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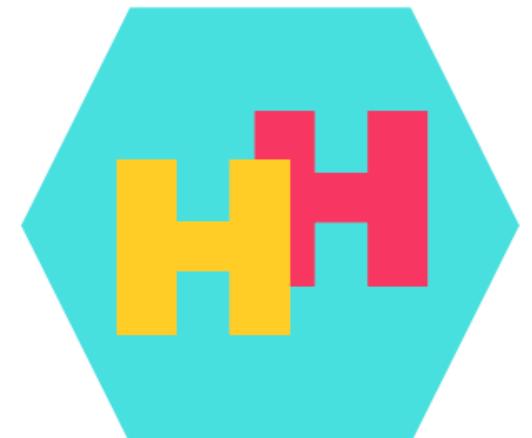
Probing the nature of electroweak symmetry breaking with Higgs boson pairs in ATLAS

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On behalf of the ATLAS Collaboration

Phenomenology 2023 Symposium, Pittsburgh

8-10 May 2023



Why look for Higgs boson pairs?

- Great progress in our understanding of the **Higgs boson** since its discovery in 2012
- But still little knowledge about the **Higgs potential**

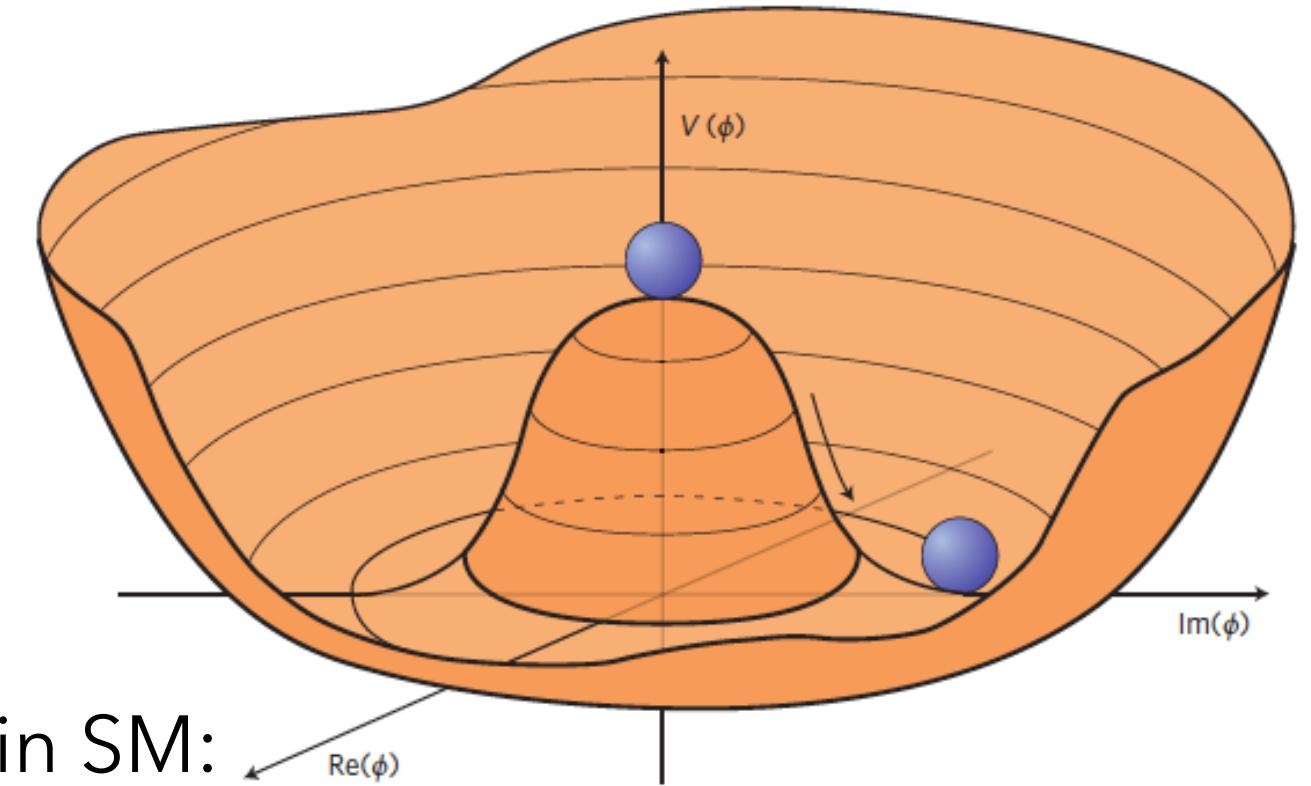
$$V(H) = \frac{1}{2}m_H H^2 + \boxed{\lambda v H^3} + \dots$$

self-interaction term

$$\lambda_{HHH} = \lambda v$$

Vacuum expectation value in SM:

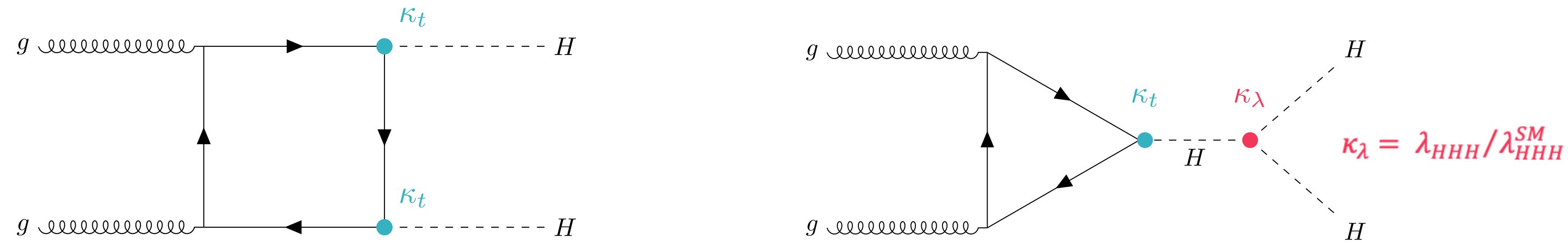
$$v = (\sqrt{2}G_F)^{-1/2} = 246 \text{ GeV}$$



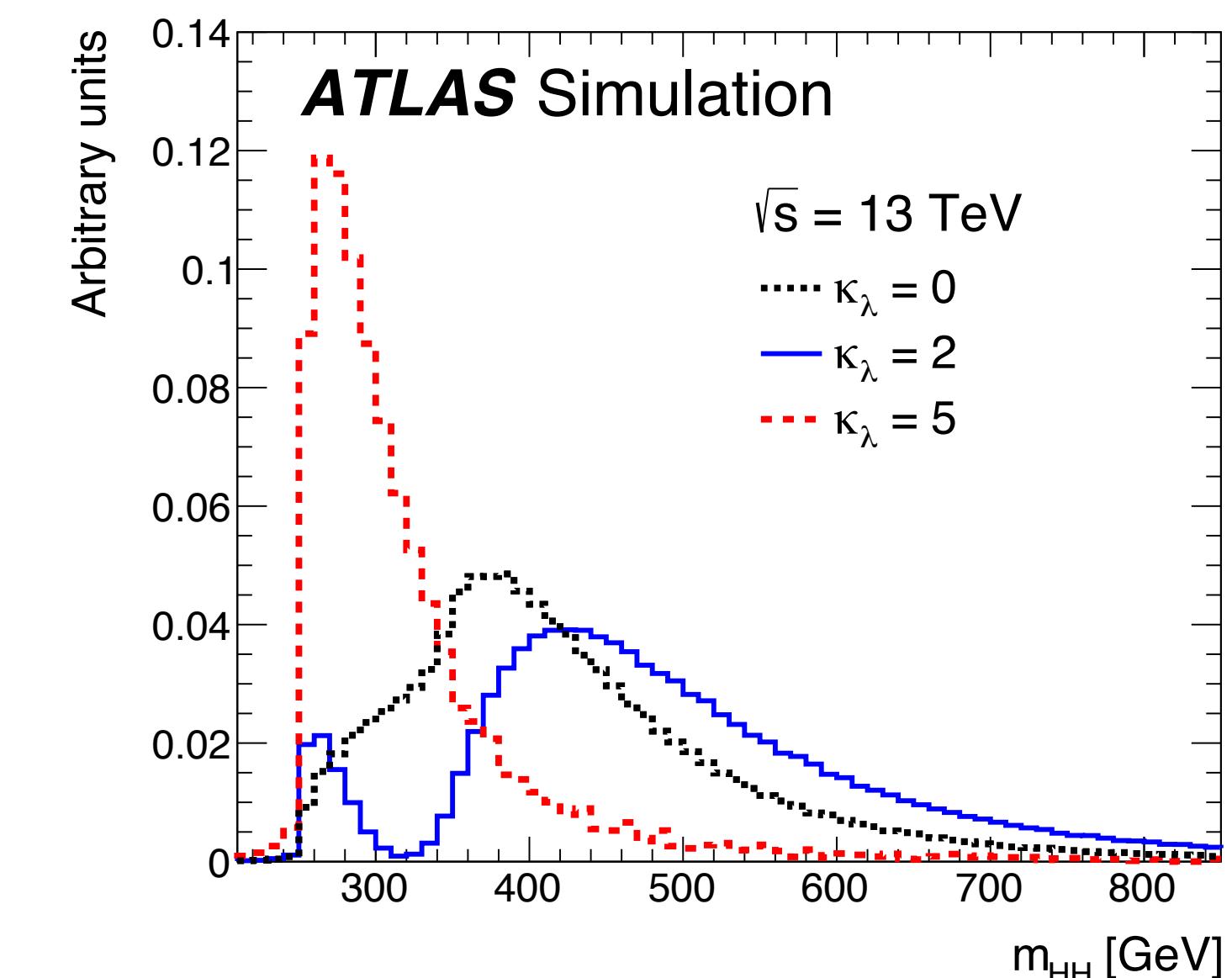
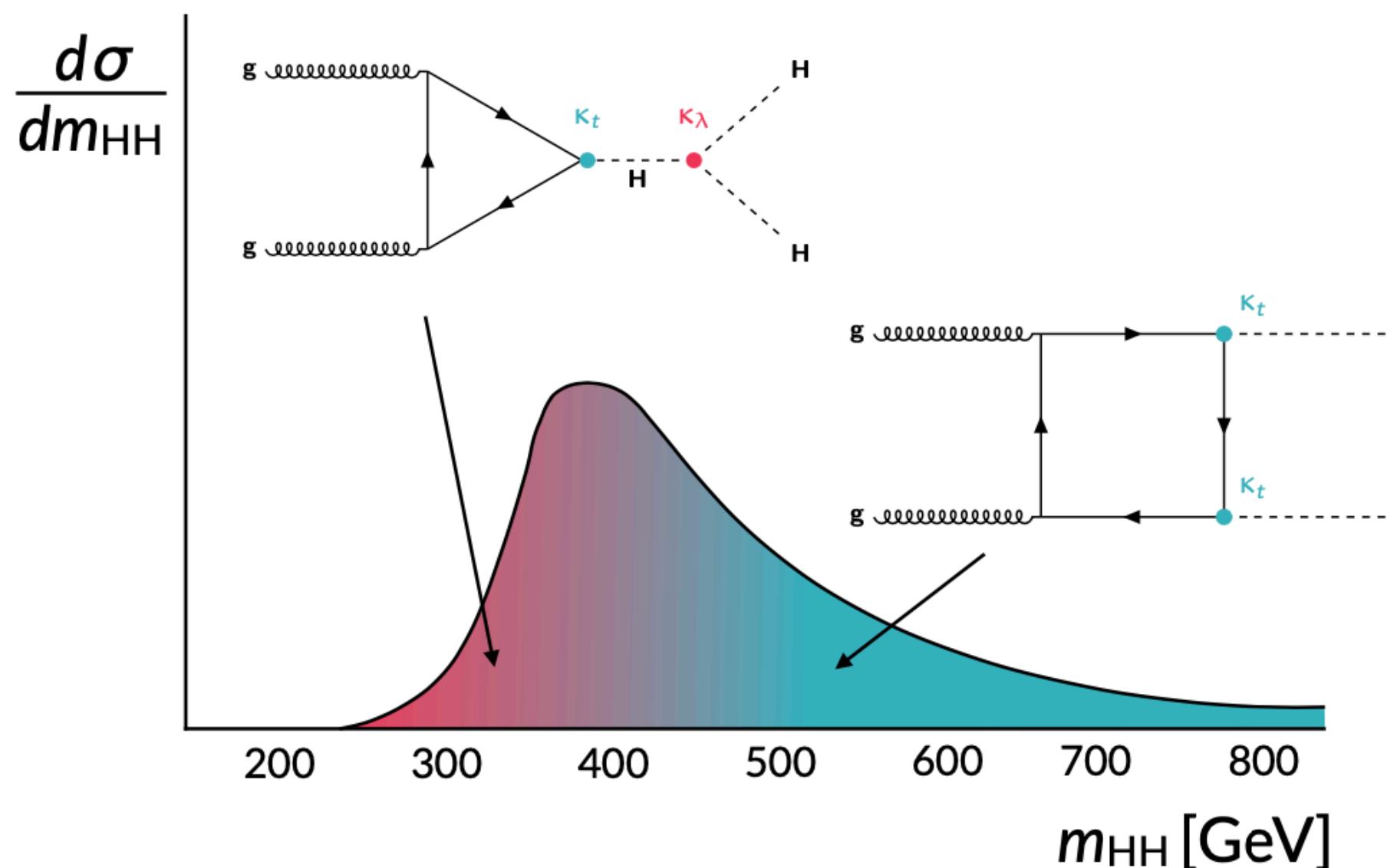
- A measurement of Higgs boson pair (HH) production could provide evidence of the **Higgs self-coupling**, which is a direct probe of the shape of the Higgs potential
- Any deviation of the self-interaction from its SM expectation is a sign of **new physics!**

Higgs boson pair production

- Gluon fusion: $\sigma_{\text{ggF}}(pp \rightarrow HH) = 31 \text{ fb}$ at 13 TeV



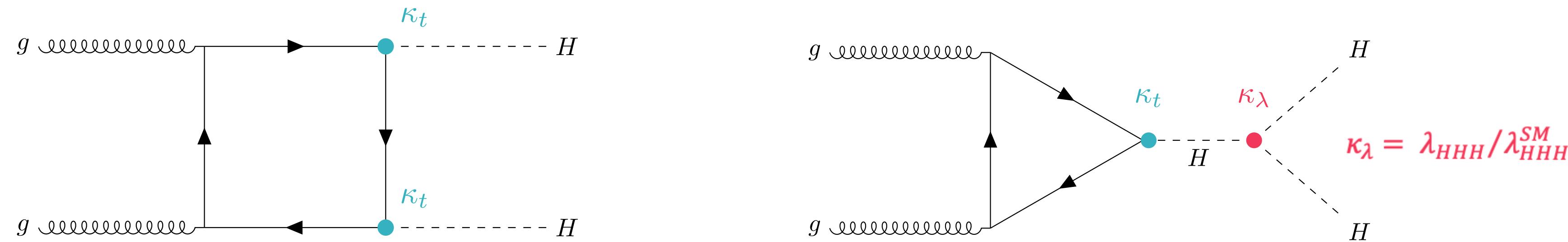
→ κ_λ variations modify the HH cross-section and kinematics w.r.t SM predictions



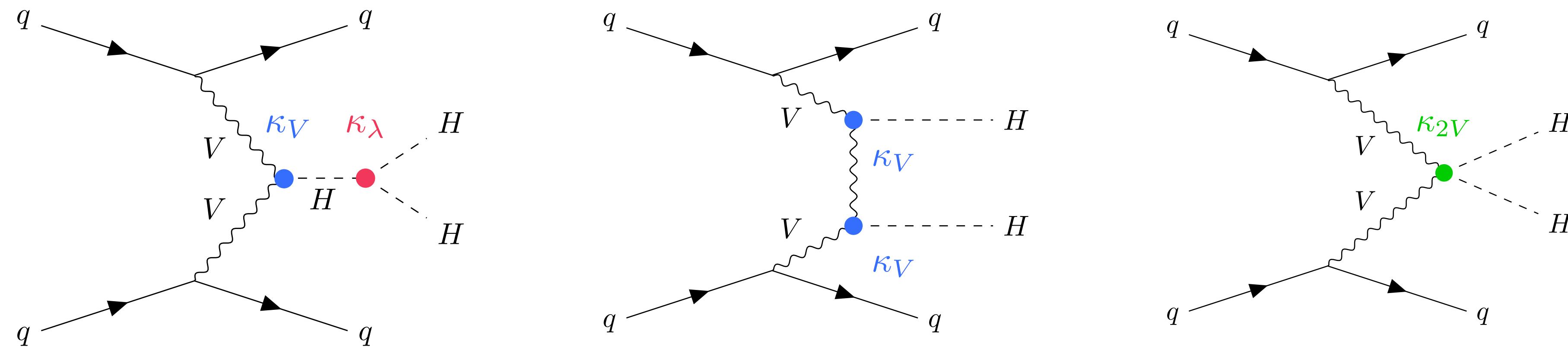
Phys. Lett. B 800 (2020) 135103

Higgs boson pair production

- Gluon fusion: $\sigma_{\text{ggF}}(pp \rightarrow HH) = 31 \text{ fb}$ at 13 TeV



- Vector-boson fusion: $\sigma_{\text{VBF}}(pp \rightarrow HH) = 1.7 \text{ fb}$ at 13 TeV



Additional sensitivity to κ_λ

Unique probe of κ_{2V} in VBF HH searches

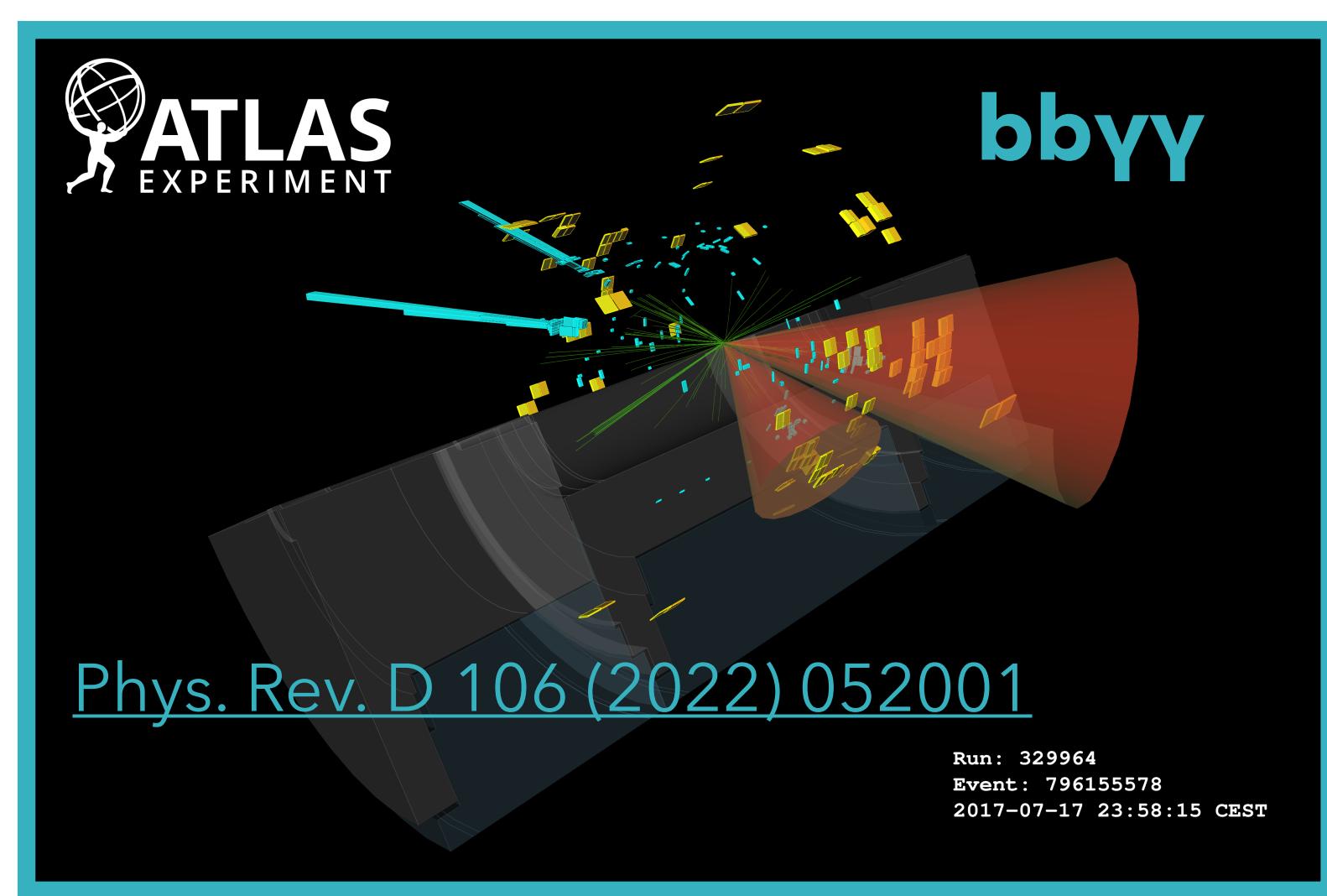
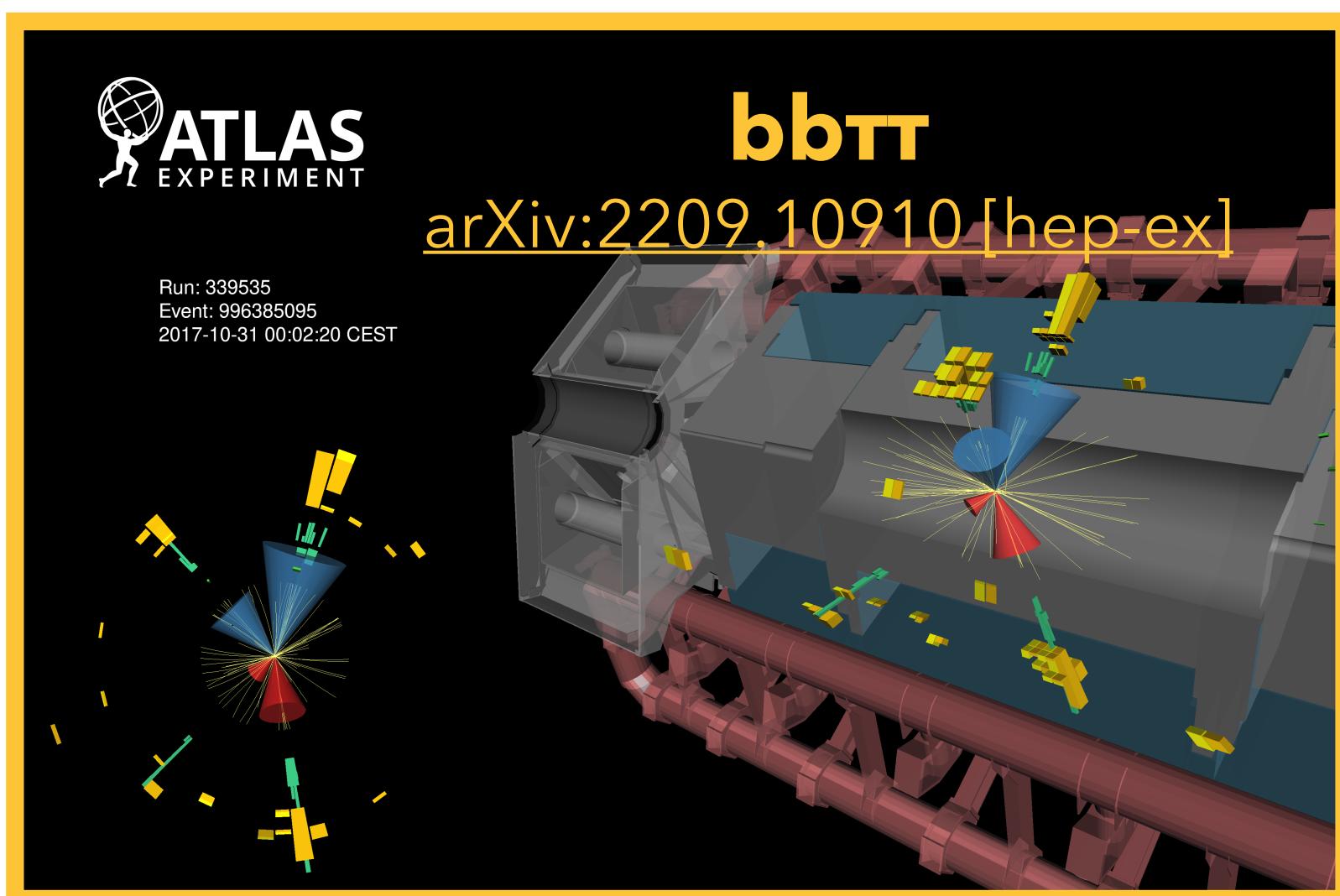
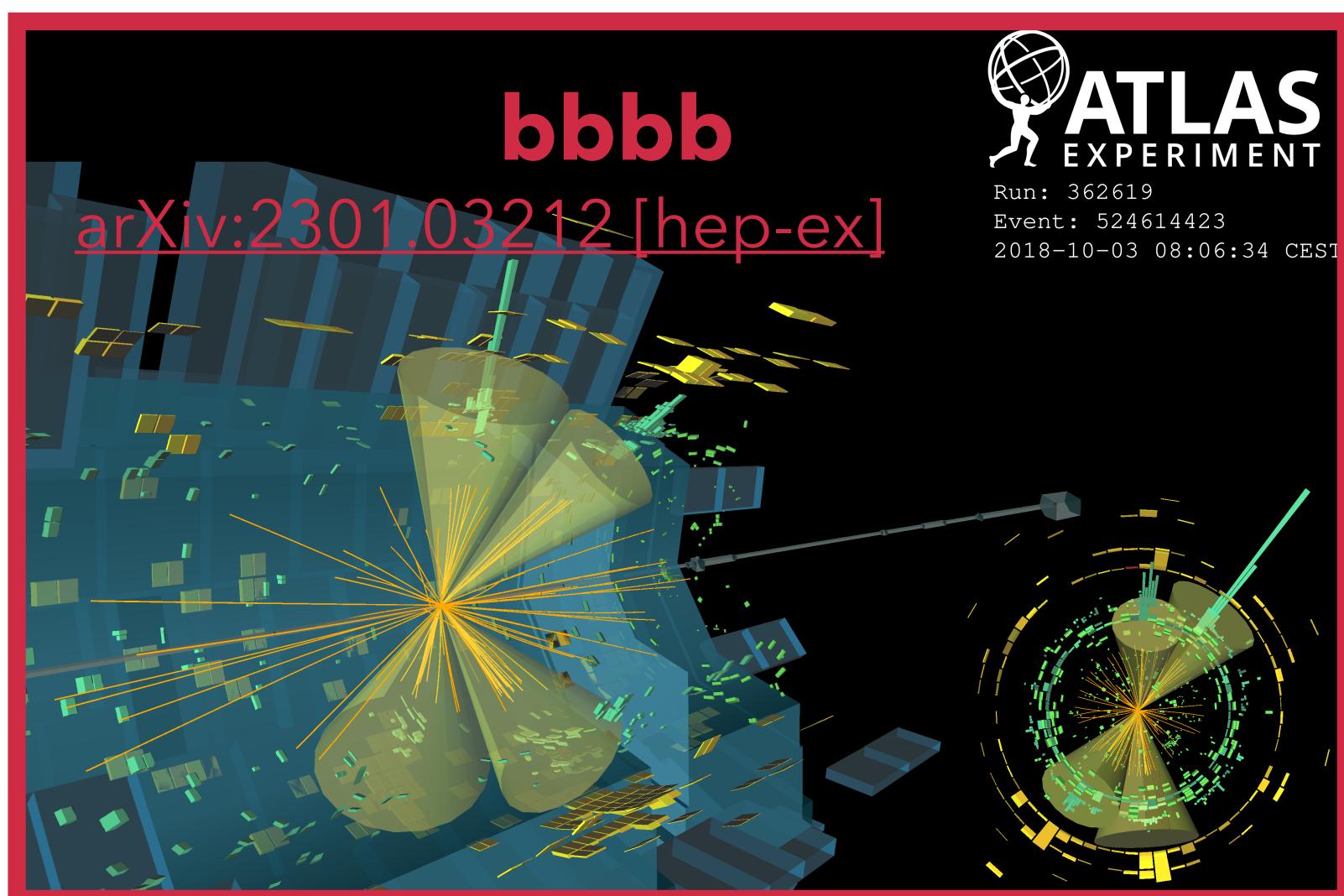
HH decay channels



Multiple decay channels, here focus on the 3 most sensitive searches and their combination:

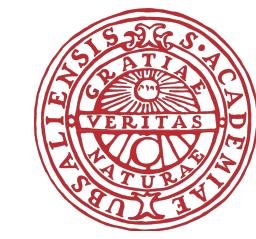
- $HH \rightarrow bbbb$: plenty of signal 👍, but challenging multi-jet background 👎
- $HH \rightarrow bb\tau\tau$: moderate signal rate, relatively clean final state ⚖️
- $HH \rightarrow bb\gamma\gamma$: very clean signature 👍, but tiny branching ratio 👎

	bb	ww	$\tau\tau$	zz	$\gamma\gamma$
bb	34%				
ww	25%	4.6%			
$\tau\tau$	7.3%	2.7%	0.39%		
zz	3.1%	1.1%	0.33%	0.069%	
$\gamma\gamma$	0.26%	0.10%	0.028%	0.012%	0.0005%

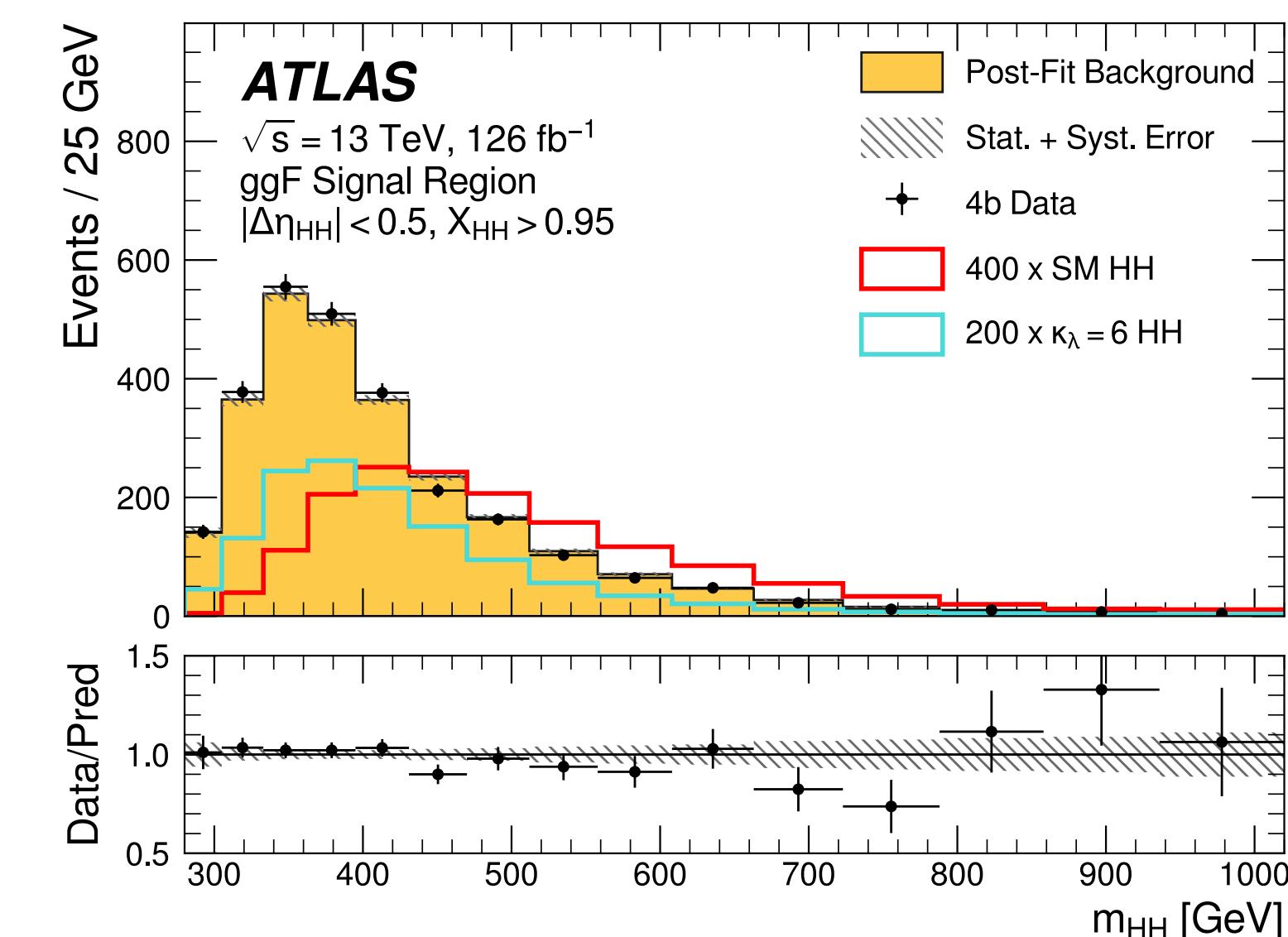
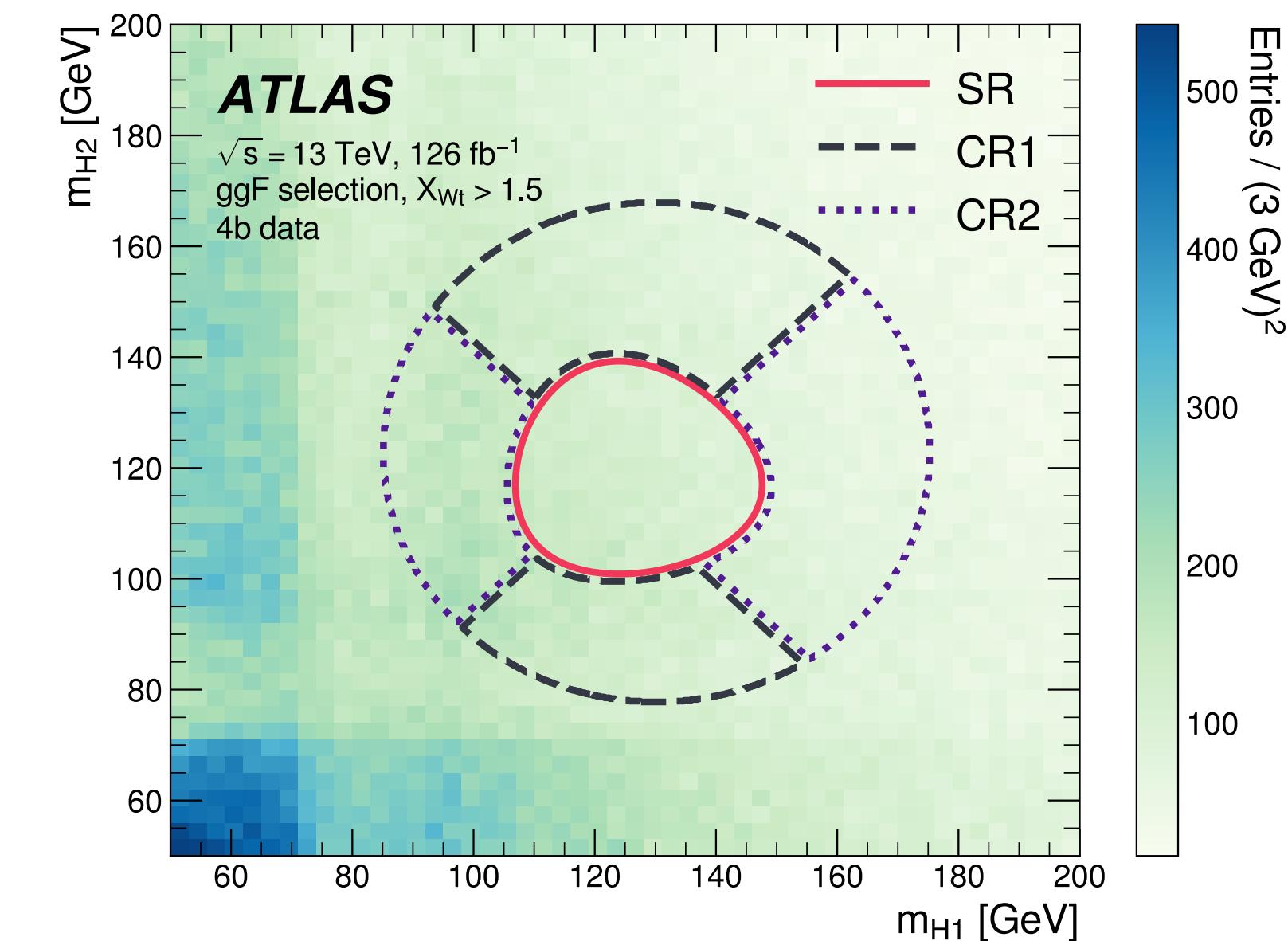


$HH \rightarrow bbbb$ search

[arXiv:2301.03212 \[hep-ex\]](https://arxiv.org/abs/2301.03212)

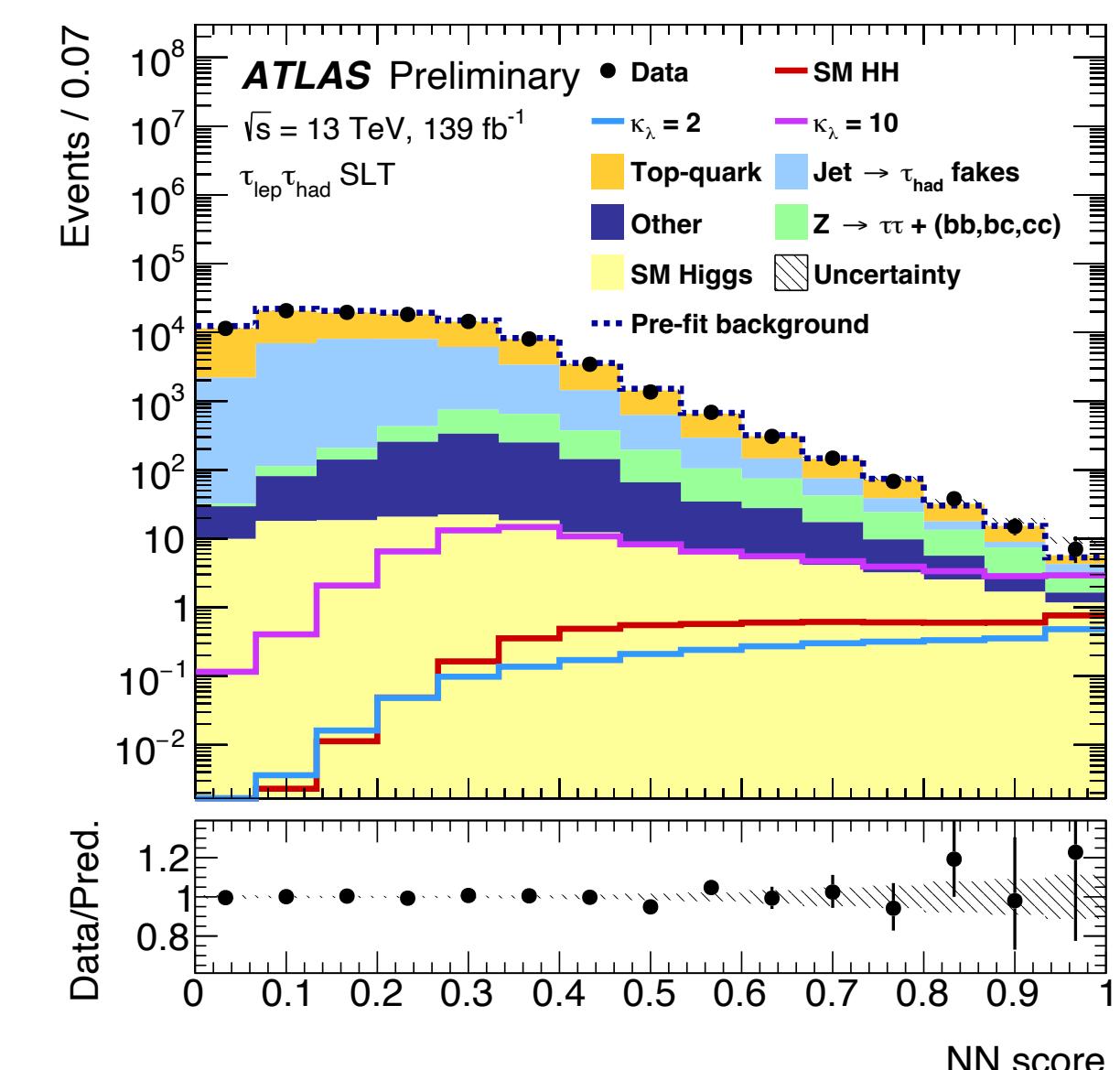
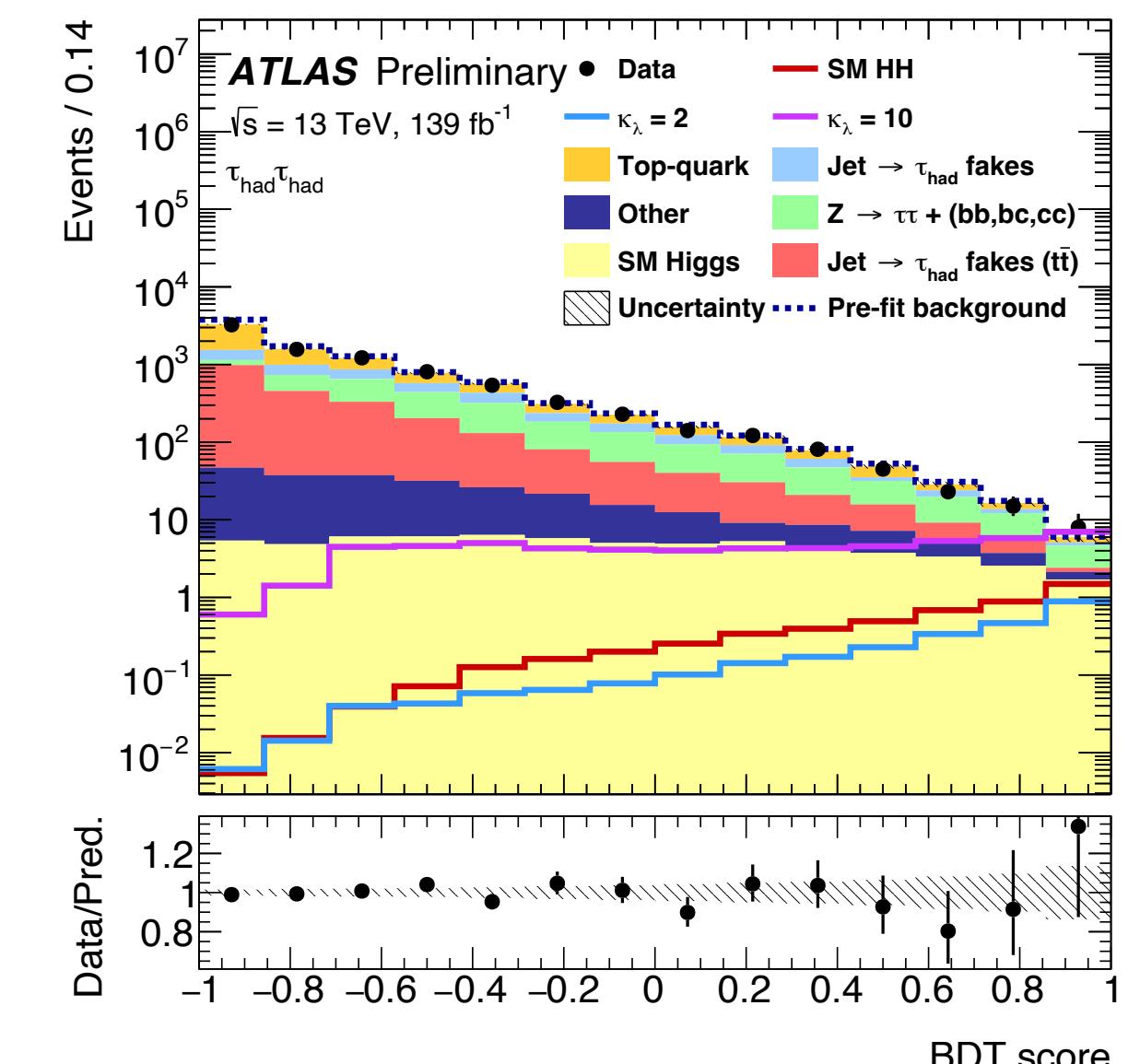


- Event selection
 - Trigger on events with at least 2 b -tagged jets
 - Two Higgs boson candidates reconstructed from 4 b -tagged jets with the smallest ΔR_{jj} pairing
 - 6 ggF and 2 VBF categories based on forward jets and HH kinematic properties
- Main backgrounds: 90% multi-jet, 10% $t\bar{t}$ events
- Data-driven background estimates
 - $2b \rightarrow 4b$ event reweighting in SR (NN-based weights from CR)
- Fit m_{HH} in all categories



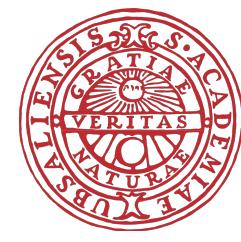
- Event selection
 - 2 sub-channels split by tau decay mode, in total 3 SRs driven by trigger
 - $b\bar{b}\tau_{\text{had}}\tau_{\text{had}} \rightarrow \text{BDT}$
 - $b\bar{b}\tau_{\text{lep}}\tau_{\text{had}}$ SLT (single-lepton triggers) $\rightarrow \text{NN}$
 - $b\bar{b}\tau_{\text{lep}}\tau_{\text{had}}$ LTT (lepton+tau triggers) $\rightarrow \text{NN}$
- Dominant backgrounds:
 - True taus from $t\bar{t}$ and $Z \rightarrow \tau\tau + \text{HF}$ (simulation)
 - Fake taus from multi-jet and $t\bar{t}$ processes (data-driven)
 - Fit MVA outputs in all categories together with m_{ll} in a $Z \rightarrow ee/\mu\mu + \text{HF}$ CR

MVA discriminants
trained with kinematic variables

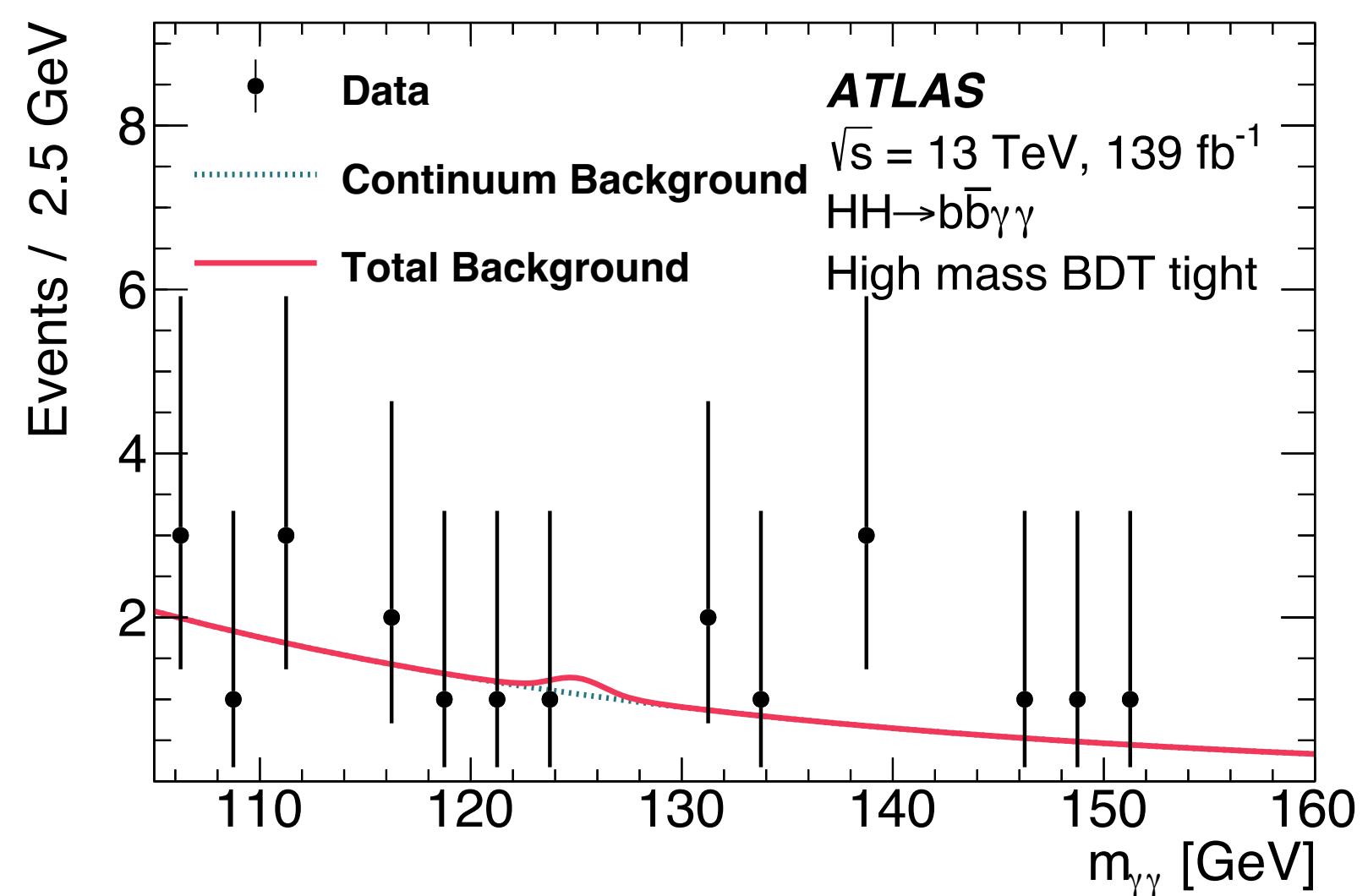
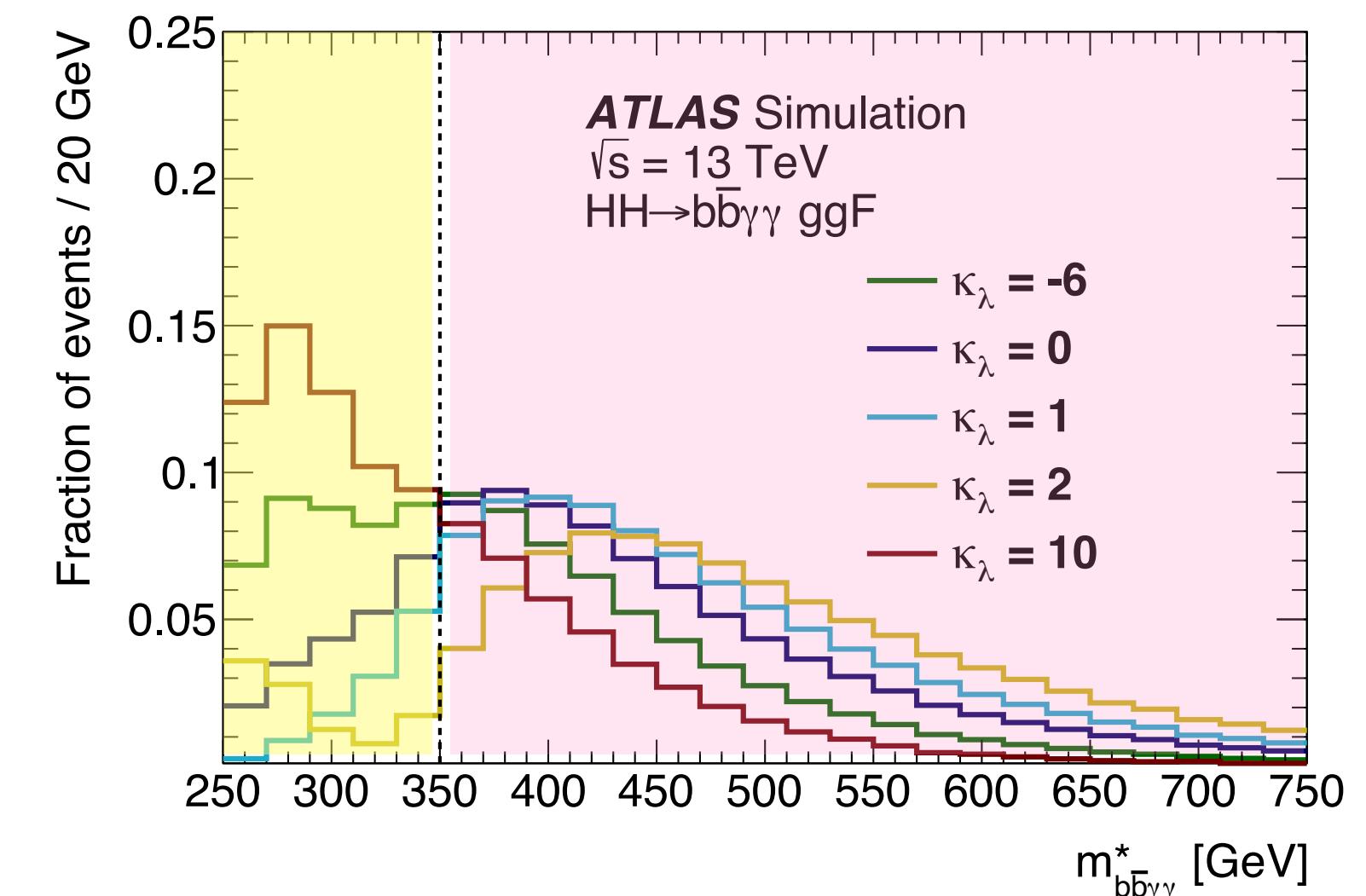


$\text{HH} \rightarrow \text{bb}\gamma\gamma$ search

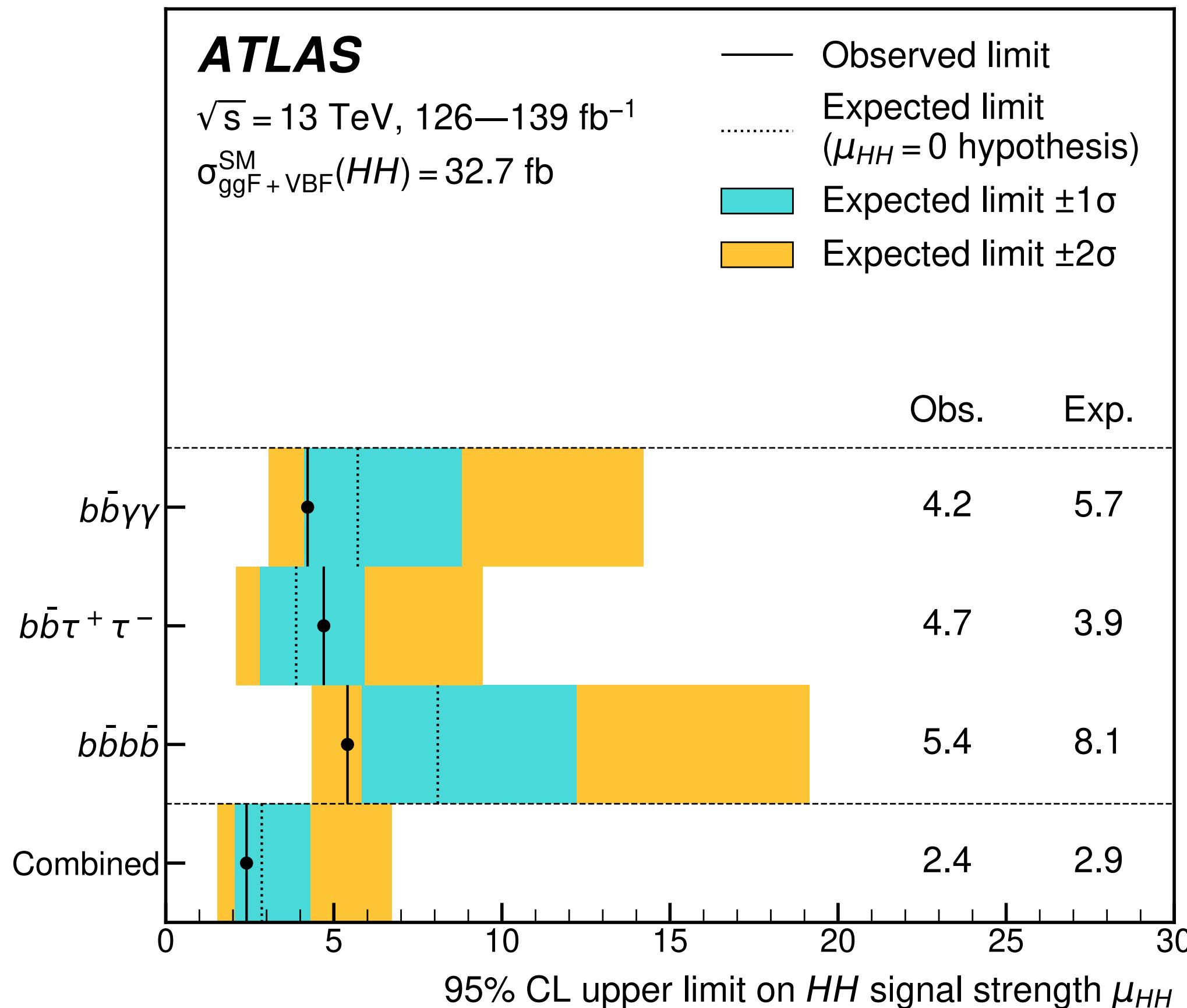
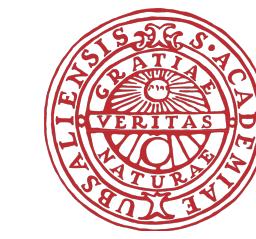
[Phys. Rev. D 106 \(2022\) 052001](#)



- Event selection
 - At least 2 photons and exactly 2 b -tagged jets
 - Split events into low- and high-mass regions to target different signal hypotheses
 - Separate BDTs in those two regions
 - Categorise into loose- and tight-BDT SRs for each mass region → 4 SRs in total
- Main backgrounds: $\gamma\gamma$ continuum and single Higgs boson
- Fit $m_{\gamma\gamma}$ in all categories



Combination: $HH \rightarrow bbbb + bb\tau\tau + bb\gamma\gamma$



- Statistical combination of the three most sensitive HH decay channels (full Run-2 data)
- Yielded the best observed limits yet on HH production

Observed (Expected) 95% CL limit: 2.4 (2.9) \times SM prediction

Factor ~ 3.4 improvement compared to 36 fb^{-1} results

([Phys. Lett. B 800 \(2020\) 135103](#))

→ improved analysis design, object reconstruction & identification

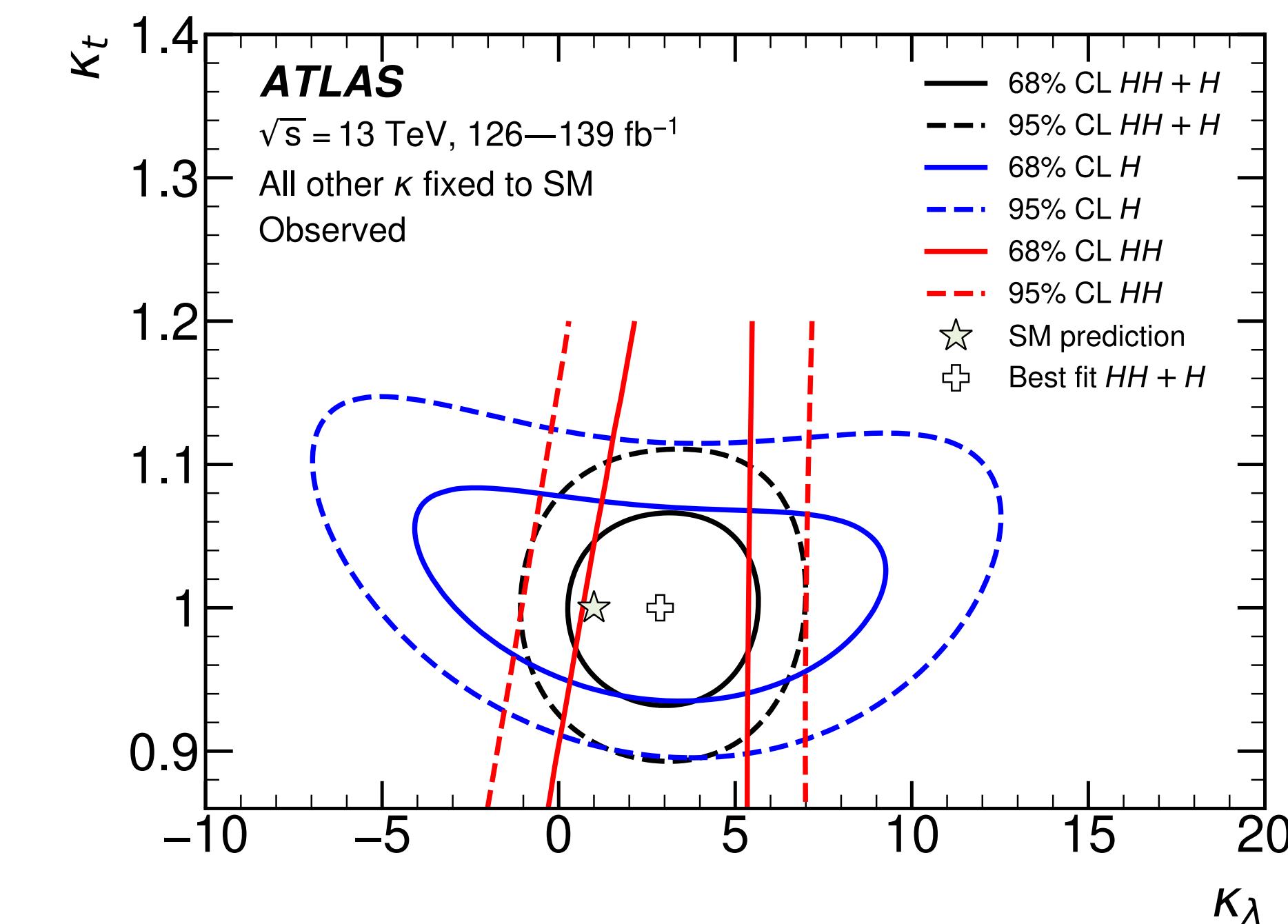
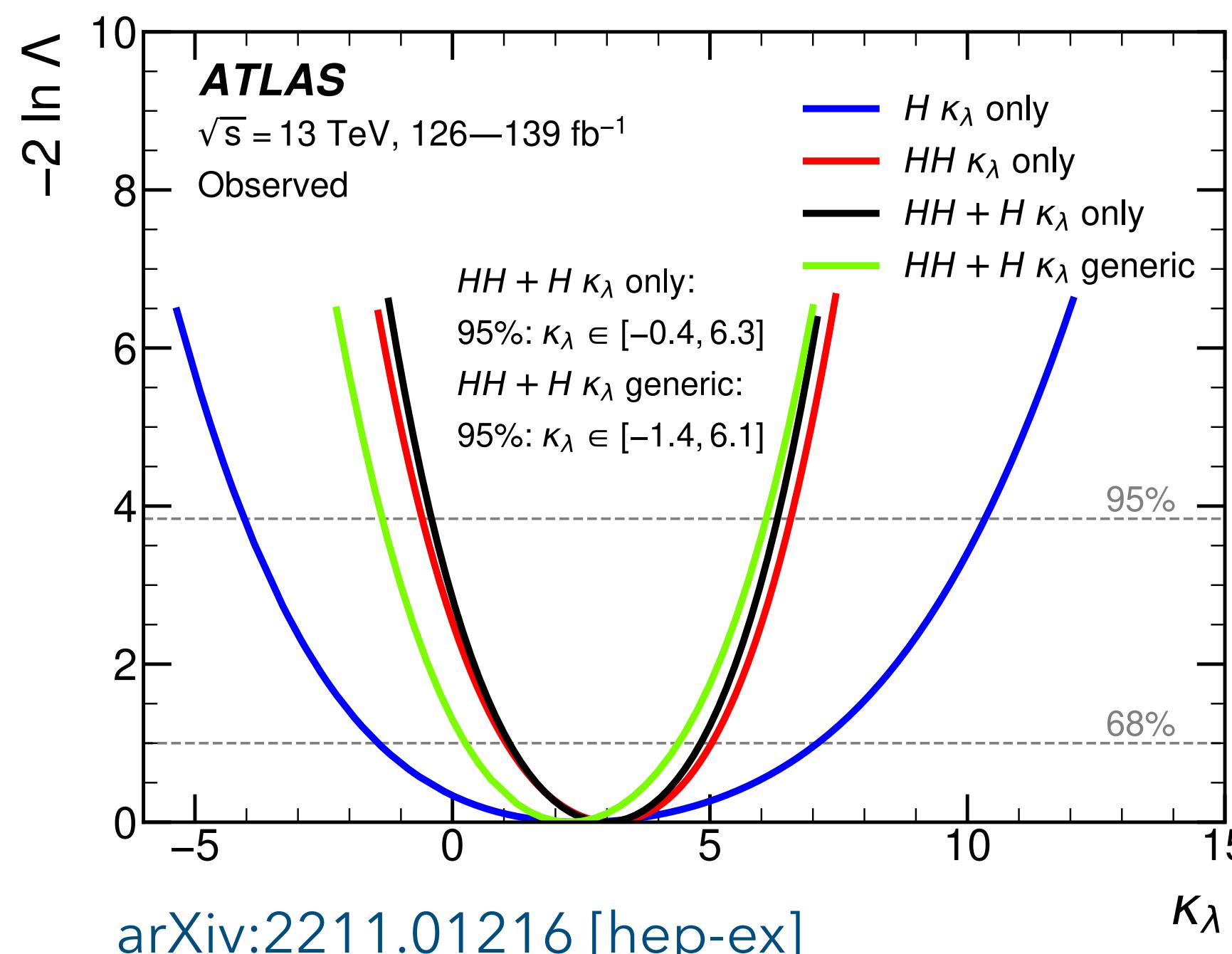
[arXiv:2211.01216 \[hep-ex\]](#)

HH+H constraints

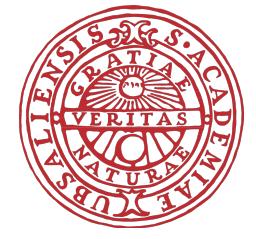
- Single Higgs boson processes also sensitive to κ_λ through NLO electroweak loop corrections



- Better constraints achieved from the combination of $HH + H$ searches
 - Combination with single-Higgs allows for a less model-dependent interpretation



HH in Effective Field Theories

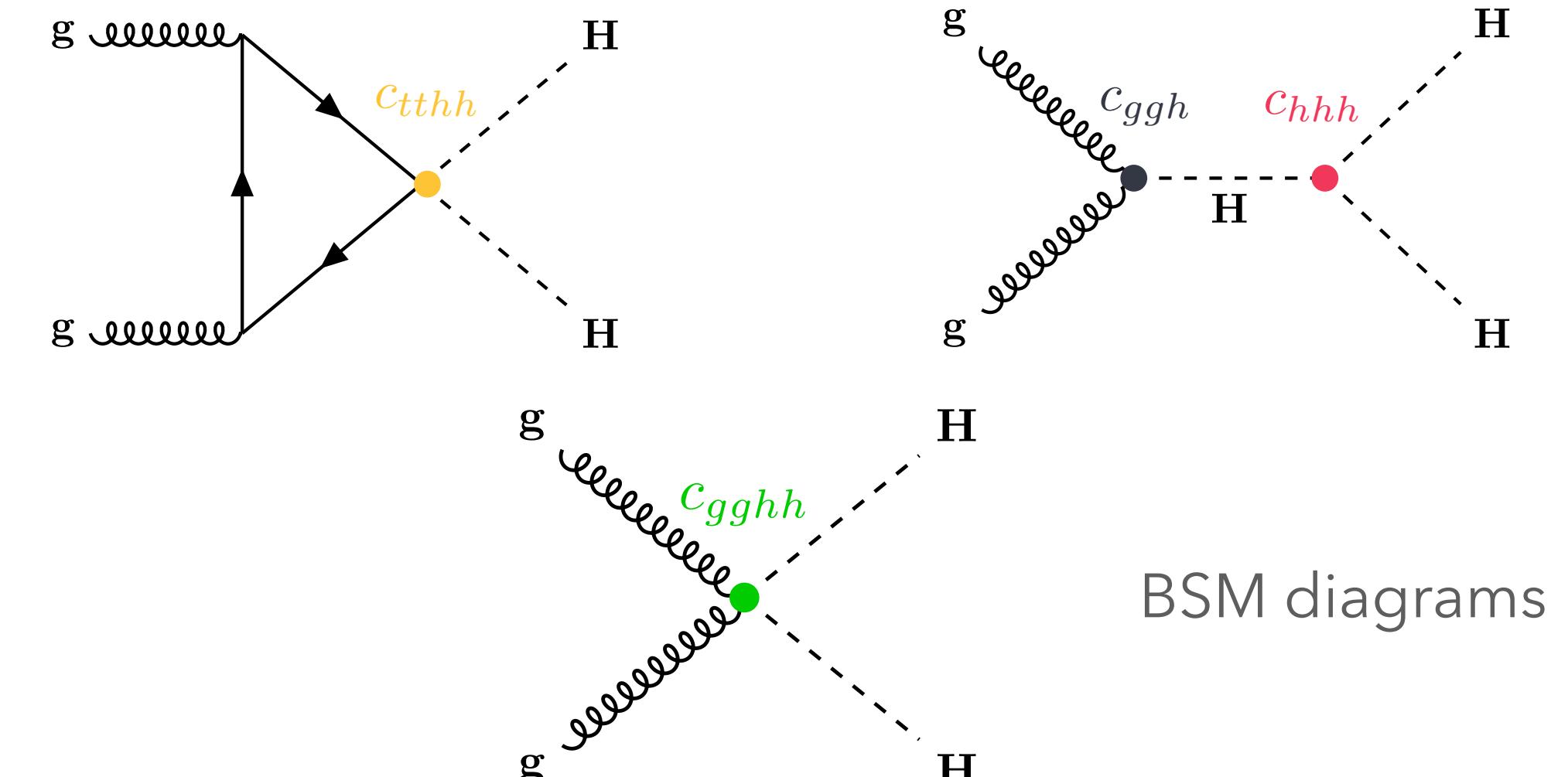


- New physics effects can be parametrised without strong model-dependence by EFTs
- Two common frameworks in HH : SMEFT and Higgs EFT (HEFT)
 - Underlying different assumptions (e.g. SMEFT uses SM symmetries)
- Effective operators can modify the $gg \rightarrow HH$ production in various ways

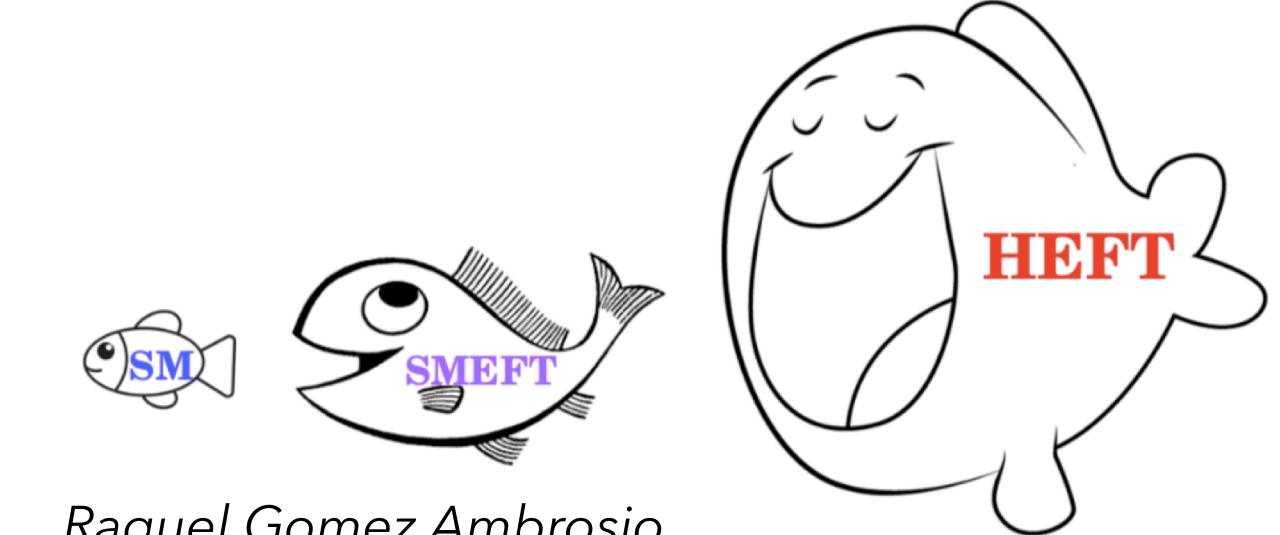
$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{n,i} \frac{c_i^{(n)}}{\Lambda^{n-4}} \mathcal{O}_i^{(n)}$$

Wilson Coefficient	Operator
c_H	$(H^\dagger H)^3$
$c_{H\square}$	$(H^\dagger H)\square(H^\dagger H)$
c_{tH}	$(H^\dagger H)(\bar{Q}\tilde{H}t)$
c_{HG}	$H^\dagger H G_A^{\mu\nu} G_A^{\mu\nu}$
c_{tG}	$(\bar{Q}\sigma^{\mu\nu} T^A t)\tilde{H} G_A^{\mu\nu}$

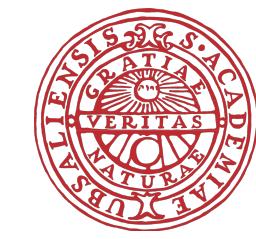
$$\mathcal{L}_{\text{HEFT}} \supset -m_t \left(\textcolor{blue}{C_{tth}} \frac{h}{v} + \textcolor{orange}{C_{tthh}} \frac{h^2}{v^2} \right) \bar{t}t - \textcolor{red}{C_{hhh}} \frac{m_h^2}{2v} h^3 + \frac{\alpha_s}{8\pi} \left(c_{ggg} \frac{h}{v} + \textcolor{green}{C_{gggh}} \frac{h^2}{v^2} \right) G_{\mu\nu}^a G^{a,\mu\nu}$$



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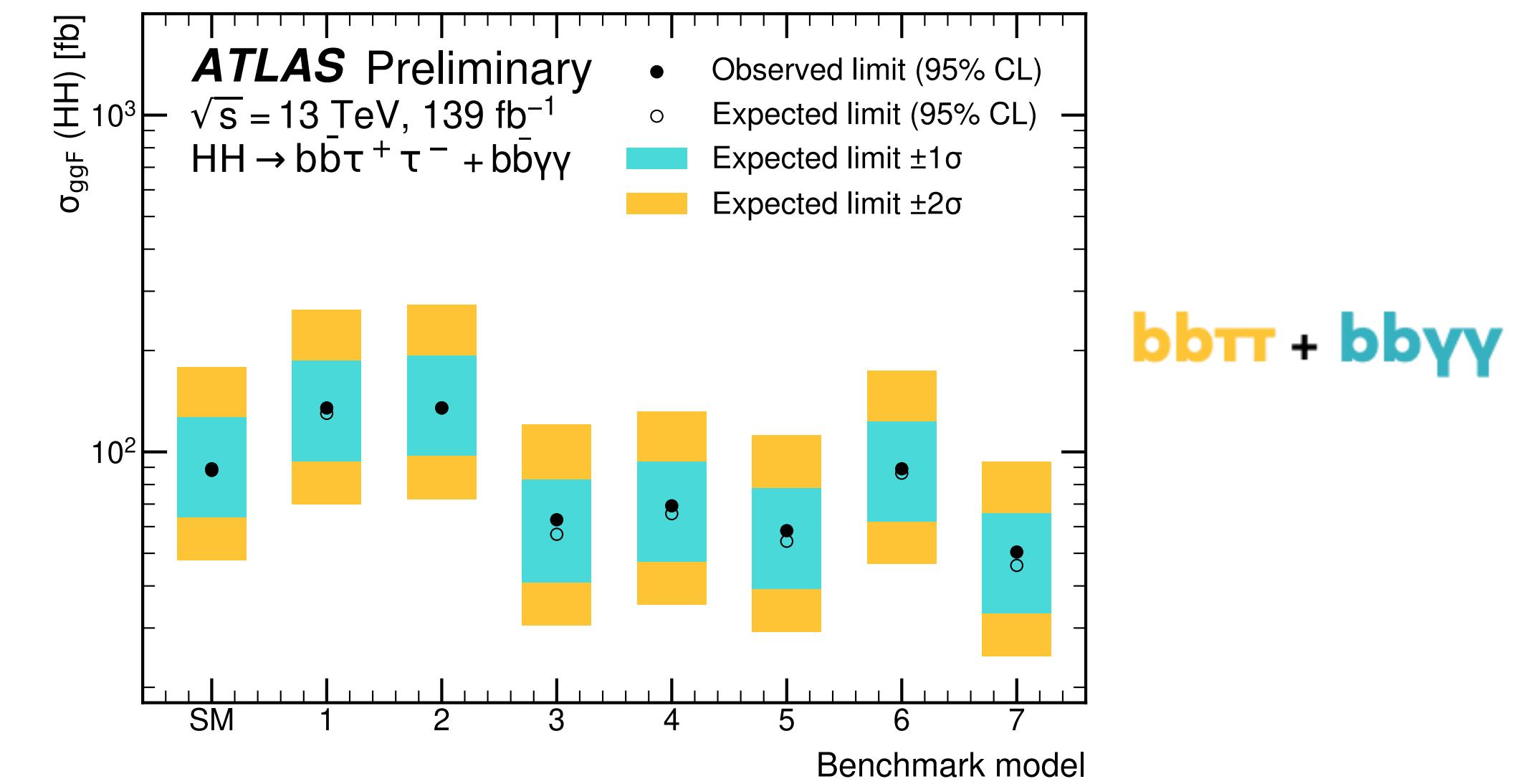
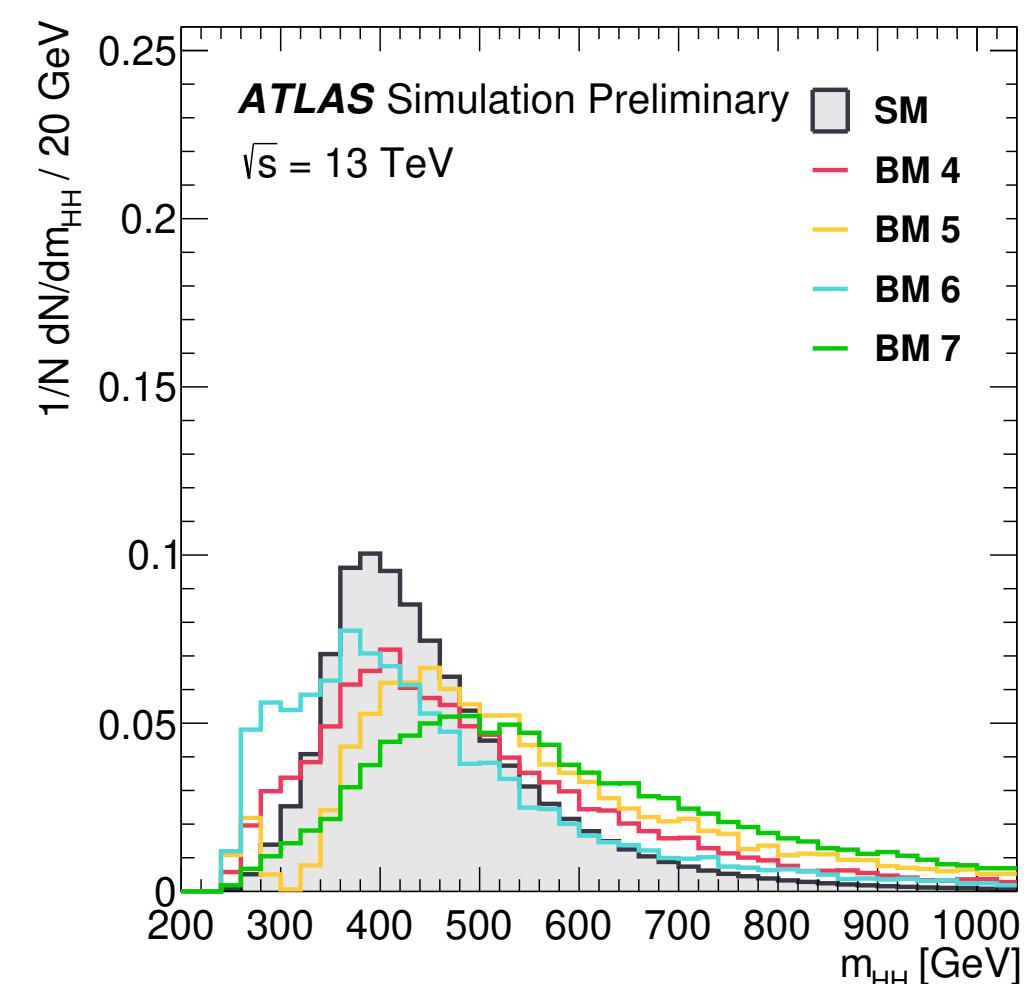
EFT re-interpretations of HH searches



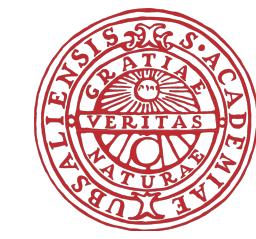
- Upper limits on seven HEFT benchmarks with characteristic m_{HH} shape features

[ATL-PHYS-PUB-2022-019](#)

Benchmark model	c_{hhh}	c_{tth}	$c_{gg\bar{h}}$	c_{gghh}	c_{tthh}
SM	1	1	0	0	0
BM 1	3.94	0.94	1/2	1/3	-1/3
BM 2	6.84	0.61	0.0	-1/3	1/3
BM 3	2.21	1.05	1/2	1/2	-1/3
BM 4	2.79	0.61	-1/2	1/6	1/3
BM 5	3.95	1.17	1/6	-1/2	-1/3
BM 6	5.68	0.83	-1/2	1/3	1/3
BM 7	-0.10	0.94	1/6	-1/6	1



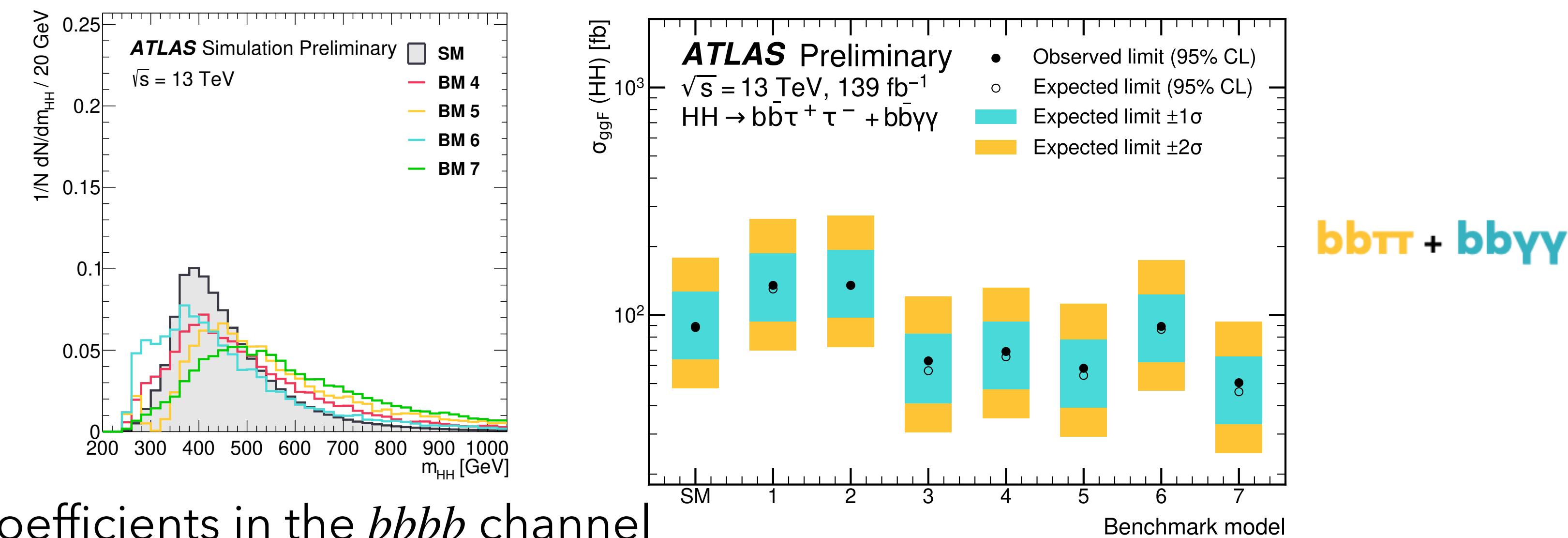
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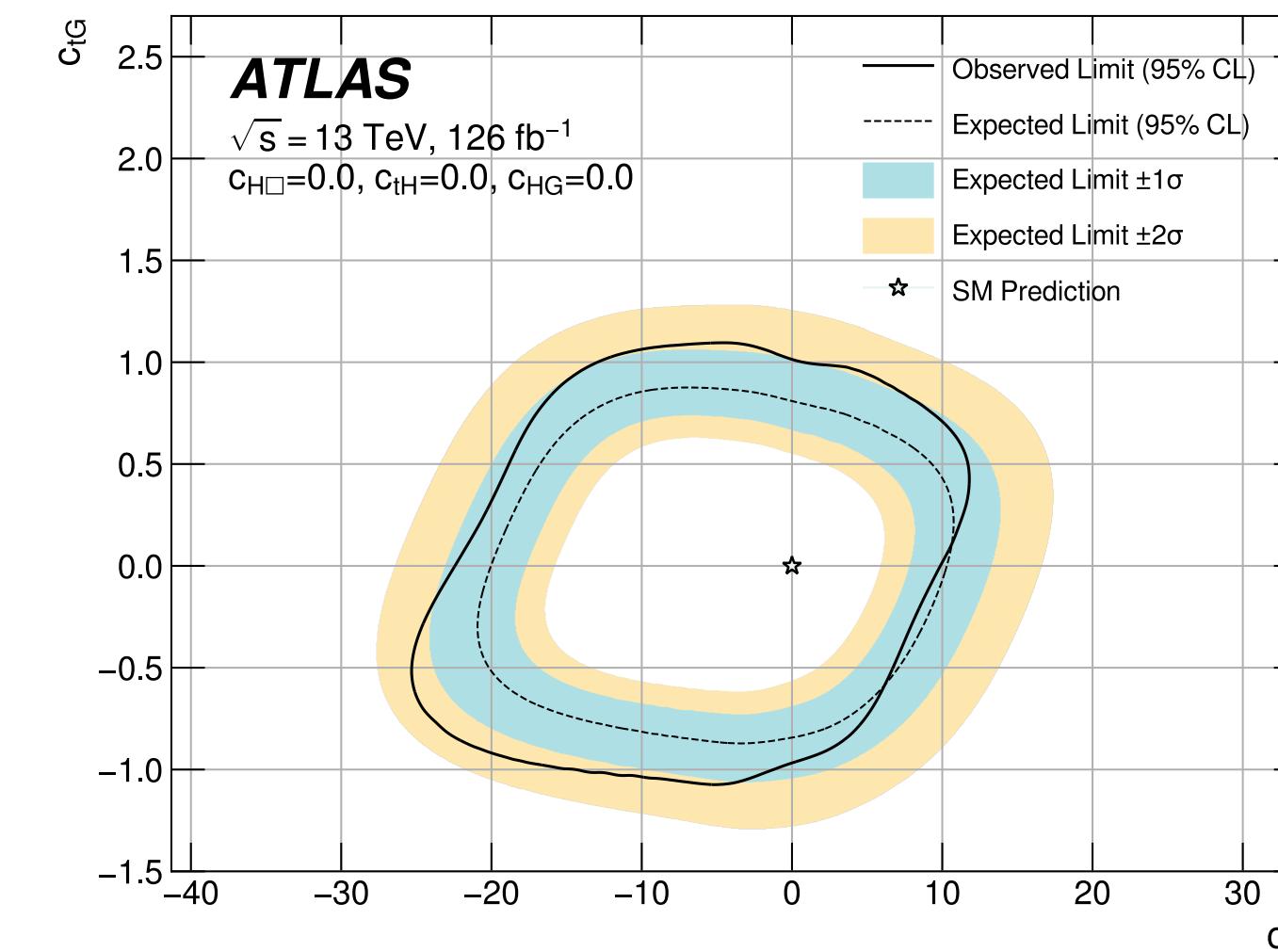
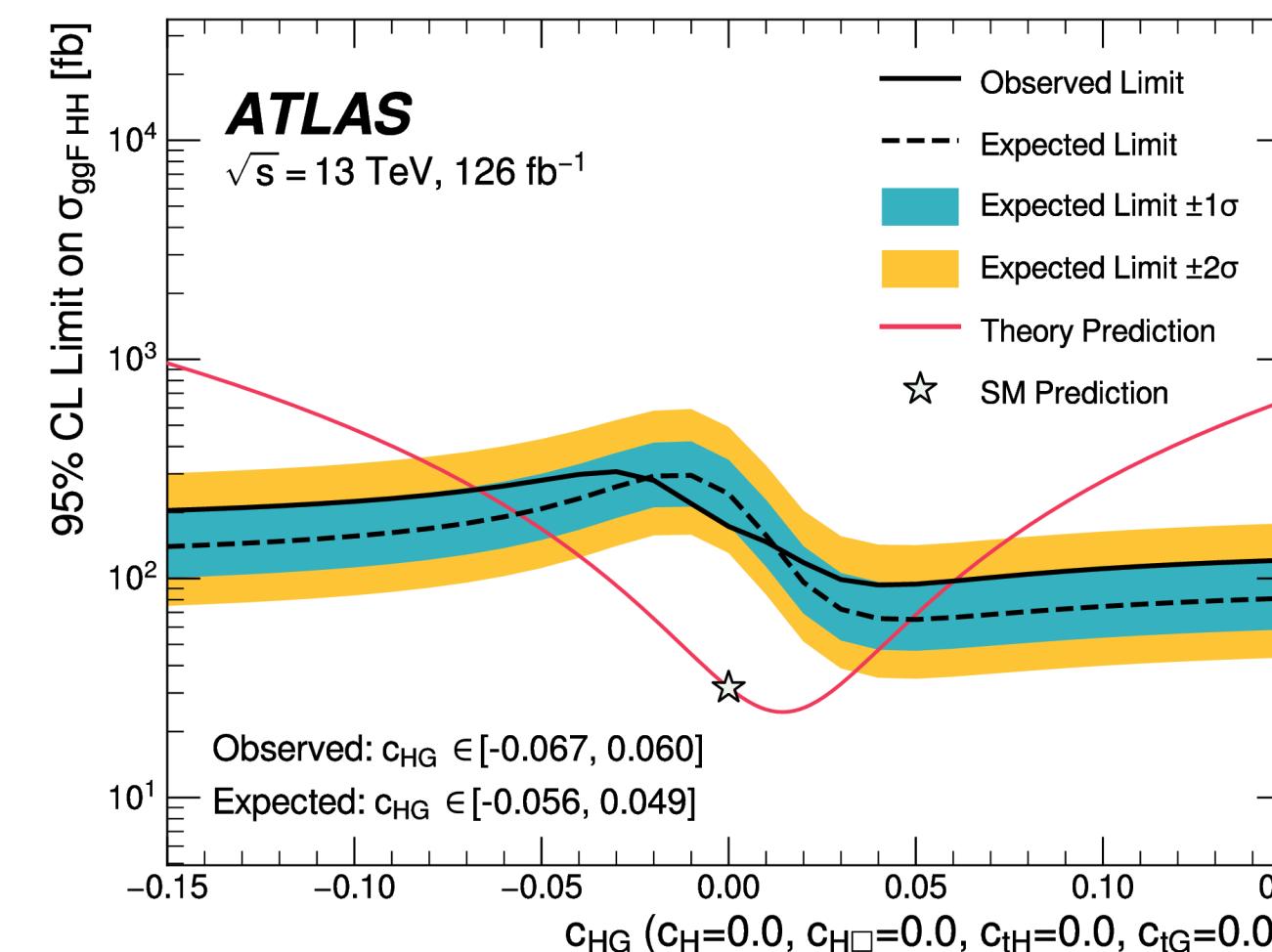
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[ATL-PHYS-PUB-2022-019](#)

Benchmark model	c_{hhh}	c_{tth}	$c_{gg\bar{h}}$	$c_{gg\bar{h}h}$	$c_{t\bar{t}hh}$
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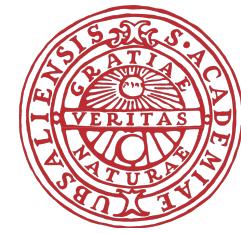
- First constraints on SMEFT Wilson coefficients in the $bbbb$ channel



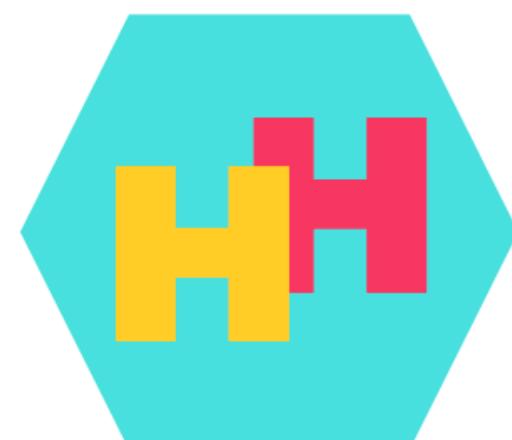
bbbb

[arXiv:2301.03212 \[hep-ex\]](#)

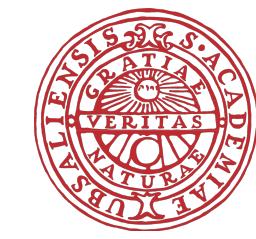
Summary



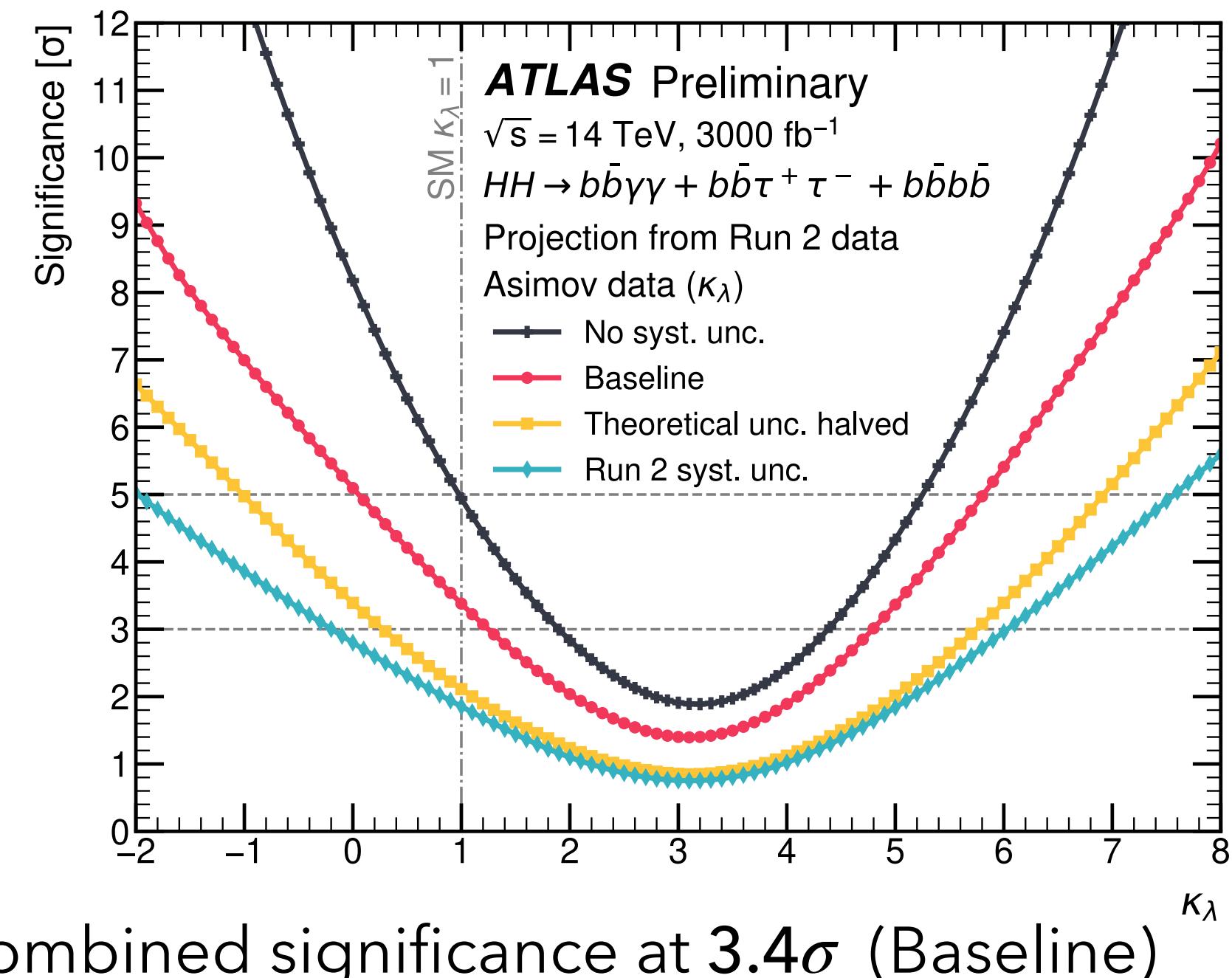
- Searching for Higgs boson pairs is the **ultimate probe of the shape of the Higgs potential** and thereby of the nature of electroweak symmetry breaking
- No single golden channel - **combinations** are the key
 - **95% CL upper limit** on μ_{HH} at **2.4** \times SM prediction
 - **$HH + H$** combination allows for more robust interpretations
- **EFT interpretations** becoming increasingly popular (HEFT and SMEFT results in ATLAS)
- Further studies performed in parallel (e.g. VBF κ_{2V} constraints, VHH search)
 - All public results can be found in <https://twiki.cern.ch/twiki/bin/view/AtlasPublic/HDBSPublicResults>



Outlook - HL-LHC prospects

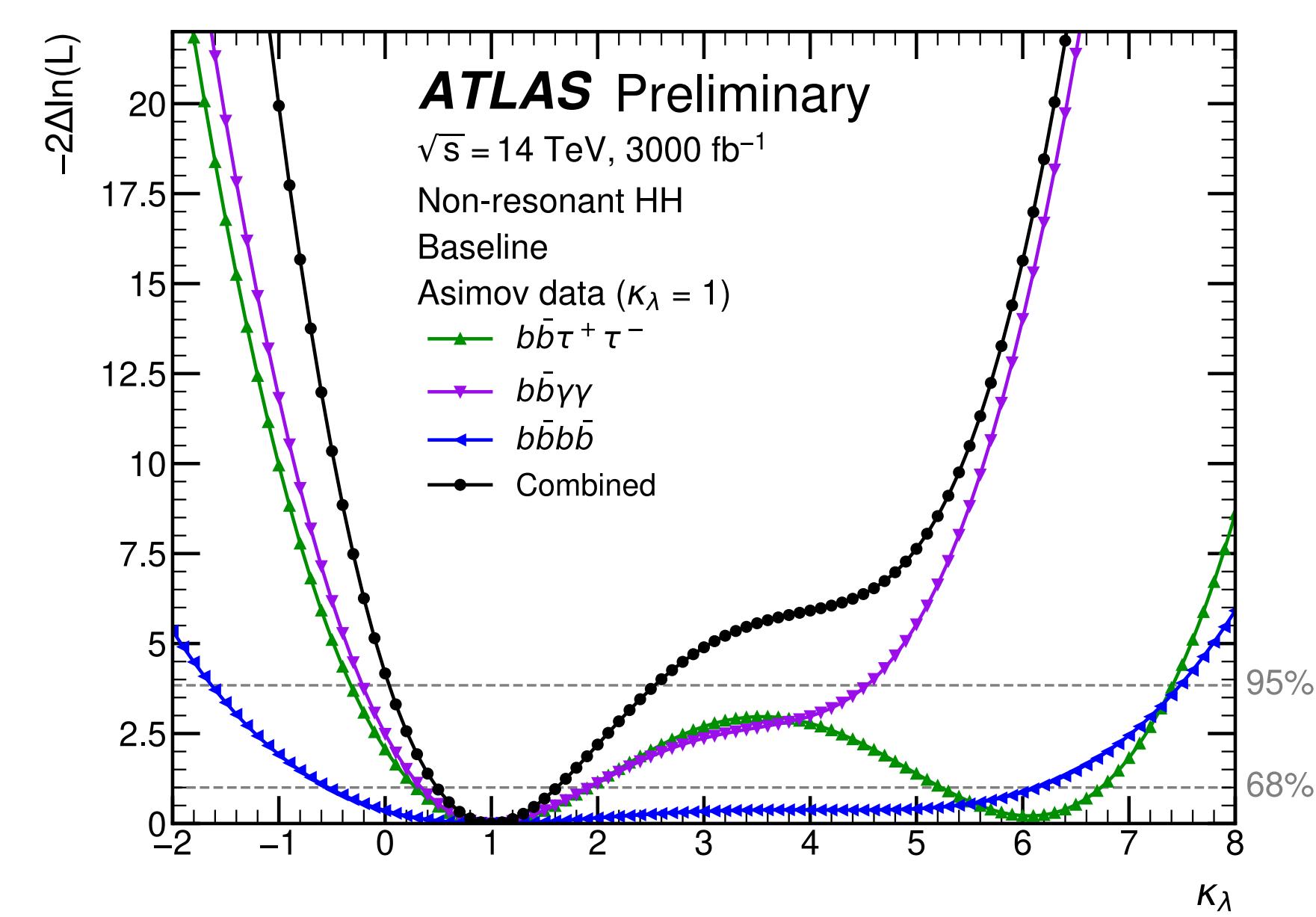


- HL-LHC critical for observing and measuring low cross-section processes like HH
- Prospect studies through extrapolation of the Run-2 results to 3000 fb^{-1} and 14 TeV
- Consider different systematic uncertainty scenarios. "Baseline" \rightarrow best guess of the future uncertainties



Combined significance at 3.4σ (Baseline)

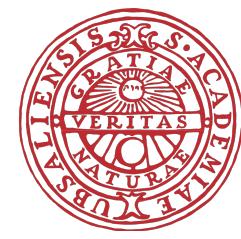
Evidence is reachable!



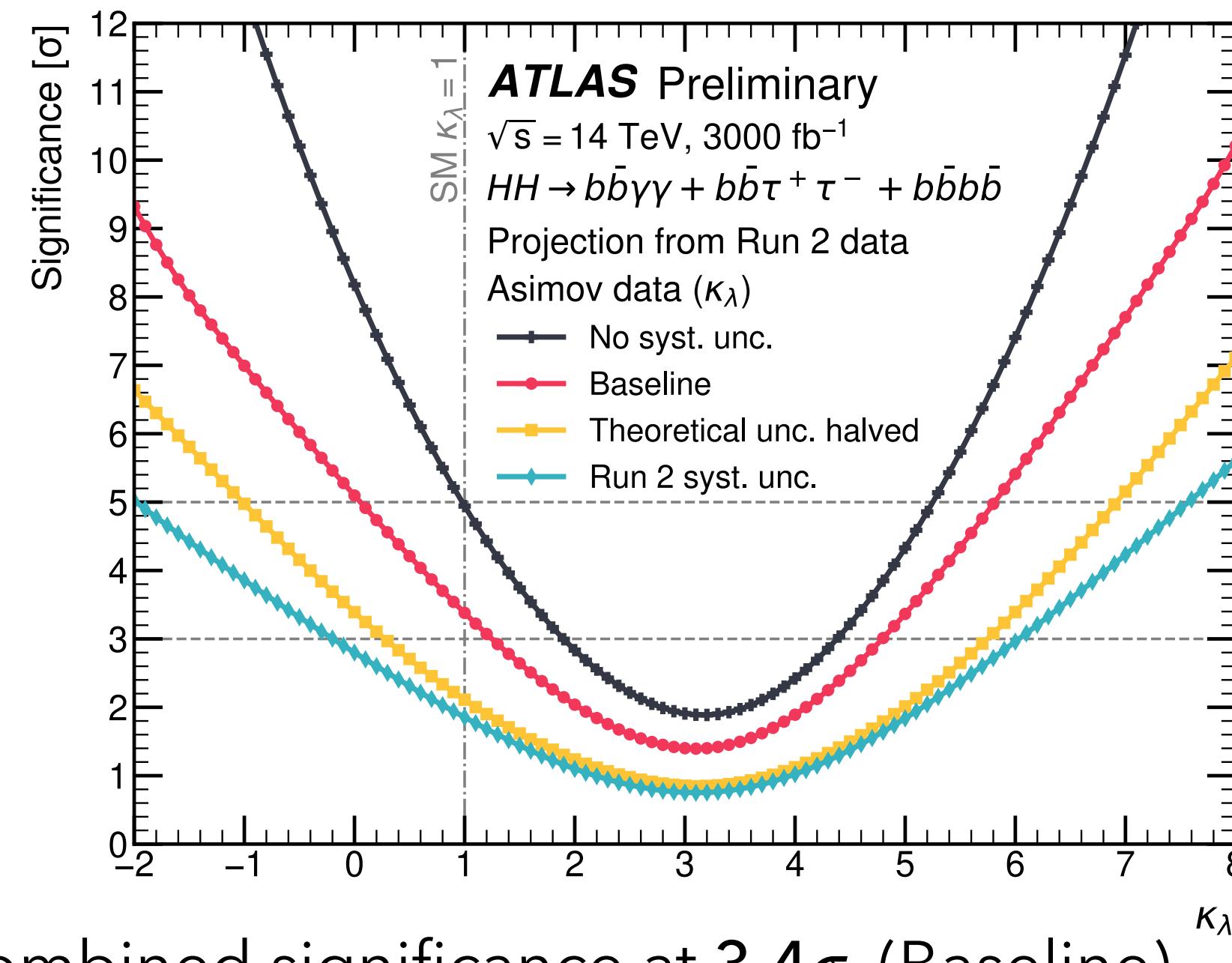
68% confidence interval for κ_λ at $[0.5, 1.6]$ (Combined - Baseline)

[ATL-PHYS-PUB-2022-053](#)

Outlook - HL-LHC prospects



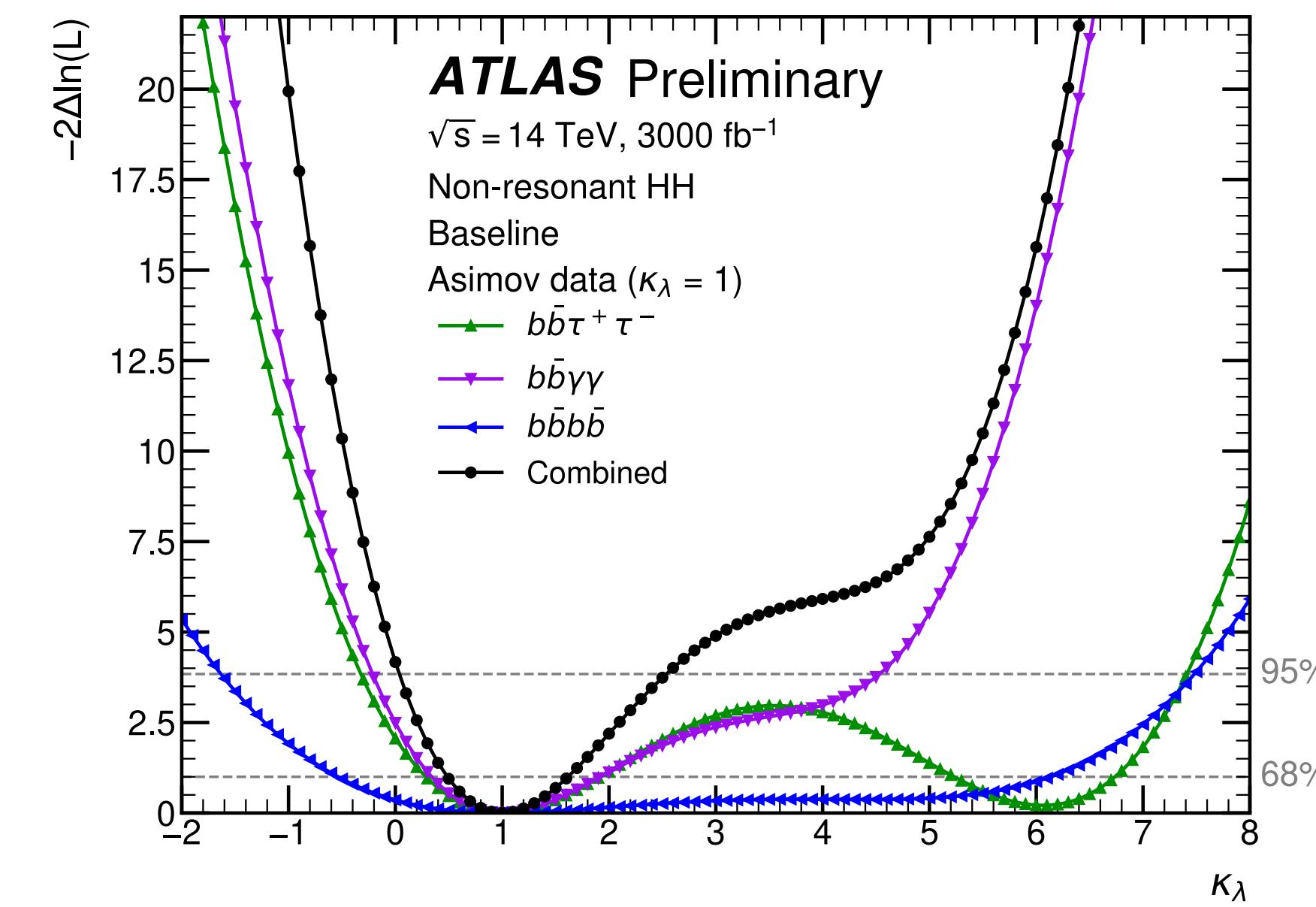
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Combined significance at 3.4σ (Baseline)

Evidence is reachable!

Observation is likely possible with an ATLAS + CMS combination

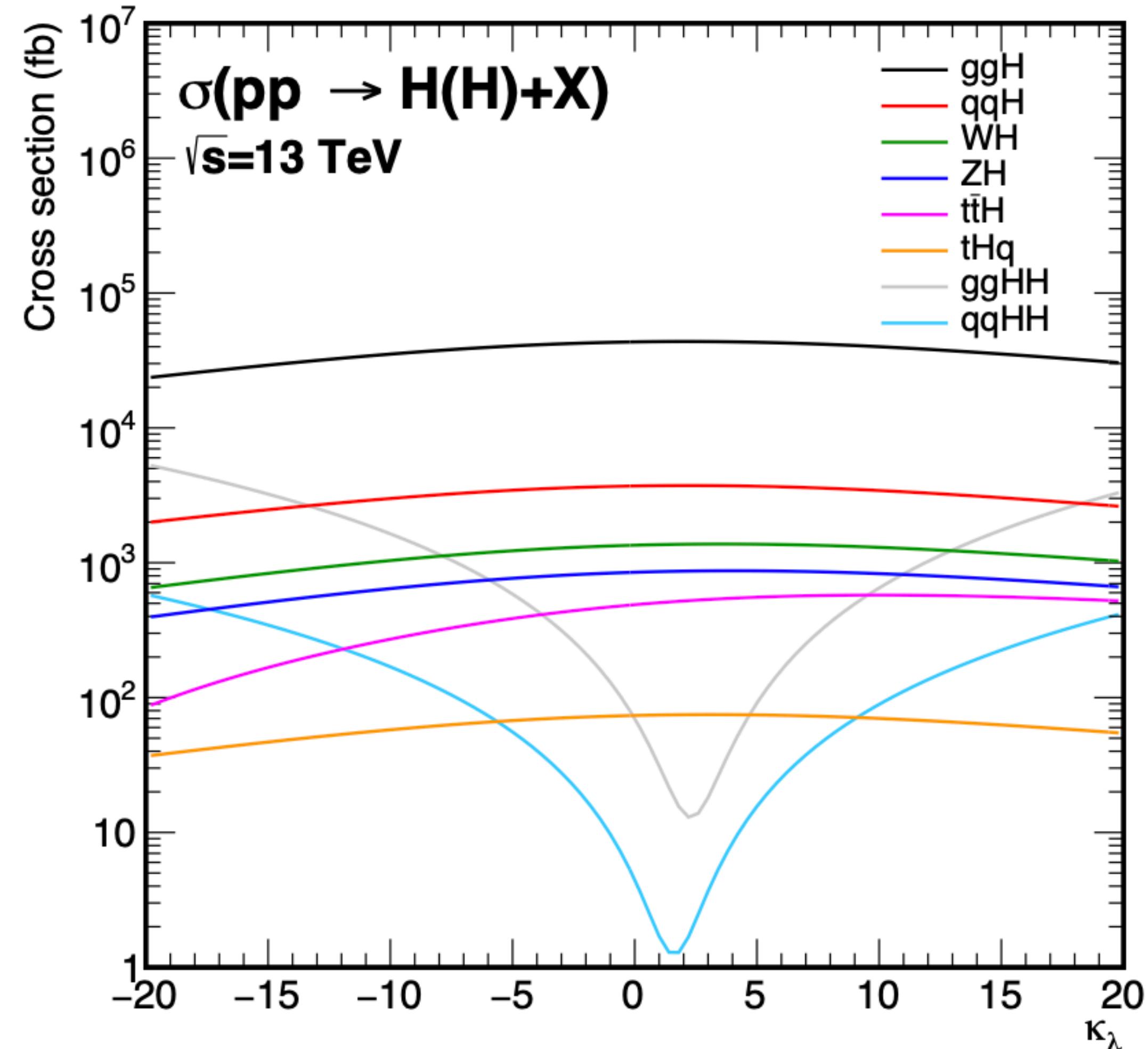


68% confidence interval for κ_λ at $[0.5, 1.6]$ (Combined - Baseline)

[ATL-PHYS-PUB-2022-053](#)

Back-up

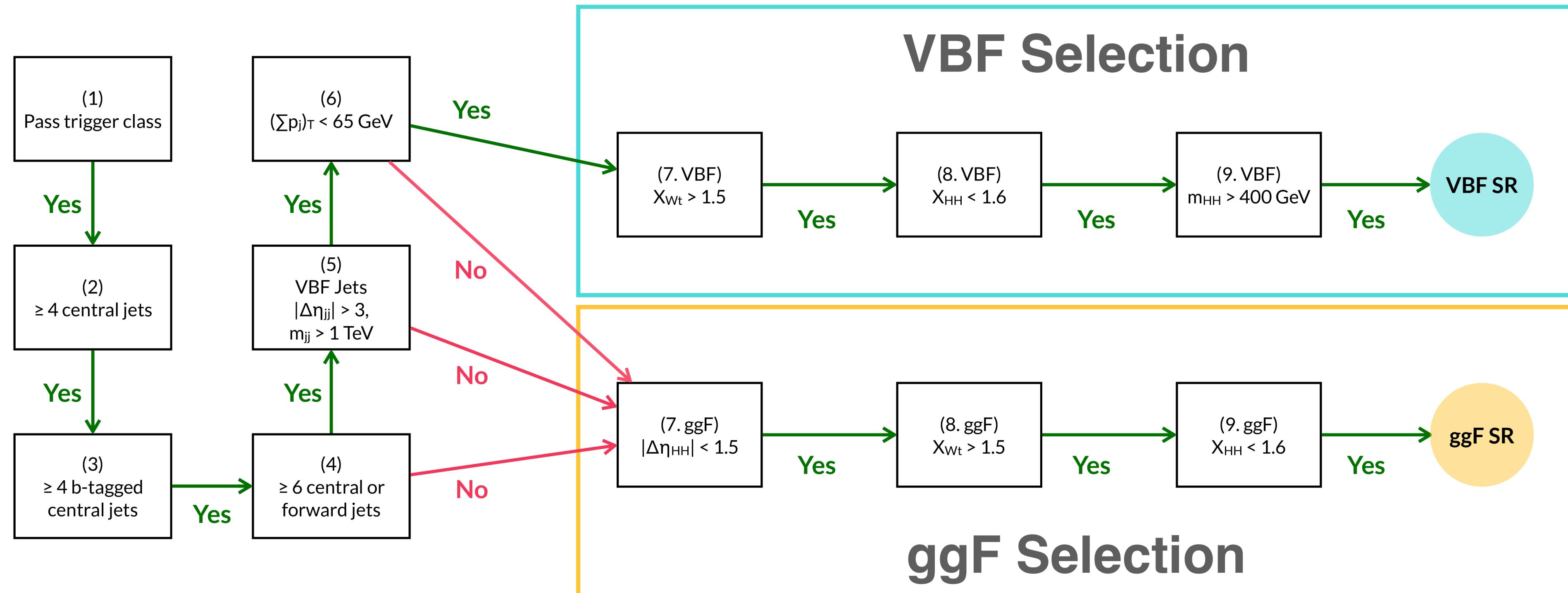
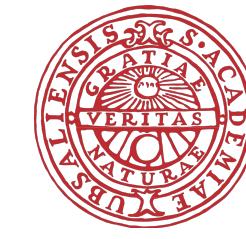
H and HH cross-sections vs. κ_λ



[LHCHWG-2022-002](#)

HH \rightarrow bbbb [event selection]

[arXiv:2301.03212 \[hep-ex\]](https://arxiv.org/abs/2301.03212)

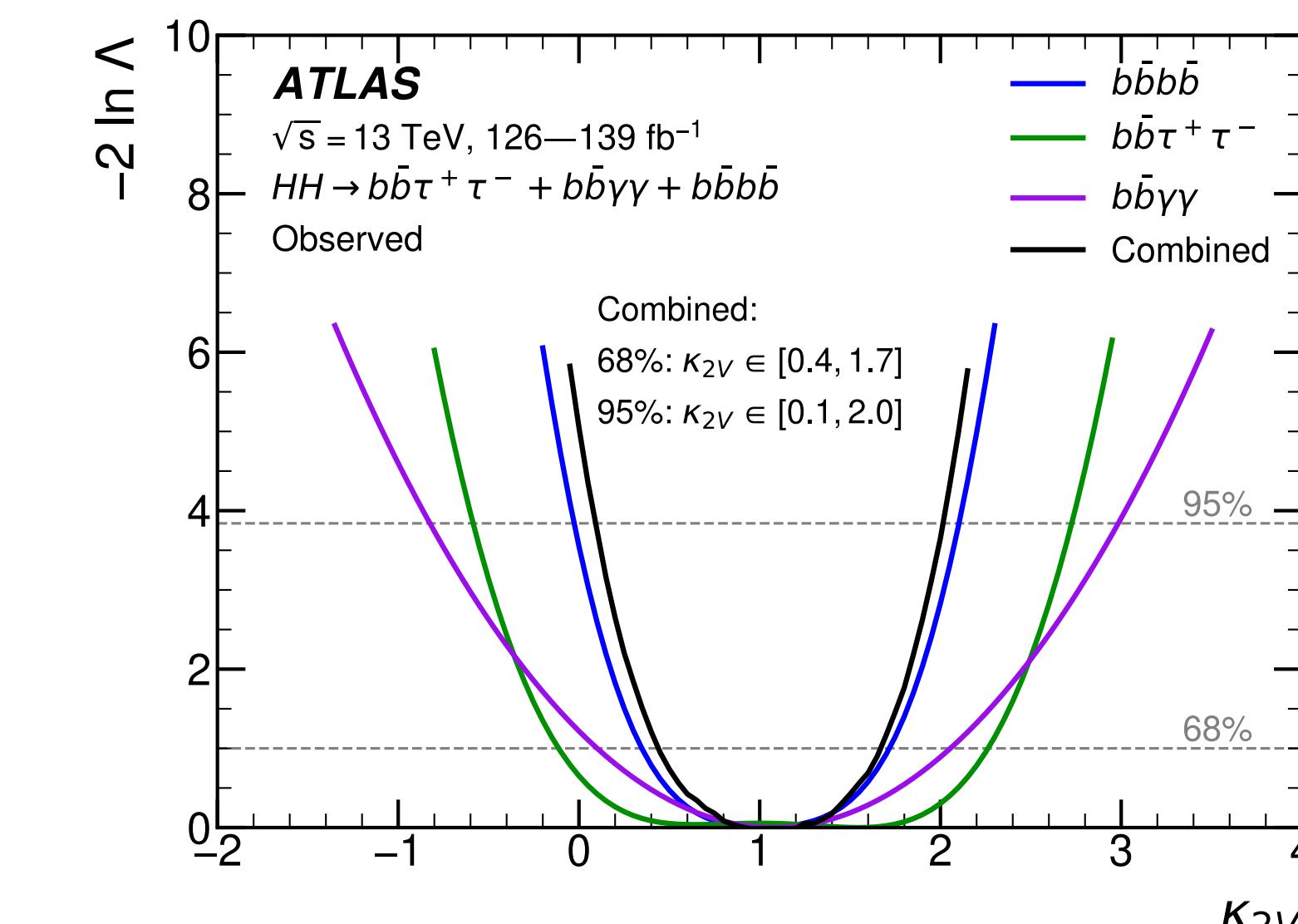
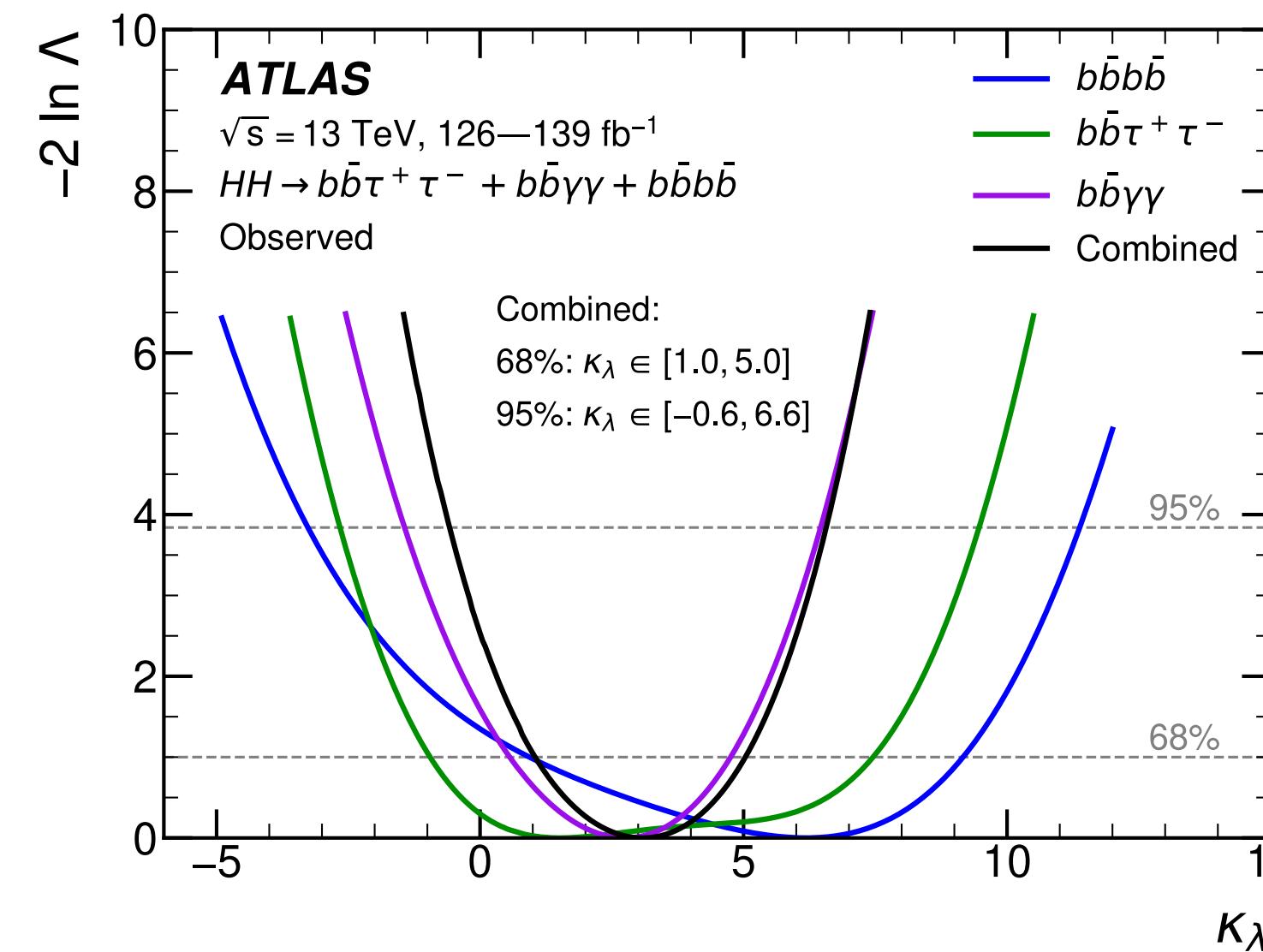
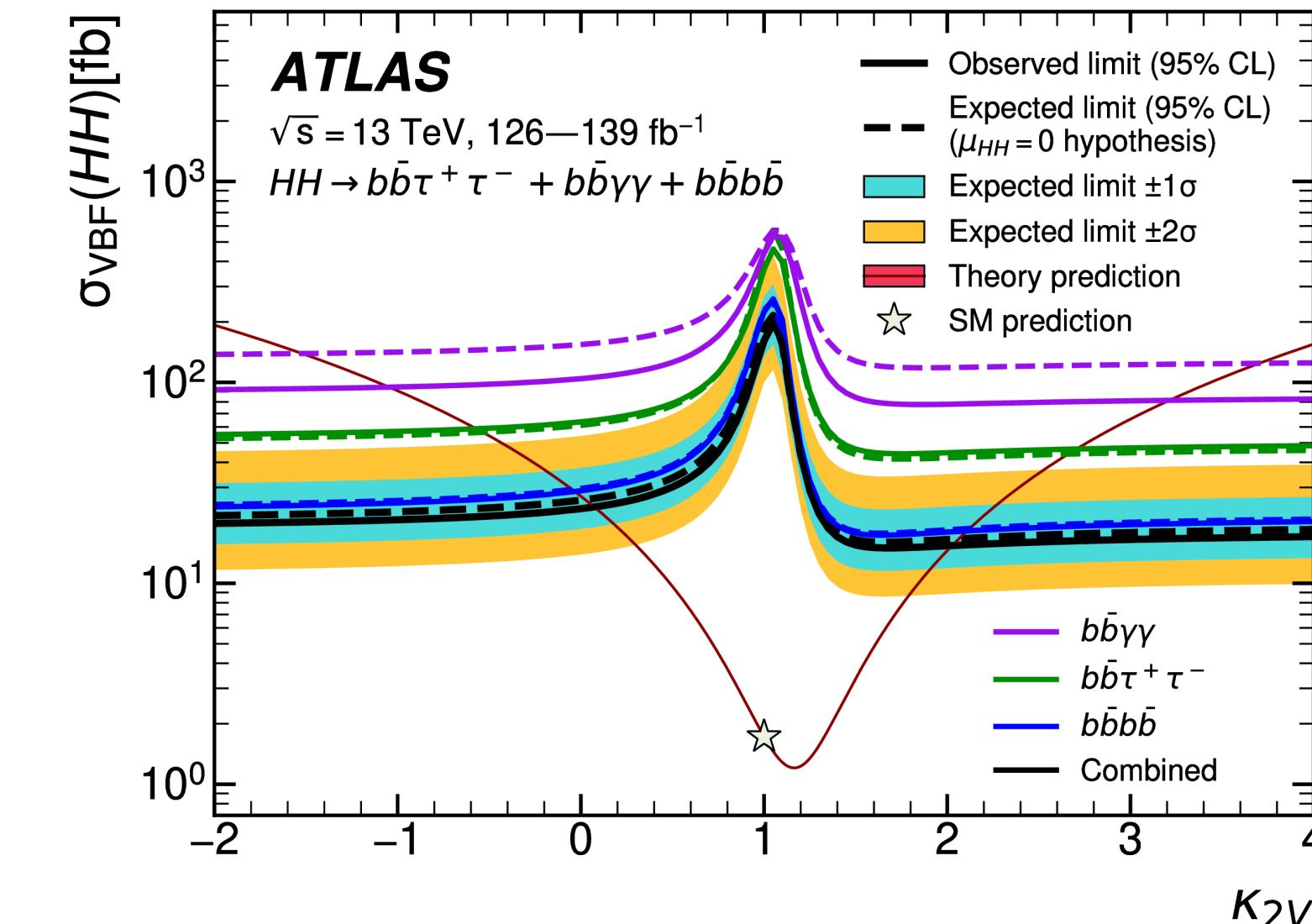
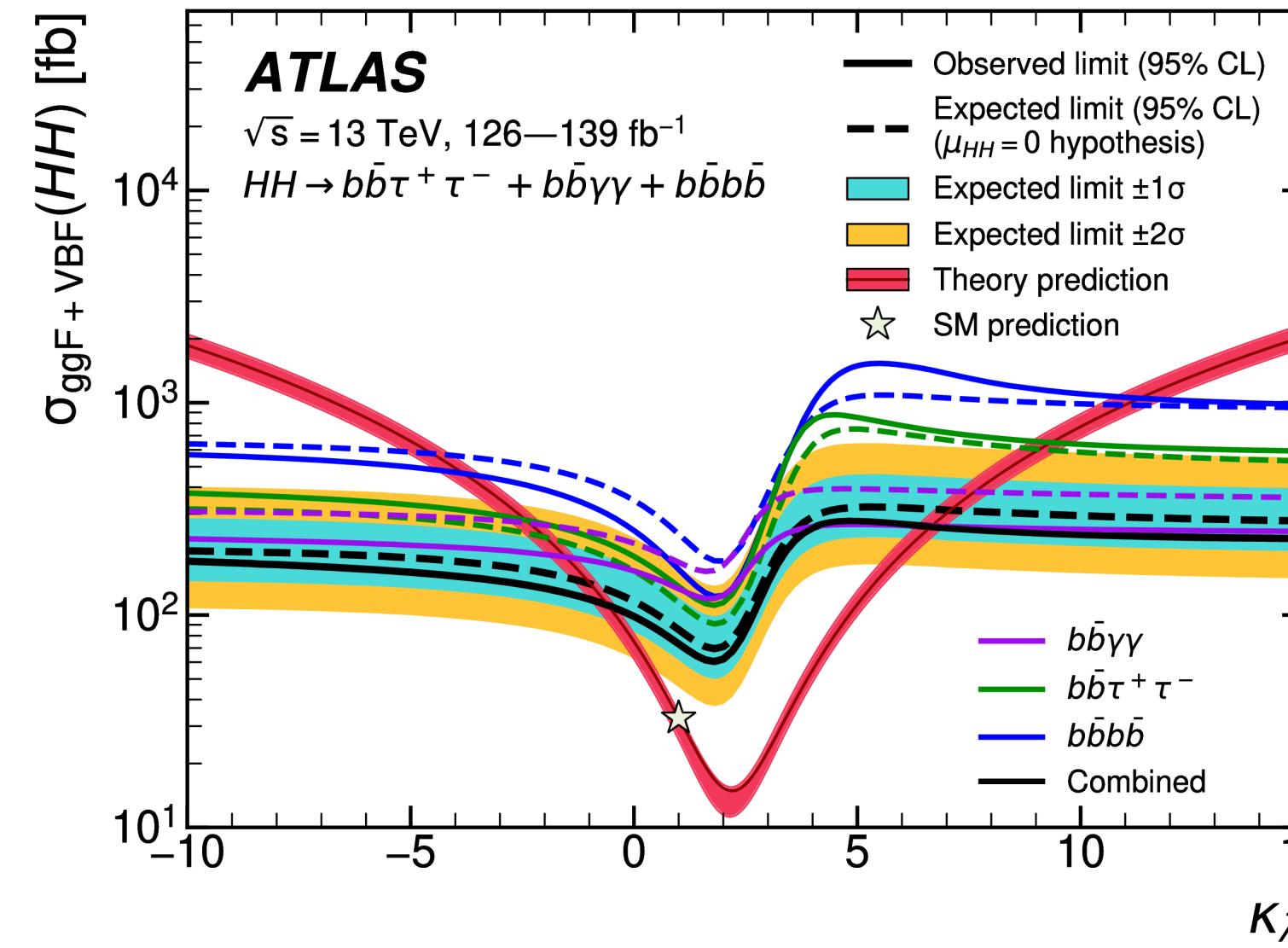


$$X_{Wt} = \min \left[\sqrt{\left(\frac{m_{jj} - m_W}{0.1 m_{jj}} \right)^2 + \left(\frac{m_{jjb} - m_t}{0.1 m_{jjb}} \right)^2} \right]$$

$$X_{HH} = \sqrt{\left(\frac{m_{H1} - 124 \text{ GeV}}{0.1 m_{H1}} \right)^2 + \left(\frac{m_{H2} - 117 \text{ GeV}}{0.1 m_{H2}} \right)^2}$$

Combined results

[arXiv:2211.01216 \[hep-ex\]](https://arxiv.org/abs/2211.01216)



EFT frameworks

SMEFT

- Canonical counting, expansion in $1/\Lambda$

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{n,i} \frac{c_i^{(n)}}{\Lambda^{n-4}} \mathcal{O}_i^{(n)}$$

- SM symmetries and fields, traditional EWSB mechanism (Higgs field: $\text{SU}(2)_L$ doublet)
- More restrictive (correlated Wilson coefficients)

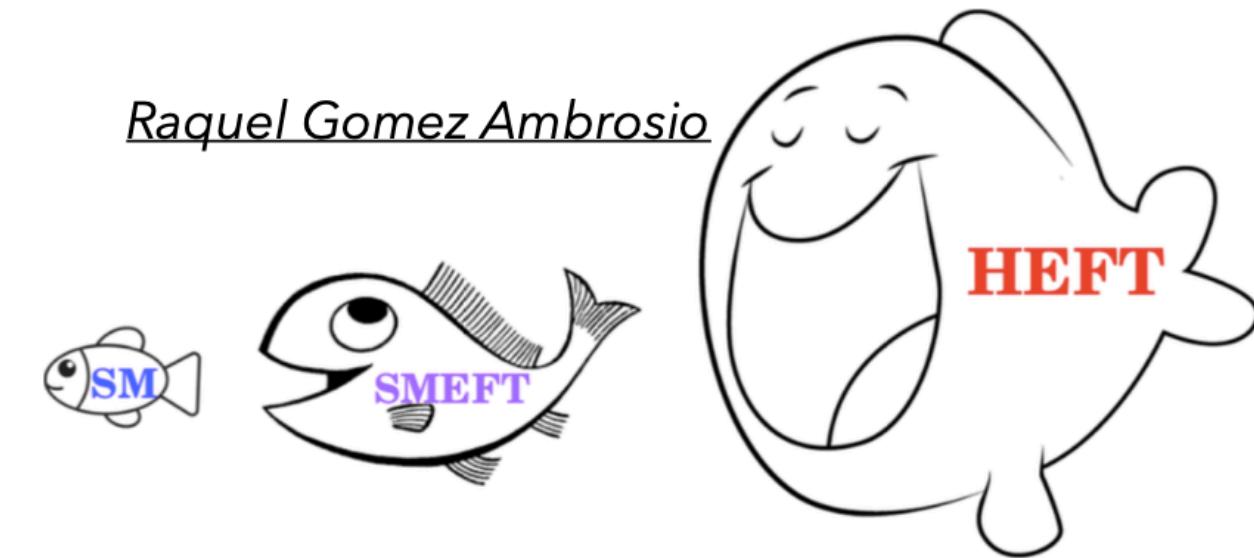
HEFT

- No power-counting like in SMEFT, more similar to chiral perturbation theory

$$\mathcal{L}_{d_\chi} = \mathcal{L}_{(d_\chi=2)} + \sum_{L=1}^{\infty} \sum_i \left(\frac{1}{16\pi^2} \right)^L c_i^{(L)} O_i^{(L)}$$

- Higgs field: EW singlet
- Much more general (independent couplings)

Raquel Gomez Ambrosio



SMEFT vs. HEFT

SMEFT

$$\Delta \mathcal{L}_{\text{Warsaw}} = \frac{C_{H,\square}}{\Lambda^2} (\phi^\dagger \phi) \square (\phi^\dagger \phi) + \frac{C_{HD}}{\Lambda^2} (\phi^\dagger D_\mu \phi)^* (\phi^\dagger D^\mu \phi) + \frac{C_H}{\Lambda^2} (\phi^\dagger \phi)^3 + \left(\frac{C_{uH}}{\Lambda^2} \phi^\dagger \phi \bar{q}_L \tilde{\phi} t_R + \text{h.c.} \right) + \frac{C_{HG}}{\Lambda^2} \phi^\dagger \phi G_{\mu\nu}^a G^{\mu\nu,a} + \frac{C_{uG}}{\Lambda^2} (\bar{q}_L \sigma^{\mu\nu} T^a G_{\mu\nu}^a \tilde{\phi} t_R + \text{h.c.})$$

$$\Delta \mathcal{L}_{\text{SILH}} = \frac{\bar{c}_H}{2v^2} \partial_\mu (\phi^\dagger \phi) \partial^\mu (\phi^\dagger \phi) + \frac{\bar{c}_u}{v^2} y_t (\phi^\dagger \phi \bar{q}_L \tilde{\phi} t_R + \text{h.c.}) - \frac{\bar{c}_6}{2v^2} \frac{m_h^2}{v^2} (\phi^\dagger \phi)^3 + \frac{\bar{c}_{ug}}{v^2} g_s (\bar{q}_L \sigma^{\mu\nu} G_{\mu\nu} \tilde{\phi} t_R + \text{h.c.}) + \frac{4\bar{c}_g}{v^2} g_s^2 \phi^\dagger \phi G_{\mu\nu}^a G^{a\mu\nu}$$

HEFT

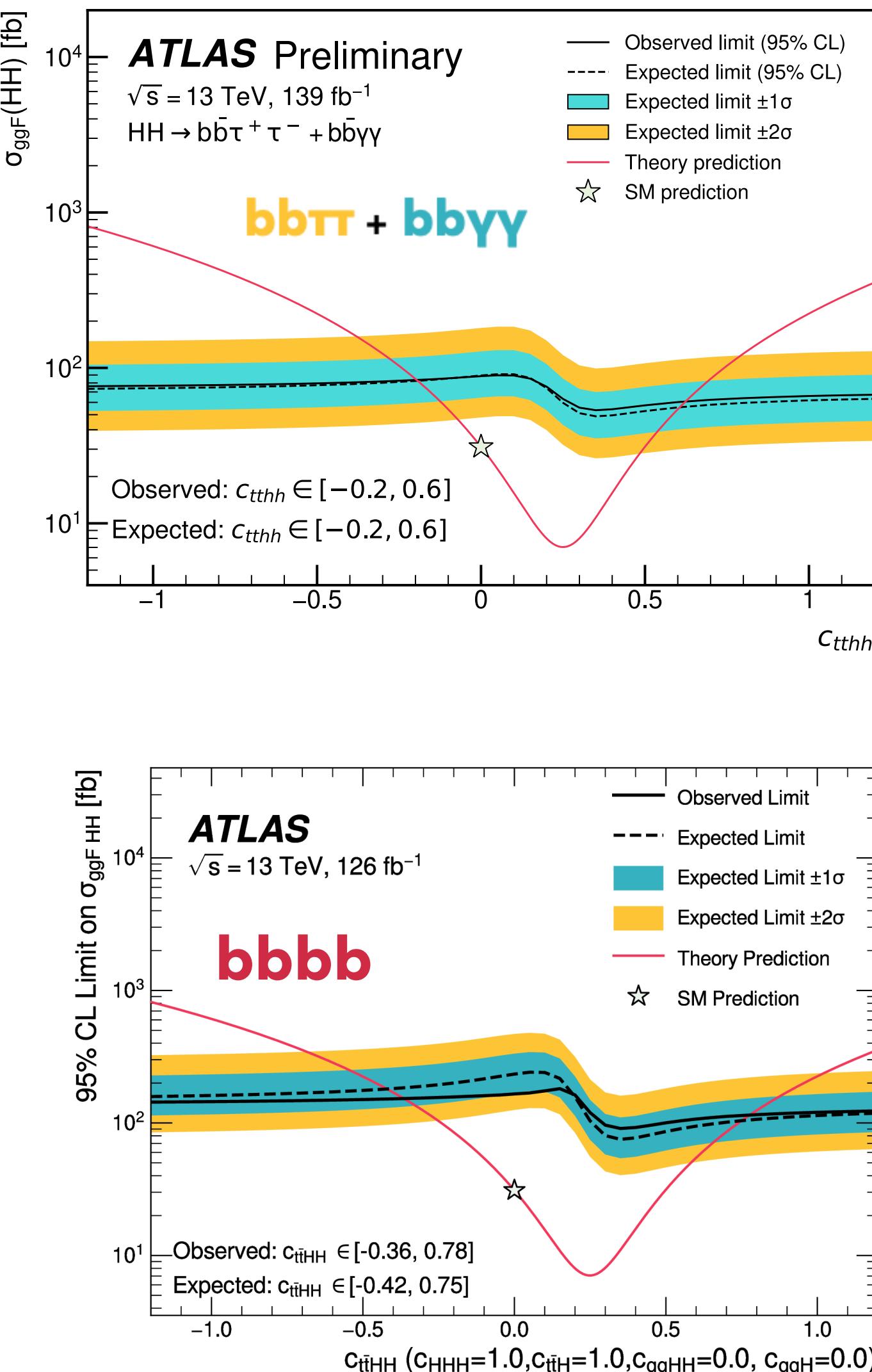
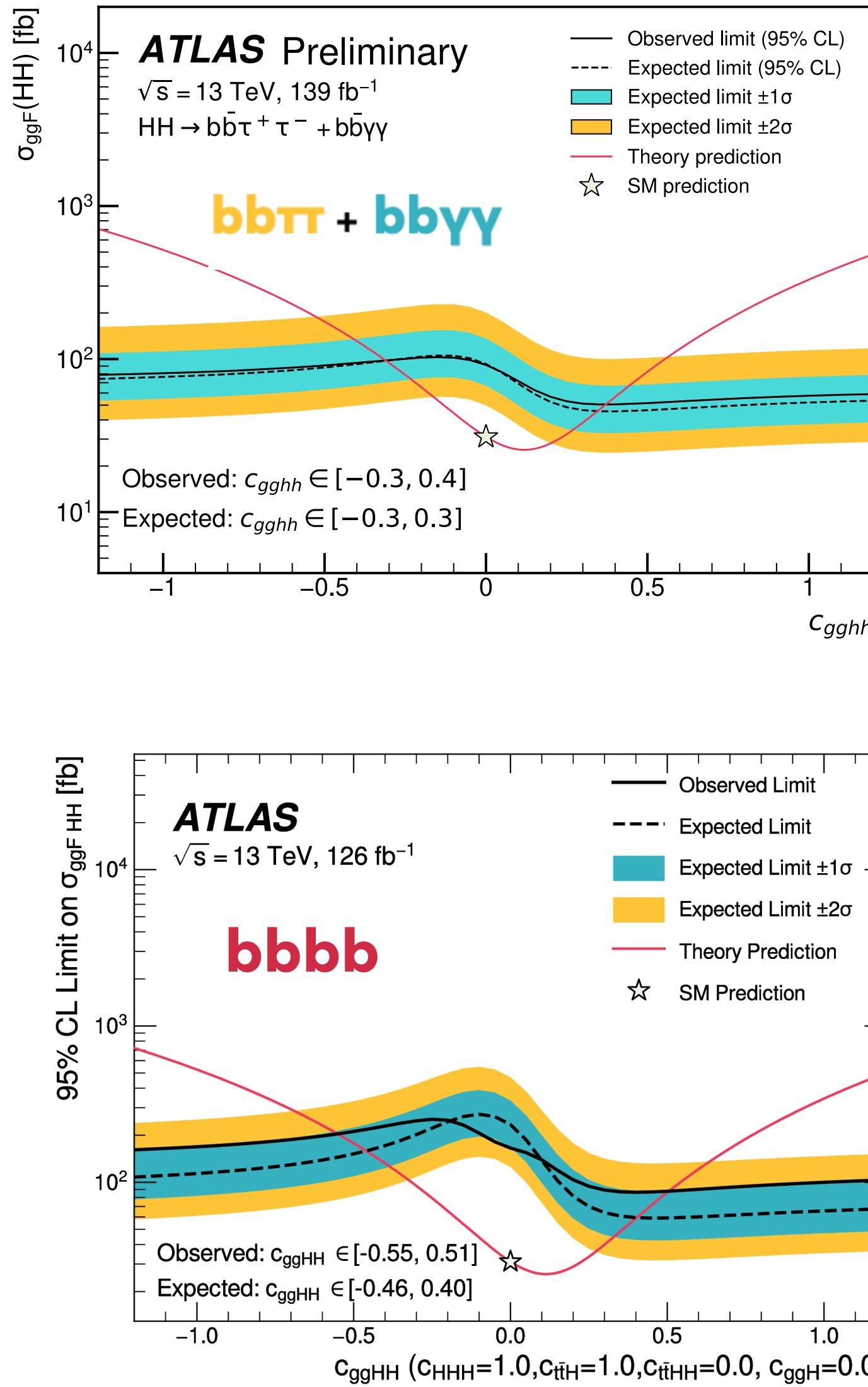
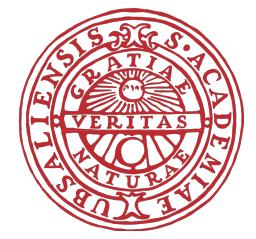
$$\Delta \mathcal{L}_{\text{HEFT}} = -m_t \left(c_t \frac{h}{v} + c_{tt} \frac{h^2}{v^2} \right) \bar{t} t - c_{hhh} \frac{m_h^2}{2v} h^3 + \frac{\alpha_s}{8\pi} \left(c_{ggh} \frac{h}{v} + c_{gghh} \frac{h^2}{v^2} \right) G_{\mu\nu}^a G^{a,\mu\nu}$$

Naive translation after field redefinition up to $\mathcal{O}(\Lambda^{-2})$ in Lagrangian ($C_{H,\text{kin}} = C_{H,\square} - \frac{1}{4}C_{HD}$)

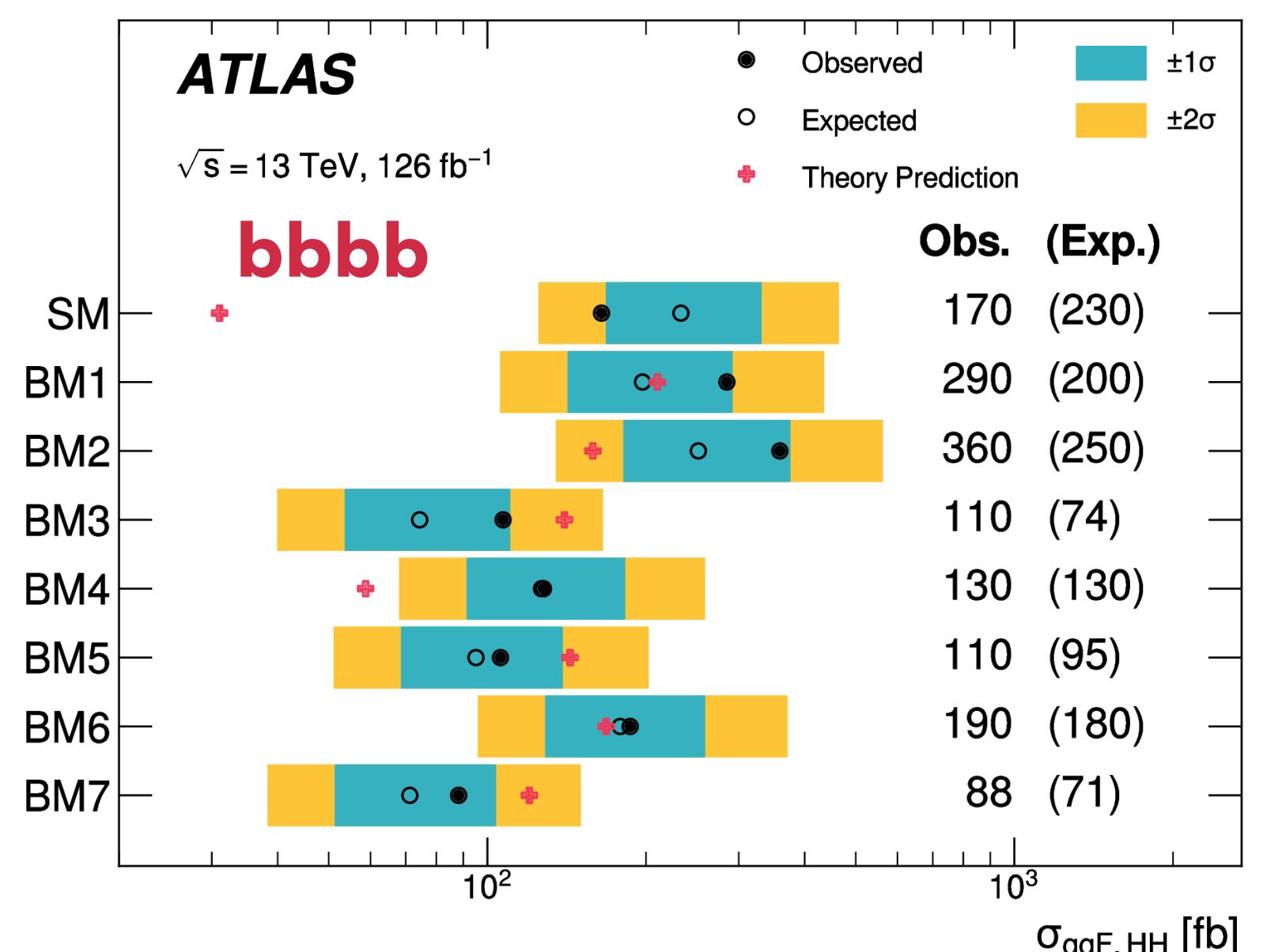
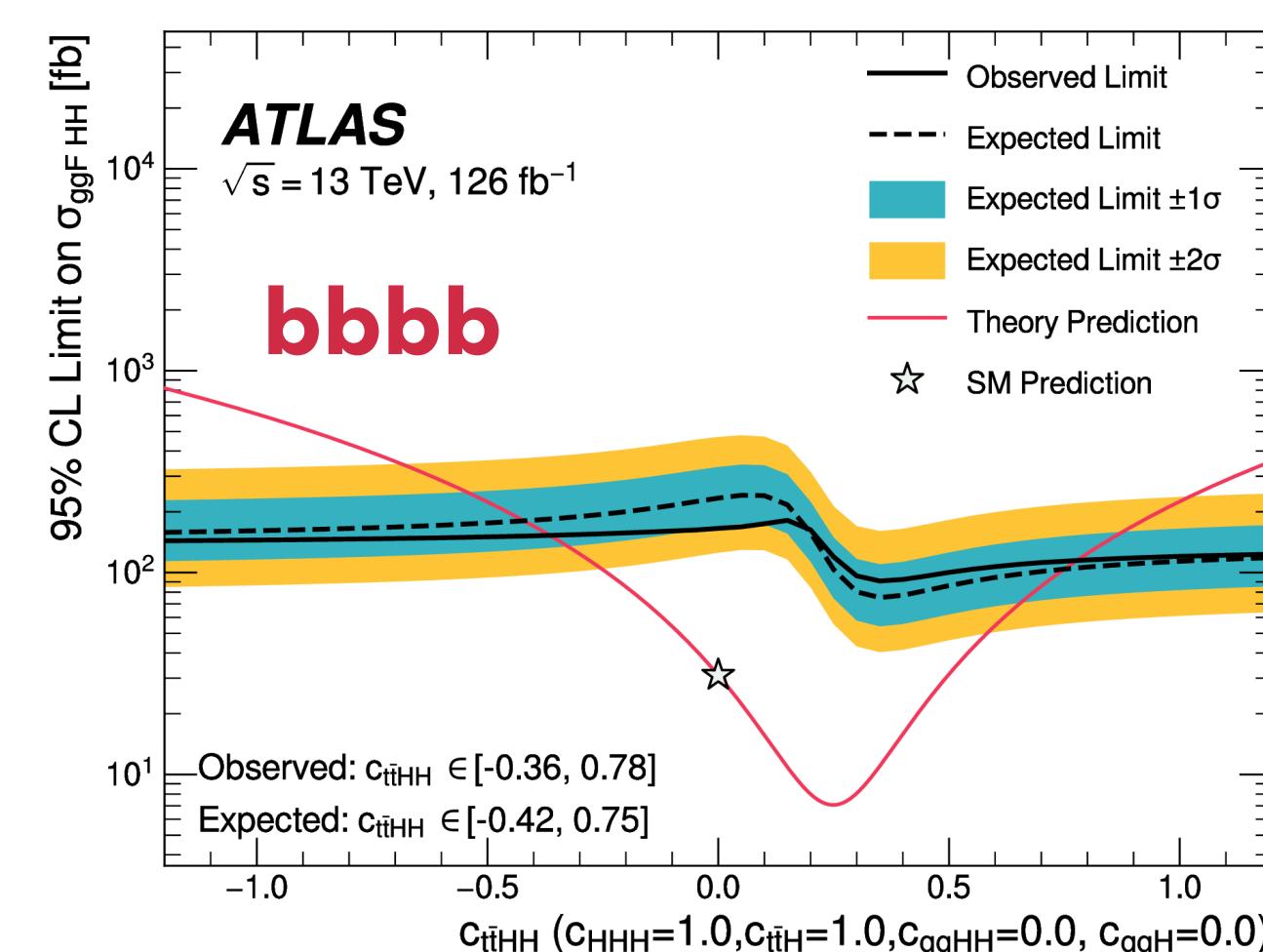
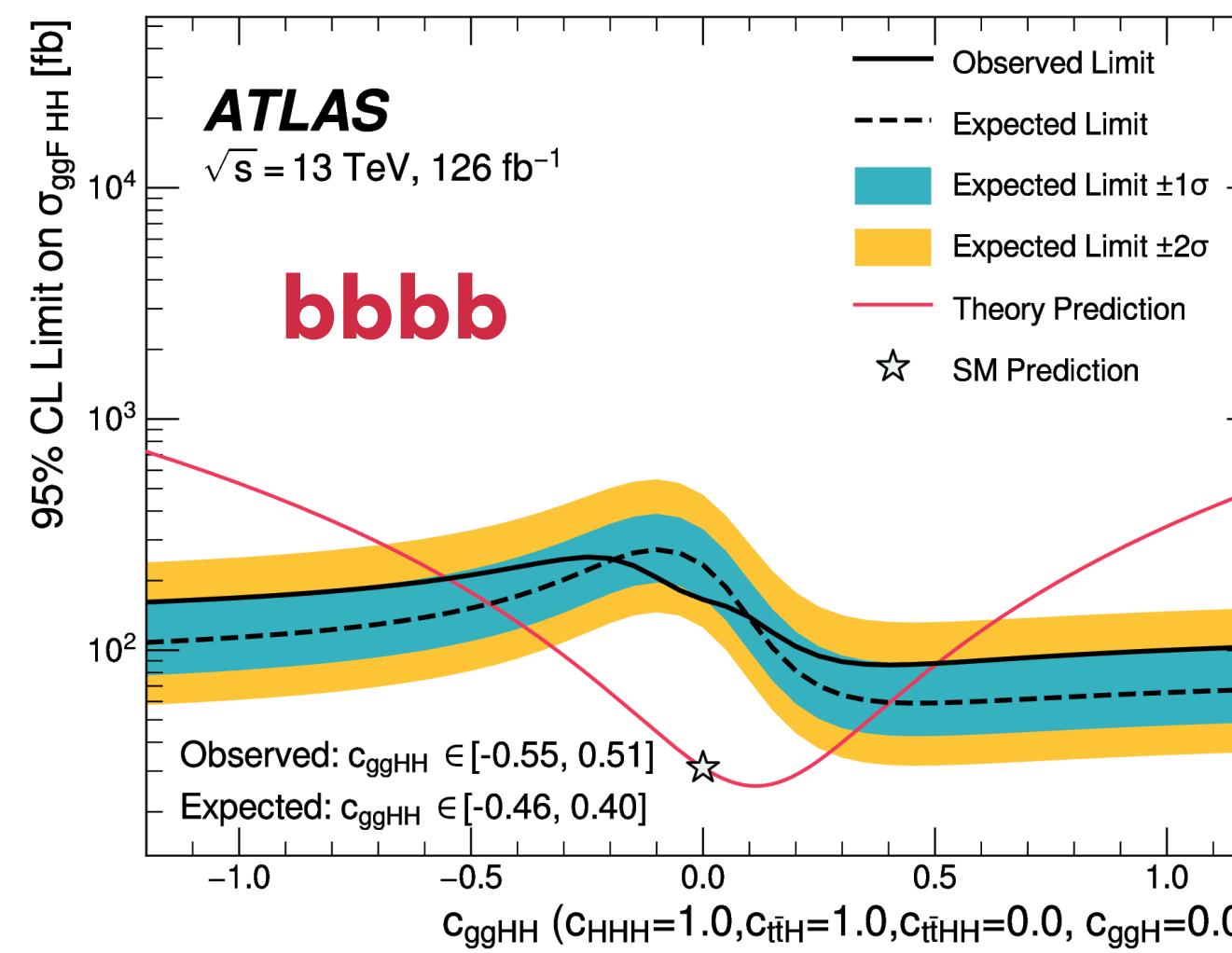
- ! Different assumptions, different EFT validity range
- ! Translation contains α_s which is a running parameter, typically evaluated at $\mu_0 = m_{hh}/2$
- ! Not generally applicable in practical calculations

HEFT	SILH	Warsaw
c_{hhh}	$1 - \frac{3}{2}\bar{c}_H + \bar{c}_6$	$1 - 2\frac{v^2}{\Lambda^2} \frac{v^2}{m_h^2} C_H + 3\frac{v^2}{\Lambda^2} C_{H,\text{kin}}$
c_t	$1 - \frac{\bar{c}_H}{2} - \bar{c}_u$	$1 + \frac{v^2}{\Lambda^2} C_{H,\text{kin}} - \frac{v^2}{\Lambda^2} \frac{v}{\sqrt{2}m_t} C_{uH}$
c_{tt}	$-\frac{\bar{c}_H + 3\bar{c}_u}{4}$	$-\frac{v^2}{\Lambda^2} \frac{3v}{2\sqrt{2}m_t} C_{uH} + \frac{v^2}{\Lambda^2} C_{H,\text{kin}}$
c_{ggh}	$128\pi^2 \bar{c}_g$	$\frac{v^2}{\Lambda^2} \frac{8\pi}{\alpha_s} C_{HG}$
c_{gghh}	$64\pi^2 \bar{c}_g$	$\frac{v^2}{\Lambda^2} \frac{4\pi}{\alpha_s} C_{HG}$

HEFT interpretations



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[arXiv:2301.03212 \[hep-ex\]](#)