

Precision Cosmological Constraints on Atomic Dark Matter

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Atomic Dark Matter

• Dark sector fermions p_D and e_D with masses m_{p_D} , m_{e_D} , oppositely charged under a dark U(1) gauge force mediated by a massless dark photon with dark fine structure constant α_D .

$$\mathcal{L}_{aDM} = A_{\mu\nu}A^{\mu\nu} + i\bar{p}_D(\not\!\!D - m_{p_D})p_D + i\bar{e}_D(\not\!\!D - m_{e_D})e_D$$

- Asymmetric, overall neutral abundance that accounts for a fraction $\mathbf{f}_{\mathbf{D}} \equiv \frac{\Omega_{aDM}}{\Omega_{DM}} \leq 1$ of the dark matter.
- Temperature ratio of the dark photons to SM photons today

 $\boldsymbol{\xi} \equiv T_D/T_{SM}$. Equivalently, $\Delta N_D \equiv \left(\frac{8}{7}\right) \left(\frac{11}{4}\right)^{\frac{1}{3}} \xi^4$, effective number of additional neutrinos at CMB.

• First studied by Kaplan et al. in 0909.0753.

Some Motivations for Atomic Dark Matter

Experiment

- Hubble tension.
- S_8 tension.
- Small scale behaviour of DM (e.g. core-cusp).

Theory

- Approximately Z₂-symmetric mirror sectors can address naturalness issues.
- e.g. Mirror Twin Higgs
- Twin/mirror sectors have their own gauge forces, and can have dark matter – dark radiation interactions.
- The SM sector has rich particle content and forces. Why shouldn't the dark matter?



Questions

• How much do current cosmological observations constrain the atomic dark matter (aDM) parameter space?

• To what extent can aDM alleviate the Hubble and S_8 tensions?

Cosmology of Atomic Dark Matter

- Early universe: dark plasma.
- Dark acoustic oscillations (DAOs) due to dark photon pressure support.
- Dark recombination at temperatures below dark hydrogen binding energy.
- Dark photons decouple from atomic dark matter.
- The DAOs and additional radiation in the early universe leave distinctive imprints on the CMB and large-scale structure (LSS) of the universe.
- First studied by Cyr-Racine et al. in 1209.5752, 1310.3278.

CMB and Matter Power Spectrum

- Complex effect on CMB power spectrum. Percent-level deviations from ΛCDM.
- Matter power spectrum: Suppression and oscillations for *k* that enter horizon before dark decoupling.
- Lower ΔN_D , higher binding energy $B_D \rightarrow$ suppression starts at higher k.
- Non-linear evolution smears out oscillations at low redshifts. [2101.12229]

Spectra computed with CLASS_aDM, at https://github.com/jp-barron/class_adm-3.1.git





Datasets: CMB, BAO, H_0 , S_8

- Planck 2018 TT, TE, EE, high- and low-l, lensing. [1907.12875]
- BAO: Measurements of D_V/r_S from 6dFGS, SDSS, BOSS. [1106.3366, 1409.3242, 1607.03155]
- Pantheon supernova dataset.
 [1710.00845]
- SH0ES measurement of H₀ using Cepheid variables and Type Ia supernovae. [2112.04510]
- Planck SZ 2013 measurement of S_8 . [1303.5080]



- 3D scans: constraining ΔN_D , f_D , and m_{e_D}
 - m_{e_D} , $\alpha_D \sim T_{dec}$.

Parameter

Prior

- Large f_D not ruled out for sufficiently large m_{e_D} .
- Constraint on ΔN_D sharply transitions from < 0.3 to much tighter as m_{e_D} is lowered.

 $\Delta N_{\rm D}$

 $0 < \Delta N_D < 1$

f_D

 $0 < f_D < 1$





H_0 and S_8



Without Direct Measurements

With Direct Measurements





Conclusion

- CLASS_aDM code can compute CMB and linear matter power spectrum for a wide range of aDM parameter values.
- We computed the current constraints on the full aDM parameter space for the first time by comparing with precision cosmological observations.
- There is no preference for atomic dark matter from CMB observations alone, but large aDM fractions can be allowed from large-scale measurements.
- When including low-redshift measurements of the Hubble constant and structure, the Hubble and S_8 tensions can be eased but not eliminated. A significant fraction of aDM that recombines near $z = 2 \times 10^4$ is preferred, along with a non-zero amount of dark radiation, in order to accommodate the local measurements.





Backup Slides

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Dark Matter: CDM or something more? 150 NGC 6503 • The cold, collisionless dark matter ĬŢŢŢĸĸĸĸĸĸĸŢĸĸĸĸĸĸĸĸŢĬ paradigm explains many observations V_c (km s⁻¹) extremely well. • Galaxy rotation curves. 50 • The Bullet cluster. • The cosmic microwave background spectrum. 20 • Large scale structure in the universe. 10 30 Freese, 2008 Radius (kpc) $flat - \Lambda CDM$ • But it's not perfect. 67.4^{+1.1} DES+BAO+B • Core vs cusp problem. Late • Hubble tension. • S_8 tension. • Could a component of dark matter have arly vs. Lat Verde. non-trivial interactions in the dark sector? Treu. 4.0aRiess. 2019 78 80 66 68 76 $H_0 \, [{\rm km \, s^{-1} \, Mpc^{-1}}]$

Dark recombination

- Dark recombination occurs after T_D falls below $B_D = \frac{\alpha_D^2 \mu_D}{2}$.
- Redshift of dark recombination controlled by B_D and ΔN_D .
- In CLASS: Use Saha equation for ionization fraction x_D above 0.999.

$$\frac{x_D^2}{1 - x_D} = \frac{1}{n_D} \left(\frac{T_{DM} m_{e_D}}{2\pi}\right)^{3/2} e^{-B_D/T_{DM}}$$

- Then, solve for $\frac{dx_D}{dt}$ with effective multilevel atom formalism using HYREC-2. [Lee and Ali-Haïmoud, 2007.14114]
 - Neglects some sub-dominant radiative transfer effects that are only modelled for parameters near SM values.







Dark decoupling

- When energy transfer rate between dark photons and dark electrons/protons drops below Hubble, they decouple.
- Dominant energy exchange mechanism is usually Thomson scattering, with cross-section

$$\sigma_{T_D} = \frac{8\pi}{3} \left(\frac{\alpha_D}{m_{e_D}}\right)^2$$

- Other processes can dominate over Thomson scattering and delay decoupling
 - For high α_D and m_{e_D}/m_{H_D} , Rayleigh scattering.
 - For low α_D and ξ , bremsstrahlung, photo-recombination/ionization*.
- Opacity of dark plasma set by Thomson, Rayleigh, photo-ionization rates.



*Photo-heating/cooling processes not yet included in CLASS-aDM due to numerical instabilities – to be added soon. Similar rate to bremsstrahlung.

Hubble and S_8 tensions

- H_0 measures the rate of expansion of the universe today.
 - Planck 2018 (CMB): $H_0 = 67.4 \pm 0.5$ km/s/Mpc
- S_8 measures the size of fluctuations in the matter density at the scale $8 h^{-1}$ Mpc.
 - Planck 2018 (CMB): $S_8 = 0.834 \pm 0.016$
 - KiDS-1000 (weak lensing): $S_8 = 0.754^{+0.027}_{-0.029}$



Dark Acoustic Oscillations



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Constraining aDM parameter space

- Used Markov chain Monte Carlo package MontePython to sample posterior distribution of the model parameters, using publicly available likelihoods for the chosen datasets.
- Generated 95% confidence level contours from posterior.
- Ran \geq 6 chains for each scan.
- Required Gelman-Rubin convergence R 1 < 0.02 for each parameter.
- CDM parameters: $\{\omega_b, \omega_{cdm}, h, \ln 10^{10}A_S, n_s, \tau_{reio}\}$
 - *h* instead of θ_s because of a bug in CLASS 3.1 shooting procedure.
- Baseline dataset: Planck + BAO + Pantheon. Then add LSS, $H_0,\,S_8$ measurements.



Results: 5D scans

- Allow all five parameters $(f_D, \Delta N_D, m_{p_D}, m_{e_D}, \alpha_D)$ to vary.
- Only ΔN_D is robustly constrained.
- Recover Planck limit of $\Delta N_{\rm eff} \leq 0.3.$
- No preference for non-zero f_D , but large fractions are possible.



Parameter	f _D	ΔN_D	m_{p_D}	m_{e_D}	α_D
Prior	$0 < f_D < 1$	$0 < \Delta N_D < 1$	$0 < \log_{10}(\frac{m_{p_D}}{\text{GeV}}) < 3$	$-4 < \log_{10}(\frac{m_{e_D}}{\text{GeV}}) < -2$	$0.005 < \alpha_D < 0.4$



H_0 and S_8 : With direct measurements

- Larger 95% confidence region.
- Best-fit aDM values pulled further than Λ CDM, with lower minimum χ^2 .
- $\Lambda CDM + \Delta N_{eff}$ unable to accommodate lower S_8 .
- DM-DR interactions necessary to allow lower S_8 .



Including local H_0 and S_8 measurements

- Preference for non-zero $\Delta N_D \sim 0.3$ and $f_D > 0.1$.
- Strong preference for redshift of dark recombination around 2×10^4 .
- No preference for particular value of m_{p_D} .

