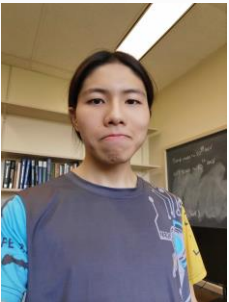


Phenomenology of wavelike vector dark matter nonminimally coupled to gravity

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@ PHENO 2023, University of Pittsburgh

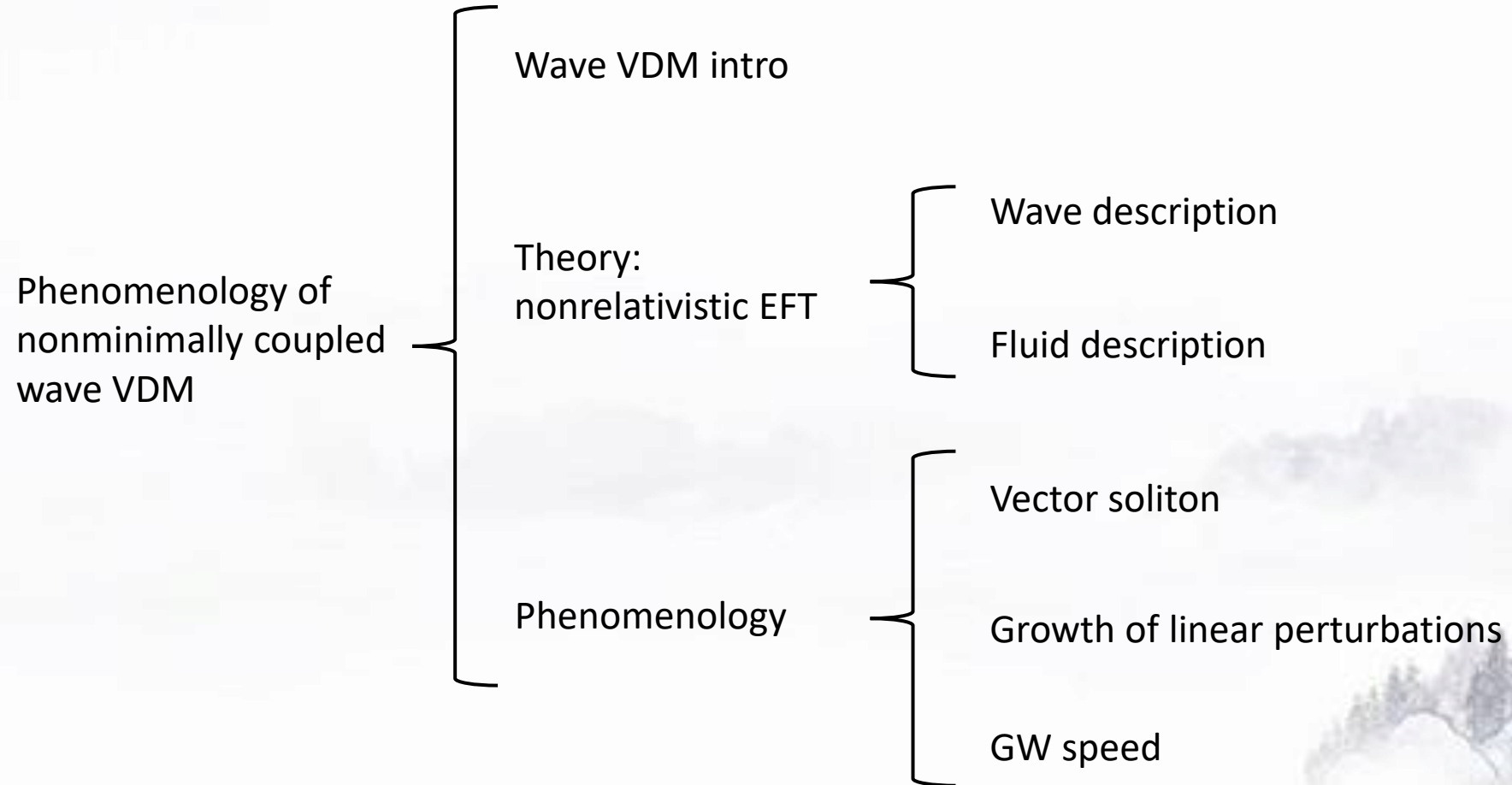
May 08, 2023

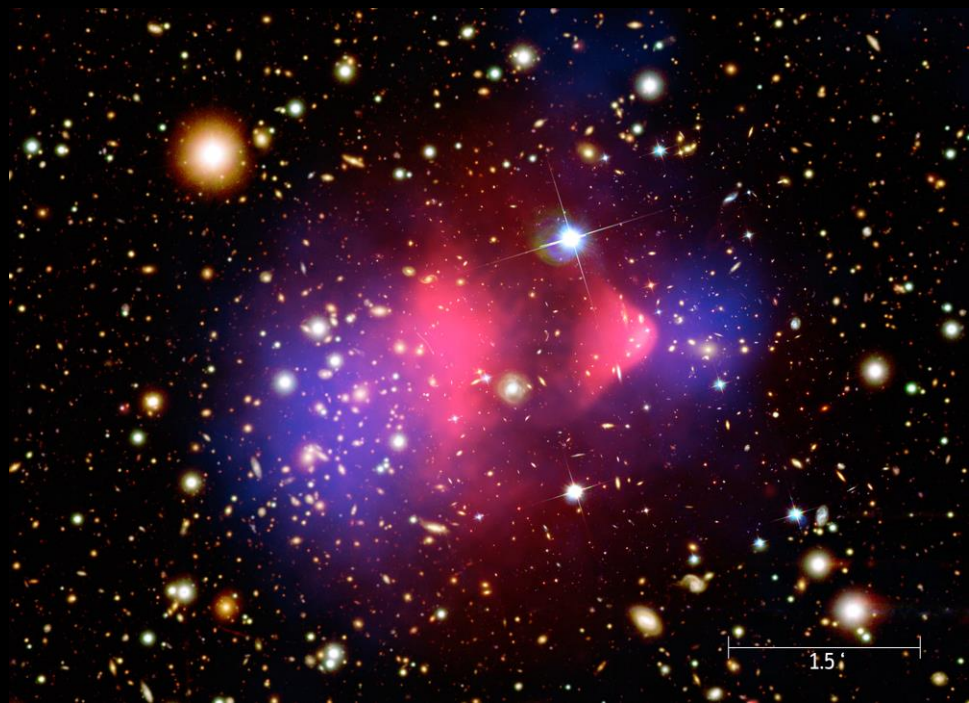
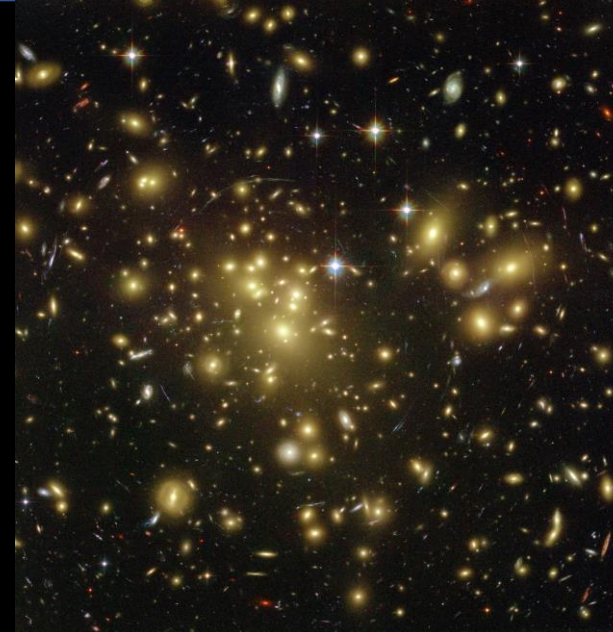
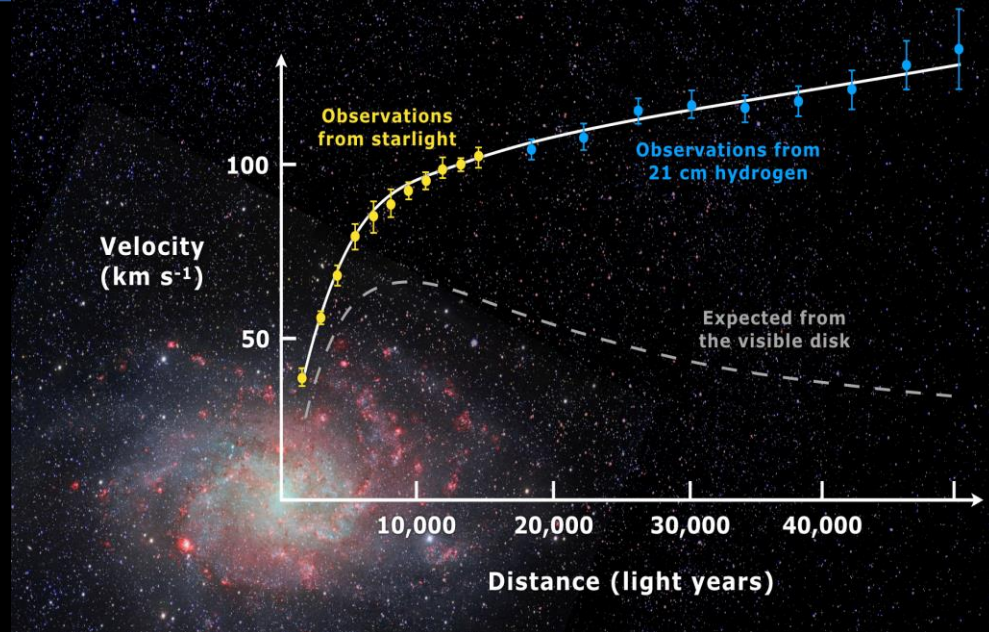


Collaborated with Siyang Ling (PhD student at Rice U).
The paper will be available in arXiv in a few hours.



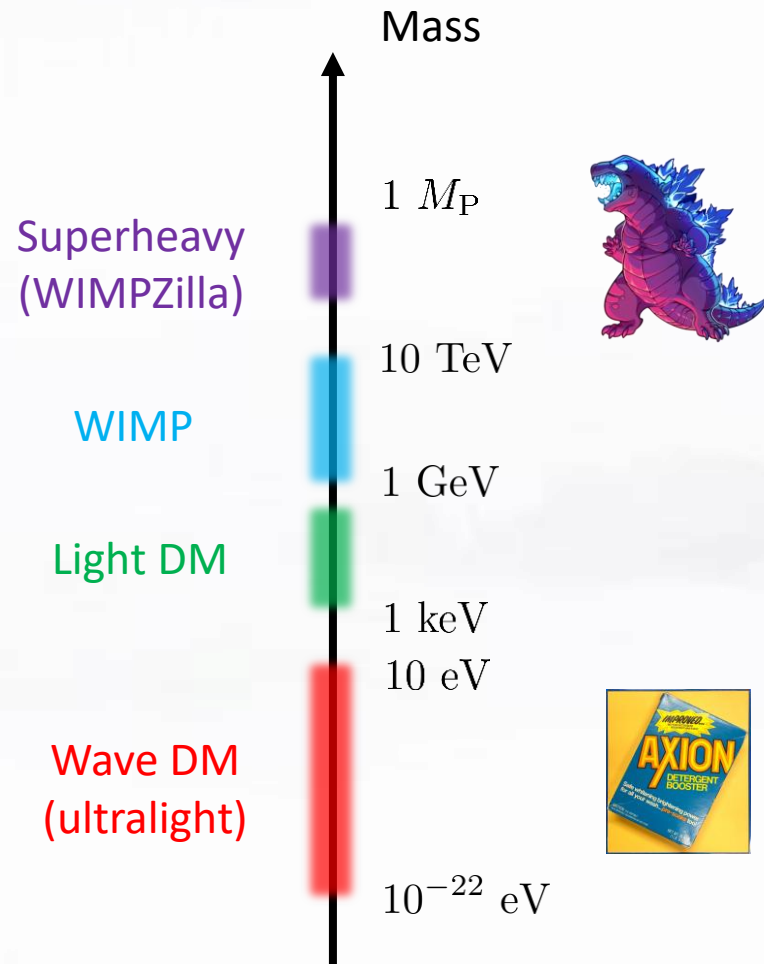
Talk structure





Rotation curves
 Low mass-to-light ratios
 Bullet cluster
 Matter power spectrum
 Galaxy cluster's virial velocity
 Gravitational lensing
 Supernovae type Ia
 Cosmic microwave background
 Baryon acoustic oscillations
 ...

DM mass landscape



Dark matter halo

Luminous matter

$$\Delta x \cdot \Delta p \sim h$$

Wave DM is composed of bosons

Structure is suppressed within de Broglie wavelength

$$\lambda_{\text{dB}} = \frac{2\pi}{mv} = 0.48 \text{ kpc} \left(\frac{10^{-22} \text{ eV}}{m} \right) \left(\frac{250 \text{ km/s}}{v} \right) \lesssim 1 \text{ kpc}$$
$$m \gtrsim 10^{-22} \text{ eV}$$

Number of particles in de Broglie volume λ_{dB}^3

$$N_{\text{dB}} \sim \left(\frac{34 \text{ eV}}{m} \right)^4 \left(\frac{250 \text{ km/s}}{v} \right)^3 \gg 1 \quad \left\{ \begin{array}{l} \text{Quantum effects are small} \\ m \ll 30 \text{ eV} \end{array} \right.$$

Wave DM cannot be fermions (Tremaine-Gunn bound)

$$n = \int \frac{d^3 p}{(2\pi)^3} f \lesssim \frac{g}{(2\pi)^3} \frac{4\pi}{3} (mv)^3 \sim \frac{\rho_\chi}{m}$$
$$m \gtrsim 20 \text{ eV}$$

Hui (2021)



VDM production mechanism

Nonminimal coupling
is needed

Kitajima & Nakayama (2023)
Nakayama (2019).

- Misalignment mechanism

-> Fixed polarization in cosmological scales

Graham, et al. (2016);
Kolb & Long (2021).

- Gravitational particle production

-> Natural avoidance of isocurvature constraints

Agrawal, et al. (2020);
Co, et al. (2019);
Bastero-Gil, et al. (2019).

- Scalar-vector oscillation

-> Stealing energy from axions/inflaton

Long & Wang (2019)

- Decay of topological defects

-> Longitudinal modes are preferred at the production,
but the subsequent evolution is complicated

Other motivations for nonminimal couplings:

- Effective theory
- Renormalization of gravity
- Rich phenomenology

For thermal freeze-out and freeze-in, see Barman, et al. (2022).

Theory

Higher-dimension operators,
suppressed by powers of X^2/M_{P}^2

$$S = \int d^4x \sqrt{-g} \left[\frac{M_{\text{P}}^2}{2} R - \frac{1}{4} X_{\mu\nu} X^{\mu\nu} - \frac{m^2}{2} X_\mu X^\mu + \frac{\xi_1}{2} R X_\mu X^\mu + \frac{\xi_2}{2} R^{\mu\nu} X_\mu X_\nu + \cdots \right]$$

Nonrelativistic EFT

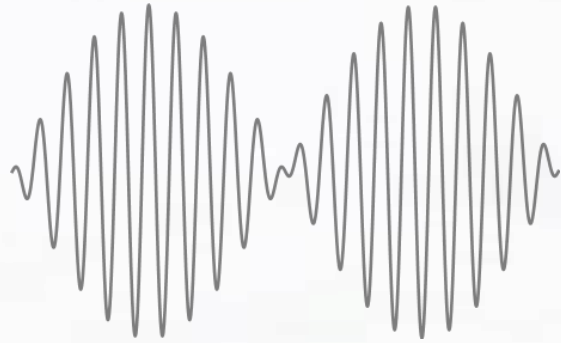
Wave description: Schroedinger-Poisson-Friedmann (SPF) equations
Convenient for VDM dynamics, simulations, quantization

Fluid description: Continuity-Euler-Friedmann equations
Convenient for structure formation

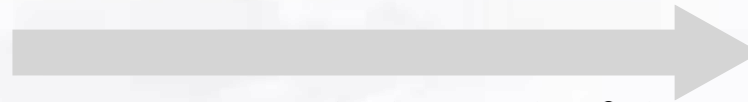
Nonrelativistic EFT

Vector field equations + Einstein equations

$$X_\mu(t, \mathbf{x}) = \frac{1}{\sqrt{2ma}} [e^{-imt} \psi_\mu(t, \mathbf{x}) + e^{imt} \psi_\mu^*(t, \mathbf{x})]$$

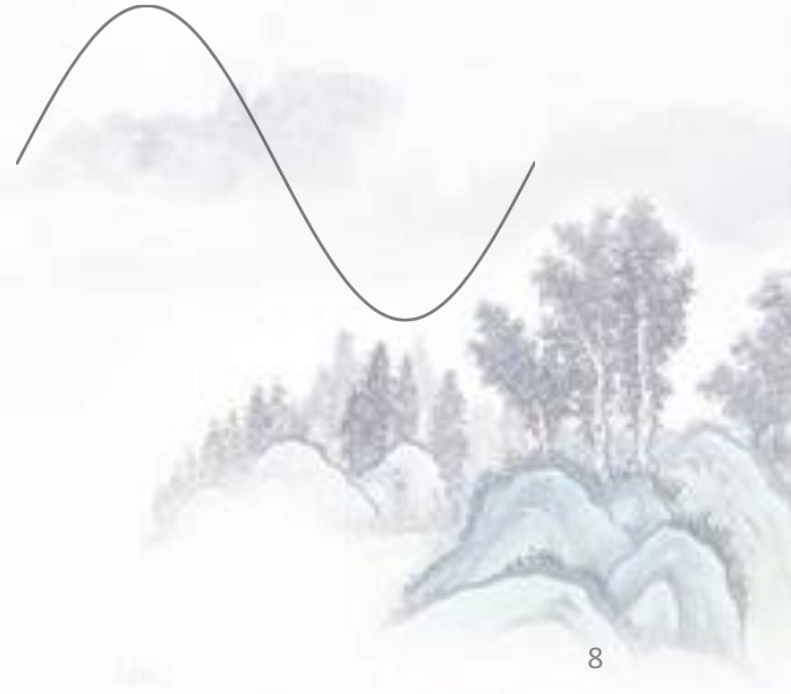


Integrate out fast modes



$$\epsilon_H \sim \frac{H}{m}, \quad \epsilon_t \sim \left| \frac{\partial_t}{m} \right|, \quad \epsilon_k \sim \left| \frac{\nabla^2}{m^2 a^2} \right|,$$
$$\epsilon_\psi \sim \frac{|\psi_i|}{a M_{\text{pl}} \sqrt{m}}, \quad \epsilon_g \sim \Phi, \quad \epsilon_\xi \sim \left| \frac{\xi R}{m^2} \right|.$$

Nonrelativistic EFT for slow modes



Wave description

Effective action:

$$S = \int d^4x \left[M_{\text{P}}^2 a (-3\dot{a}^2 + \Phi \nabla^2 \Phi - 6a\ddot{a}\Phi) \right. \\ \left. + \frac{1}{2} a^2 m \left| \psi_0 - \frac{i}{a^2 m} \nabla \cdot \psi \right|^2 + i \dot{\psi} \cdot \psi^* + \frac{1}{2a^2 m} (\nabla^2 \psi) \cdot \psi^* - m \Phi |\psi|^2 \right. \\ \left. + \frac{|\psi|^2}{2a^2 m} [2\xi_1 (\nabla^2 \Phi + 3\dot{a}^2 + 3a\ddot{a}) + \xi_2 (\nabla^2 \Phi + 2\dot{a}^2 + a\ddot{a})] \right] \times [1 + \mathcal{O}(\epsilon)]$$

Multicomponent
SPF equations:

$$i\partial_t \psi_i = -\frac{\nabla^2}{2ma^2} \psi_i + m\Phi_N \psi_i + 2m\Phi_\xi \psi_i - \frac{(2\xi_1 + \xi_2)\nabla^2 \Phi_\xi}{2ma^2} \psi_i , \\ \frac{\nabla^2}{a^2} \Phi_N = 4\pi G(\rho - \bar{\rho}) , \quad \Phi_\xi = -\frac{(2\xi_1 + \xi_2)2\pi G}{m^2} \rho , \quad \Phi = \Phi_N + \Phi_\xi \\ H^2 = \frac{8\pi G}{3} \bar{\rho} , \quad \rho = \frac{1}{a^3} m |\psi|^2 \equiv \sum_{i=1}^3 \rho_i ,$$

Effective self-interactions

Energy density for each field component

Fluid description

Madelung transformation

$$\psi_i = \sqrt{\frac{\rho_i a^3}{m}} e^{i\theta} , \quad \mathbf{v}_i \equiv \frac{1}{ma} \nabla \theta_i$$

Continuity and Euler equations

$$\begin{aligned} \dot{\rho}_i + 3H\rho_i + \frac{1}{a} \nabla \cdot (\rho_i \mathbf{v}_i) &= 0 , \\ \dot{\mathbf{v}}_i + H\mathbf{v}_i + \frac{1}{a} (\mathbf{v}_i \cdot \nabla) \mathbf{v}_i &= -\frac{1}{a} \nabla \left(\Phi_N + \Phi_{Q,i} + 2\Phi_\xi - \frac{(2\xi_1 + \xi_2)}{2m^2 a^2} \nabla^2 \Phi_\xi \right) \\ \Phi_{Q,i} &\equiv -\frac{1}{2a^2 m^2} \frac{\nabla^2 \sqrt{\rho_i}}{\sqrt{\rho_i}} \end{aligned}$$

Phenomenology

Phenomenology

Mass-radius relation of polarized vector solitons

Potential impacts on galaxy density profiles

Growth of linear perturbations

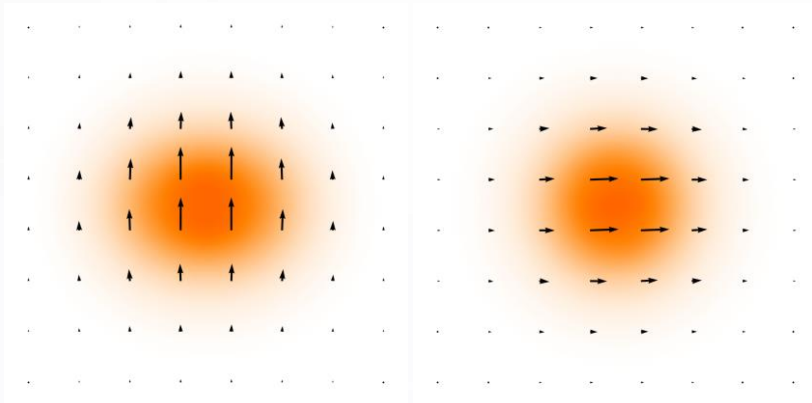
The success of cold DM at large scales should be retained

Correction to the GW speed

Constraints are obtained from GW 170817 & GRB 179817A and lack of gravitational Cherenkov radiation

Mass-radius relation of vector solitons

Ground-state vector solitons



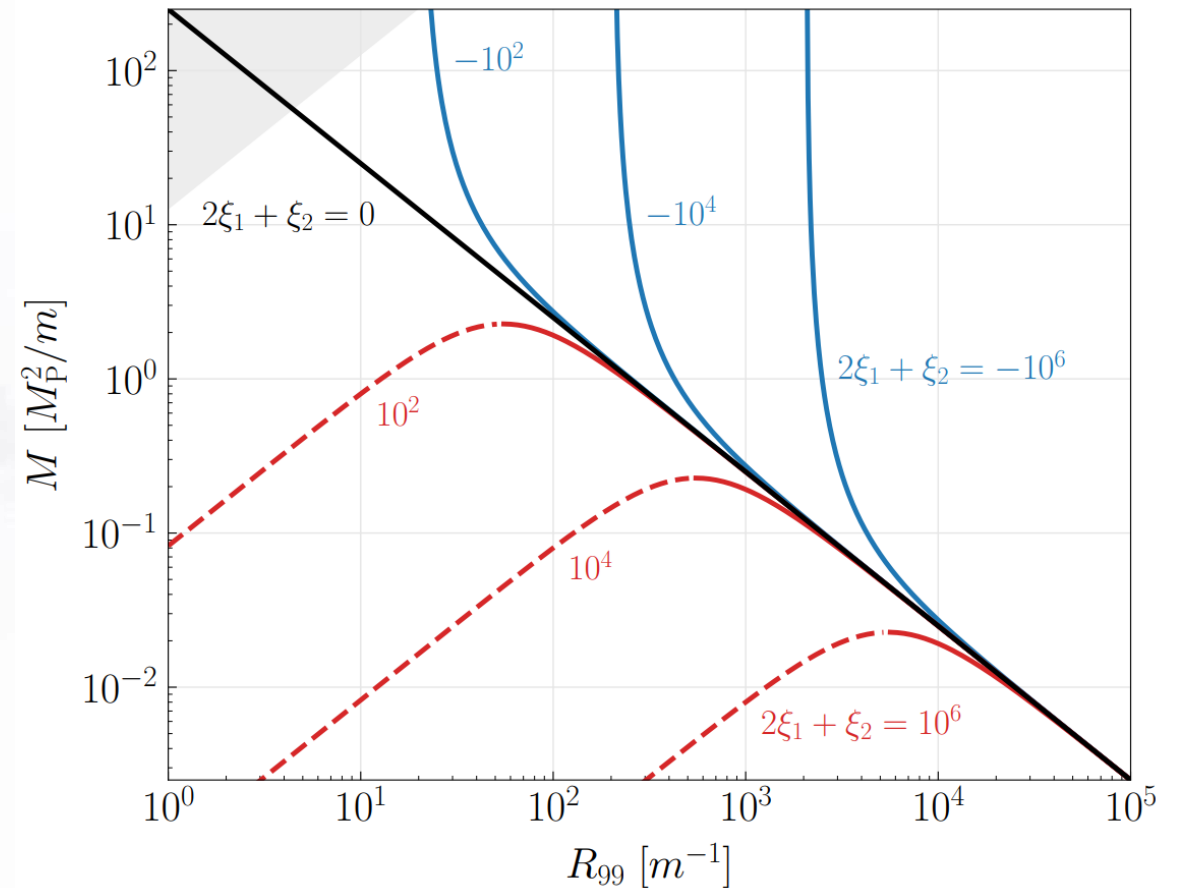
$S = 0$

$S = \hbar N$

Also see:
Jain & Amin (2022)
Zhang, Jain & Amin (2022)



$$M = \frac{\beta_1 (mR)^3}{(mR)^4 + \beta_2 (2\xi_1 + \xi_2)(mR)^2 + \beta_3 (2\xi_1 + \xi_2)^2} \frac{M_P^2}{m}$$



Growth of linear perturbations

$$\delta_i \equiv (\rho_i - \bar{\rho}_i)/\bar{\rho}_i, \quad \delta_i \propto \Delta_1, \quad (\partial_t^2 + 2H\partial_t + \Omega^2)\Delta_1 = 0$$

Wave DM

$$\Omega^2 = -4\pi G\bar{\rho} \left[\underbrace{1 - \frac{m^2}{4\pi G\bar{\rho}} \frac{k^4}{4a^4m^4}}_{\text{Cold DM, which is consistent with matter power spectrum observation for } k < k_{\text{obs}} \sim 10^3 k_{\text{eq}}} + (2\xi_1 + \xi_2) \frac{k^2}{a^2m^2} + (2\xi_1 + \xi_2)^2 \frac{k^4}{4a^4m^4} \right]$$

Cold DM, which is consistent with matter power spectrum observation for $k < k_{\text{obs}} \sim 10^3 k_{\text{eq}}$

$$\xrightarrow[\text{for } k < k_{\text{obs}}]{\text{Require } \Omega^2 \approx -4\pi G} |2\xi_1 + \xi_2| \ll \frac{a_{\text{eq}}^2 m^2}{k_{\text{obs}}^2} = 10^{10} \left(\frac{m}{10^{-20} \text{eV}} \right)^2 \left(\frac{10^{-28} \text{eV}}{H_{\text{eq}}} \right)^2$$

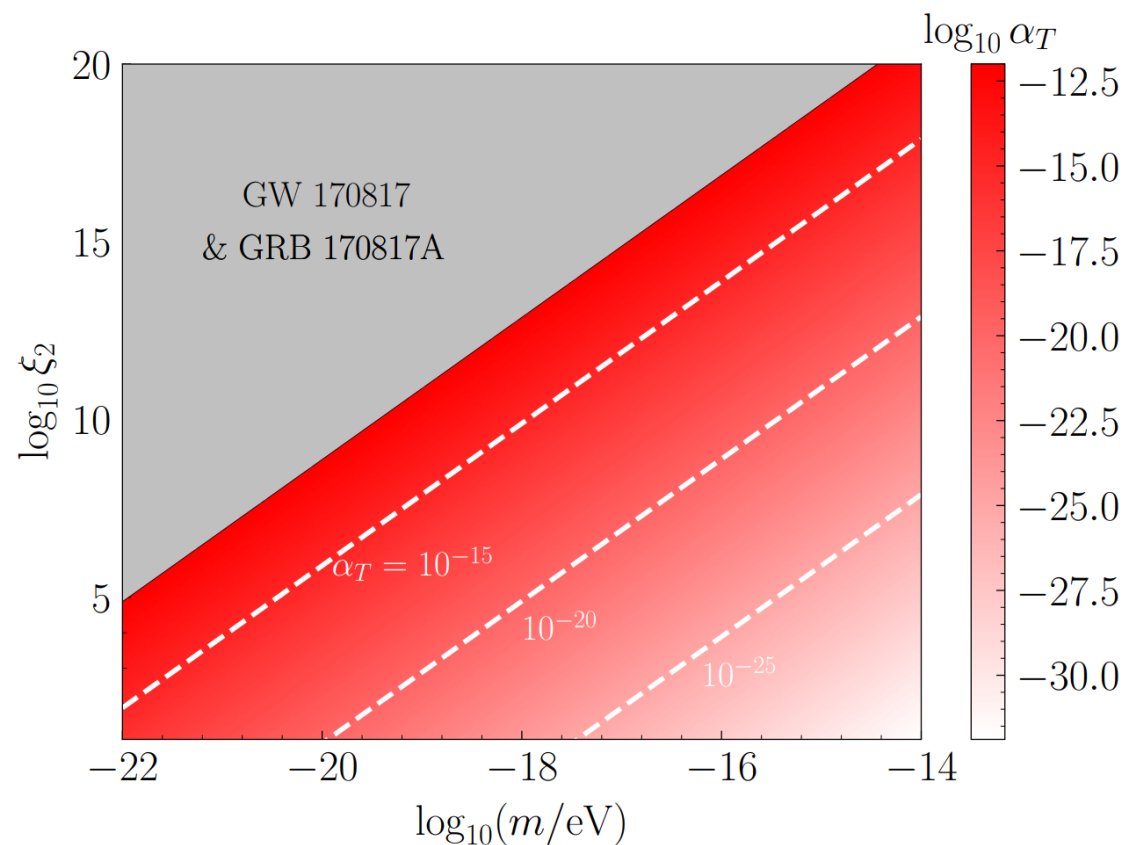
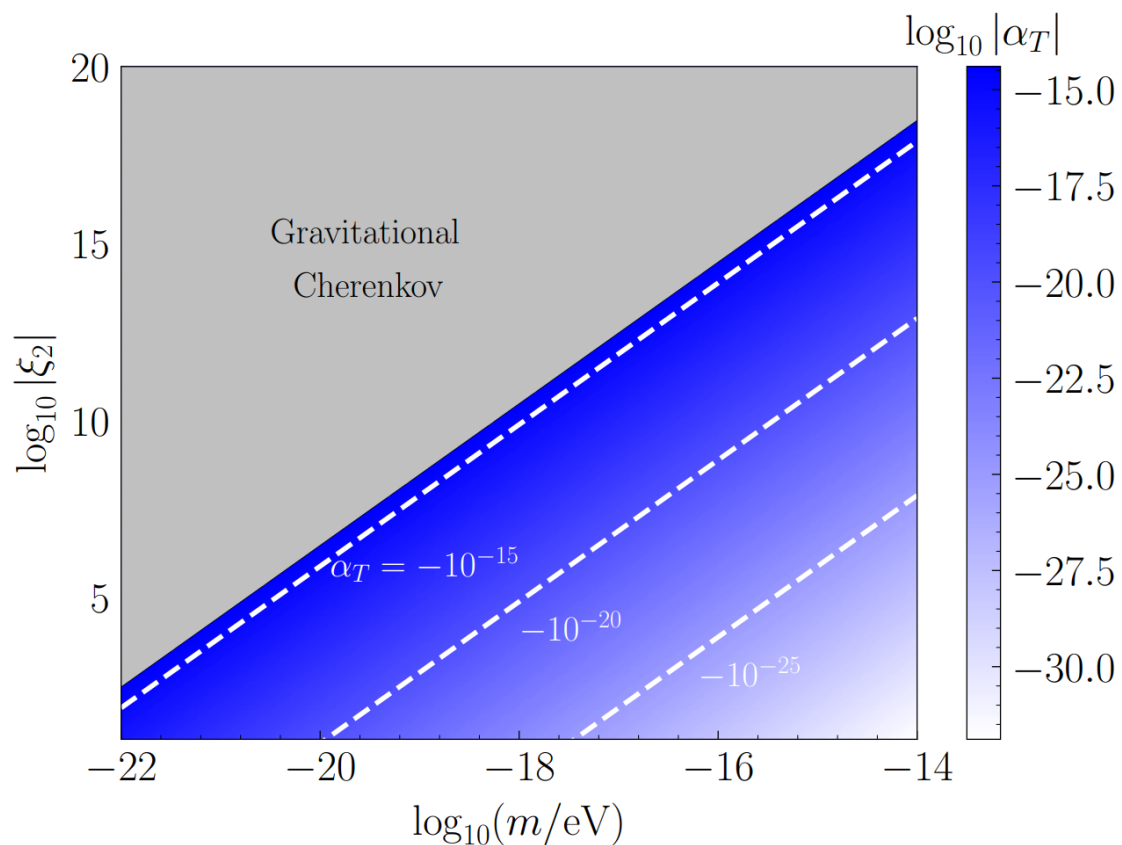
$$\left\{ \begin{array}{l} \text{For } \frac{m^2}{4\pi G\bar{\rho}} \ll (2\xi_1 + \xi_2)^2 \\ \text{For } \frac{m^2}{4\pi G\bar{\rho}} \gg (2\xi_1 + \xi_2)^2 \end{array} \right. \left\{ \begin{array}{l} \rightarrow \text{cold DM at large scales} \\ \rightarrow \text{enhanced small-scale structure than wave DM} \end{array} \right.$$

$$\rightarrow \text{wave DM, with Jeans scale } k_J \approx k_{J,0} \left[1 + (2\xi_1 + \xi_2) \sqrt{\frac{\pi G\bar{\rho}}{m^2}} \right]$$

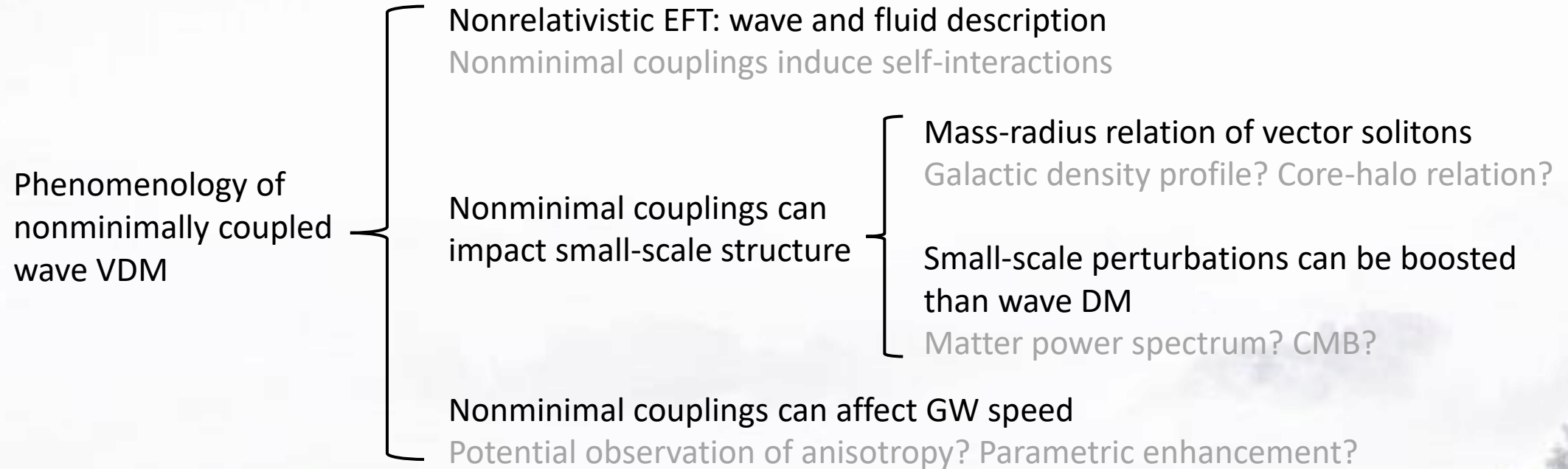
Constraints on GW speed

Energy density for component
in the direction of GWs

$$S^{(2)} = \frac{1}{2} \sum_{\lambda} \int d^4x M_*^2 \left[\dot{h}_{\lambda}^2 - c_T^2 (\nabla h_{\lambda})^2 \right] , \quad \alpha_T \equiv c_T^2 - 1 = \frac{\xi_2 \rho_{\hat{n}}}{m^2 M_{\text{P}}^2} \simeq \frac{\xi_2 \rho}{3m^2 M_{\text{P}}^2}$$



Summary



Constraint on nonminimal couplings

$$\frac{|\xi_1|}{m^2} \ll 10^{50} \text{eV}^{-2}, \quad -3 \times 10^{46} \text{eV}^{-2} \lesssim \frac{\xi_2}{m^2} \lesssim 8 \times 10^{48} \text{eV}^{-2}$$