

Nucleon TMDPDFs from Lattice QCD

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Jinchen He

University of Maryland, College Park

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Parton Model

Any hadron with large momentum can be considered as a composition of a number of point-like "partons". (By Richard Feynman in 1969)

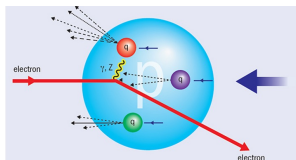


Figure 1: DIS with parton model, from CERN.

$$\sigma_{\text{DIS}} \propto \left| \begin{array}{c} l \quad l' \\ q \\ p \end{array} \right|^2 \approx \left| \begin{array}{c} k \approx \xi P \\ p \end{array} \right|^2 \otimes \left| \begin{array}{c} l \quad l' \\ q \\ \xi P \end{array} \right|^2$$

Figure 2: DIS factorization. [1]

What are TMDPDFs

Parton Distribution Functions: the probability to find a parton (quark or gluon) in a hadron as a function of momentum fraction x .

Transverse-Momentum-Dependent PDFs: the distribution densities to find a parton carrying a longitudinal momentum fraction x and transverse momentum k_T in a hadron.

TMDPDFs can be regarded as a tomography of the nucleon, which reveals the nucleons 3D internal structure.

$$\sigma_{\text{DY}} \propto \left| \text{Diagram 1} \right|^2 \approx \left| \text{Diagram 2} \right|^2 \otimes \left| \text{Diagram 3} \right|^2 \otimes \left| \text{Diagram 4} \right|^2$$

Figure 3: Drell Yan factorization.[1]

Why Care about TMDPDFs

Understanding the inner structure of hadrons.

TMD processes are the most important processes in high energy collisions, like SIDIS on EIC.

TMDPDFs are important inputs for experiments on TMD processes.

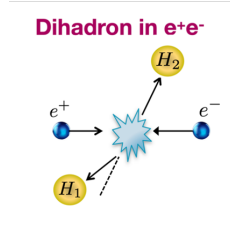
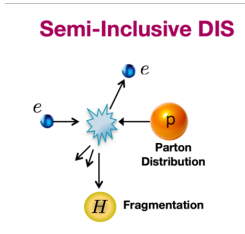
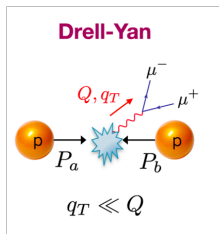


Figure 4: TMD processes. [1]

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Lattice QCD

Path integral formalism

$$\langle \hat{O} \rangle = \int \mathcal{D}A \mathcal{D}\psi \mathcal{D}\bar{\psi} \hat{O} e^{iS_{\text{QCD}}(A, \psi, \bar{\psi})}$$

$$\langle \hat{O} \rangle = \int \mathcal{D}A \mathcal{D}\psi \mathcal{D}\bar{\psi} \hat{O} e^{-S_{\text{QCD}}^E(A, \psi, \bar{\psi})}$$

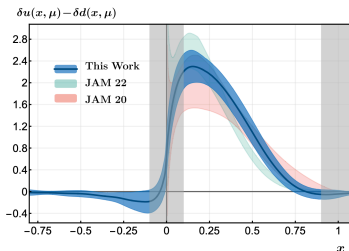
Monte Carlo Sampling

$$P(A, \psi, \bar{\psi}) = \frac{e^{-S_{\text{QCD}}^E(A, \psi, \bar{\psi})}}{\int \mathcal{D}A \mathcal{D}\psi \mathcal{D}\bar{\psi} e^{-S_{\text{QCD}}^E(A, \psi, \bar{\psi})}}$$

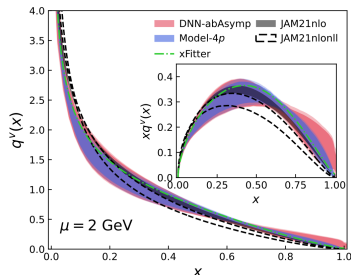
$$\langle \hat{O} \rangle = \langle \hat{O} \rangle_{\text{Sample average}}$$

Lattice QCD

Lattice QCD is a powerful non-perturbative method



(a) Proton transversity PDF. [2]



(b) Pion valance PDF. [3]

Figure 5: Achievements on Lattice QCD.

Large Momentum Effective Theory(LaMET)

Quasi distribution: Equal time, P_z -dependent.

Light-cone distribution: Separate on time, universal.

$$(\Delta t)^2 - (\Delta x)^2 - (\Delta y)^2 - (\Delta z)^2 = 0$$

$$(\Delta \tau)^2 + (\Delta x)^2 + (\Delta y)^2 + (\Delta z)^2 = 0$$

$$\tilde{f}_\Gamma(x, b_\perp, \zeta_z, \mu) \sqrt{S_I(b_\perp, \mu)} = H_\Gamma\left(\frac{\zeta_z}{\mu^2}\right) e^{\frac{1}{2} \ln\left(\frac{\zeta_z}{\mu^2}\right) K(b_\perp, \mu)} f(x, b_\perp, \mu, \zeta) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{\zeta_z}, \frac{M^2}{(P^z)^2}, \frac{1}{b_\perp^2 \zeta_z}\right)$$

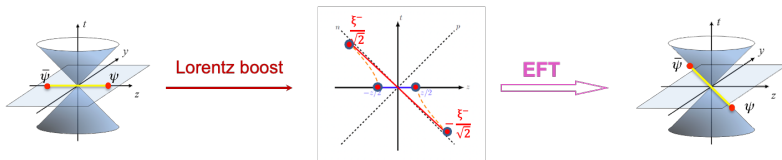


Figure 6: Large Momentum Effective Field Theory.

Recipe

Correlator on lattice.

Ground state fit.

Renormalization

x-distribution.

LaMET matching.

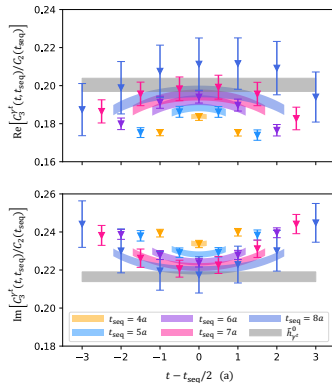
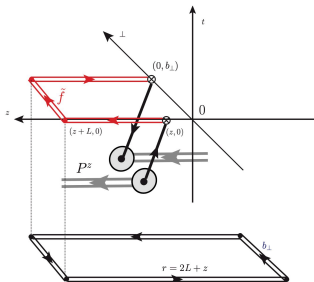


Figure 7: Ground state fit.[4]

Renormalization

$$\tilde{h}_{\Gamma}(z, b_{\perp}, P^z, a, \mu) = \lim_{L \rightarrow \infty} \frac{\tilde{h}_{\Gamma}^0(z, b_{\perp}, P^z, a, L)}{\sqrt{Z_E(2L + z, b_{\perp}, a)} Z_O(1/a, \mu, \Gamma)} \quad (1)$$



Z_E : divergence from Wilson links

Z_O : divergence from quark-gauge link vertices

Figure 8: Non-perturbative renormalization.[5]

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Comparison with Phenomenological Results

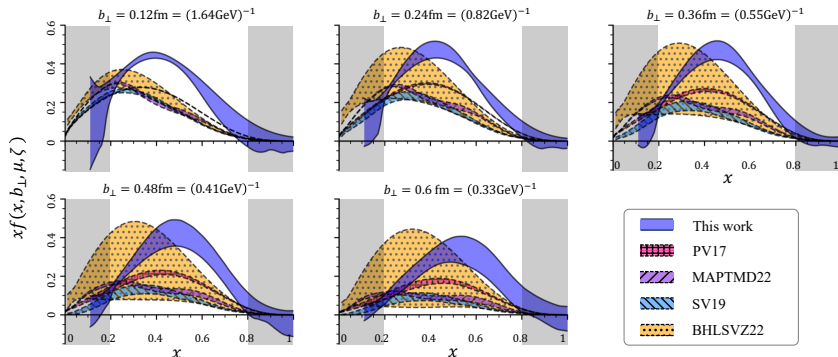


Figure 9: Final results on x-dependence. [4][6][7][8][9]

Comparison with Phenomenological Results

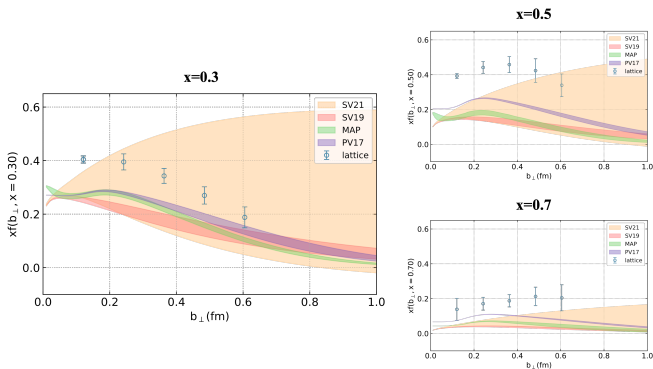


Figure 10: Final results on b -dependence. [4][6][7][8][9]

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Improvement

Direction	Status	Prospect
Lattice spacing	$a = 0.12 \text{ fm}$	Continuum limit
Hadron momentum	$P_z = 2.58 \text{ GeV}$	Larger γ factor
Transverse behavior	$b = 0.6 \text{ fm}$	QCD Confinement
More		Higher power correction

Other TMDPDFs

Leading Quark TMDPDFs



Nucleon Spin



Quark Spin

		Quark Polarization		
		Un-Polarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1 = \text{Unpolarized}$		$h_1^\perp = \text{Boer-Mulders}$
	L		$g_{1L} = \text{Helicity}$	$h_{1L}^\perp = \text{Worm-gear}$
	T	$f_{1T}^\perp = \text{Sivers}$	$g_{1T}^\perp = \text{Worm-gear}$	$h_1 = \text{Transversity}$ $h_{1T}^\perp = \text{Pretzelosity}$

Figure 11: Leading quark TMDPDFs. [1]

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Thank You

Unpolarized Transverse-Momentum-Dependent Parton Distributions of the Nucleon from Lattice QCD (Lattice Parton Collaboration (LPC))



Jin-Chen He,^{1,2,3} Min-Huan Chu,^{1,2} Jun Hua,^{4,5} Xiangdong Ji,³ Andreas Schäfer,⁶
Yushan Su,³ Wei Wang,^{2,*} Yi-Bo Yang,^{7,8,9,10} Jian-Hui Zhang,^{11,12} and Qi-An Zhang^{13,†}

Figure 12: e-Print: 2211.02340. [4]

Back Up: Lattice Setup

2+1+1 flavors of HISQ action by MILC ($a=0.12$ fm);

Valance pion mass: 310 MeV, 220 MeV;

Gamma structure: γ^t and γ^z ;

Hadron momentum: 1.72 GeV, 2.15 GeV, 2.58 GeV.

Back Up: Uncertainties

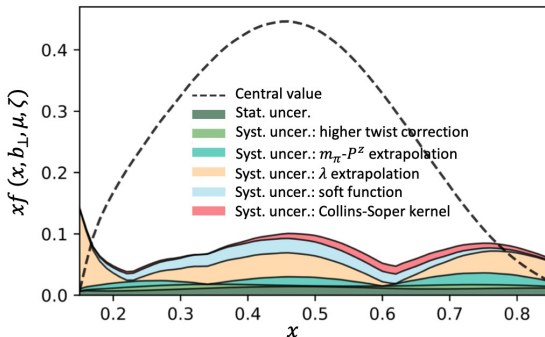


Figure 13: Comparison between different uncertainties. [4]

Back Up: Definition

PDF

$$q(x, \epsilon) \equiv \int \frac{d\xi^-}{4\pi} e^{-ixP^+\xi^-} \langle P | \bar{\psi}(\xi^-) \gamma^+ U(\xi^-, 0) \psi(0) | P \rangle,$$

$$U(\xi^-, 0) = P \exp \left(-ig \int_0^{\xi^-} d\eta^- A^+(\eta^-) \right).$$

TMDPDF

$$\tilde{f}_{i/p}^{0(u)}(x, \mathbf{b}_T, \epsilon, \tau, xP^+)$$

$$\equiv \int \frac{db^-}{2\pi} e^{-ib^-(xP^+)} \left\langle p(P) \left| \left[\bar{\psi}_i^0(b^\mu) W_\square(b^\mu, 0) \frac{\gamma^+}{2} \psi_i^0(0) \right]_\tau \right| p(P) \right\rangle,$$

$$\tilde{S}_{n_a n_b}^0(b_T, \epsilon, \tau) = \frac{1}{N_c} \langle 0 | \text{Tr}[W(b_T)] | 0 \rangle.$$