

Parity-violating Signals from Cosmological Collider

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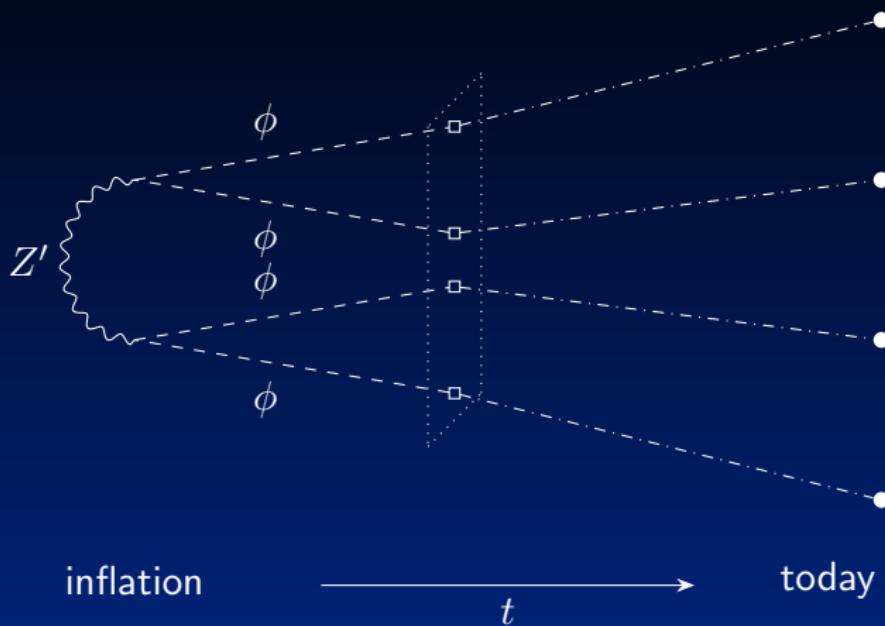
work in progress with Lian-Tao Wang, Zhong-Zhi Xianyu, Yi-Ming Zhong

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Inflation as a Probe for New Particles



$$H_I \lesssim 10^{13} \text{ GeV}$$

$$E \sim 3 \text{ K}$$

Why Parity-violating Interaction?

- Potentially low background

mock data based on structure formation model, gravity and baryonic effects on small scale should be dominantly parity-preserving

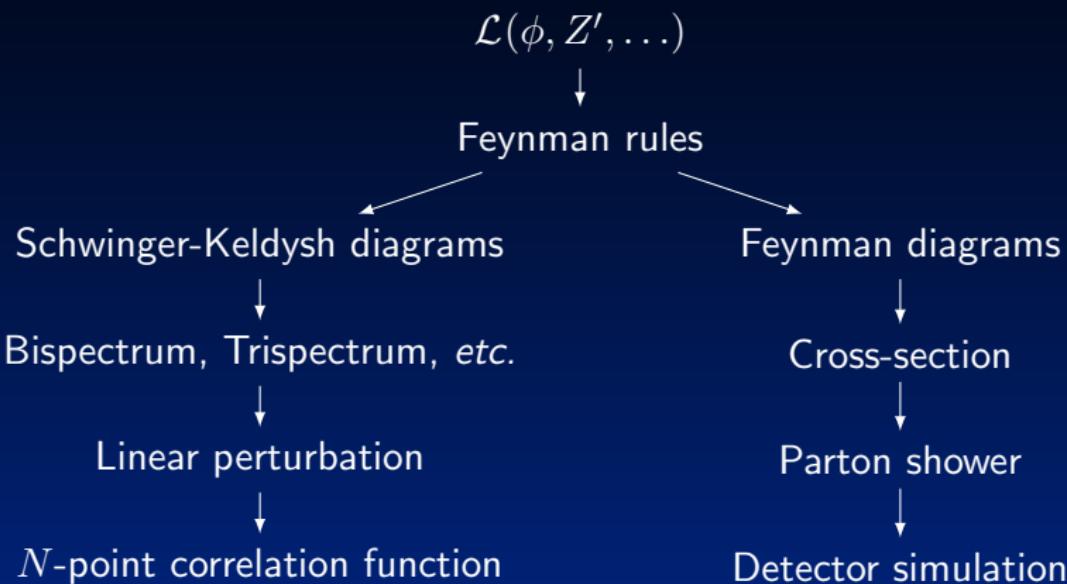
[Hou, Slepian, Cahn: 2206.03625], [Philcox: 2206.04227], ...

- Theoretically motivated

to produce large signal due to chemical potential enhancement

[Wang, Xianyu: 2004.02887], [Qin, Xianyu: 2208.13790], [Creque-Sarbinowski, Alexander, Kamionkowski, Philcox: 2303.04815], ...

Cosmological Collider vs. Particle Collider



Spin-1 Model: Chemical Potential

Abelian Higgs model with coupling to some background $\theta(t)$.

$$\begin{aligned} \mathcal{L} \supset & \sqrt{-g} \left[\frac{1}{4} Z_{\mu\nu} Z^{\mu\nu} + D_\mu \mathcal{H}^* D^\mu \mathcal{H} \right] + \frac{c_0 \theta(t)}{4} \epsilon^{\mu\nu\rho\sigma} Z_{\mu\nu} Z_{\rho\sigma} \\ & \sim \frac{1}{2} \mathbf{Z} \left(\frac{d^2}{d\tau^2} + k^2 + a^2 m_Z^2 \right) \mathbf{Z} + Z_i (-i \underbrace{c_0 \dot{\theta}}_{\triangleq c} a \epsilon^{ijk} k_j) Z_k \end{aligned} \quad (1)$$

with $m_Z \sim H_I$.

Spin-1 Model: Equation of Motion

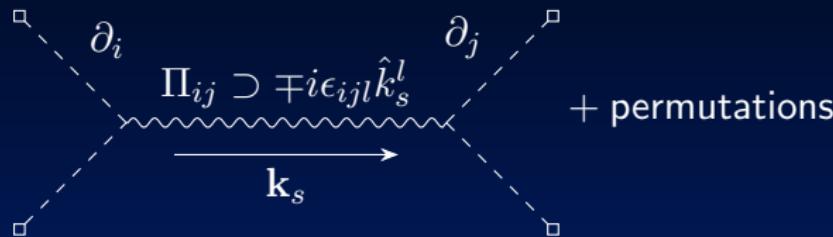
$$\begin{aligned}\mathbf{Z}_- &= \epsilon_- \cdot \exp\left(+\frac{\pi c}{2H_I}\right) \cdot f_-(m_Z, c_0\dot{\theta}; ik\tau) e^{i\mathbf{k} \cdot \mathbf{x}} \\ \mathbf{Z}_L &= \epsilon_L \cdot \quad \quad \quad 1 \quad \quad \quad \cdot g(m_Z; ik\tau) e^{i\mathbf{k} \cdot \mathbf{x}} \quad (2)\end{aligned}$$

$$\mathbf{Z}_+ = \epsilon_+ \cdot \exp\left(-\frac{\pi c}{2H_I}\right) \cdot f_+(m_Z, c_0\dot{\theta}; ik\tau) e^{i\mathbf{k} \cdot \mathbf{x}}$$

Chemical potential c enhances a particular polarization.

Spin-1 Model: Interaction with Inflaton

$$\mathcal{L}_{\text{int}} \propto a(\tau) \delta^{ij} \phi' \partial_i \phi Z_j. \quad (3)$$

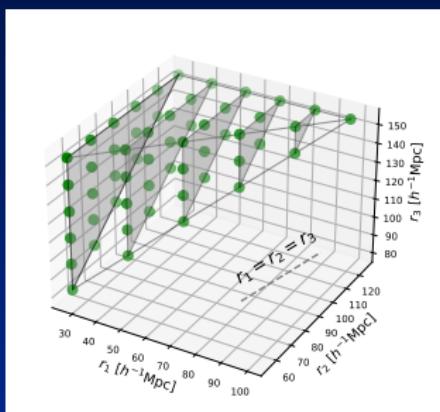


$$\langle \phi^4 \rangle_{\text{full, PO}} = \mathcal{J}(k_1, \dots, k_4, k_s) \left[i (\hat{\mathbf{k}}_1 \times \hat{\mathbf{k}}_2) \cdot \hat{\mathbf{k}}_s \right] + \text{permutations} \quad (4)$$

$$\langle \phi^4 \rangle_{\text{toy, PO}} = \left[i (\hat{\mathbf{k}}_1 \times \hat{\mathbf{k}}_2) \cdot \hat{\mathbf{k}}_s \right] + \text{permutations} \quad (5)$$

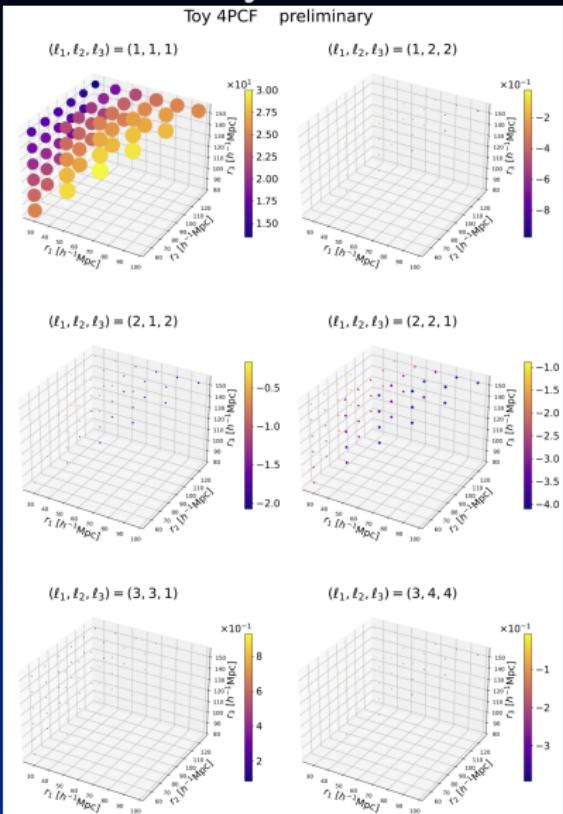
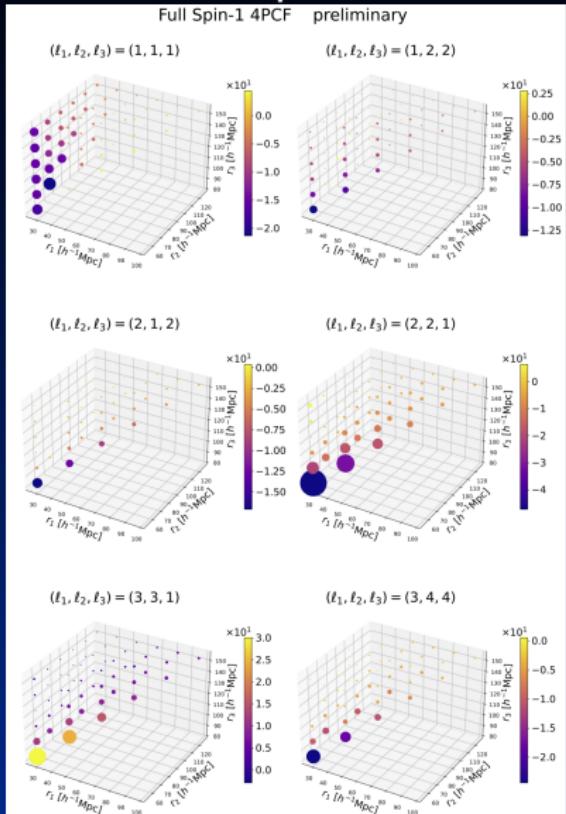
Reading Position-space Correlation Coefficients

$$\begin{aligned} \langle \delta_g^4 \rangle &= \langle \delta_g^4 \rangle (r_1, r_2, r_3, \theta_1, \theta_2, \theta_3), \\ \implies \langle \delta_g^4 \rangle &= \sum_{\ell} \zeta_{\ell_1, \ell_2, \ell_3}(r_1, r_2, r_3) \underbrace{\left[C_m^{\ell} Y_{\ell_1}^{m_1}(\theta_1) Y_{\ell_2}^{m_2}(\theta_2) Y_{\ell_3}^{m_3}(\theta_3) \right]}_{\triangleq \mathcal{P}_{\ell_1, \ell_2, \ell_3}(\theta_1, \theta_2, \theta_3)} \end{aligned} \quad (6)$$



$$(\ell_1, \ell_2, \ell_3) = (1, 1, 1) \sim (\hat{\mathbf{r}}_1 \times \hat{\mathbf{r}}_2) \cdot \hat{\mathbf{r}}_3.$$

Full Spin-1 Model vs. Contact Toy Model



Summary

- Inflation offers new probes for heavy particles
- Primordial parity-violating interactions can show up in the large-scale structure survey data
- We have established a pipeline to compute the position-space correlation function from models
- Full spin-1 model shows non-trivial correlation coefficients compared to contact toy model

What's next:

- Relation between equilateral (or local) templates and full model?
- Efficient tools to perform the Fourier transform from momentum space to position space?
- Unfolding position-space data to momentum space?
- Key features to distinguish various full models in either position space or momentum space?

Why Trispectrum?

The only parity-odd invariant tensor is ϵ_{ijk} .

$$\langle \phi^N \rangle \propto \epsilon_{ijl} (k_1^i k_2^j k_3^l).$$

Momentum conservation \implies at least 4 external particles, i.e.
 $\langle \phi^4 \rangle$. Alternatively,

2-point correlation

$\mapsto \#.$ 3-point correlation



4-point correlation



Some LSS Redshift Surveys

- BOSS and eBOSS (from SDSS)
- DES (CTIO)
- WiggleZ (AAT)
- KiDS (VLT)
- Euclid (ESA)
- Vera Rubin Observatory (LSST)
- Roman Space Telescope (WFIRST)
- ...

Spin-1 Model: Interaction with Inflaton

$$\begin{aligned} -\mathcal{L} &\supset \sqrt{-g} \left[\frac{c_1}{\Lambda} \partial_\mu \phi (\mathcal{H}^* D^\mu \mathcal{H}) + \frac{c_2}{\Lambda^2} (\partial \phi)^2 |\mathcal{H}|^2 + \text{h.c.} \right] \\ &\supset -\underbrace{\text{Im} \left\{ \frac{c_1 \dot{\phi}_0 m_Z}{\Lambda} \right\}}_{\triangleq \rho_{1,Z}} \frac{a^2}{\dot{\phi}_0} \eta^{\mu\nu} \partial_\mu \phi Z_\nu h - a^3 \underbrace{\frac{c_2 v \dot{\phi}_0}{\Lambda^2} \phi' h}_{\triangleq \rho_2}. \end{aligned} \quad (7)$$

Just A Fourier Transform?

$$\langle \delta_g^4 \rangle (\mathbf{r}_1, \dots, \mathbf{r}_4) \sim \int \prod_{i=1}^4 \left(\frac{dk_i}{2\pi} \right)^3 e^{i\mathbf{k}_i \cdot \mathbf{r}_i} \langle \phi^4 \rangle (\mathbf{k}_1, \dots, \mathbf{k}_4) \quad (8)$$

For instance, 10 grids per dimension, 10^{12} grids in total!

Just A Fourier Transform?

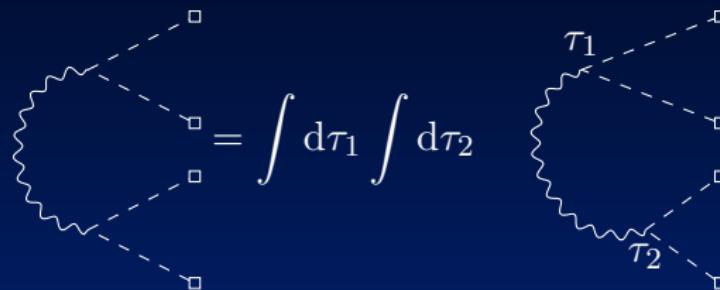
Instead, if we can “separate” $\langle \phi^4 \rangle(\mathbf{k}_1, \dots, \mathbf{k}_4) = f_1(\mathbf{k}_1) \dots f_4(\mathbf{k}_4)$
[Lee, Dvorkin: 2001.00584], [Smith, Senatore, Zaldarriaga: 1502.00635], then

$$\langle \delta_g^4 \rangle(\mathbf{r}_1, \dots, \mathbf{r}_4) \sim \prod_{i=1}^4 \left[\int \left(\frac{d\mathbf{k}_i}{2\pi} \right)^3 e^{i\mathbf{k}_i \cdot \mathbf{r}_i} f_i(\mathbf{k}_i) \right]. \quad (9)$$

Then, only $4 \times 10^3 \ll 10^{12}$ grids are needed. However,
 $\langle \phi^4 \rangle = f_1 \dots f_4$ is a very strict requirement.

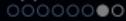
Hints from Diagrams

$$\mathcal{L}_{\text{int}} \propto a(\tau) \eta^{\mu\nu} \phi' \partial_\mu \phi Z_\mu. \quad (3)$$

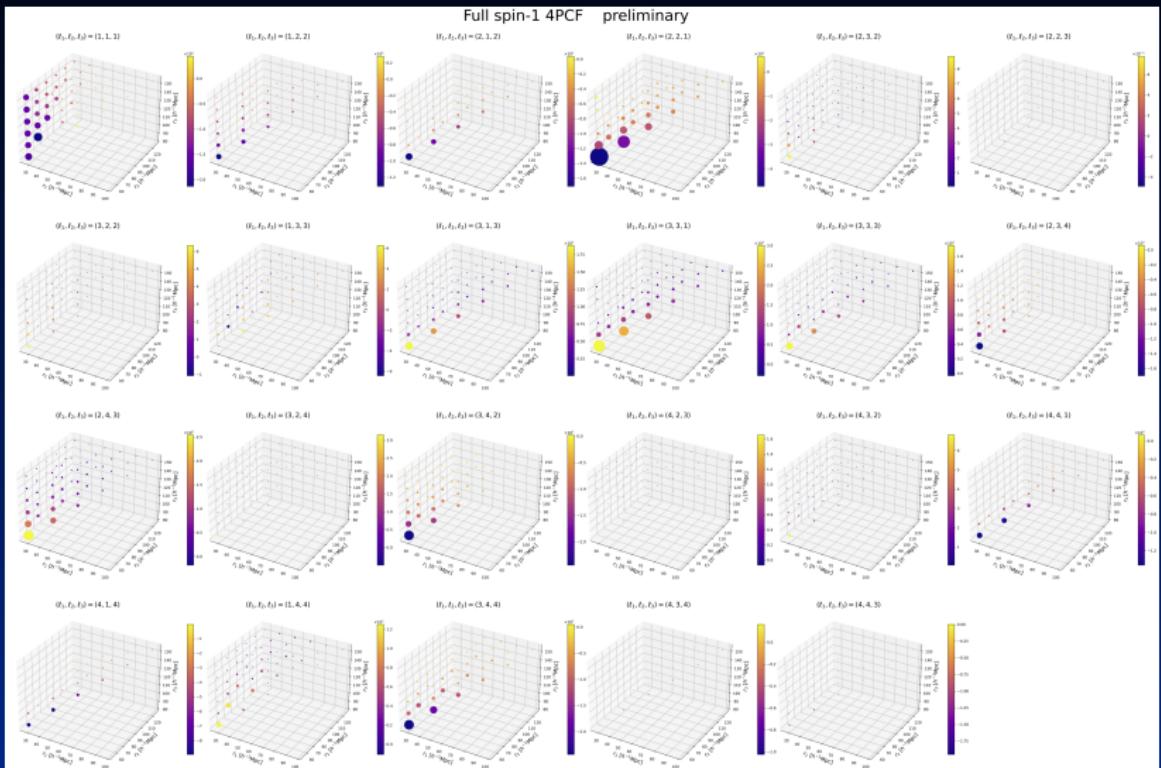


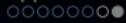
$$\frac{\tau}{k} \sim \exp(ik\tau) \cdot \frac{(1 + ik\tau)}{k^3}$$

After Wick rotation, we only need to sample over small τ .



Full Spin-1 Result





Equilateral Result

