

Neutrino Annihilation Signatures from Inelastic DM in Neutron Stars

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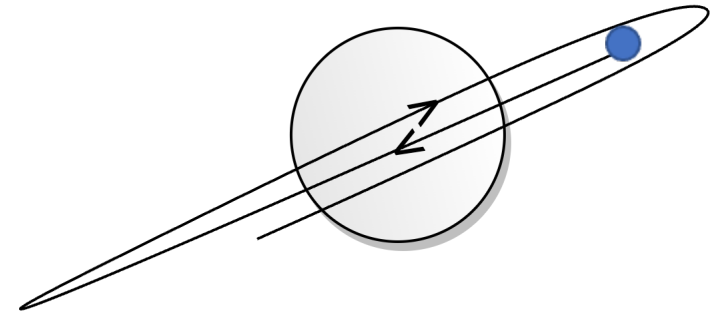
DM Capture in NS

DM can recoil against neutrons in the NS \longrightarrow losing energy

For capture, the dark matter must lose enough of its energy through collisions with scattering sites in the star

For a mass range of $\text{GeV} < m_\chi < \text{PeV}$

a single scatter is enough for the DM particle to lose enough energy to get completely captured.



$$C \propto \pi R^2 \frac{M_{NS}}{m_\chi} \text{Min} \left[1, \frac{\sigma}{\sigma_{sat}} \right]$$

DM Capture in NS

DM can recoil against neutrons in the NS \longrightarrow losing energy

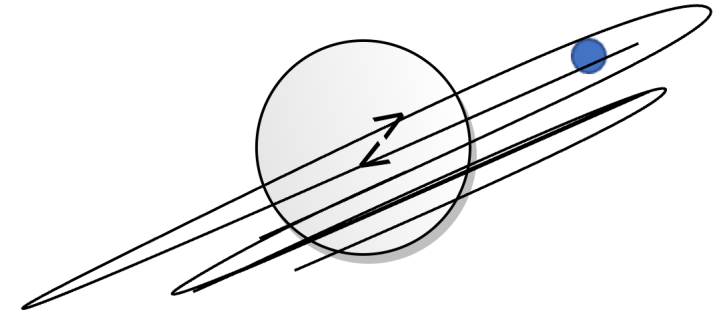
For capture, the dark matter must lose enough of its energy through collisions with scattering sites in the star

For heavy dark matter, $m_\chi > \text{PeV}$

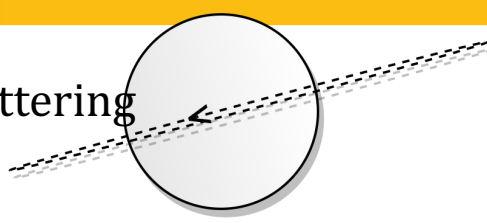
Heavier dark matter has more initial kinetic energy in the halo.

\therefore it needs to lose more energy to be captured, i.e., it needs to scatter more.

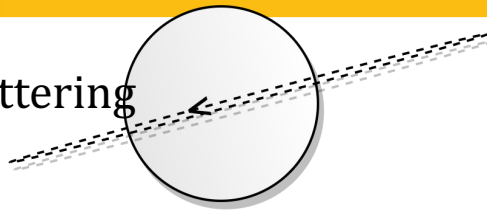
Multiscatter capture becomes more important in this case .



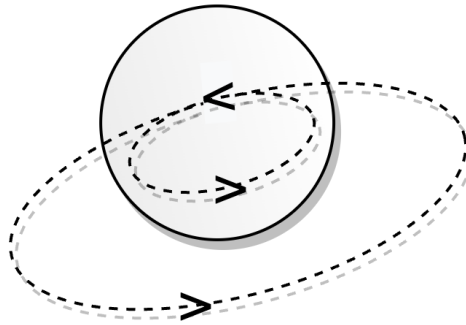
DM is captured as it loses energy by scattering



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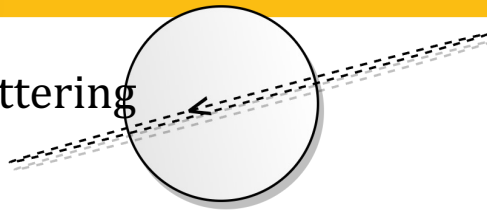


First stage of thermalization:
DM continues in a closed orbit
till its orbit is now contained
within the NS



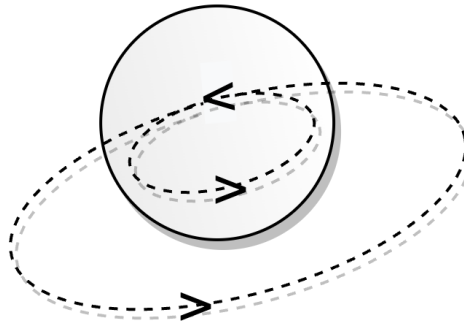
The timescale on which this happens is called $\tau_{1,therm}$

DM is captured as it loses energy by scattering



First stage of thermalization:

DM continues in a closed orbit till its orbit is now contained within the NS

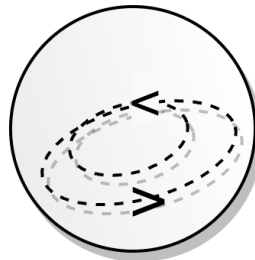


The timescale on which this happens is called $\tau_{1,therm}$

Second stage of thermalization:

DM moves completely inside the star on the orbit which shrinks to the thermal radius,

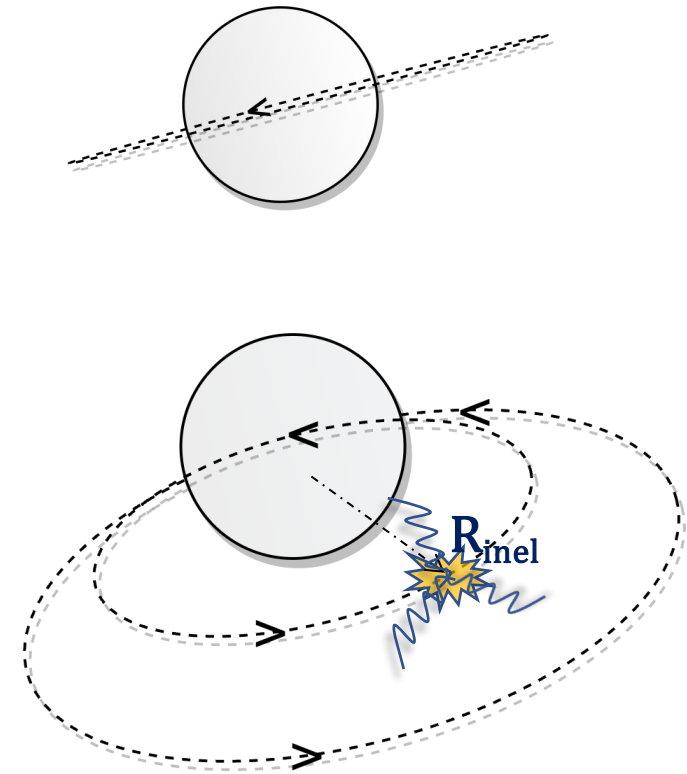
i.e., $T_x = T_{NS}$

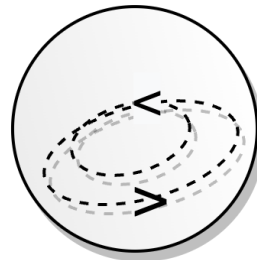
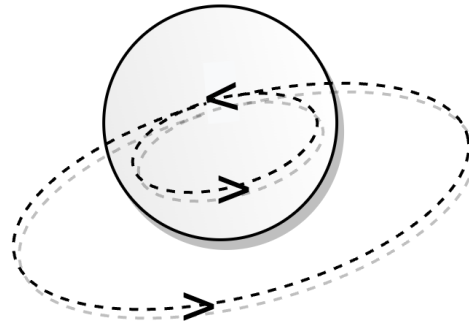
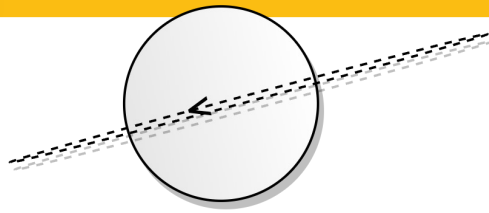


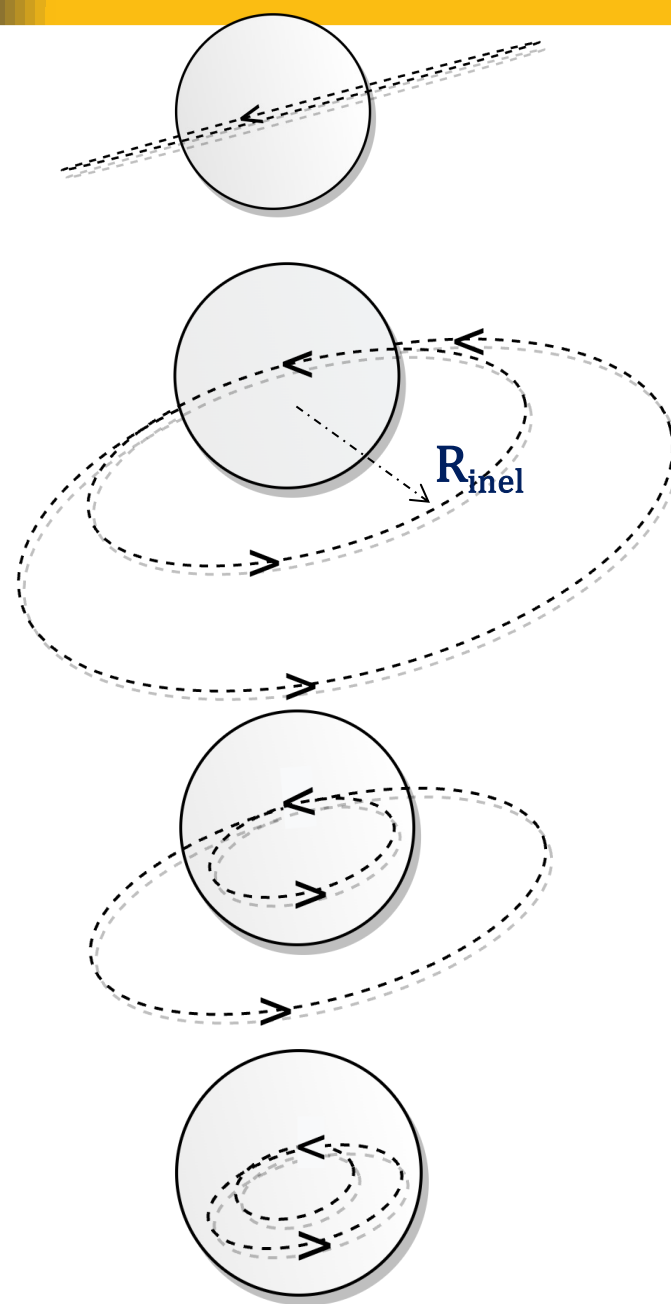
The timescale on which this happens is called $\tau_{2,therm}$

Inelastic DM

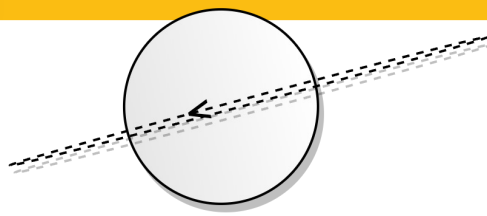
- We investigate the possibility of Inelastic DM thermalizing **outside** NS before it eventually settles in the NS.
- Depending on the DM model, it might be possible to detect this partial DM annihilation outside the NS, in the form of **neutrinos or gamma ray signatures**.





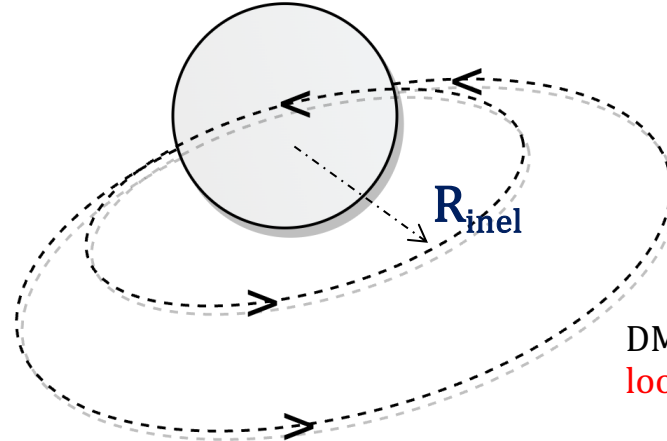


τ_{eq}
Rate of Capture = Rate of Annihilation



In the case of inelastic, $\tau_{eq} \ll \tau_{1,therm}$

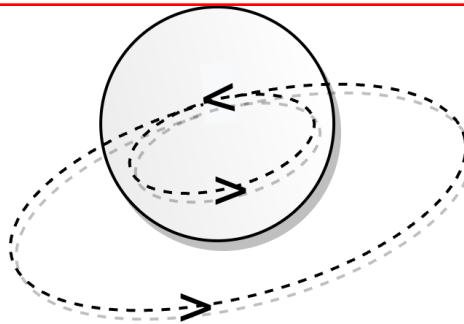
$\tau_{inel,therm}$
Time to thermalize to $R_{inel} > R_{NS}$



In the case of $\sigma_{inelastic}$, this is stage 2
&
 $\tau_{inel,therm} \ll \tau_{1,therm}$

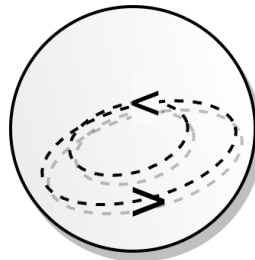
DM predominantly scatters inelastically off nuclei **with negligible loop level elastic scattering**

$\tau_{1,therm}$
Time to thermalize to $R < R_{NS}$



In the case of $\sigma_{elastic}/\sigma_{inelastic}$,
this is stage 2/stage 3

$\tau_{2,therm}$
Time to thermalize to $T_\chi = T_{NS}$

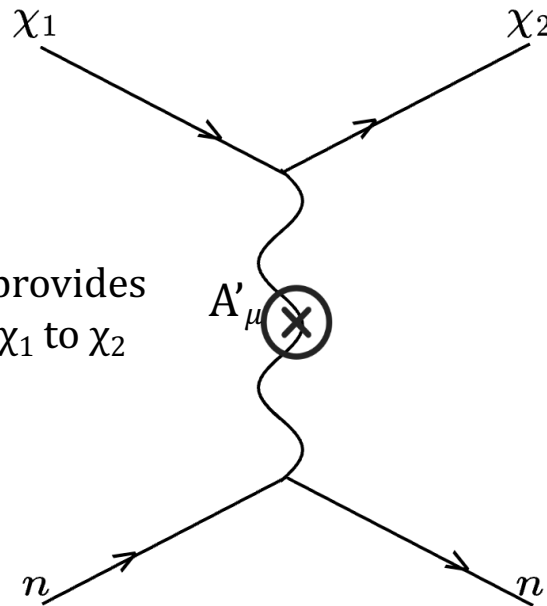


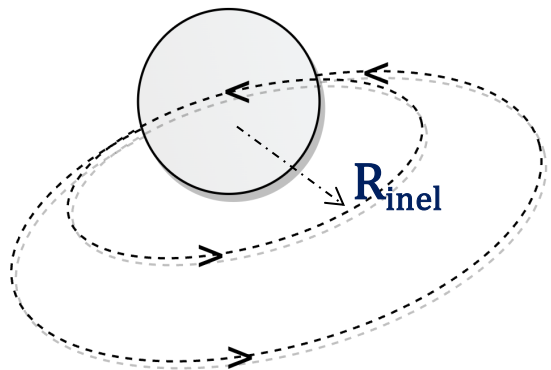
Inelastic Dark Matter

Inelastic dark matter models feature a DM particle χ_1 with mass m_χ

and a slightly heavier state χ_2 with mass $m_\chi + \delta m$, where $\delta m \ll m_\chi$

NS can accelerate DM to $\sim 0.7c$ during infall. This provides sufficient kinetic energy to allow up-scattering of χ_1 to χ_2

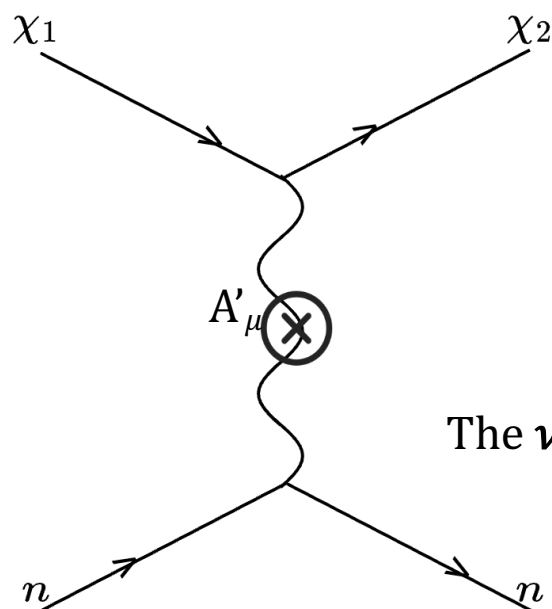




Inelastic Dark Matter

δ imposes an energy threshold below which interactions are suppressed.
(i.e. R_{inel} depends on δ)

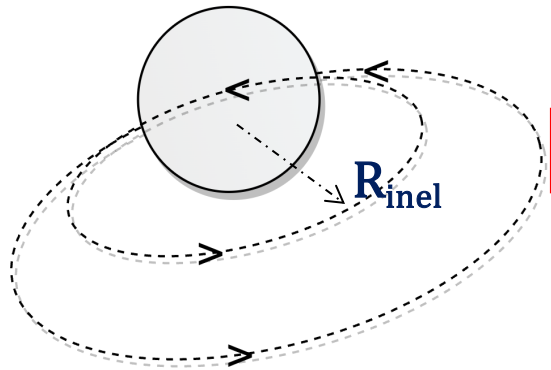
For $m_\chi > \text{GeV} \rightarrow \delta \in [50\text{-}300 \text{ MeV}]$



DM can partly thermalize and annihilate outside NS, and can produce SM particles

The ν (or γ) flux can be detected by ν (or γ) detectors

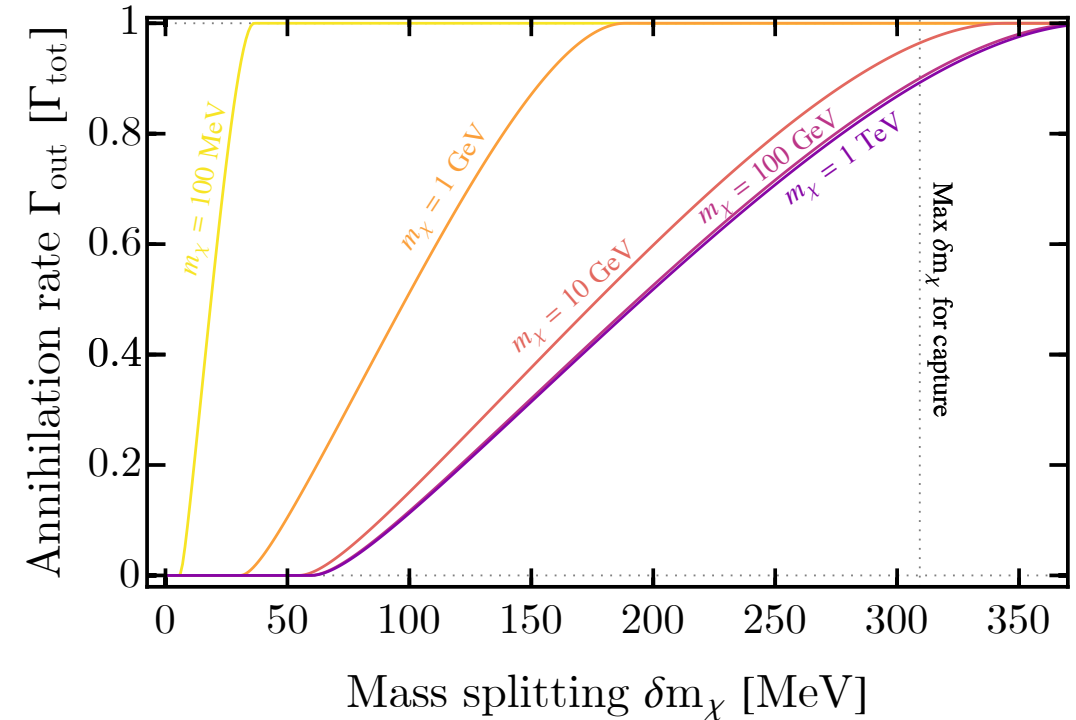
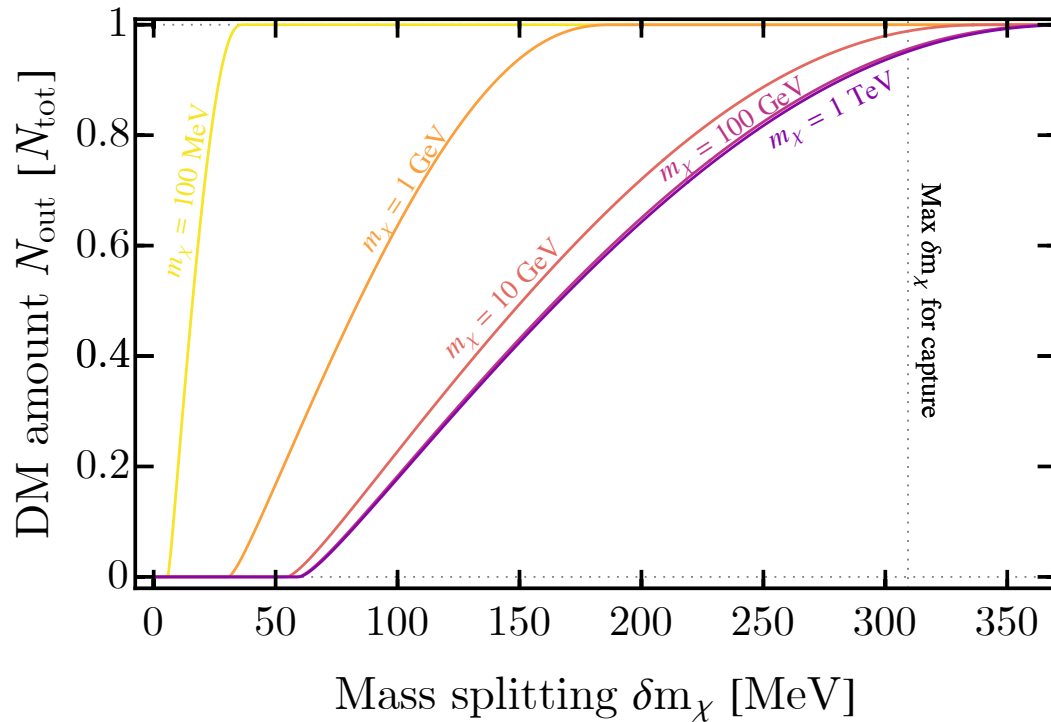
Timescales for DM thermalization



Timescale		
τ_{eq}	Time for capture-annihilation equilibrium	$\sim 9 \text{ hours}$
$\tau_{1,inel}$	Time to thermalize to R_{inel}	$\sim 12 \text{ hours}$
$\tau_{1,el}$	Time to thermalize to $R \leq R_{NS}$	$\sim 100 \text{ years}$
τ_2	Time to thermalize to $T_{NS} = T_\chi$	$\sim 10^5 \text{ years}$

Thermalization timescales for $m_\chi \sim \mathcal{O}(TeV)$, $\delta = 100 \text{ MeV}$, $\sigma_{inel} = 2 \times 10^{-45} \text{ cm}^2$, $\sigma_{el} \simeq 10^{-50} \text{ cm}^2$

Amount of DM annihilating outside NS



Fraction of the total captured and thermalized DM outside the neutron star comes from the condition

$$E_{\text{CM}} \gtrsim m_n + m_\chi + \delta m_\chi$$

Where E_{CM} is the center-of-mass energy of a dark matter particle transiting a neutron star

n_{NS} & ρ_{DM} Models

- ν flux produced by a single NS is too faint to be detected.
- $\sim 10^9$ NS predicted in the Milky Way with densest population in galactic center

\therefore Adopting NS distribution models

- Surface density of disk :

[arXiv: 0908.3182v1](#)

$$\log \Sigma(R) = a_0 + a_1 R + a_2 R^2 + a_3 R^3 + a_4 R^4$$

- Number density distribution :

[arXiv: 2101.12213](#)

$$n_{\text{NS}} = 5.98 \times 10^3 \left(\frac{r}{1 \text{ pc}} \right)^{-1.7} \text{ pc}^{-3}; \quad 0.1 \text{ pc} < r < 2 \text{ pc},$$

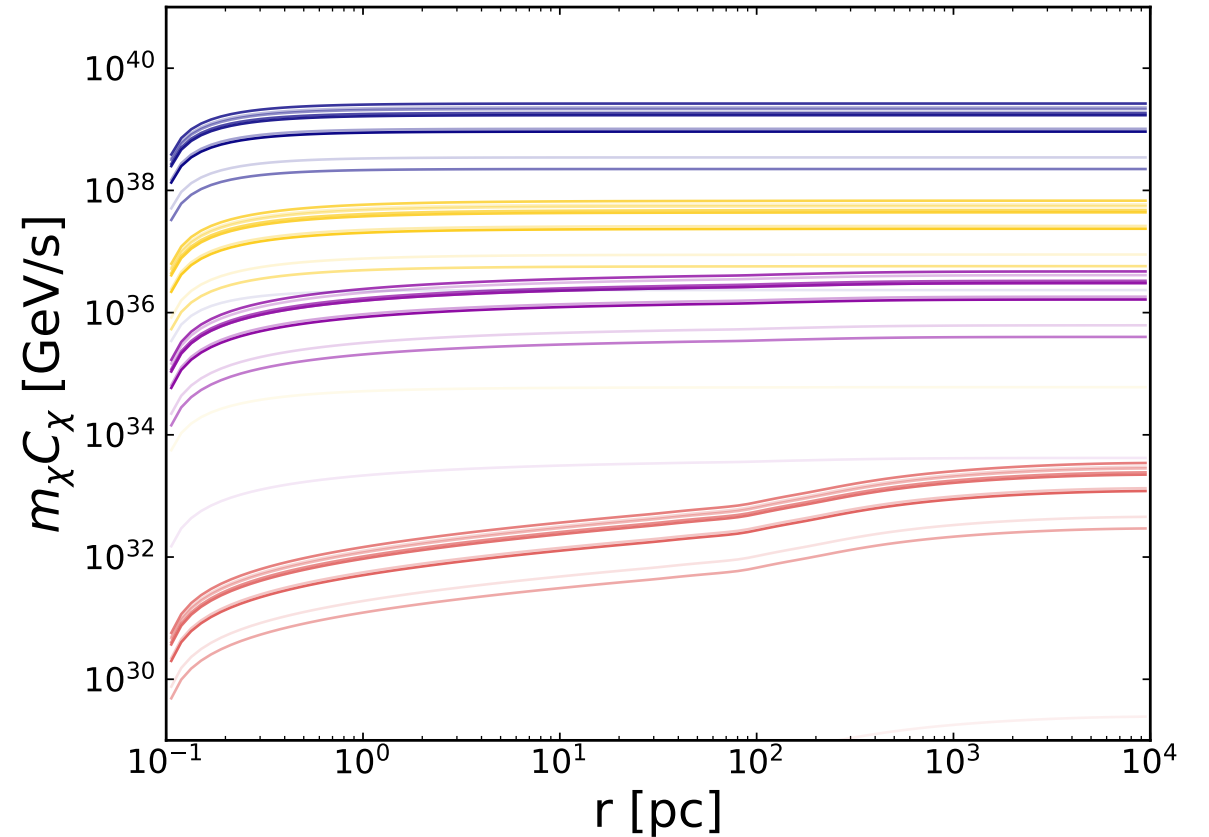
$$= 2.08 \times 10^4 \left(\frac{r}{1 \text{ pc}} \right)^{-3.5} \text{ pc}^{-3}; \quad r > 2 \text{ pc}$$

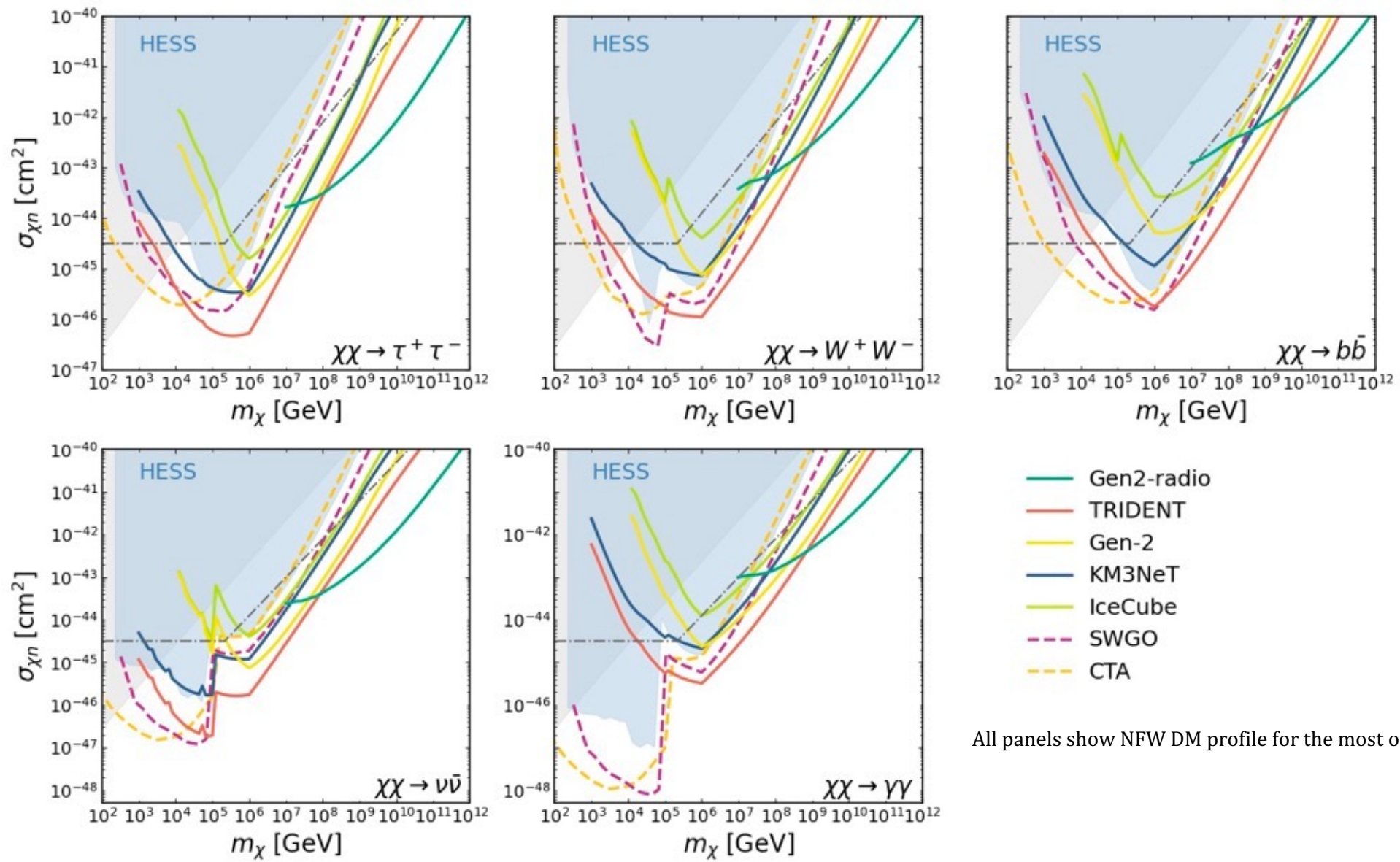
ν flux signal

Then ν flux for a distribution of NS models + DM halo density profiles –

$$\Gamma_{tot} = \int_{r_{min}}^{r_{max}} \Gamma_{ann}(r) \rho_{NS}(r) 4\pi r^2 dr,$$

$$\phi_{\nu/\gamma}(E_{\nu/\gamma}) = \frac{\Gamma_{out}}{4\pi D^2} \frac{dN_{\nu/\gamma, ch}}{dE_{\nu/\gamma}}(E_{\nu/\gamma})$$





All panels show NFW DM profile for the most optimistic NS distribution

References

- N. Sartore, E. Ripamonti, A. Treves, and R. Turolla, Galactic neutron stars in space and velocity distributions in the disk and in the halo, *Astronomy and Astrophysics* 510 (2009).
- R. K. Leane, T. Linden, P. Mukhopadhyay, and N. Toro, Celestial-body focused dark matter annihilation throughout the Galaxy, *Phys. Rev. D* 103, 075030 (2021), arXiv:2101.12213 [astro-ph.HE].
- M. Baryakhtar, J. Bramante, S. W. Li, T. Linden, and N. Raj, Dark Kinetic Heating of Neutron Stars and an Infrared Window on WIMPs, SIMPs, and Pure Higgsinos, *Phys. Rev. Lett.* 119, 131801 (2017), arXiv:1704.01577 [hep-ph].
- J. F. Acevedo, J. Bramante, R. K. Leane, and N. Raj, Warming Nuclear Pasta with Dark Matter: Kinetic and Annihilation Heating of Neutron Star Crusts, *JCAP* 03, 038, arXiv:1911.06334 [hep-ph].
- M. Baryakhtar, J. Bramante, S. W. Li, T. Linden, and N. Raj, Dark Kinetic Heating of Neutron Stars and An Infrared Window On WIMPs, SIMPs, and Pure Higgsinos, *Phys. Rev. Lett.* 119, 131801 (2017), arXiv:1704.01577 [hep-ph].
- J. Bramante, A. Delgado, and A. Martin, Multiscatter stellar capture of dark matter, *Phys. Rev. D* 96, 063002 (2017), arXiv:1703.04043 [hep-ph].
- M. Baryakhtar, J. Bramante, S. W. Li, T. Linden, and N. Raj, Dark Kinetic Heating of Neutron Stars and an Infrared Window on WIMPs, SIMPs, and Pure Higgsinos, *Phys. Rev. Lett.* 119, 131801 (2017), arXiv:1704.01577 [hep-ph].
- D. Bose, T. N. Maity, and T. S. Ray, Neutrinos from captured dark matter annihilation in a galactic population of neutron stars, *JCAP* 2022 (5), 001, arXiv:2108.12420
- A. D. Avrorin et al. (Baikal-GVD), Baikal-GVD: status and prospects, *EPJ Web Conf.* 191, 01006 (2018), arXiv:1808.10353 [astro-ph.IM].
- S. Adrian-Martinez et al. (KM3Net), Letter of intent for KM3NeT 2.0, *J. Phys. G* 43, 084001 (2016), arXiv:1601.07459 [astro-ph.IM].
- M. G. Aartsen et al. (IceCube-Gen2), IceCube-Gen2: the window to the extreme Universe, *J. Phys. G* 48, 060501 (2021), arXiv:2008.04323 [astro-ph.HE].
- M. Agostini et al. (P-ONE), The Pacific Ocean Neutrino Experiment, *Nature Astron.* 4, 913 (2020), arXiv:2005.09493 [astro-ph.HE].
- Z. P. Ye et al., Proposal for a neutrino telescope in South China Sea, (2022), arXiv:2207.04519 [astro-ph.HE].
- J. A. Aguilar et al. (RNO-G), Design and Sensitivity of the Radio Neutrino Observatory in Greenland (RNO-G), *JINST* 16 (03), P03025, arXiv:2010.12279 [astro-ph.IM].
- J. Alvarez-Muniz et al. (GRAND), The Giant Radio Array for Neutrino Detection (GRAND): Science and Design, *Sci. China Phys. Mech. Astron.* 63, 219501 (2020), arXiv:1810.09994 [astro-ph.HE].
- A. V. Olinto et al. (POEMMA), The POEMMA (Probe of Extreme Multi-Messenger Astrophysics) observatory, *JCAP* 06, 007, arXiv:2012.07945 [astro-ph.IM].
- P. Allison et al., Measurement of the real dielectric permittivity ϵ_r of glacial ice, *Astropart. Phys.* 108, 63 (2019), arXiv:1712.03301 [astro-ph.IM].
- T. K. Gaisser, Spectrum of cosmic-ray nucleons, kaon production, and the atmospheric muon charge ratio, *Astropart. Phys.* 35, 801 (2012), arXiv:1111.6675 [astro-ph.HE].
- R. Enberg, M. H. Reno, and I. Sarcevic, Prompt neutrino fluxes from atmospheric charm, *Phys. Rev. D* 78, 043005 (2008), arXiv:0806.0418 [hep-ph].
- A. D. Martin, M. G. Ryskin, and A. M. Stasto, Prompt neutrinos from atmospheric $\bar{c}c$ and $b\bar{b}$ production and the gluon at very small x , *Acta Phys. Polon. B* 34, 3273 (2003), arXiv:hep-ph/0302140.
- R. Abbasi et al. (IceCube), The IceCube high-energy starting event sample: Description and flux characterization with 7.5 years of data, *Phys. Rev. D* 104, 022002 (2021), arXiv:2011.03545 [astro-ph.HE].
- R. Abbasi et al. (IceCube), Improved Characterization of the Astrophysical Muon–neutrino Flux with 9.5 Years of IceCube Data, *Astrophys. J.* 928, 50 (2022), arXiv:2111.10299 [astro-ph.HE].
- J. Aalbers et al. (LZ), First Dark Matter Search Results from the LUX-ZEPLIN (LZ) Experiment, (2022), arXiv:2207.03764 [hep-ex].
- C. Amole et al. (PICO), Dark Matter Search Results from the Complete Exposure of the PICO-60 C3F8 Bubble Chamber, *Phys. Rev. D* 100, 022001 (2019), arXiv:1902.04031 [astro-ph.CO].



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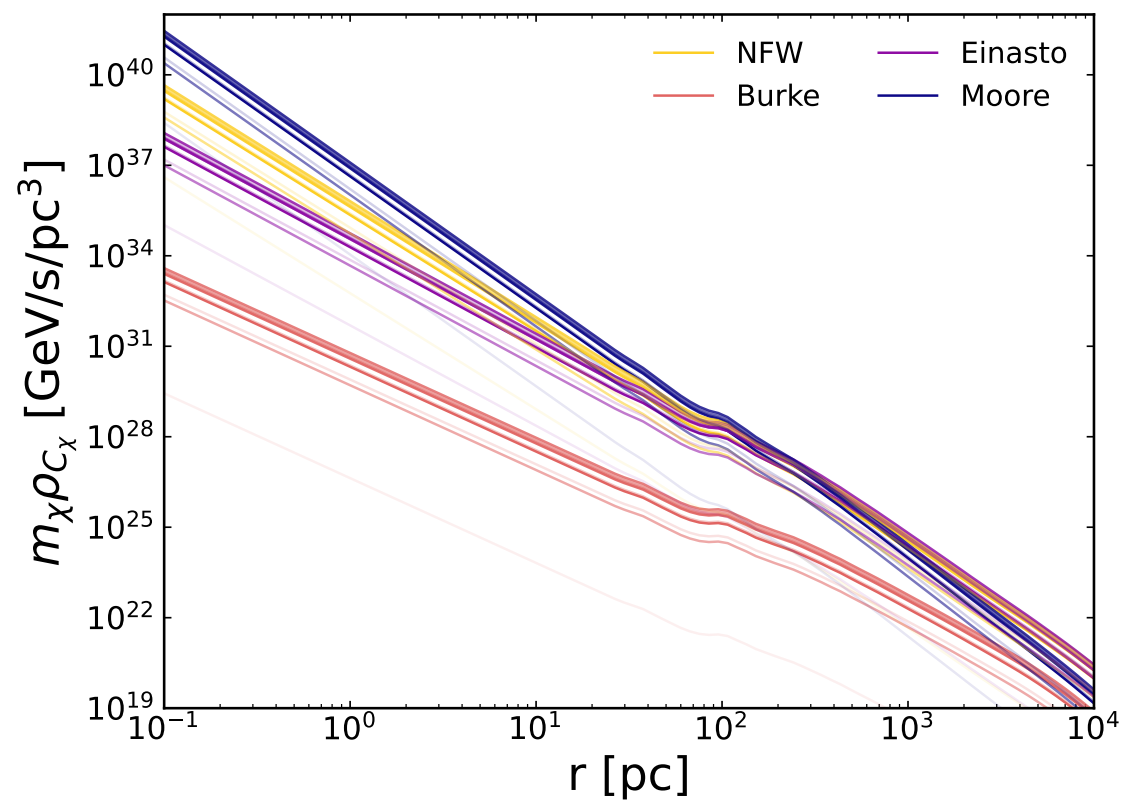
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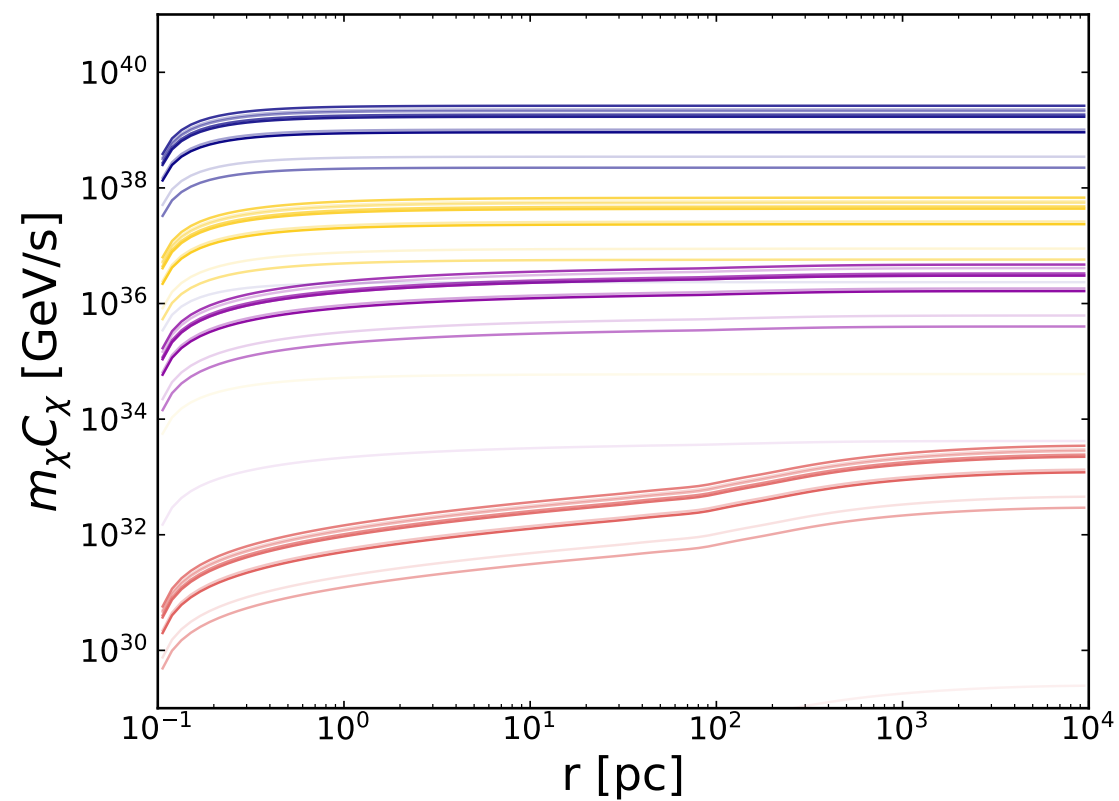
Questions?

Backup Slides

Mass capture rate densities



Mass capture rates



η_{NS} & ρ_{DM} Models

DM density profiles : $\rho_{\text{NFW}}(r) = \frac{\rho_0}{\left(\frac{r}{r_s}\right) \left[1 + \left(\frac{r}{r_s}\right)\right]^2}$

$$\rho_{\text{Moore}}(r) = \frac{\rho_0}{\left(\frac{r}{r_s}\right)^\gamma \left[1 + \left(\frac{r}{r_s}\right)\right]^{3-\gamma}}$$

$$\rho_{\text{Bur}}(r) = \frac{\rho_0}{\left(1 + \frac{r}{r_s}\right) \left[1 + \left(\frac{r}{r_s}\right)^2\right]}$$

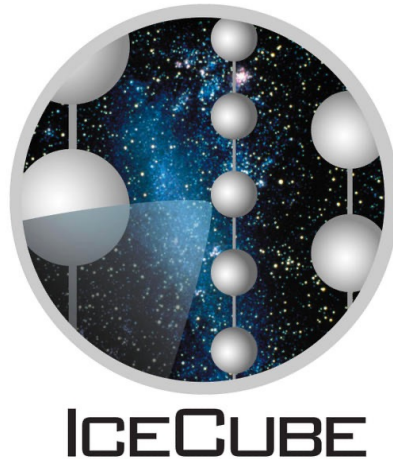
$$\rho_{\text{Einasto}}(r) = \rho_0 \exp \left[- \left(\frac{2}{\alpha} \right) \left\{ \left(\frac{r}{r_s} \right)^\alpha - 1 \right\} \right]$$

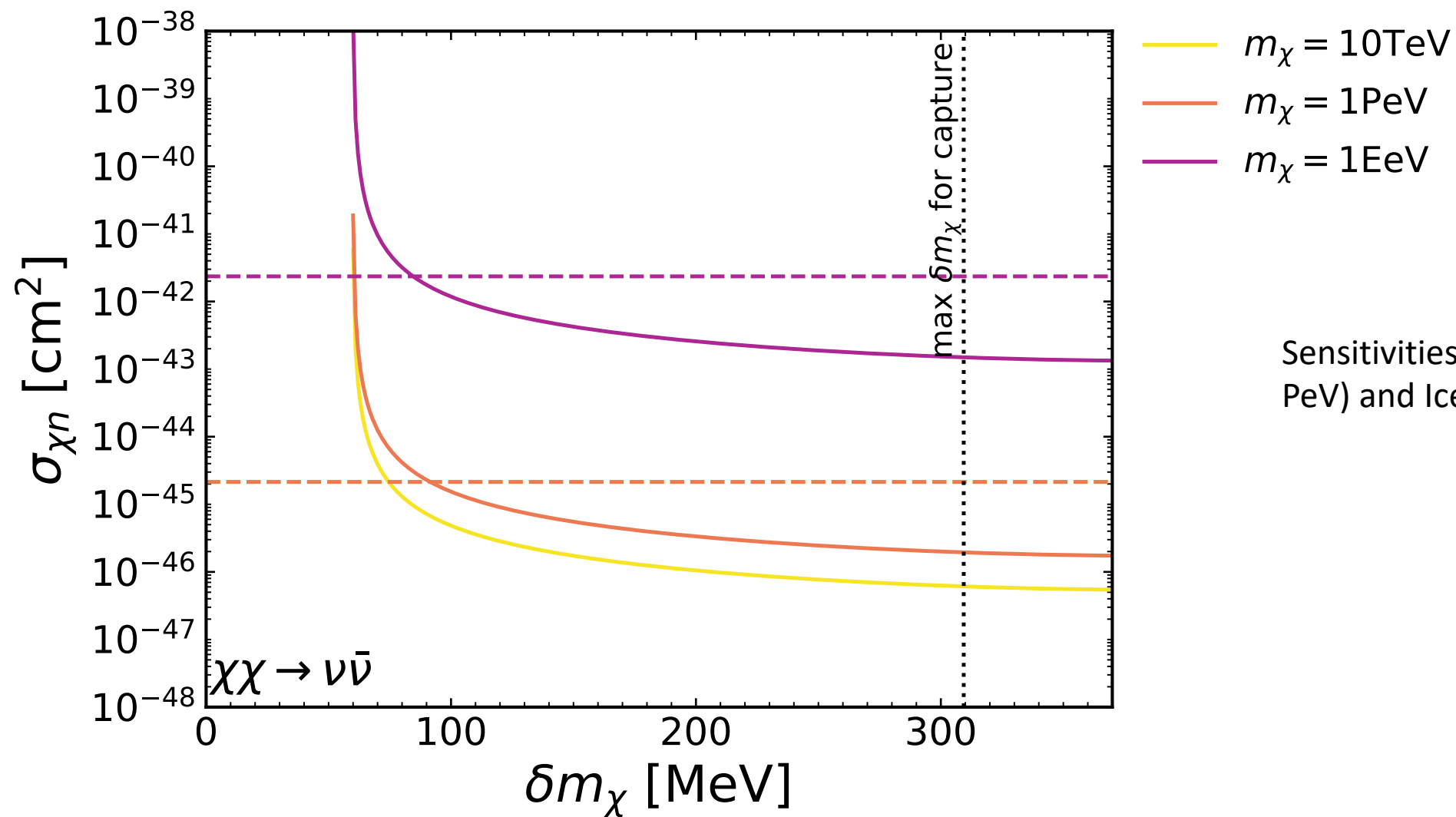
Then ν flux for a distribution of NS models + DM halo density profiles –

$$\phi = \int_r d^3r \eta_{\text{NS}}(r) \dot{m} \left(\frac{\rho_\chi(r)}{0.42 \text{ GeV cm}^{-3}} \right) \left(\frac{230 \text{ km s}^{-1}}{v_\chi(r)} \right)$$

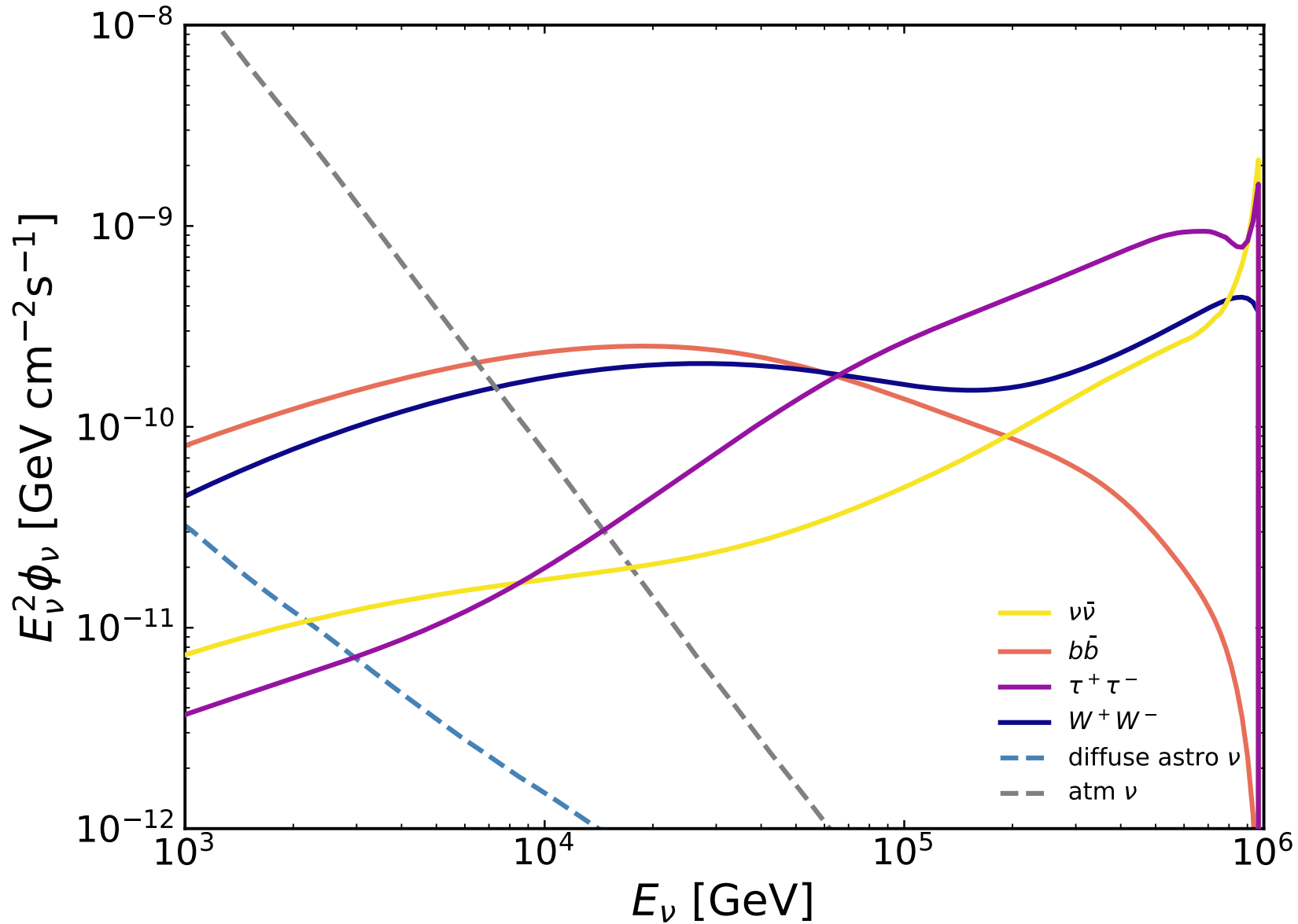
Note: halo velocities have been adopted from MW RC [[Sofue Y. 2020](#)]

Telescope Sensitivities & Detection Prospects





Sensitivities of TRIDENT (10 TeV and 1 PeV) and IceCube-Gen2 radio (1EeV).



DM=1 PeV with the background fluxes , the capture rate used is the maximum.