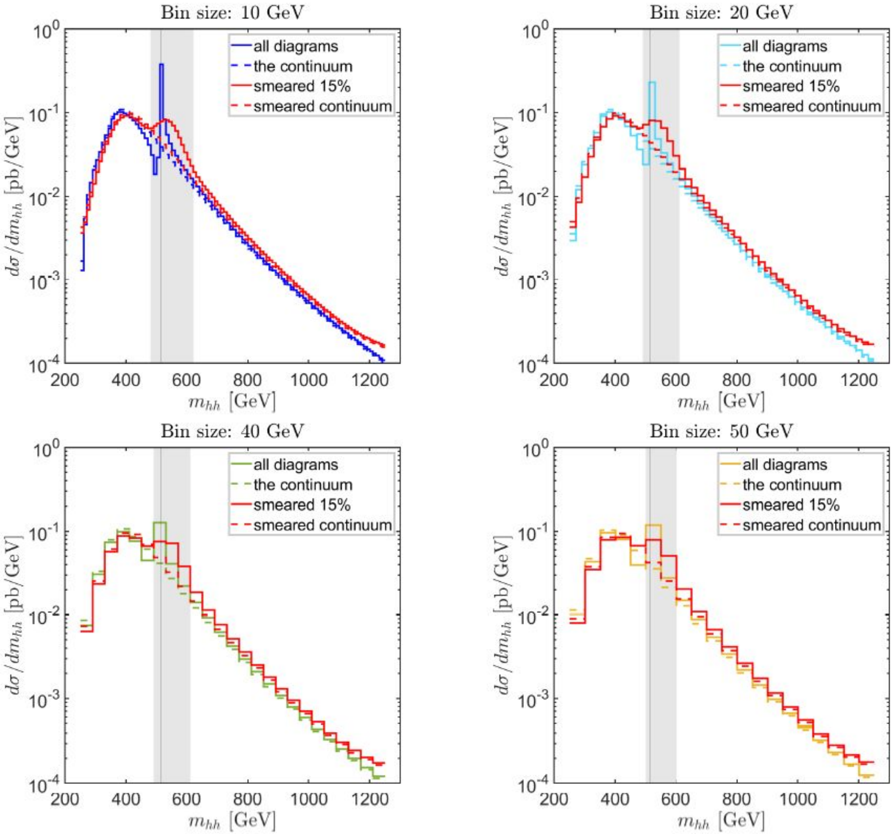


- Differential cross section measurements are affected by the finite resolution of particle detectors → observed spectrum is "**smeared**".
- We try to mimic this effect by artificially smearing the theoretical prediction introducing **Gaussian uncertainties** in the invariant mass.





Experimental challenges: binning (15 % smearing)



- We define a value for the 'significance of the signal' according to the excess of the number of events. Assuming: $\mathcal{L} = 6000 \text{ fb}^{-1}$.

$$R := \frac{\sum_i (N^R - N^C)}{\sqrt{\sum_i N^C}}$$

Bin size	R
10 GeV	3.79
20 GeV	3.99
40 GeV	3.87
50 GeV	4.04

- Window definition:

$(N^R - N^C) >$ (bin size)/50

→ Smearing dilutes more the resonance than binning.

Conclusion

Invariant mass distributions give information about:

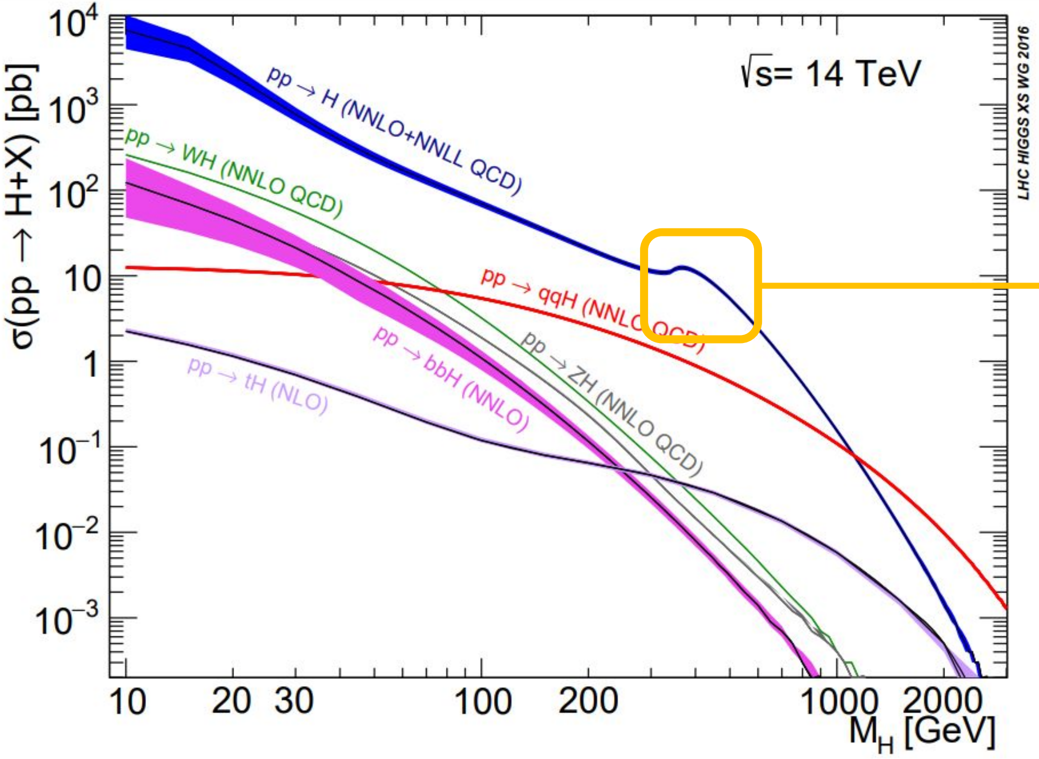
1. Deviations of κ_λ that can be seen in the low m_{hh} threshold,

2. **resonant production** that can be embedded in BSM models:

- **mass** of the intermediate Higgs boson $\longrightarrow$ **position** of the resonance.
- **total decay width** of the resonance $\longrightarrow$ **height** of the resonance.
- relative sign of the **couplings** $\longrightarrow$ **structure** of the resonance.

These effects may be (partially) washed out by experimental precision (**smearing**).

Backup: Single Higgs production

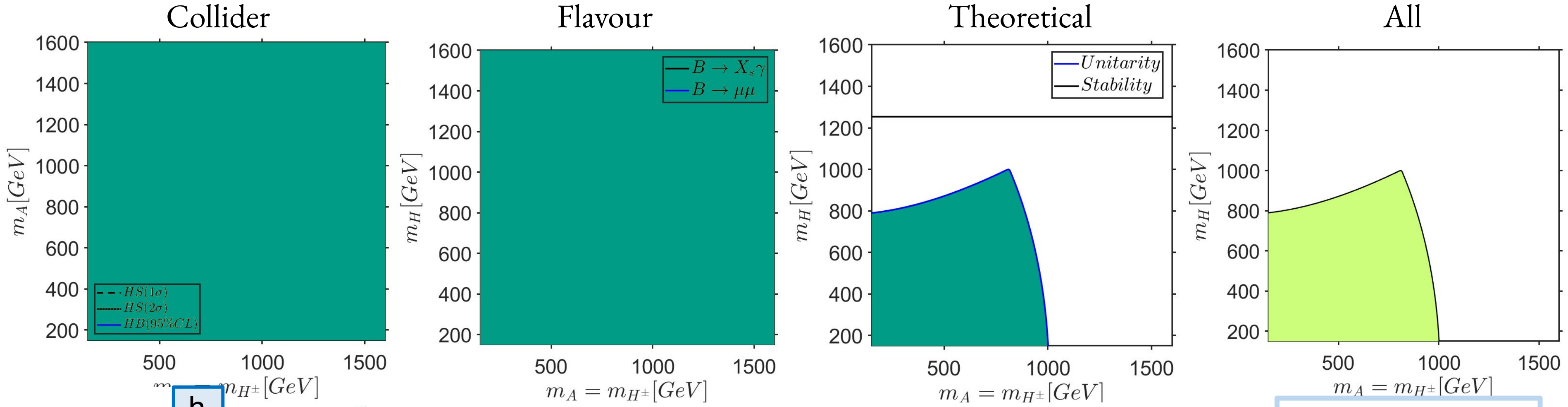


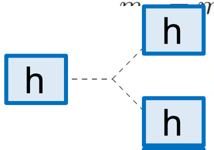
Top pair threshold ⟶ gives a hint on the results for Higgs pair production

[LHC Higgs Working Group: CERN Yellow Report 4]

Constraints and Feynman rules

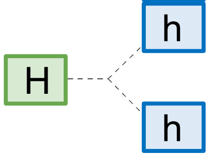
BP: Type I, $\cos(\beta-\alpha) = 0.2$, $\tan\beta = 10$, $m_{12}^2 = m_H^2 \cos^2\alpha / \tan\beta$





$$\lambda_{hhh} = \frac{1}{2v^2}\left\{m_h^2 s_{\beta-\alpha}^3 + \left(3m_h^2 - 2\bar{m}^2\right) c_{\beta-\alpha}^2 s_{\beta-\alpha} + 2\cot 2\beta \left(m_h^2 - \bar{m}^2\right) c_{\beta-\alpha}^3\right\}$$

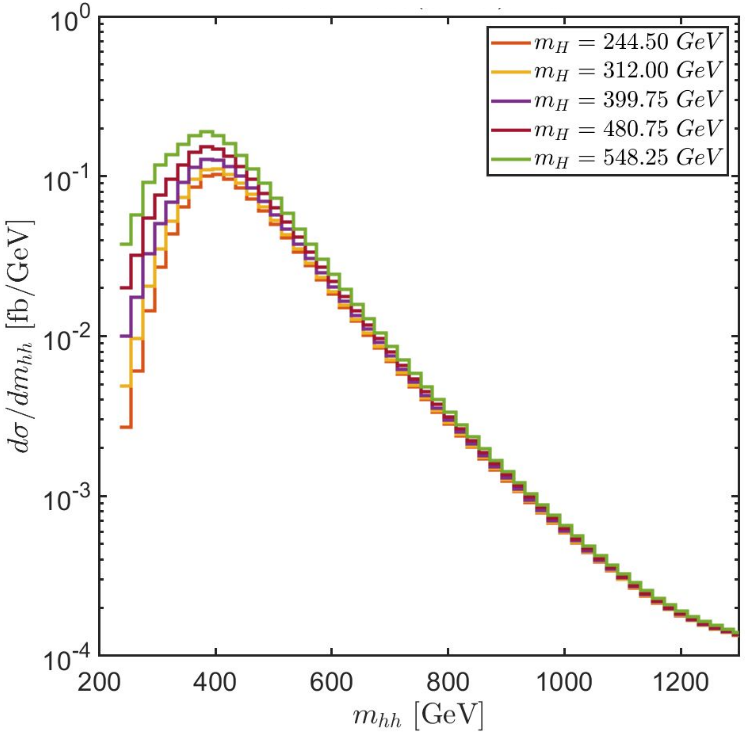
$$\bar{m}^2 = \frac{m_{12}^2}{sin(\beta)cos(\beta)}$$



$$\lambda_{hhH} = \frac{-c_{\beta-\alpha}}{2v^2}\left\{\left(2m_h^2 + m_H^2 - 4\bar{m}^2\right) s_{\beta-\alpha}^2 + 2\cot 2\beta \left(2m_h^2 + m_H^2 - 3\bar{m}^2\right) s_{\beta-\alpha} c_{\beta-\alpha} - \left(2m_h^2 + m_H^2 - 2\bar{m}^2\right) c_{\beta-\alpha}^2\right\}$$

Smearing applied on the invariant mass distributions

BP: Type I, $\cos(\beta-\alpha) = 0.1$, $\tan\beta = 10$, $m_{12}^2 = m_H^2 \cos^2\alpha / \tan\beta$, $m_H = m_A = m_{H^\pm}$



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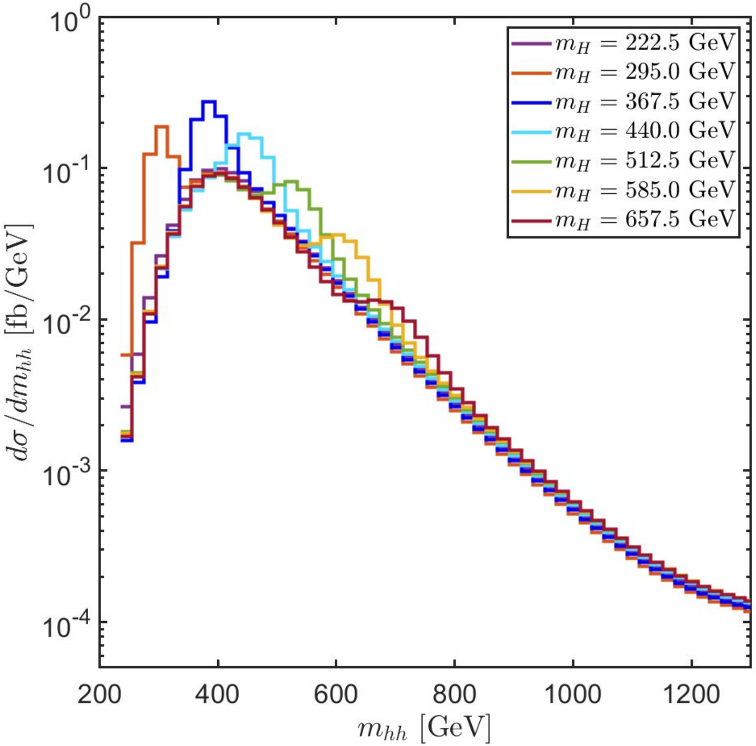
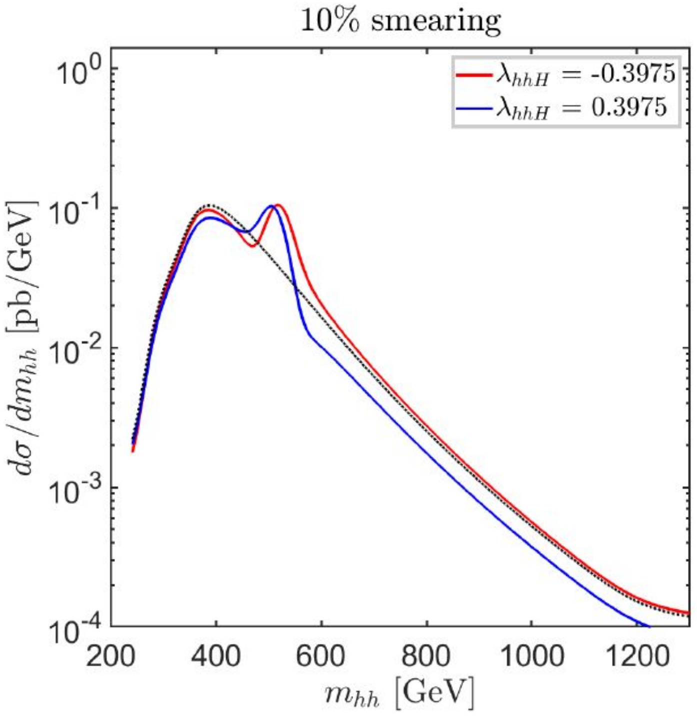
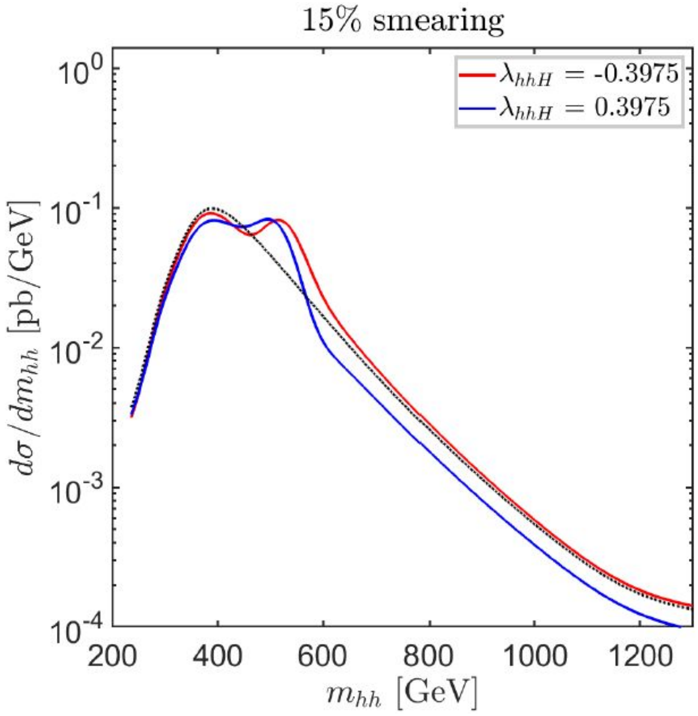


Smearing applied on the structure of the resonance





Binning applied on the structure of the resonance

For 15 % smearing: