



Detecting Ultralight Dark Photon Dark Matter Using Optomechanical Sensors

2305.XXXX

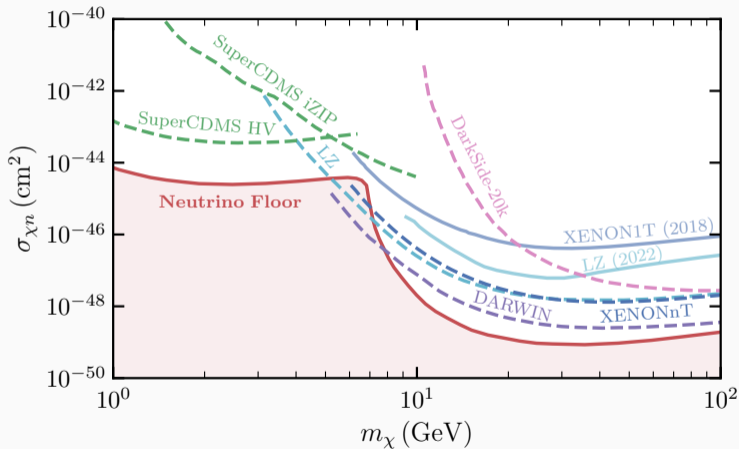
Pheno 2023

Dorian Amaral¹ Erqian Cai¹ Christopher Tunnell¹

May 8th, 2023

¹Rice University

The WIMP Era in Distress



The WIMP Era in Distress (According to ChatGPT)

*WIMPs retreat, unseen,
Alternatives now emerge,
Mysteries persist.*

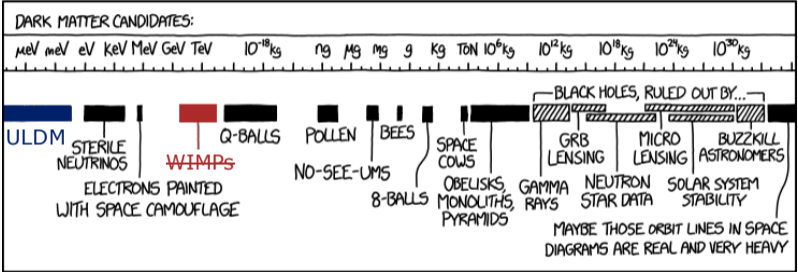
— ChatGPT

The WIMP Era in Distress (According to ChatGPT)

WIMPs retreat, unseen,
Alternatives now emerge,
Mysteries persist.

— ChatGPT

Beyond the WIMP Paradigm

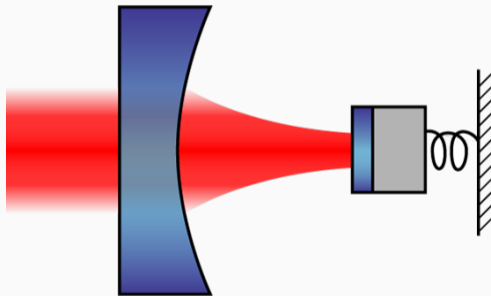


Adapted from xkcd

- Dark photon from a new gauged $U(1)_{B-L}$ symmetry is a popular BSM extension
- Characterised by two parameters: coupling (g_{B-L}) and mass (m_{DM})
- Ultralight dark photons lead to an ever-present, **oscillating** background **dark electric field**
- This leads to a very small differential acceleration between materials with different charge-mass ratios!

The Optomechanical Sensor

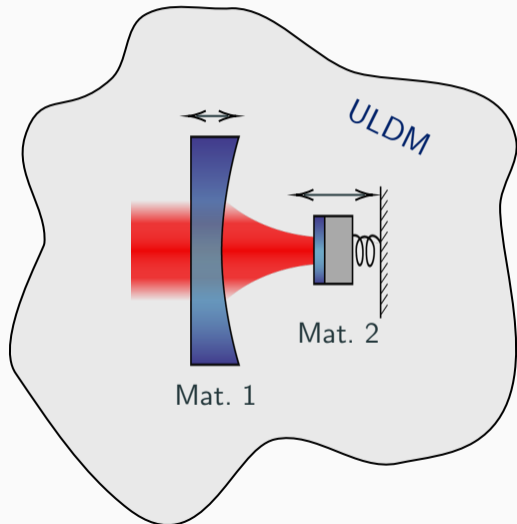
- Sensors use coupling between optical and mechanical modes to measure quantities very precisely
- Incredible developments due to gravitational wave efforts!
- Important parameters:
 - Mechanical resonant frequency (ω_0)
 - Temperature of surroundings (T)
 - Sensor mass (m_{sensor})
 - Mode decay rates



Detecting Ultralight Dark Matter

- Sensor constantly immersed in **oscillating dark electric field**
- If mirrors made of different materials, get **differential acceleration**
- This is our measurable!

$$\Delta a \sim \underbrace{g_{B-L}}_{\text{Coupling}} \underbrace{\left| \frac{N_1}{A_1} - \frac{N_2}{A_2} \right|}_{\text{Diff. } q-m \text{ ratio}} \underbrace{a_0}_{\text{DM Frequency}} \cos(\omega_{\text{DM}} t)$$

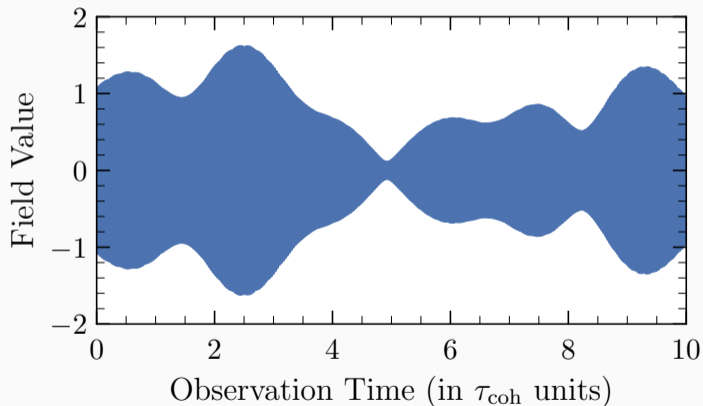


Our Goal

- Cavity accelerometers have been used to draw limits on $B - L$ dark photon DM
Peter W. Graham et al. **1512.06165**, Daniel Carney et al. **1908.04797**,
Jack Manley et al. **2007.04899**
- **However**, a likelihood-led treatment incorporating stochastic field properties and DM signal shape is lacking
- Want to develop this treatment in contact with experimentalists in **Windchime Collaboration!**

How can optomechanical sensors help us to learn more about the nature of ULDM?

The Stochastic Field: Long Observation Time



- Net field a superposition of many partial waves oscillating at slightly different frequencies

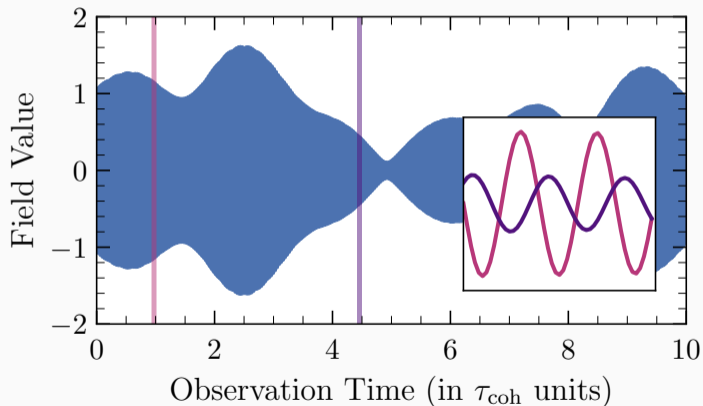
$$\omega_{\text{DM}} = m_{\text{DM}} + \frac{1}{2} m_{\text{DM}} v^2$$

- Leads to 'beat-like' effect on timescales

$$\tau_{\text{coh}} = \frac{1}{m_{\text{DM}} v^2} \sim \frac{10^6}{m_{\text{DM}}}$$

- For $T_{\text{obs}} \gg \tau_{\text{coh}}$, our sensor samples many coherent patches

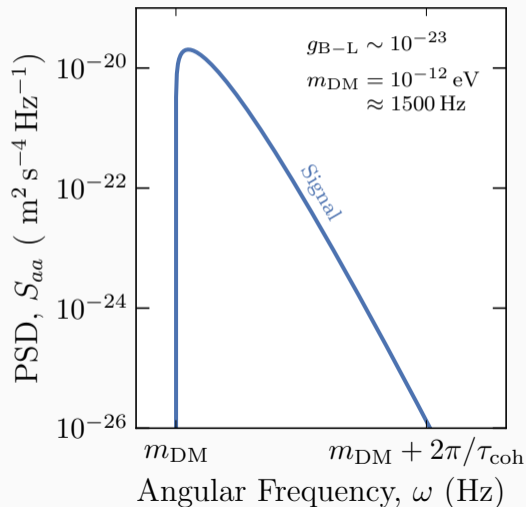
The Stochastic Field: Short Observation Time



- For $T_{obs} \ll \tau_{coh}$, our sensor samples one coherent patch
- We observe a clean sinusoid in time series
- But now the amplitude is **stochastic!**

Focus here on $T_{obs} \gg \tau_{coh}$

ULDM in Frequency Space

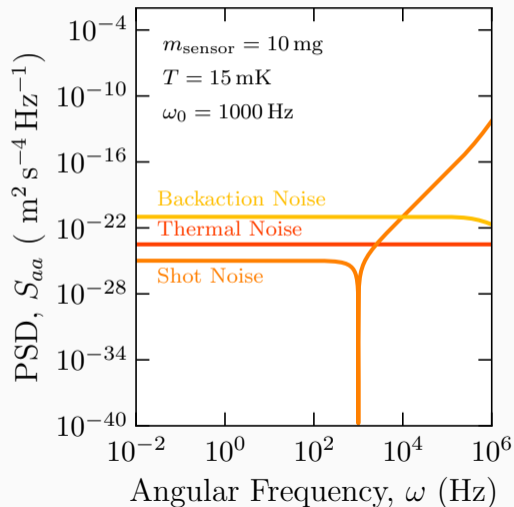


- Fourier space is best place to do our inferencing
- We consider the **power spectral density**, PSD (power per frequency)
- In $T_{\text{obs}} \gg \tau_{\text{coh}}$ case, signal appears as a **broadened peak** at $\omega \approx m_{\text{DM}}$
- Each frequency bin is exponentially distributed

Joshua W. Foster et al. [1711.10489](#)

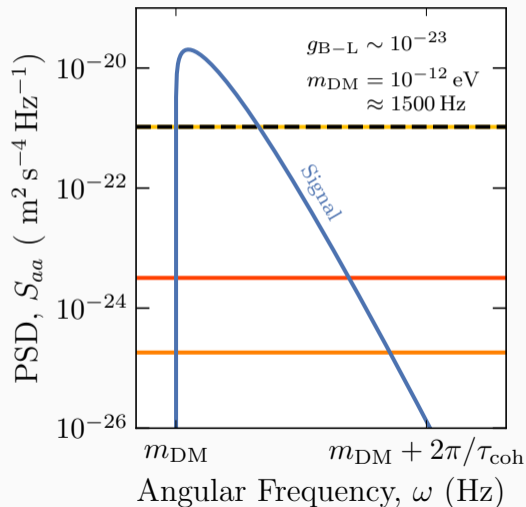
Likelihood(data) \sim Exp(\langle signal \rangle)

ULDM in Frequency Space: Signal vs Background



- We have to distinguish this peak from many background sources:
 - Thermal noise
 - Shot noise
 - Backaction noise

ULDM in Frequency Space: Signal vs Background



- We have to distinguish this peak from many background sources:
 - Thermal noise
 - Shot noise
 - Backaction noise

Frequentist Limit Setting

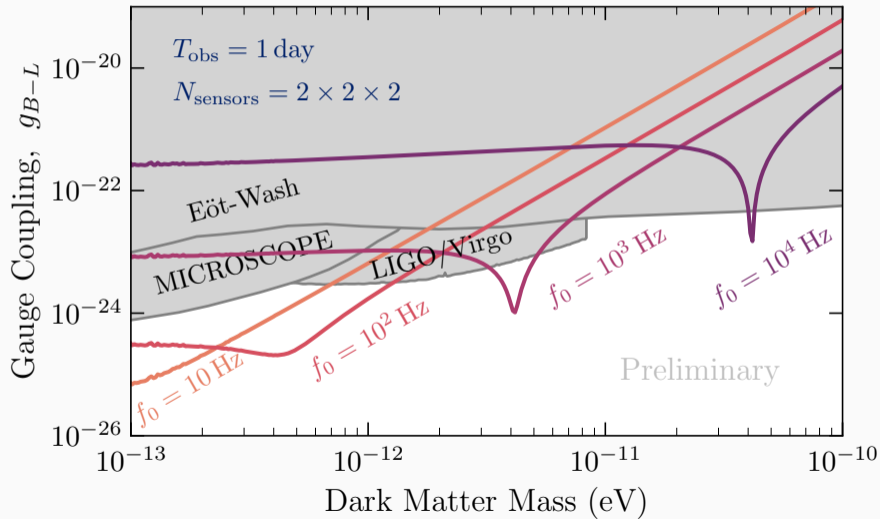
1. Assume we see nothing: 'generate' Asimov background data set
2. For a given m_{DM} , construct signal given a $g_{B-L} \neq 0$
3. Consider log-likelihood ratio test statistic

$$q = -2 \ln \left[\frac{\mathcal{L}(\text{data} | g_{B-L}, m_{\text{DM}})}{\mathcal{L}(\text{data} | g_{B-L} = 0, m_{\text{DM}})} \right]$$

(Likelihood for axions developed by [Joshua W. Foster et al. 1711.10489](#)):

4. Use q to exclude coupling at desired confidence level! (90%)

Limits



- WIMPs are facing an **existential crisis**
- Ultralight $B - L$ dark photons are a **well-motivated DM candidate**
- Optomechanical sensors at up to the task of measuring effects of dark electric field
- A **small** array could quickly **exclude new parameter space!**

Optomechanical sensors will form powerful probes of ultralight dark photon dark matter!

$$A_x \sim \sum_{i=1}^{N_{\text{waves}}} \frac{\sqrt{\rho_{\text{DM}}}}{m_{Z'}} \text{Re} \left(e^{im_{\text{DM}}(1+\frac{1}{2}v_i^2)t+\theta_i} \right) \hat{\mathbf{e}} \cdot \hat{\mathbf{x}}$$

$$v \sim f_{\text{SHM}}(\mathbf{v})$$

$$\theta \sim U(0, 2\pi)$$

$$\hat{\mathbf{e}} \sim \text{Unit Sphere}$$

$$S_{aa}^{\text{Th}} \equiv \frac{4k_B T \gamma}{m_s}$$

$$S_{aa}^{\text{SN}}(\omega) \equiv \frac{\hbar \kappa L^2}{2\omega_L P_L} |\chi_c(\omega)|^{-2} |\chi_m(\omega)|^{-2}$$

$$S_{aa}^{\text{BA}}(\omega) \equiv \frac{2\hbar \omega_L P_L}{m_s^2 L^2 \kappa} |\chi_c(\omega)|^2$$

$$|\chi_m(\omega)|^{-2} = (\omega^2 - \omega_0^2)^2 + \gamma^2 \omega^2$$

$$|\chi_c(\omega)|^{-2} = \frac{\omega^2 + \kappa^2/4}{\kappa}.$$

$$q = 2 \sum_k \left[S_{aa}^k \left(\frac{1}{\lambda_k(\boldsymbol{\theta})} - \frac{1}{\lambda_k(\hat{\boldsymbol{\theta}})} \right) + \ln \left(\frac{\lambda_k(\boldsymbol{\theta})}{\lambda_k(\hat{\boldsymbol{\theta}})} \right) \right]$$

$$p \equiv \int_{q^{\text{lim}}}^{\infty} f(q) \, dq = 1 - \alpha,$$

$$f(q) = \frac{1}{2} \delta(0) + \frac{1}{2} \chi_{k=1}^2(q).$$