Vector boson dark matter in a Classically Conformal U(1) extension of the Standard Model

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#### Classically conformal extended SM Radiative Symmetry breaking as origin of SM Higgs potential

• SM Higgs potential requires  $-m^2$  term for EW symmetry breaking (EWSB); usually added by hand.

$$V = \boxed{-m^2(\Phi^{\dagger}\Phi)} + \lambda_{\Phi}(\Phi^{\dagger}\Phi)^2, \quad m^2 > 0$$
 (1)

- Alternative origin: Radiative symmetry breaking via Coleman-Weinberg mechanism (Coleman & Weinberg, 1973).
- Extend SM minimally with hidden U(1) gauge group (dubbed U(1)<sub>H</sub>), containing a Higgs scalar Φ.
- Implement CW mechanism for  $U(1)_H$  by imposing classical conformality (Iso, Okada, & Orikasa, 2009).



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# Coleman-Weinberg Mechanism

Radiative Symmetry breaking as origin of SM Higgs potential

• Hidden  $U(1)_H$  sector scalar effective potential of the form

$$V_{\phi} = \lambda_{\phi} \left(\Phi^{\dagger}\Phi\right)^{2} + V_{1loop}$$
$$= \frac{1}{4}\lambda_{\phi}\phi^{4} + \frac{\beta_{\phi}}{8}\phi^{4} \left(ln\left[\frac{\phi^{2}}{v_{\phi}^{2}}\right] - \frac{25}{6}\right), \text{ where } \phi = \sqrt{2}\text{Re}\left[\Phi\right]$$
(2)

 $\bullet$  Combined SM Higgs and  $\Phi$  scalar potential is

$$V = \lambda_{h} \left( H^{\dagger} H \right)^{2} - \left[ \lambda_{mix} \left( H^{\dagger} H \right) \left( \Phi^{\dagger} \Phi \right) \right] + \lambda_{\phi} \left( \Phi^{\dagger} \Phi \right)^{2} + V_{1loop} \quad (3)$$

- Radiative symmetry breaking in  $U(1)_H$  sector at  $\langle \phi 
  angle = {\it v}_\phi$
- With  $\lambda_{mix} > 0$ ,  $\langle \phi \rangle = v_{\phi}$  generates negative SM Higgs mass squared term, driving EW symmetry breaking.



#### **Coupling Analysis**

Conventional: 
$$V = \frac{\lambda_h}{4}(h^2 - v_h^2)^2 + \frac{\lambda_\phi}{4}(\phi^2 - v_\phi^2)^2 + \frac{\lambda_{mix}}{4}(h^2 - v_h^2)(\phi^2 - v_\phi^2)$$
  
Conformal:  $V = \frac{\lambda_h}{4}h^4 + \frac{\lambda_\phi}{4}\phi^4 + \frac{\beta_\phi}{8}\phi^4 \left(\ln\left[\frac{\phi^2}{v_\phi^2}\right] - \frac{25}{6}\right) - \frac{\lambda_{mix}h^2\phi^2}{4}$ 

• Mass-squared matrices defined as

$$M_{sq} = \begin{pmatrix} \partial_h^2 V & \partial_h \partial_\phi V \\ \partial_\phi \partial_h V & \partial_\phi^2 V \end{pmatrix} \Big|_{h=v_h, \phi=v_\phi} = \begin{pmatrix} m_h^2 & M^2 \\ M^2 & m_\phi^2 \end{pmatrix}$$

• Diagonalize  $M_{sq}$  to find mixing of eigenstates:

$$h = h_1 \cos(\theta) + h_2 \sin(\theta)$$
  
$$\phi = -h_1 \sin(\theta) + h_2 \cos(\theta)$$

In our analysis, we set  $M_{h_1} > 2M_{h_2}, \theta \ll 1 \Rightarrow [h \sim h_1, \phi \sim h_2]$ 

## **Coupling Analysis**

• Express potentials in terms of observables and extract couplings.

For  $M_{h_1} > 2M_{h_2}, \theta \ll 1$ , conventional system coupling goes as

$$g_{h_{1}h_{2}h_{2}} \simeq -\frac{M_{h_{1}}^{2}}{2v_{\phi}} \left(1 + 2\frac{M_{h_{2}}^{2}}{M_{h_{1}}^{2}}\right) \theta \quad \text{for} \quad \theta \ll \frac{v_{h}}{v_{\phi}}, \qquad (4)$$

$$g_{h_{1}h_{2}h_{2}} \simeq \frac{M_{h_{1}}^{2}}{2v_{h}} \left(1 + 2\frac{M_{h_{2}}^{2}}{M_{h_{1}}^{2}}\right) \theta^{2} \quad \text{for} \quad \frac{v_{h}}{v_{\phi}} \lesssim \theta, \qquad (5)$$

while conformal system coupling goes as

$$g_{h_1h_2h_2} \simeq -\frac{M_{h_2}^2}{2v_h} \left(1 - 4\frac{M_{h_2}^2}{M_{h_1}^2}\right) \theta^2.$$
 (6)

• Combination of cancellation of lower order  $\theta$  terms and unique structure leads to coupling suppression in the conformal system.



5/10

### **Coupling Analysis**



• Using  $M_{h_1} = 125$  GeV,  $M_{h_2} = 25$  GeV,  $v_h = 246$  GeV,  $v_{\phi} = 10^4$  GeV, and  $|\theta| = 0.1$ :

Conventional system:  $g_{h_1h_2h_2} = 0.424$ CW system:  $g_{h_1h_2h_2} = -0.0107$ 



Numerical Analysis: Br  $(h_1 \rightarrow h_2 h_2)$ 

- Gray regions excluded by LHC (ATLAS, 2020) and LEP-II (for  $M_{h_2} = 25$ GeV) (LEP-II, 2003)
- Prospective ILC search reach indicated by blue region for anomalous Higgs decay (Liu, Wang, Zhang, 2017) and red region for anomalous coupling (Barklow et. al., 2018).



Figure: Conventional (dashed) and Conformal (solid) branching ratios.  $M_{h_2} = 10$  (red), 25 (black), and 50 (blue) GeV.

# $U(1)_H$ vector boson Dark Matter

- Consider Z', the gauge boson of U(1)<sub>H</sub>, as DM candidate
- Four main annihilation processes for  $Z'Z' \rightarrow h_2h_2$
- Reproducing observed DM relic density  $\Omega_{DM}h^2 = 0.12$  (Planck 2018) requires  $\langle \sigma v_{rel} \rangle \sim 1$  pb



Figure:  $Z'Z' \rightarrow h_2h_2$  DM annihilation process diagrams



# $U(1)_H$ vector boson Dark Matter

- Non-relativistic approx. for preliminary results.
- Red, Black, Blue lines correspond to non-excluded parameter space for  $\theta$  below LEP-II bounds for  $M_{h_2} = 10$ (red), 25 (black), and 50 (blue) GeV, respectively.
- $\langle \sigma v_{rel} \rangle \sim 1 \text{pb}$  satisfied along purple curve. Conformal model satisfies this condition and reproduces  $\Omega_{DM} h^2 = 0.12$  at intersection points.



Figure: Gauge coupling  $g_H$  vs. Z' (DM) mass, for select values of  $m_{h_2}$ .  $\Omega_{DM}h^2 = 0.12$  is reproduced along the purple line.

## Summary

- Classical conformal structure & Coleman-Weinberg mechanism as origin of EW Symmetry breaking.
  - ► Radiative symmetry breaking in U(1)<sub>H</sub> sector induces negative SM Higgs mass term, driving EW symmetry breaking
- Unique structure of conformal potential greatly affects Higgs physics/phenomenology
  - ▶ Most notably, trilinear coupling  $g_{h_1h_2h_2}$  suppression in CW model vs. conventional model
  - ▶ Models distinguishable by precision measurement of anomalous Higgs coupling alongside (non-)observation of anomalous Higgs decay  $h_1 \rightarrow h_2 h_2 \rightarrow b \bar{b} b \bar{b}$  at future  $e^+e^-$  colliders (ILC)
- With Z' as DM candidate,  $\Omega_{DM}h^2 = 0.12$  can still be satisfied for appropriate choice of scalar mass  $m_{h_2}$ .

