Unfolding Higgs-top CP structure

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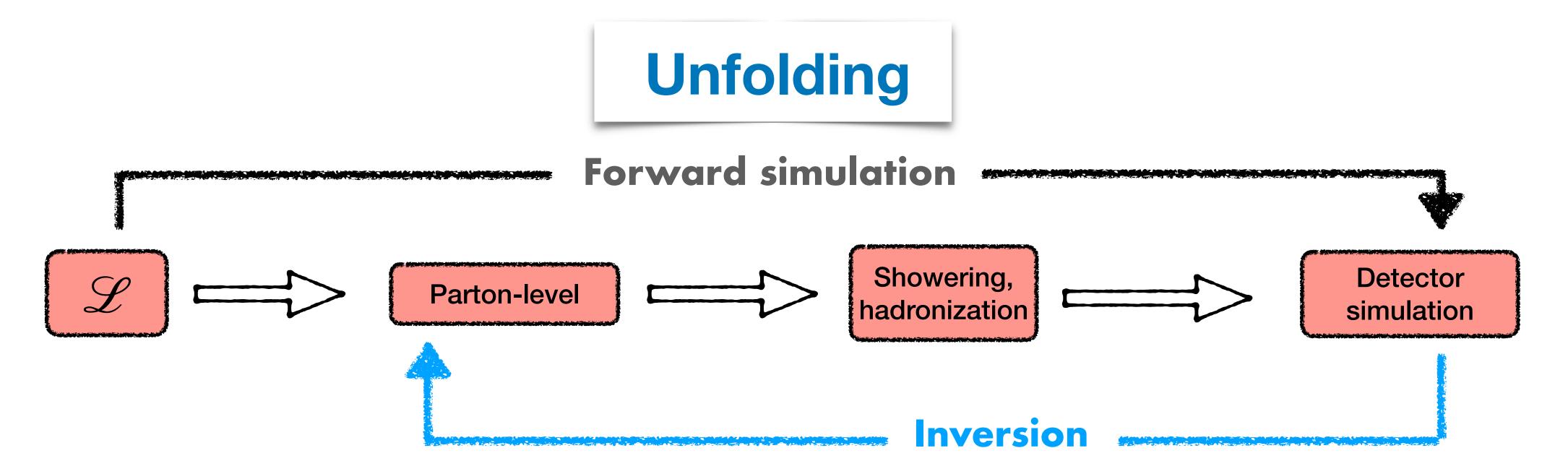
With

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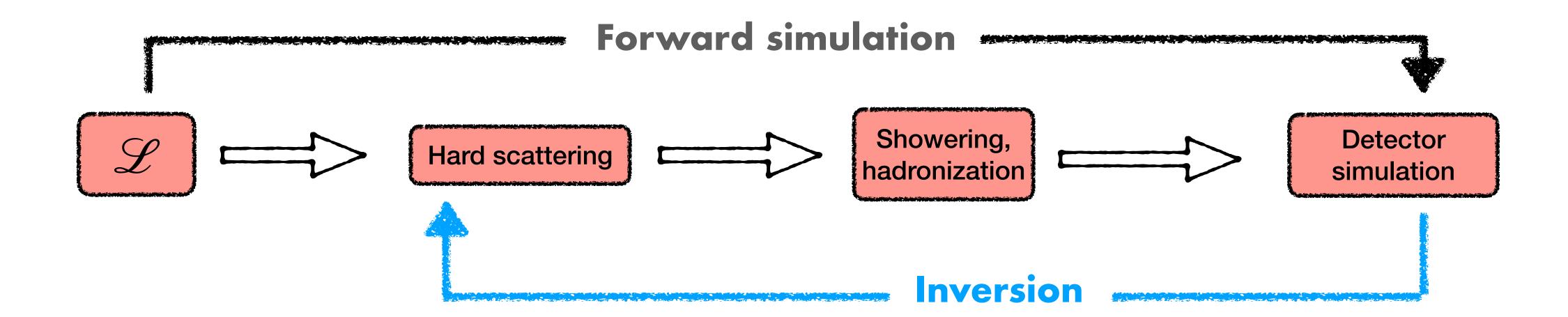


- Conventional LHC analysis involves comparing measured data with MC events simulated under NP hypothesis.
 - Reconstructed LHC events present a convoluted version of the `true' underlying physics.
 - Forward simulation chain can be highly resource intensive.

Invert simulation chain o apply on measured data o reconstruct phase space at the parton-level

→ compare new physics hypotheses at the parton-level.

Unfolding



The unfolding model is required to be:



Able to invert multi-dimensional d.o.f.

Possible with machine learning based generative models.

[Talk by R. Winterhalder]

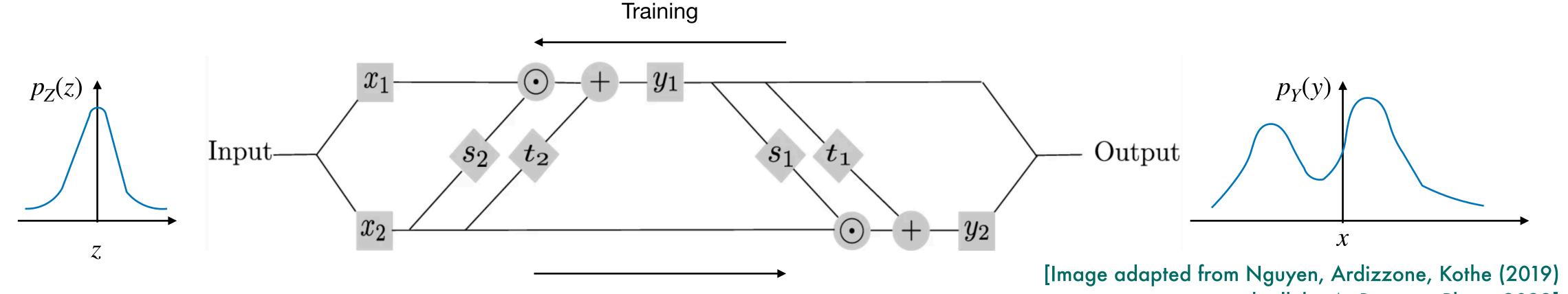
- ***** Generative Adversarial networks
- * Variational Autoencoders
- *Normalising flows

[Bellagente, Butter, Kasieczka, Plehn, Winterhalder (2020)]
[Bellagente, Butter, Kasieczka, Plehn, Rousselot, Winterhalder, Ardizzone, Kothe (2020)]
[Andreassen, Komiske, Metodiev, Nachman, Thaler (2020)]
[Komiske, McCormack, Nachman (2021)]

Normalising flows

- Series of invertible layers that transform complex (Y) to simple probability distributions (Z).
- Tractable Jacobian $J: p_Y(y) = p_Z(g(y))J$ Z = g(Y) g: Invertible function
- Fast sampling and density estimation.

[Talk by T. Heimel]



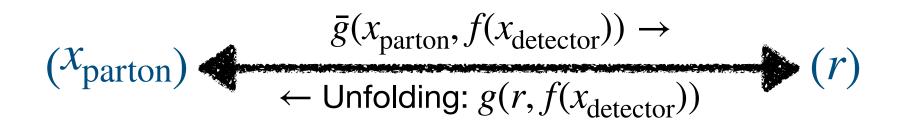
Inversion and talk by A. Butter at Pheno-2022]

Bijective map between parton-level and detector-level phase space

Forward simulation:
$$\bar{g} \rightarrow \begin{pmatrix} x_{\text{detector}} \\ r_{\text{parton}} \end{pmatrix}$$
 Forward simulation: $\bar{g} \rightarrow \begin{pmatrix} x_{\text{detector}} \\ r_{\text{detector}} \end{pmatrix}$

[Bellagente, Butter, Kasieczka, Plehn, Rousselot, Winterhalder, Ardizzone, Kothe (2020)]

Conditional INN



• Generate probability distributions at the parton-level, given detector-level events x_{detector}

 \longrightarrow Conditional on x_{detector} $\longrightarrow x_{\text{parton}}$ mapped with r

[Bellagente, Butter, Kasieczka, Plehn, Rousselot, Winterhalder, Ardizzone, Kothe (2020)]

Target phase space for unfolding can be chosen flexibly to include:

| iet radiation | Particle decays | Particle decay

Unfolding

- We use the Bayesian version of cINN
 - Stable network predictions
 - Allows the estimation of training-related uncertainties.

and talk by A. Butter at Pheno-2022]

[Butter, Heimel, Hummerich, Krebs, Plehn, Rousselot, Vent (2021)]

[Image adapted from Nguyen, Ardizzone, Kothe (2019)

Case study: CP measurement in Higgs-top interactions

- New sources of CPV interactions can explain the baryon asymmetry in the universe.
- ullet CPV in hVV interactions is extensively tested at the LHC

[Englert, Goncalves, Mawatari, Plehn (2012)] [Djouadi, Godbole, Mellado, Mohan (2013)] [Anderson, Bolognesi, Caola, Gao et al. (2013)]

- CPV in $hf\bar{f}$ couplings manifest at tree-level:
 - \rightarrow desirable choice: $ht\bar{t}$

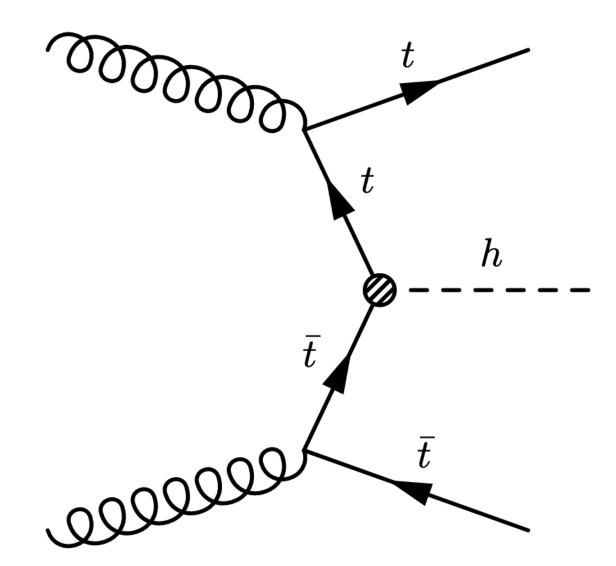
$$\mathcal{L} = -\frac{m_t}{v} \kappa_t h \bar{t} (\cos \alpha + i \gamma_5 \sin \alpha) t$$

- $pp \rightarrow h$ (+ jets): indirect constraints.

 [Brod, Haisch, Zupan (2013)]

 [Dolan, Harris, Jankowiak, Spannowsky (2014)]
- $pp \to t\bar{t}h$: opportunity to directly probe α and κ_t

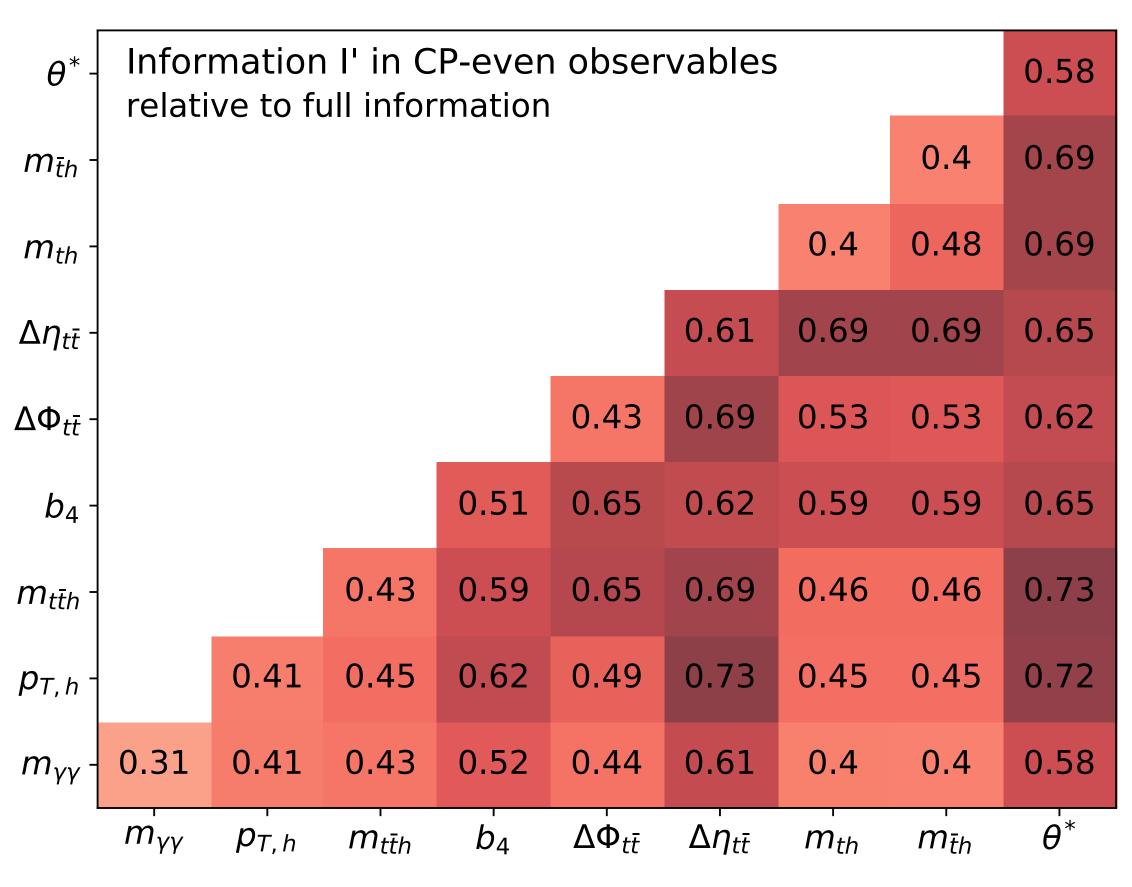
[Boudjema, Godbole, Guadagnoli, Mohan (2015)]
[Buckley, Goncalves (2016)]
[Azevedo, Onofre, Filthaut, Goncalo (2017)]



Current limit (ATLAS: 2004.04545): $|\alpha| < 43^0 \text{ at } 95\% \text{ CL}$

Improved statistics @ HL-LHC paves the pathway for precision studies.

$t\bar{t}(h \to \gamma\gamma)$ @ HL-LHC



Importance matrix at the **non-linear level** computed using Fisher Information metric

[RKB, Goncalves, Kling (2021)]

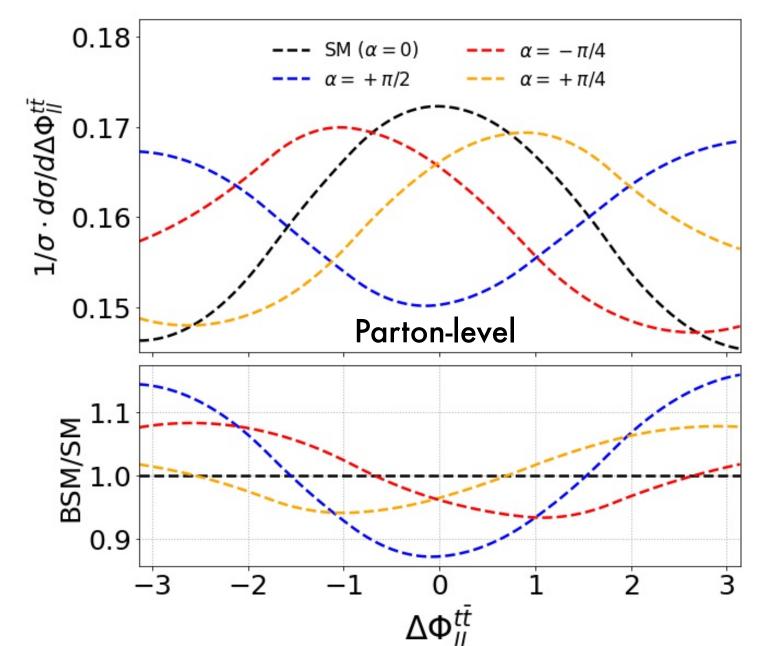
Sensitive only to non-linear new physics effects.

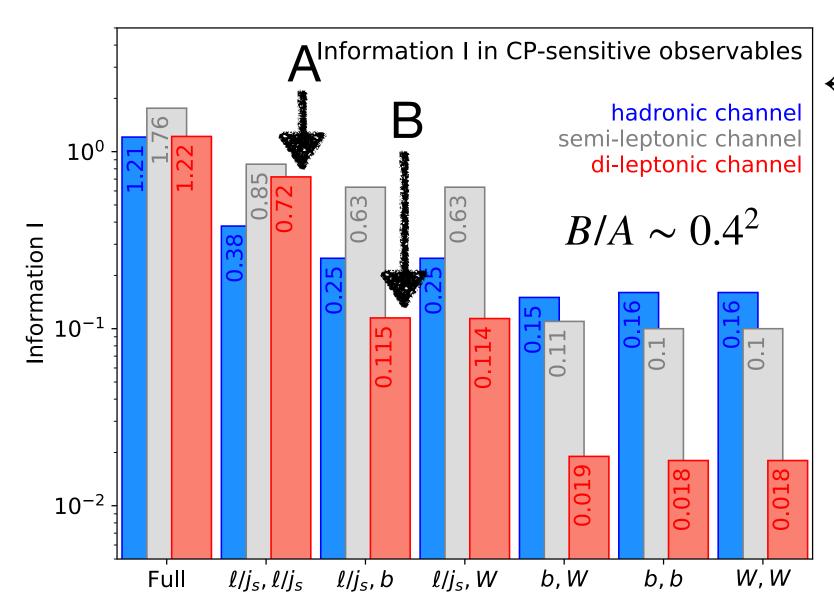
CP-odd observables

- Short lifetime for t $(10^{-25} s) \rightarrow$ Spin correlations can be traced back from their decay products.
- CP-odd observables constructed from antisymmetric tensor products

$$\epsilon(p_t, p_{\bar{t}}, p_i, p_j) \sim \epsilon_{\mu\nu\rho\sigma} p_t^{\mu} p_{\bar{t}}^{\nu} p_i^{\rho} p_j^{\sigma}$$
:

$$\Delta \phi_{ij}^{t\bar{t}} = \mathbf{sgn} \left[\vec{p}_t \cdot (\vec{p}_i \times \vec{p}_j) \right] \arccos \left[\frac{\vec{p}_t \times \vec{p}_i}{|\vec{p}_t \times \vec{p}_i|} \cdot \frac{\vec{p}_t \times \vec{p}_j}{|\vec{p}_t \times \vec{p}_j|} \right]$$





 \leftarrow Spin correlations scale with the spin analysing power β_i .

[Mileo, Kiers, Szynkman, Crane, Gegner (2016); Goncalves, Kong, Kim (2018)]; RKB, Goncalves, Kling (2021)]

$$\frac{1}{\Gamma} \frac{d\Gamma}{d\cos\xi_i} = \frac{1}{2} \left(1 + \beta_i P_t \cos\xi_i \right)$$
 Fisher Info = $\mathbb{E} \left[\frac{\partial \log p(\mathbf{x} \mid \kappa_t, \alpha)}{d\alpha} \frac{\partial \log p(\mathbf{x} \mid \kappa_t, \alpha)}{d\alpha} \right]$

Kinematic reconstruction efficiency is limited at the detector level

Use Machine learning techniques to maximize the extraction of NP information from depleted CP-odd and CP-even observables.

Event parametrization

- In the training data, event information at the parton level can be parametrised through the 4-momentum of the final state particles
 - → may include redundant d.o.f.
- Reconstruction of sharp intermediate mass peaks can be improved by adding maximum mean discrepancy between the truth and generated distributions in the loss function.

[Butter, Plehn, Winterhalder (2019)]

Affects only the target distributions

[Bellagente, Butter, Kasieczka, Plehn, Rousselot, Winterhalder, Ardizzone, Kothe (2020)]

Avoids large model dependence

Complications in training and performance limitations.

- Alternative approach:
- → invariant mass features can be learnt directly with appropriate phase-space parametrization.
- → also provide direct access to the most important CP-even and CP-odd observables.

$$\vec{p}_{t\bar{t}}, m_{t_{\ell}}, |\vec{p}_{t_{\ell}}^{\text{CS}}|, \theta_{t_{\ell}}^{\text{CS}}, \phi_{t_{\ell}}^{\text{CS}}, m_{t_{h}},$$

$$\operatorname{sign}(\Delta \phi_{\ell \nu}^{t\bar{t}}) m_{W_{\ell}} |\vec{p}_{\ell}^{t\bar{t}}|, \theta_{\ell}^{t\bar{t}}, \phi_{\ell}^{t\bar{t}}, |\vec{p}_{\nu}^{t\bar{t}}|$$

$$\operatorname{sign}(\Delta \phi_{du}^{t\bar{t}}) m_{W_{h}}, |\vec{p}_{d}^{t\bar{t}}|, \theta_{d}^{t\bar{t}}, \Delta \phi_{\ell d}^{t\bar{t}}, |\vec{p}_{u}^{t\bar{t}}|$$

Unfolding $t\bar{t}h$ events

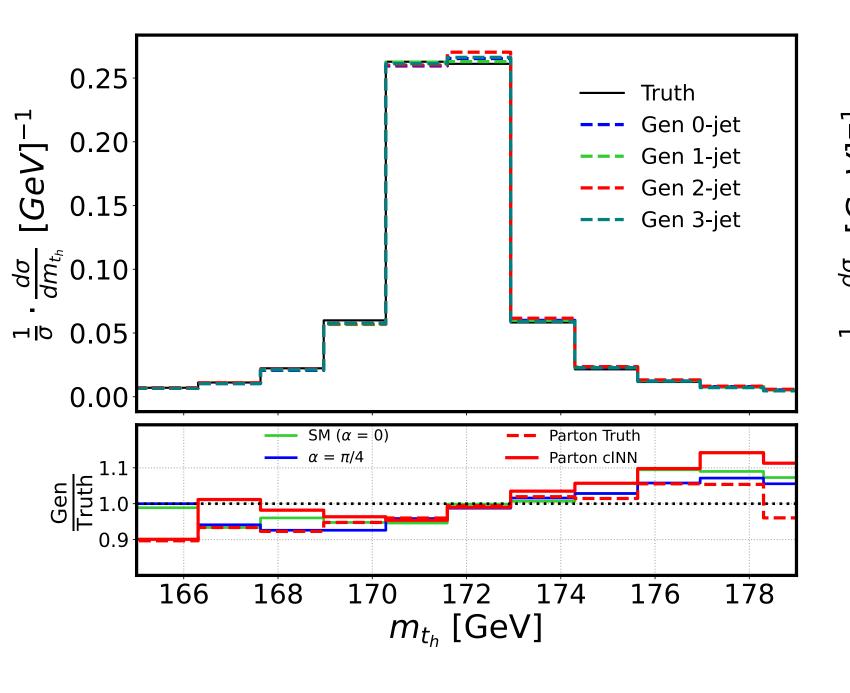
Training dataset: Parton-level: $(t \to \ell \nu b)(\bar{t} \to jj\bar{b})(h \to \gamma \gamma)$

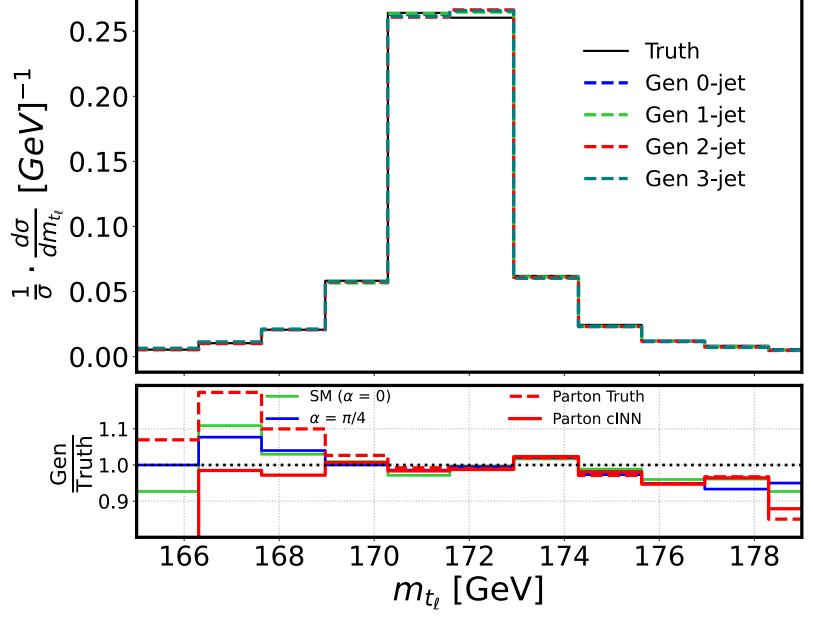
Detector-level: $1\ell + 2b + 2\gamma + MET + 2$ jet

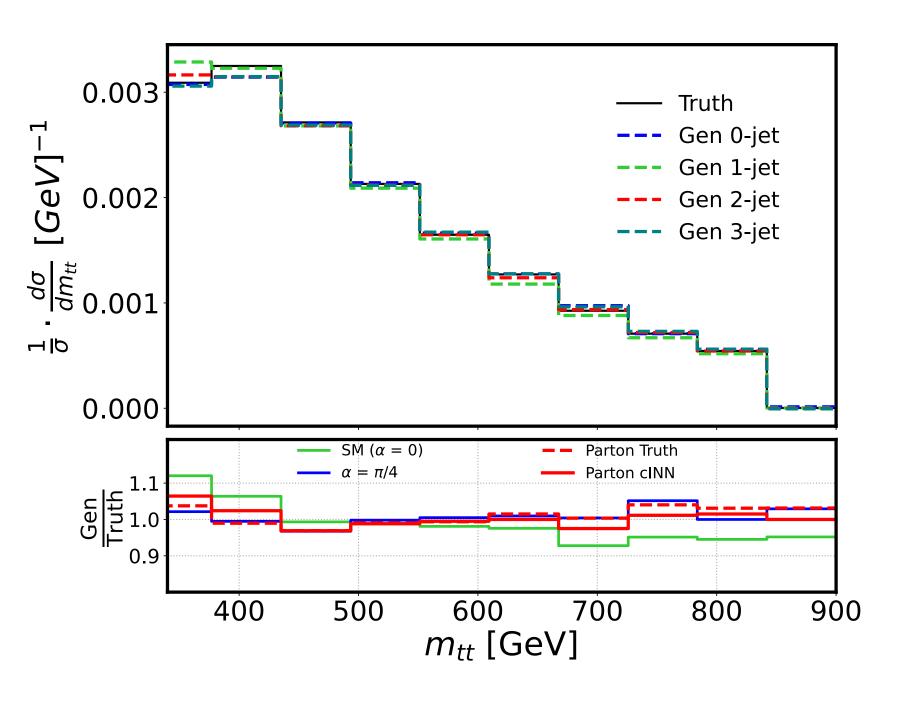
≤ 3 jet inclusive

≤ 4 jet inclusive

≤ 5 jet inclusive

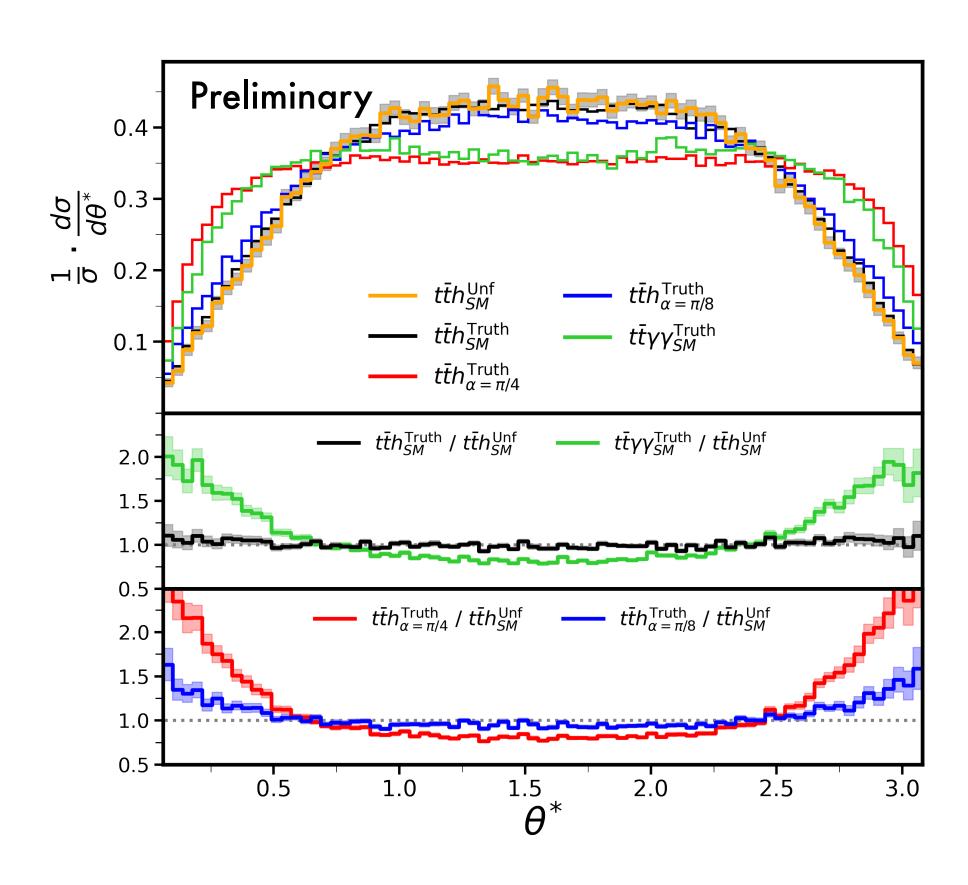


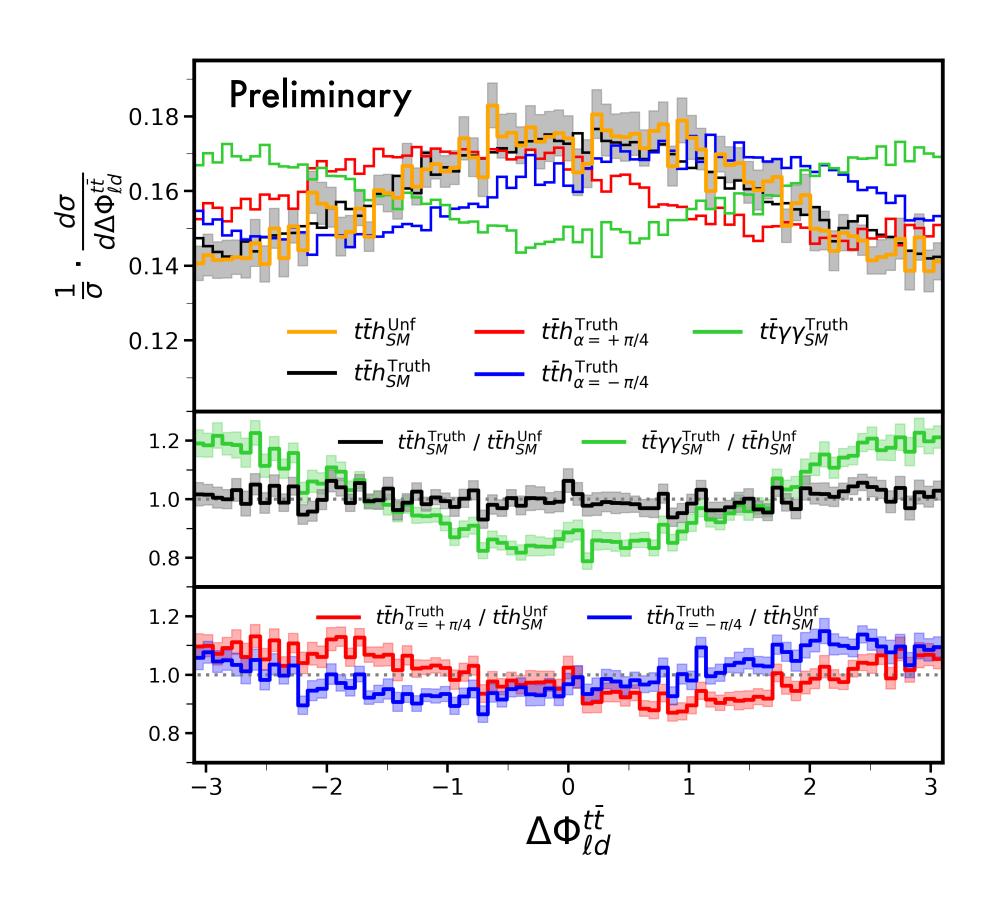




Unfolding $t\bar{t}h$ events

Trained on SM $pp \to (t \to \ell \nu b)(\bar t \to jj\bar b)(h \to \gamma\gamma)$ + with upto 4 additional jets





Outlook

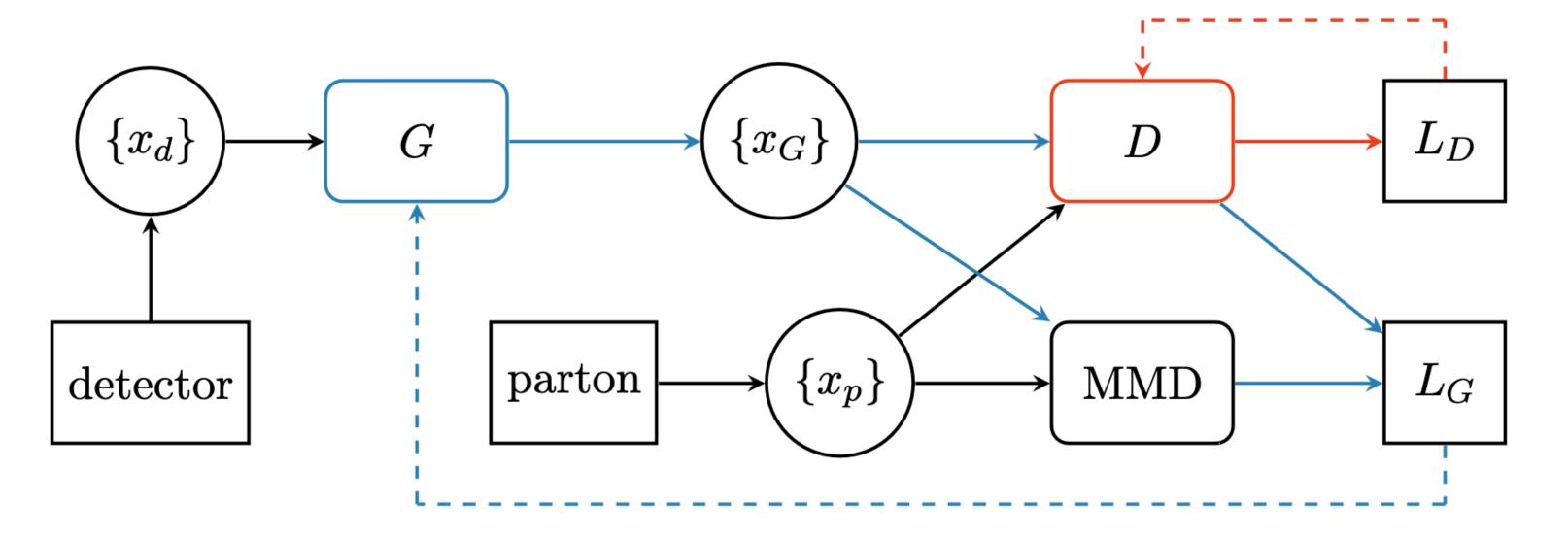
- Full inversion to the hard-scattering level ..
- Towards precision studies with unfolding ..
- Explore model dependence of the unfolding setup ...

Back up slides

Unfolding with GANs

[Bellagente, Butter, Kasieczka, Plehn, Winterhalder(2019)]

[Butter, Plehn, Winterhalder(2019)]



[Image from Bellagente, Butter, Kasieczka, Plehn, Winterhalder (2019)]

$$L_{\text{Discriminator}} = \langle -log D(x) \rangle_{x \sim P_p} + \langle -log (1 - D(x)) \rangle_{x \sim P_G}$$

$$L_{\text{Generator}} = \langle -log D(x) \rangle_{x \sim P_G}$$

Event parametrization

• For $pp \to (t_\ell \to W_\ell b_\ell)$ $(\bar{t}_h \to udb_h) h$, a typical parametrisation could be:

$$\left\{ m_t, p_{T,t}, \eta_t, \boldsymbol{\phi}_t, m_W, \eta_W^t, \boldsymbol{\phi}_W^t, \eta_{\ell,u}^W, \boldsymbol{\phi}_{\ell,u}^W, \boldsymbol{\phi}_{\ell,u}^W \right\}$$

- Can reproduce the intermediate mass distributions.
- Limited spin-correlation information.
- Alternatively, provide direct access to the most important CP-even and CP-odd observables:

$$\vec{p}_{t\bar{t}}, m_{t_{\ell}}, |\vec{p}_{t_{\ell}}^{\mathsf{CS}}|, \theta_{t_{\ell}}^{\mathsf{CS}}, \phi_{t_{\ell}}^{\mathsf{CS}}, m_{t_{h}},$$

$$\operatorname{sign}(\Delta \phi_{\ell \nu}^{t\bar{t}}) m_{W_{\ell}} |\vec{p}_{\ell}^{t\bar{t}}|, \theta_{\ell}^{t\bar{t}}, \phi_{\ell}^{t\bar{t}}, |\vec{p}_{\nu}^{t\bar{t}}|$$

$$\operatorname{sign}(\Delta \phi_{du}^{t\bar{t}}) m_{W_{h}}, |\vec{p}_{d}^{t\bar{t}}|, \theta_{d}^{t\bar{t}}, \Delta \phi_{\ell d}^{t\bar{t}}, |\vec{p}_{u}^{t\bar{t}}|$$