

Probing BSM physics with Solar and Atmospheric Neutrinos in Dark Matter Experiments

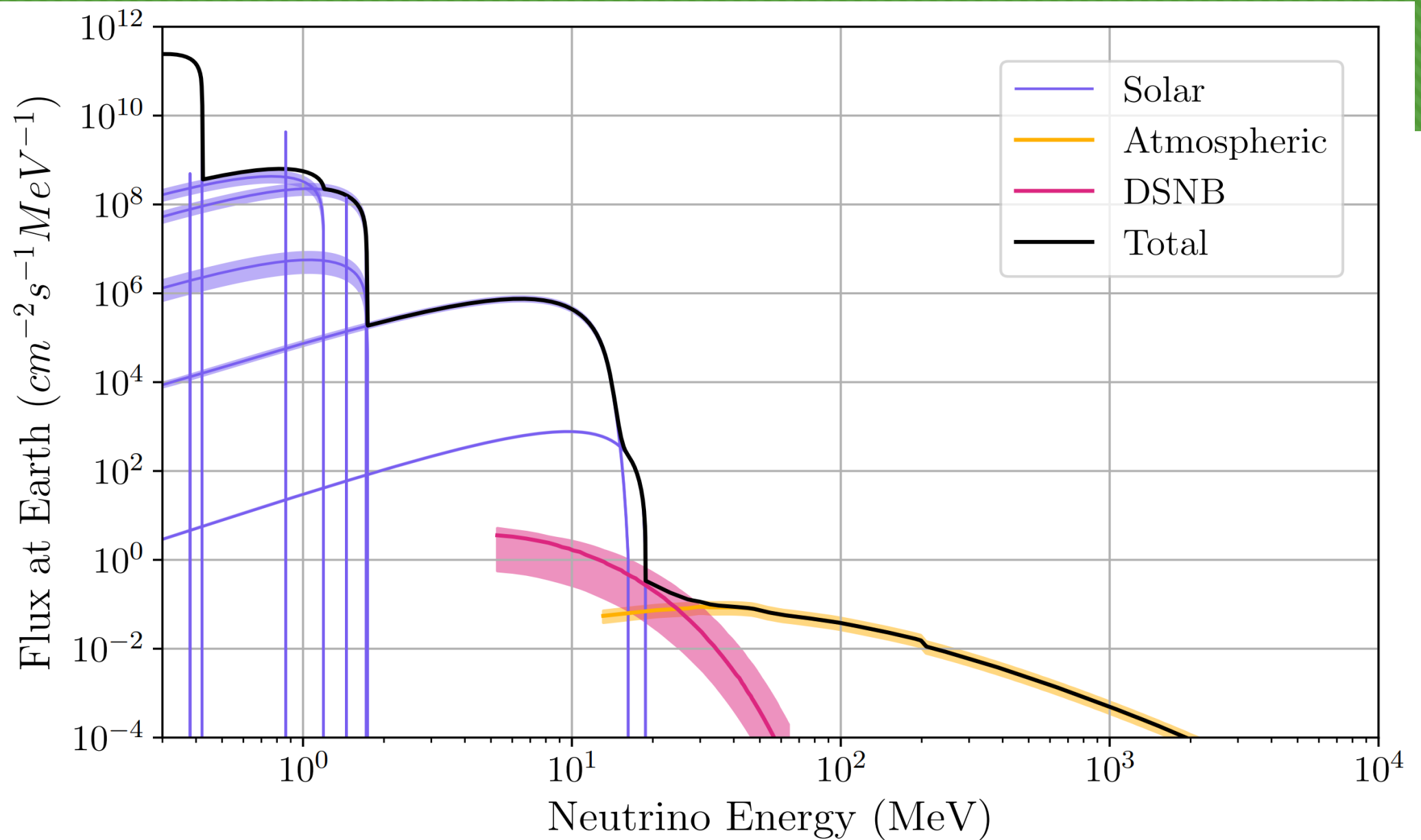
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Phenomenology 2023 Symposium — Pittsburg, PA



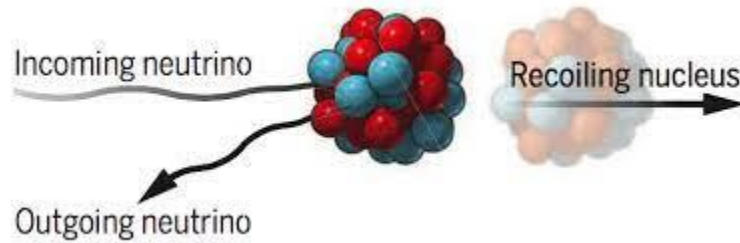
Based on work with Tien-Tien Yu and Volodymyr Takhistov

The Neutrino Flux



Neutrino Scattering

Coherent Scattering:



Differential recoil rate:

$$\frac{dR}{dE_R} = N_T \int_{E_\nu^{min}} \frac{d\sigma}{dE_R} \frac{dN_\nu}{dE_\nu} dE_\nu$$

A little kinematics:

$$E_\nu^{min} = \frac{1}{2} \left(E_R + \sqrt{E_R^2 + 2E_R m_N} \right)$$

SM cross-section:

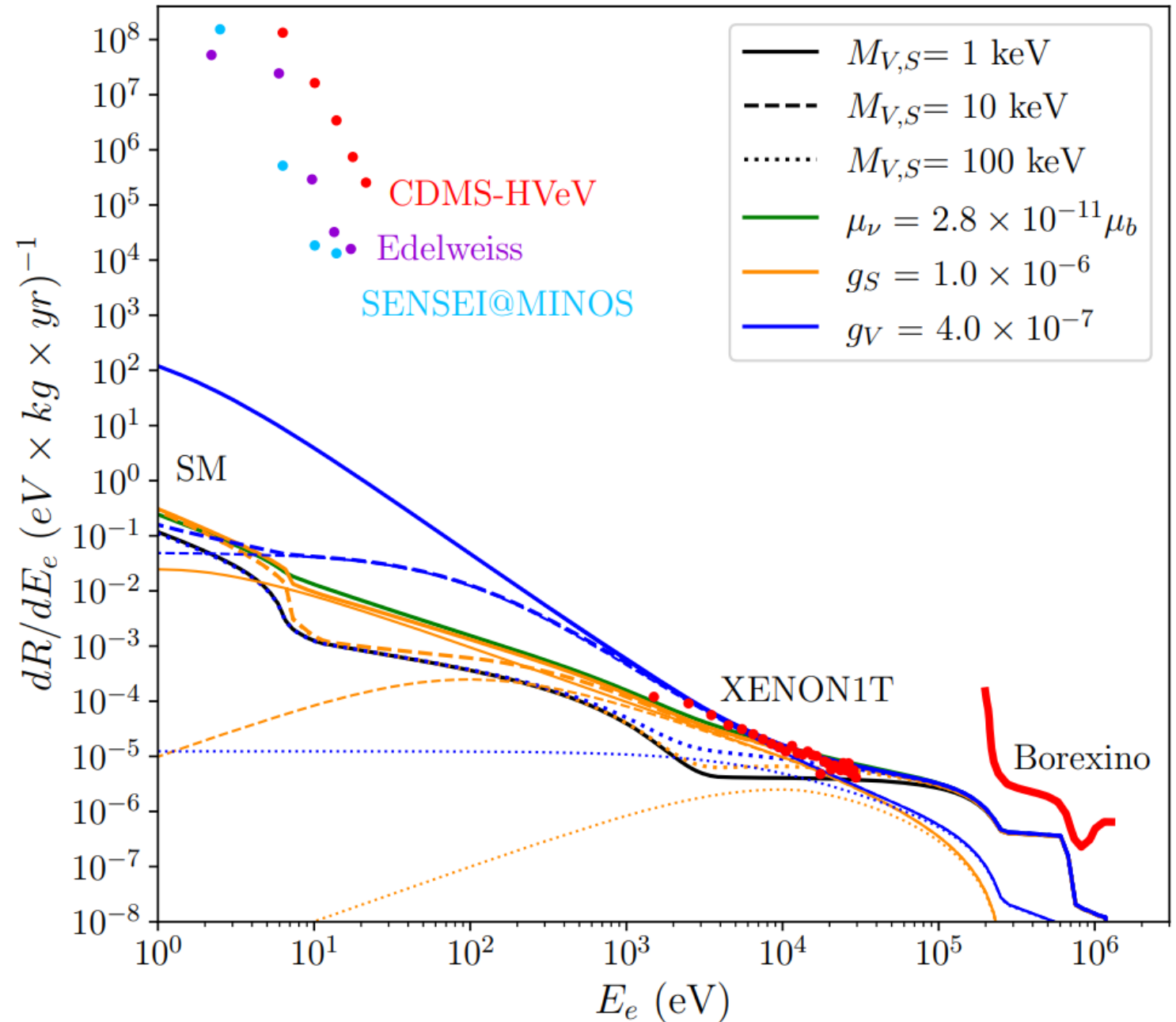
$$\frac{d\sigma}{dE_N} = \frac{G_F^2}{4\pi} Q_v^2 m_N \left(1 - \frac{m_N E_N}{2E_\nu^2} \right) F^2(E_N)$$

Recoil Energy

Scattering events produce a spectrum of recoils

Light vector mediators are most visible at low energy

Looking at higher energy has advantages for scalar interactions



Scalar NSI

- Two dimension 4 operators

$$\mathcal{L} \supset \phi (g_{\nu S} \bar{\nu}_R \nu_L + g_{qS} \bar{q} q)$$

- Dimension 6 operator

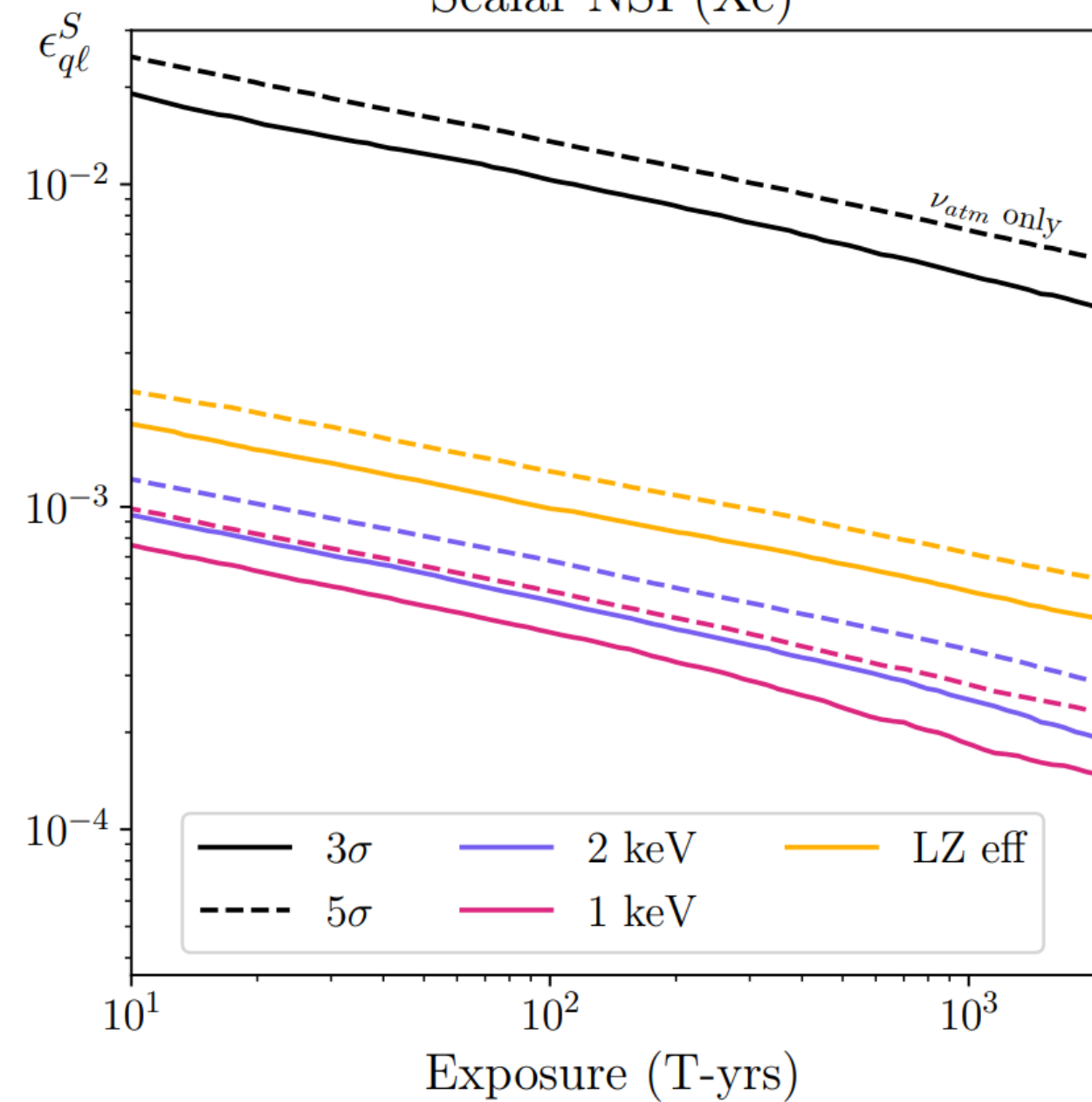
$$\mathcal{L} \supset G_F \sum_{q,\ell} \epsilon_{q\ell}^S \bar{\nu}_\ell \nu_\ell \bar{q} q \quad \epsilon_{q\ell}^S = \frac{g_{\nu_\ell S} g_{qS}}{M_S^2 G_F}$$

- High mass cross section

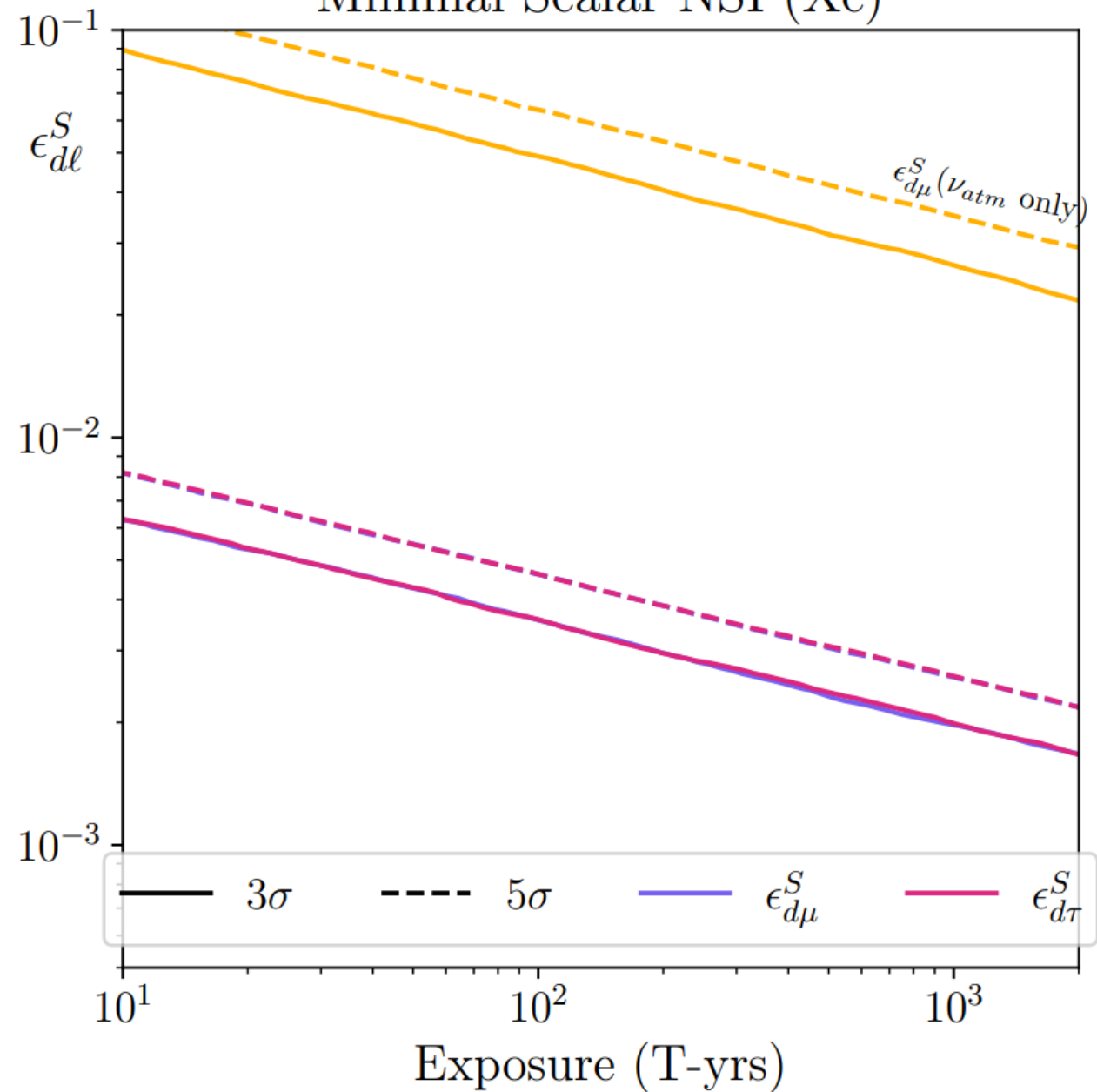
$$\left. \frac{d\sigma_N}{dE_R} \right|_S = \frac{F^2(E_R) q_s^2 E_R m_N^2}{4\pi E_\nu^2 (2E_R m_N + m_S^2)^2} \Longrightarrow \frac{G_F^2 \epsilon_{q\ell}^2 q_S^2 E_R m_N^2}{4\pi E_\nu^2}$$

$$q_S \approx 14.7A - 3.3Z$$

Scalar NSI (Xe)



Minimal Scalar NSI (Xe)



UV Motivation: Leptoquarks

- GUT models (unification of quarks and leptons)
- Possible solutions to:
 - Neutrino masses
 - $(g - 2)_\mu$
 - Flavor anomalies
- LHC constraints currently dominant (from pair production and Drell-Yan)

A Two Leptoquark Model

- Based on the Zee Model of neutrino mass generation
- $S_1 = (\bar{3}, 1, 1/3)$

$$\mathcal{L}_{S_1} \supset \lambda_L^{i\ell} S_1 \bar{Q}_i^c \epsilon L_\ell + \lambda_R^{i\ell} S_1^* \bar{u}_R^i \bar{\ell}_R + \tilde{\lambda}_R^{i\ell} S_1^* \bar{d}_R^i \bar{\nu}_R^\ell + h.c. - \frac{1}{2} m_{S_1}^2 S_1^2$$

- $\tilde{R}_2 = (\bar{3}, 2, 1/6)$

$$\mathcal{L}_{\tilde{R}_2} \supset \lambda^{i\ell} \bar{d}_R^i \tilde{R}_2 \epsilon L_\ell + \tilde{\lambda}^{i\ell} \bar{Q}_i^c \tilde{R}_2^\dagger \bar{\nu}_R^\ell + h.c. - \frac{1}{2} m_{\tilde{R}_2}^2 \tilde{R}_2^2$$

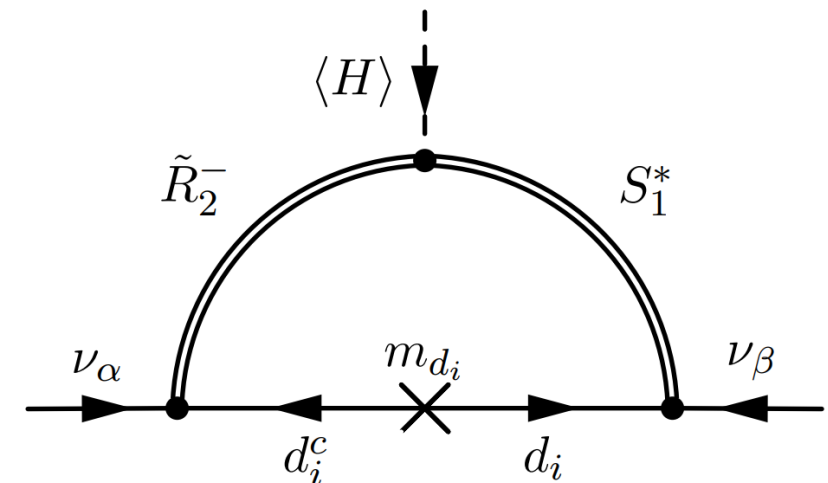
Neutrino Masses

- Expanding SU(2) doublets, dropping unnecessary terms, adding a higgs coupling

$$\mathcal{L} \supset \lambda^{i\ell} \left(\nu_\ell \bar{d}_i R^- - \ell \bar{d}_i R^+ \right) + \lambda_L^{i\ell} \left(\nu_\ell d_i - \ell \bar{u}_i^c \right) S_1 - \mu R^- H^0 S_1^* + h.c.$$

- R^- and S_1^* mix via Higgs $\rightarrow m_{1,2}$ in the new basis mixed by angle θ

$$M_\nu = \frac{3 \sin(2\theta)}{32\pi^2} \log \left(\frac{m_1^2}{m_2^2} \right) (\lambda M_d \lambda_L^T + \lambda_L M_d \lambda^T)$$



Leptoquarks in the IR

- Integrate out the LQ

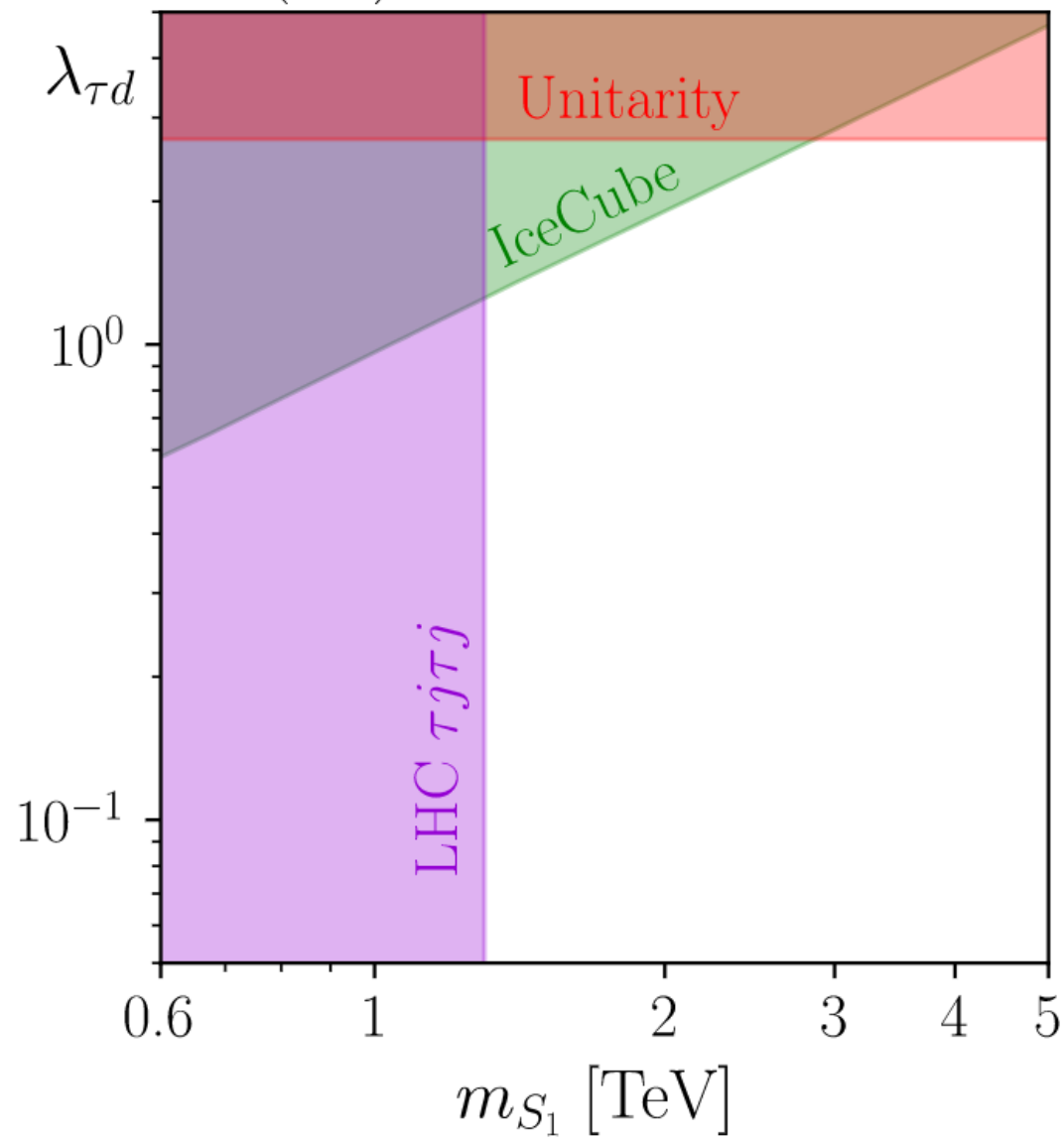
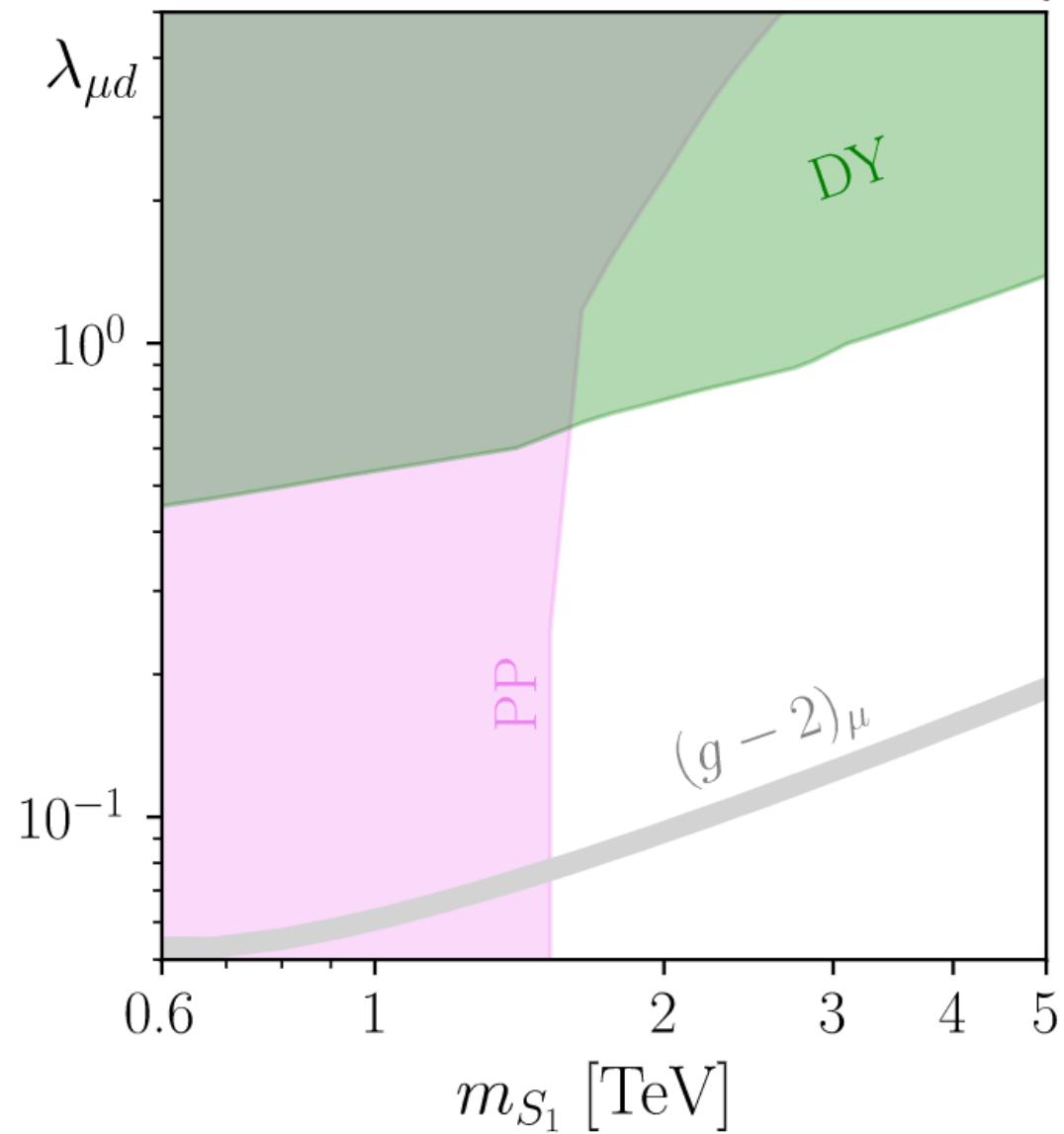
$$\mathcal{L} \supset \frac{\lambda_{d\ell}^L \lambda_{d\ell}^R}{2m_{S_1}^2} (\bar{d}_L^c \nu_{L\ell}) (\bar{d}_R^c \nu_{R\ell}) \implies \epsilon_{d\ell}^{S_1} = \lambda_{d\ell}^2 / 4G_F m_{S_1}^2$$

- New effective nuclear charge:

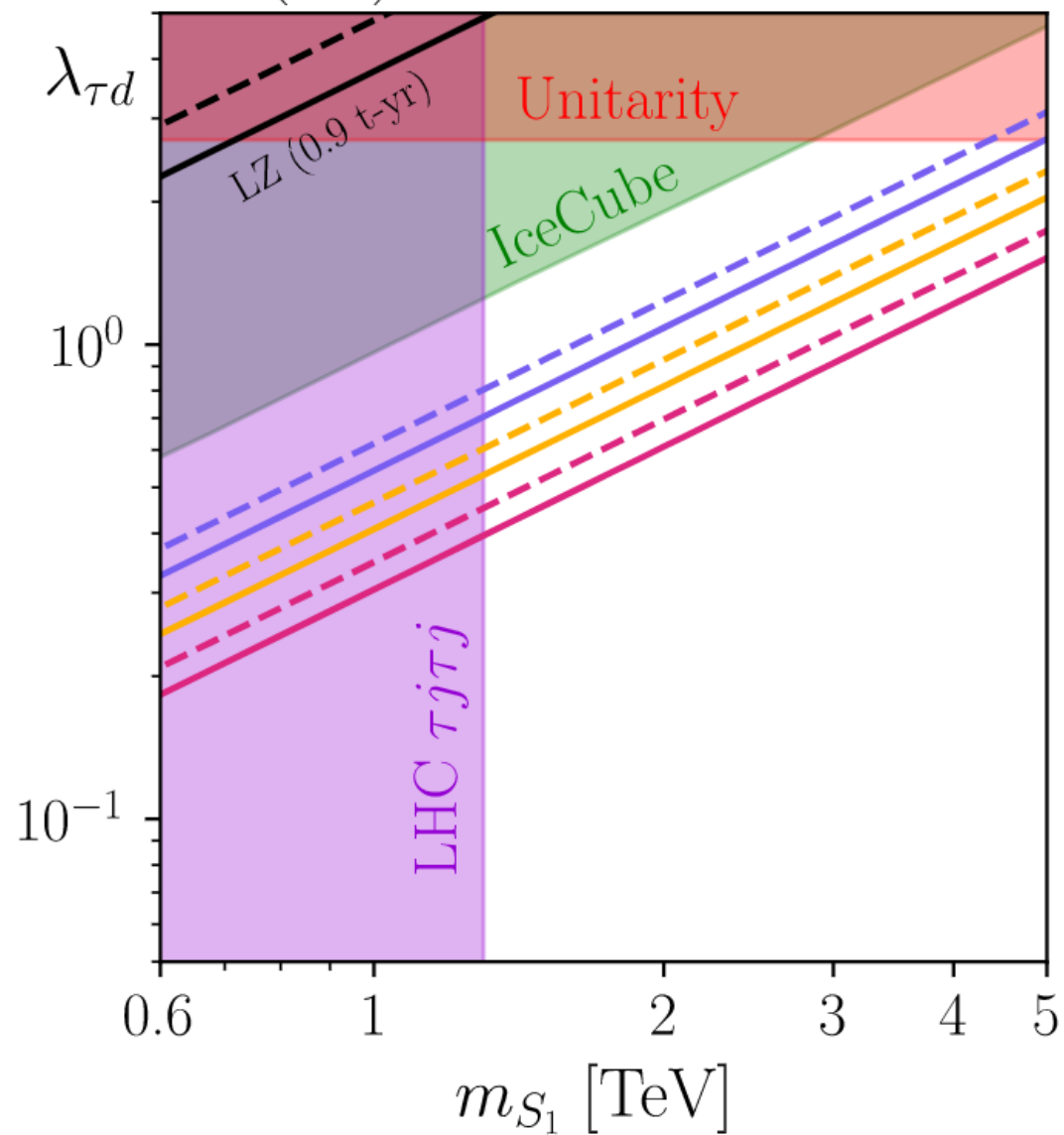
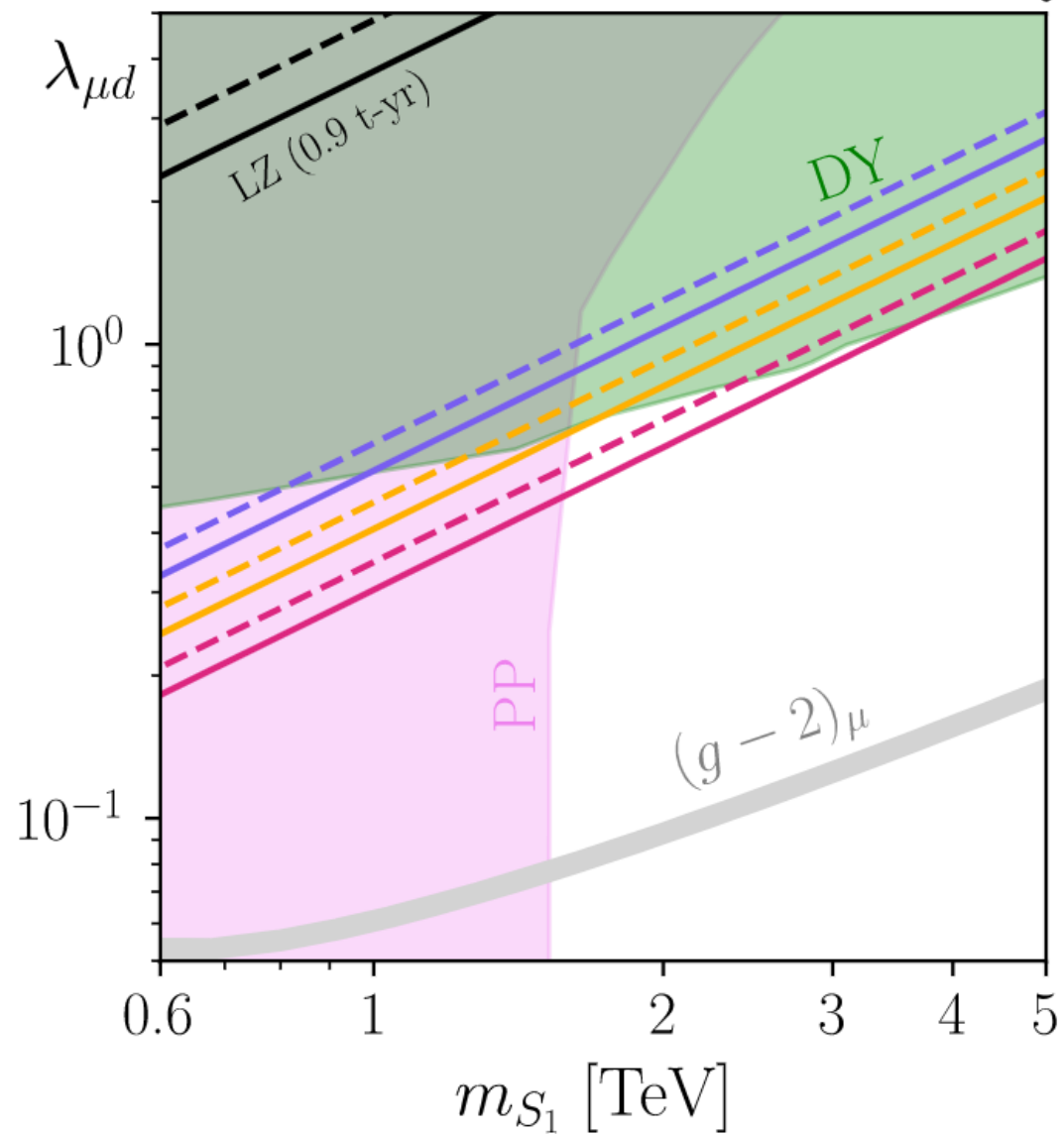
$$q_s = \frac{1}{m_d} (Z m_p f_d^p + N m_n f_d^n) \approx 3.74A - 0.75Z$$

- Careful which neutrinos we use

Scalar LQ Yukawa (Xe)



Scalar LQ Yukawa (Xe)



3σ
 5σ
 10 T-yrs
 100 T-yrs
 1000 T-yrs

Summary

- Non-Standard Neutrino interactions can produce measurable signals in Dark Matter experiments
- Such interaction can be mediated by heavy Leptoquarks
 - Future DM experiments can improve constraints on certain Leptoquarks
- Removing backgrounds and lowering thresholds is critical to realize these constraints

Thank You