Lepton Specific Extended Higgs Model

Matthew Knauss

Bernardo Gonçalves, Marc Sher Phys. Rev. D 107, 095001 (2023) arXiv:2301.08641







Motivation

- 2 Higgs Doublet Models (2HDMs) are a simple extension to Standard Model
 - SUSY
 - Axions
 - CP-Violation
 - Muon-specific 2HDM (1705.01469)
- N Doublet Models
 - Can include more doublets
 - "Private" Higgs
 - Extend muon-specific 2HDM



2HDMs

• 2HDM Potential

$$V_{2} = m_{11}^{2} \Phi_{1}^{\dagger} \Phi_{1} + m_{22}^{2} \Phi_{2}^{\dagger} \Phi_{2} + m_{12}^{2} \left(\Phi_{1}^{\dagger} \Phi_{2} + \Phi_{2}^{\dagger} \Phi_{1} \right)$$
$$V_{4} = \lambda_{1} \left(\Phi_{1}^{\dagger} \Phi_{1} \right)^{2} + \lambda_{2} \left(\Phi_{2}^{\dagger} \Phi_{2} \right)^{2} + \lambda_{3} \left(\Phi_{1}^{\dagger} \Phi_{1} \right) \left(\Phi_{2}^{\dagger} \Phi_{2} \right)$$
$$+ \lambda_{4} \left(\Phi_{1}^{\dagger} \Phi_{2} \right) \left(\Phi_{2}^{\dagger} \Phi_{1} \right) + \frac{\lambda_{5}}{2} \left[\left(\Phi_{1}^{\dagger} \Phi_{2} \right)^{2} + \left(\Phi_{2}^{\dagger} \Phi_{1} \right)^{2} \right]$$

- Yukawa couplings <u>not</u> yet defined
 - Different models



arXiv:1106.0034

Flavor Changing Neutral Current

• General form of Yukawa couplings for Q = -1/3 quarks is:

$$\mathcal{L}_Y = y_{ij}^1 \bar{\Psi}_i \Psi_j \Phi_1 + y_{ij}^2 \bar{\Psi}_i \Psi_j \Phi_2$$

- Yukawa couplings will not, in general, be simultaneously diagonalizable
 - Yukawa couplings will not be flavor diagonal
 - Leads to flavor-changing neutral currents (FCNC)
 - Can lead to processes such as $K \overline{K}$ mixing at tree level
- FCNC can be allowed but our model avoids it
 - By the Paschos-Glaschow-Weinberg theorem, all fermions with same quantum numbers must couple to the same Higgs doublet
 - For 2HDM, we have 3 groups: Q = 2/3 RH quarks, Q = -1/3 RH quarks, and RH leptons
 - For our 4HDM, we have 4 groups: quarks (q), electron (e), muon (μ), and tau (τ)



2HDM Models

- There are four possible coupling assignments which have no FCNC at tree-level
- Different Yukawa couplings between models
- All models have same neutral Higgs to W/Z boson couplings

Model	u_R^i	d_R^i	e_R^i
Type I	Φ_2	Φ_2	Φ_2
Type II	Φ_2	Φ_1	Φ_1
Type X	Φ_2	Φ_2	Φ_1
Type Y	Φ_2	Φ_1	Φ_2



Alignment Limit

- We can define two important parameters:
 - Angle which diagonalizes scalar mass-squared matrix: $\boldsymbol{\alpha}$

$$\begin{pmatrix} h_{125} \\ h_2 \end{pmatrix} = R_{\alpha} \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix}$$

- Angle which diagonalizes pseudoscalar and charged scalar mass-squared matrices: $\boldsymbol{\beta}$
 - Rotates into Higgs basis

$$\begin{pmatrix} G^{0} \\ A \end{pmatrix} = R_{\beta} \begin{pmatrix} \chi_{1} \\ \chi_{2} \end{pmatrix} \qquad \begin{pmatrix} G^{+} \\ H^{+} \end{pmatrix} = R_{\beta} \begin{pmatrix} \phi_{1}^{+} \\ \phi_{2}^{+} \end{pmatrix} \qquad \begin{pmatrix} h \\ H \end{pmatrix} = R_{\beta} \begin{pmatrix} \phi_{1} \\ \phi_{2} \end{pmatrix}$$

• Alignment Limit: SM Higgs is lightest neutral scalar



Our Model

• Scalar Potential:

$$V_{2} = \sum_{i} m_{ii}^{2} \Phi_{i}^{\dagger} \Phi_{i} + \sum_{i \neq j} m_{ij}^{2} \Phi_{i}^{\dagger} \Phi_{j}$$

$$V_{4} = \sum_{i} \lambda_{1}^{i} \left(\Phi_{i}^{\dagger} \Phi_{i}\right)^{2} + \frac{1}{2} \sum_{i \neq j} \left[\lambda_{3}^{ij} \left(\Phi_{i}^{\dagger} \Phi_{i}\right) \left(\Phi_{j}^{\dagger} \Phi_{j}\right) + \lambda_{4}^{ij} \left(\Phi_{i}^{\dagger} \Phi_{j}\right) \left(\Phi_{j}^{\dagger} \Phi_{i}\right) + \lambda_{5}^{ij} \left(\Phi_{i}^{\dagger} \Phi_{j}\right)^{2}\right]$$

$$\Phi_{i} = \begin{pmatrix} \phi_{i}^{+} \\ (v_{i} + \phi_{i} + i\chi_{i})/\sqrt{2} \end{pmatrix}$$

$$\lambda_{5}^{ij} = \left(\lambda_{5}^{ji}\right)^{*}$$

$$i, j \in \{q, e, \mu, \tau\}$$



No $q\tau$ - $e\mu$ Mixing

- $q\tau$ and $e\mu$ sectors are decoupled through absence of mixing
 - Results in block-diagonal diagonalization matrices
 - Diagonalization procedure similar to 2HDM
- In Type X model: $v_q^2 + v_{\tau}^2 = (246 \text{GeV})^2$
- In this model: $v_q^2 + v_{\tau}^2 < (246 \text{GeV})^2$
- Yukawa couplings will increase affecting production and decay of 125 GeV Higgs boson



VEV Analysis

- Compare production/decay to SM values
- Require μ_X to be consistent with unity within 20% at 95% CL

$$r \equiv \left(v_q^2 + v_\tau^2\right)^{1/2} / v$$

$$\mu_X \equiv \frac{\sigma \left(pp \to H \right) BR \left(H \to X \right)}{\sigma \left(pp \to H \right)_{SM} BR \left(H \to X \right)_{SM}}$$

 $X = gg, \mu\mu, \tau\tau, \bar{c}c, \bar{b}b, \bar{t}t, \gamma\gamma, \gamma Z, WW, ZZ$





Perturbation Theory

- Model produces extra massless scalars
 - Extra SU(2) symmetry
- Must be non-zero off-diagonal terms

$$\Phi_q = V_{11}h_1 + V_{12}h_2 + \dots$$

$$V_{11} = 1 - \frac{1}{2} \left(\lambda_5^{q\mu} v_q v_\mu\right)^2 \left[\left(\frac{c_{34}s_{12}}{m_{h_1}^2 - m_{h_3}^2}\right)^2 + \left(\frac{c_{12}s_{34}}{m_{h_1}^2 - m_{h_4}^2}\right)^2 \right]$$

- Using perturbation theory, V_{11} decreases which counteracts smaller v_q in Yukawa couplings
- $\lambda_5^{q\mu}$ has minimum value but neutral scalar mass (m_{h_3}) can be large



Aligned Model

- 125 GeV Higgs decays are consistent with SM
 - Multi-doublet models must be near alignment limit
- Model can be described using 31 parameters
 - 18 rotation angles (6 scalar, 6 pseudoscalar, 6 charged)
 - 12 masses (4 scalar, 4 pseudoscalar, 4 charged)
 - 1 vev (246 GeV)

$$\alpha_{1j} = \beta_j$$



Computational Technique

- Probing a 31-dimensional parameter space can be computationally expensive
- Generate random set of angles such that the alignment limit is obeyed
- Generate random masses
 - Still require 125 GeV scalar Higgs, 0 GeV pseudoscalar Higgs, and 0 GeV charged Higgs
- Calculate vevs
- Calculate Lagrangian parameters such that perturbativity is maintained
- Check BFB conditions
- Check experimental bounds



Neutral Higgs Decay



- Heavy neutral Higgs branching ratios
 - $t\bar{t}$ decay channel opens at 350 GeV
- Heavy neutral Higgs decays into leptons can probe this model
 - Current searches look at tauonic decays
 - Should also search for muonic and electronic decays

Charged Higgs Decay



- Heavy charged Higgs branching ratios
 - $t\overline{b}$ channel opens at 180 GeV
- Electronic decay branching ratio can be large when muonic decay branching ratio is near 1
- Possible to have electronic decay branching ratio larger than muonic decay branching ratio by order of magnitude





- We presented a 4HDM extension of the SM
 - Extends muon-specific model
- No $q\tau$ - $e\mu$ Mixing:
 - Place limit of 85% on $\left(v_q^2 + v_\tau^2\right)^{1/2}/v$
- Aligned Model:
 - Neutral and charged Higgs decays can have substantial electronic and muonic decays
 - Heavy Higgs decays can be searched through muonic or electronic branching ratios



EXTRA: Oblique Parameters

$$\begin{split} \frac{\alpha S}{4s_w^2 c_w^2} &= \left[\frac{\delta \Pi_{ZZ}(M_Z^2) - \delta \Pi_{ZZ}(0)}{M_Z^2} \right] - \frac{(c_w^2 - s_w^2)}{s_w c_w} \delta \Pi'_{Z\gamma}(0) - \delta \Pi'_{\gamma\gamma}(0), \\ \alpha T &= \frac{\delta \Pi_{WW}(0)}{M_W^2} - \frac{\delta \Pi_{ZZ}(0)}{M_Z^2}, \\ \frac{\alpha U}{4s_w^2} &= \left[\frac{\delta \Pi_{WW}(M_W^2) - \delta \Pi_{WW}(0)}{M_W^2} \right] - c_w^2 \left[\frac{\delta \Pi_{ZZ}(M_Z^2) - \delta \Pi_{ZZ}(0)}{M_Z^2} \right] \\ &- s_w^2 \delta \Pi'_{\gamma\gamma}(0) - 2s_w c_w \delta \Pi'_{Z\gamma}(0). \end{split}$$

arXiv:hep-ph/9306267v3



EXTRA: Conditions checked

- BFB Satisfied
- Perturbativity for lambda and Yukawas
- S,T are in acceptable range
- Charged Higgs < 80 GeV
- Charged scalar constributions to Higgs diphoton compatible with experiment
- Bounds from new physics contributions to B meson oscillations, K mesons are compatible with experiment
- $b \rightarrow s\gamma$ from charged Higgs are acceptable
- Heavy neutral Higgs decaying into tau pairs
- LHC direct searches for heavy charged Higgs



EXTRA: Lambda Extraction

$$\begin{split} \lambda_1^i =& \frac{1}{2v_i^3} \left(v_i M_{s,ii}^2 + \sum_{j \neq i} v_j m_{ij}^2 \right) \,, \\ \lambda_3^{ij} =& \frac{1}{v_i v_j} \left(M_{s,ij}^2 - 2M_{c,ij}^2 + m_{ij}^2 \right) \,, \\ \lambda_4^{ij} =& \frac{1}{v_i v_j} \left(2M_{c,ij}^2 - M_{p,ij}^2 - m_{ij}^2 \right) \,, \\ \lambda_5^{ij} =& \frac{1}{v_i v_j} \left(M_{p,ij}^2 - m_{ij}^2 \right) \,, \end{split}$$

