



Lepton Flavor Portal Matter



Ricardo Ximenes

Based on 2211.09918 (PLB) and 2303.12983
with Lisa Everett, Shu Tian Eu, and George
Wojcik

Physics Department, UW-Madison

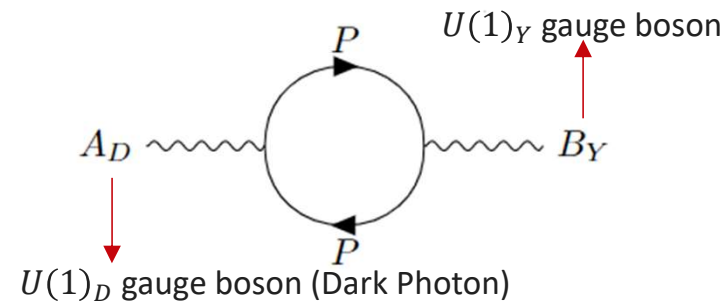
Phenomenology Symposium 2023



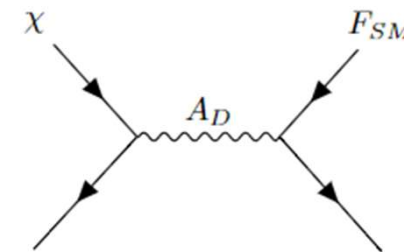
Vector Portal/Kinetic Mixing DM

- Vector portal matter (PM)_{1,2} /kinetic mixing (KM) DM setup:

	G_{SM}	$U(1)_D$
SM	R_{SM}	0
DM (χ)	$\mathbf{1}_0$	$\neq 0$
PM (P)	R_{PM}	$\neq 0$



- KM $\sim \frac{\epsilon}{2c_W} B_{\mu\nu} A_D^{\mu\nu}$ and $\epsilon \ll 1 \Rightarrow Q_D^{SM} \sim \epsilon Q_{em}$
- For $m_{DM, A_D} \sim 0.1 - 1$ GeV and $\epsilon \sim 10^{-(3-5)}$



We can recreate the correct relic abundance!

- Nature of PM?
 - Precision electroweak constraints and $H \rightarrow gg, \gamma\gamma \Rightarrow$ Must be **Vector-like (VL)**!
 - VL PM must be unstable. Without adding more particles, the only possibility is $F_{PM} \rightarrow F_{SM} h_D \Rightarrow$ PM is either a VL Quark (VLQ) or **VL Lepton (VLL)**. In this work we focus on VLL.

¹B. Holdom, Phys. Lett. B **166**, 196 (1986)

²B. Holdom, Phys. Lett. B **178**, 65 (1986)



PM/KM DM and coming up!

- Collider searches for VLL PM suggests $M_{PM} \geq 1$ TeV.
- For a set of VL PM, kinetic mixing coefficient $\epsilon \propto \sum_i Q_{Y_i} Q_{D_i} \log \frac{m_i^2}{\mu^2}$.
- Finite and calculable ϵ ? $\sum_i Q_{Y_i} Q_{D_i} = 0$.
- $\epsilon \sim 10^{-(3-5)}$ can be achieved with PM fields coming in pairs with $Q_D = \pm 1$, same Q_Y and $O(1)$ mass splitting.
- The discrepancy between the measured value of a_μ and the theoretical expectation of the SM at (significance of 4.2σ) is currently placed at^{1,2,3}

$$\Delta a_\mu = a_\mu^{exp} - a_\mu^{SM} = (251 \pm 59) \times 10^{-1}$$

- Assuming the discrepancy is due to new physics, we show that Δa_μ can be addressed in a minimal PM/KM model agnostic to the DM details.
- We present an embedding of $U(1)_D$ into a larger dark group of $SU(2)_A \otimes SU(2)_B \otimes Z_2$ and explore its rich scalar and fermion sectors.

¹b. Abi *et al.* (Muon $g-2$), Phys. Rev. Lett. **126**, 141801 (2021), arXiv:2104.03281 [hep-ex]

²g. W. Bennett *et al.* (Muon $g-2$), Phys. Rev. D **73**, 072003 (2006), arxiv:hep-ex/0602035

³t. Aoyama *et al.*, Phys Rept. **887**, 1 (2020), arXiv:2006.04822 [hep-ph]



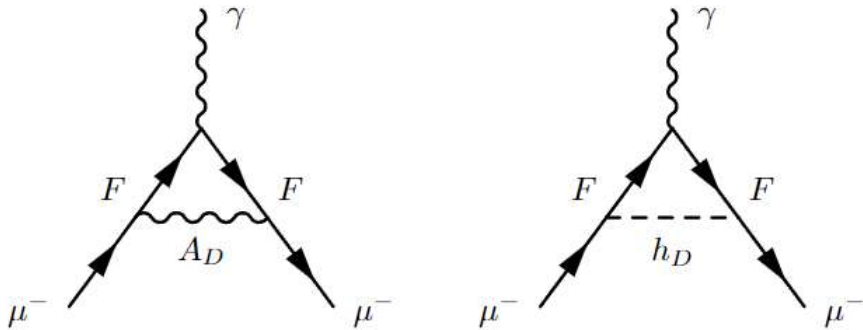
Minimal Construction

Field	$SU(2)_L \times U(1)_Y$	Q_D
$l_L = (v_L^\mu, \mu_L)^T$	$(2, -\frac{1}{2})$	0
μ_R	$(1, -1)$	0
$L_{L,R}^\pm = (N_{L,R}^\pm, L_{L,R}^\pm)^T$	$(2, -\frac{1}{2})$	± 1
$E_{L,R}^\pm$	$(1, -1)$	± 1
$S = v_S + h_D/\sqrt{2}$	$(1, 0)$	+1

$$(\bar{\mu}_L \quad \bar{L}_L^+ \quad \bar{E}_L^+ \quad \dots) \begin{pmatrix} m_\mu & y_{SL}^+ v_S & 0 & \dots \\ 0 & M_L^+ & e^{i\phi_{LE}^+} \frac{y_{LE}^+}{y_\mu} m_\mu & \dots \\ y_{SE}^+ v_S & e^{i\phi_{EL}^+} \frac{y_{EL}^+}{y_\mu} m_\mu & M_E^+ & \dots \\ \dots & \dots & \dots & \dots \end{pmatrix} \begin{pmatrix} \mu_R \\ L_R^+ \\ E_R^+ \\ \vdots \\ \dots \end{pmatrix}$$

- Five chirality-flipping masses:

$$m_\mu, e^{i\phi_{LE}^\pm} \frac{y_{LE}^\pm}{y_\mu} m_\mu, e^{i\phi_{EL}^\pm} \frac{y_{EL}^\pm}{y_\mu} m_\mu$$



$$\Delta a_\mu^{BSM} \sim \Re \left(\frac{c_L c_R}{24\pi^2} \frac{m_F}{m_\mu} \right) \left(\frac{m_\mu}{M_{BSM}} \right)^2$$

$$\Delta a_\mu^{PM} \approx -\frac{m_\mu^2}{16\pi^2} \sum_{i=+,-} \frac{y_{SL}^i y_{SE}^i}{M_L^i M_E^i} \frac{y_{LE}^i}{y_\mu} \cos \phi_{LE}^i$$

- Independent of parameters which govern the phenomenology of the dark sector!

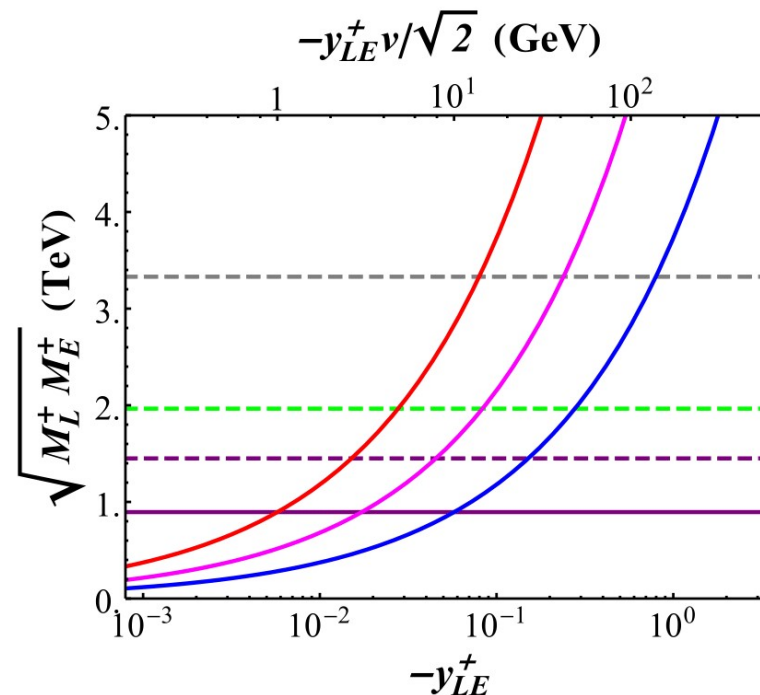


PM Collider Production

$$\Delta a_\mu^{PM} \approx -\Delta a_\mu \left(\frac{y_{LE}^+ / y_\mu}{36} \right) \left(\frac{1 \text{ TeV}}{M_E^+ / y_E^+} \right) \left(\frac{1 \text{ TeV}}{M_L^+ / y_L^+} \right)$$

- Plot of $\sqrt{M_L^+ M_E^+}$ as a function of y_{LE}^+ to reproduce Δa_μ :

- Continuous lines indicates the plots for $y_L^+ y_E^+ = 0.3, 1, 3$.
- Vertical lines indicate LHC, HL-LHC, HE-LHC, *hh*-FCC constraints on M_E^+ .
- Observed anomaly can be accounted for with $\frac{y_{LE}^+}{y_\mu} \sim O(10)$, $M_{L,E}^+ \sim O(\text{TeV})$.





$SU(2)_A \otimes SU(2)_B \otimes Z_2$ Model

- Why embed $U(1)_D$ into a larger group?
 - Justify a priori arbitrary $U(1)_D$ charge assignments.
 - Natural reason for the flavor symmetry.
- Anomaly-free embedding with finite KM for $SU(2)_A \otimes SU(2)_B \otimes Z_2$

$$SU(2)_A \times SU(2)_B \times Z_2 \xrightarrow{\langle \Phi \rangle \sim \text{TeV}} U(1)_D \times Z_2 \xrightarrow{\langle \Delta_{A,B} \rangle \sim \text{GeV}} Z'_2$$

Fields	$SU(2)_L \times U(1)_Y$	$SU(2)_A$	$SU(2)_B$	Z_2
Ψ_L, Ψ_R	$\left(2, -\frac{1}{2}\right), (1, -1)$	2	2	+1
V_L, V_R	$(1, -1), \left(2, -\frac{1}{2}\right)$	1	3	+1
S_L, S_R	$\left(2, -\frac{1}{2}\right), (1, -1)$	1	1	-1
$\Phi \sim \text{TeV}$	(1, 0)	2	2	+1
$\Delta_A \sim \text{GeV}$	(1, 0)	3	1	-1
$\Delta_B \sim \text{GeV}$	(1, 0)	1	3	-1



Scalar Sector

$$SU(2)_A \times SU(2)_B \times Z_2 \xrightarrow{\langle \Phi \rangle \sim \text{TeV}} U(1)_D \times Z_2 \xrightarrow{\langle \Delta_{A,B} \rangle \sim \text{GeV}} Z'_2$$

- Scalar potential for this model is quite complicated ($8 + 3 + 3 = 14$ variables).
- There are 6 Goldstone modes (broken generators of $SU(2)_A \otimes SU(2)_B$) and $14 - 6 = 8$ massive states.

Scalars	Mass	Z'_2	Q_D
h_1, h_2, h_4	$M_{h_1, h_2, h_4} \sim \text{TeV}$	+1	0
h^\pm	$M_h^\pm \sim \text{TeV}$	+1	± 1
h_5, h_6	$M_{h_5, h_6} \sim \text{TeV}$	-1	0
h_3	$m_{h_3} \sim \text{GeV}$	+1	*
Gauge Bosons	Mass	Z'_2	Q_D
Z_D	$M_{Z_D} \sim \text{TeV}$	+1	0
W_l^\pm	$M_{W_l} \sim \text{TeV}$	-1	± 1
W_h^\pm	$M_{W_h} \sim \text{TeV}$	-1	± 1
A_D	$m_{A_D} \sim \text{GeV}$	+1	*



Fermion Sector

Fermions	Mass	Z'_2	Q_D
e_b	m_b^{SM}	+1	0
L^0, E^0	$M_L^0, M_E^0 \sim \text{TeV}$	+1	0
e_a	m_a^{SM}	-1	0
L^\pm, E^\pm	$M_L^\pm, M_E^\pm \sim \text{TeV}$	-1	± 1

- The Z'_2 symmetry provides charged Lepton flavor protection.
- Z'_2 – positive block: One SM Lepton and Two heavy VLL.
- Z'_2 – Negative block: One SM Lepton and Two heavy VLL PM.
- Two scenarios of interest:

SM Field	Scenario A	Scenario B
Muon	e_a	e_b
Tau	e_b	e_a



Coming Up!

- Shu Tian's Flavor Portal Matter-II talk will discuss in scenarios A and B:
 - Muon $g - 2$ anomaly: Contribution from new diagrams and bigger parameter space.
 - Rich phenomenology: Diboson, VLL and VLL PM collider production.

