

Baryogenesis in Mirror Twin Higgs

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Motivation

Hierarchy problem, no SUSY found at LHC.

Mirror Twin Higgs (MTH) evades LHC bounds through a color neutral top partner. A hidden sector is generated through \tilde{S}_2 symmetry.¹

Asymmetric reheating mechanism ensures that J_{eff} bounds are satisfied.²

Baryogenesis in MTH: \tilde{S}_2 symmetry implies twin baryogenesis.

Atomic Dark Matter (ADM) has interesting astrophysical implications.^{3,4}

Main question: what does SM baryogenesis imply about mirror sector baryogenesis, and hence atomic dark matter content of universe?

¹Ši ; P- <Vb>Oiq i Kbp>- ^@pi O-q^SWf| CEVg -qt Š= CECE| I v

²Ši ; P- <Vb>] i ; q; S>di Ti Gb†>- ^@pi O-q^SWf| CEug -qt Š= cvcciCE_ul

³, i KP- Y- sS- ^@[i [<k ~S^ f| CEDg -qt Š= cuc| iCEuu_

⁴Ti G- ^>, i V- z<>Xi p- ^@ Y- ^@[i pCC<C f| CE{g -qt Š= c{ CEicl | c

Mirror Twin Higgs Review

The hidden sector (denoted B) is a mirror copy of the SM sector (denoted A).

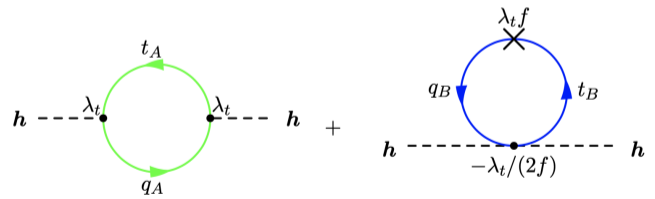
$$r \} (3), \quad r \} (2), \quad \} (1), \quad \tilde{S}_2 \quad r \} (3)_3 \quad r \} (2)_3 \quad \} (1)_3$$

\tilde{S}_2 symmetry ensures cancellation of quadratic divergences.

Twin vev higher than SM vev $H > f$, O_3 heavier than O_1 .

$3 < H/f < 10$ satisfies Higgs coupling constraints and avoids sub-percent tuning.

$J_{\text{eff}} \approx 6 \rightarrow$ need asymmetric reheating or fraternal twin Higgs (hard \tilde{S}_2 breaking to remove light degrees of freedom)⁵



$$G_{\text{SM}} \sim \mathcal{O}(1) \text{ [} \lambda_t \sim \lambda_t \text{]} \quad \text{[} \lambda_t \sim \lambda_t \text{]} \quad \text{[} \lambda_t \sim \lambda_t \text{]} \quad \text{[} \lambda_t \sim \lambda_t \text{]}$$

⁵] i ; q S >, i V - z > [i r z q s s V > - ^ @ p i r ~ ^ @ q \ f | C E l g - q t S = c l C E i C E { c C E

Jh>

AMi`Q/m+2b`B;?i@?LM/M/LM2mi`BMQb

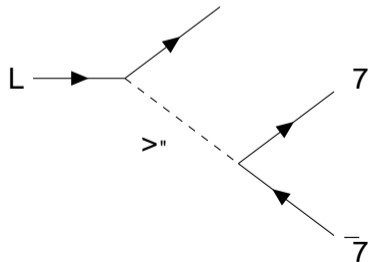
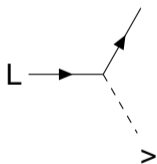
:2M2`i2b+iBp2M2kmi`B^{kk}BM^{kk}BM^{kk}#bi?b2+iQ`b

L_RBb HQM;@HBp2/Kbb2B;2Mbi i2 i? i7`22x2b Qmi`2

T2`BQ/Q7Kii2`/QKBMiBQM-/2+vb7i2`irQb2+iQ`b

KQ`2i?M"

"`M+?BM;`iBQBMiQ^{pp}7^{kk}BZ+iQH/2+vb`2i?`22@#Q/v



b v K K 2 i ` B + _ 2 ? 2 i B M ; U S ` 2 H B K B M ` v V

1 M 2 ` ; v / 2 M b B i v ? B b i Q ` v 7 Q ` Q M 2 @ ; 2 M 2 ` L 2 7 + Q M i Q m ` b 7 Q ` 7 f p 4 8 B M i ? 2 T ` M
i B Q M ` 2 ? 2 i B M ; X
Q 7 ` 2 ? 2 i Q M K b b M / ` 2 ? 2 i B M ; i 2 K

A M i 2 ` 2 b i B M ; : 2 o B Q M < R 9 y y 2 o , y R / m 2 i Q // B i B Q M H i r Q
i ? ` 2 2 @ # Q / v / 2 + v U M Q i / B b + m b b 2 / ? 2 ` 2 # m i b F K 2 7 Q ` /
K L > : 2 o b Q r B H H ? p 2 i Q ; 2 M 0 7 Q 2 i ? 2 H B ; ? i 2 b i + i B p 2 M 2 r

" `vQ;2M2bBb>BM

J BM B/2 , L /2+ vb iQ ;2M2` i2 # `vQM MmK#2` pBQH iB
qAJS # `vQ;2M2bBbVX

h?2 `2?2 iQM K v Q` K v MQi #2 i?2 ;2M2` iQMX

h?2 ;2M2` iQM ? b iQ #2 +QmTH2/ iQ [m `M B KM;2M +QH Q`

K;2M K > 8y.y2o7`QK G>* #QmM/ 7Q` +QH Q`2/ b+ H `X

JQ/2H

AMi`Q/m+2 7QHHRBM; BMi2` +iBQMb BM i?2 G ;` M;B M

$$L \quad v \approx GL \quad L \quad | \quad 0 \quad . \quad . \quad y$$

$$v \approx "GL" \quad "L" \quad "I" \quad 0 \quad . \quad . \quad y$$

q2 // T `iB+H2 + Q.Mi2 Mi

| JmHiBTaH(2)i | a(l)k | l(R | a(l)j) | a(l)k | l(R) |
|--------------|-------|------|--------|-------|------|
| j | R | @kfj | R | R | y |
| R | R | y | j | R | @kfj |
| L | R | y | R | R | y |
| L | R | y | R | R | y |

J bb 2B;2Mbi i2b M/ 2B;2Mp Hm2b

$$L = \rho \frac{R}{k} (L \quad L"); \quad J = K_L \quad K "$$

Cases we study

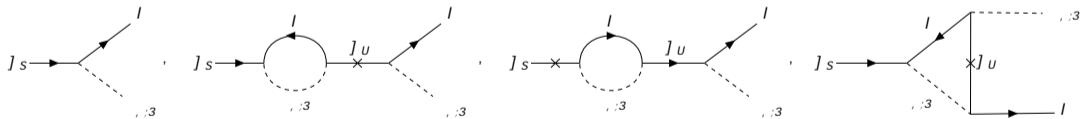
CASE 1: reheaton = generator = J_1

J_2 interferes, need $\lambda_{J_2} \leq \lambda_{J_1}$

If $\lambda_{J_2} > \lambda_{J_1}$, SM baryogenesis suppressed unless $\lambda_{J_1} \sim \lambda_{J_2}$ (further study ongoing for the resonant baryogenesis limit)

If $\lambda_{J_2} < \lambda_{J_1} < \lambda_{J_3}$, SM baryogenesis unsuppressed, twin baryogenesis suppressed

If $\lambda_{J_2} \sim \lambda_{J_3} < \lambda_{J_1}$, both unsuppressed.



Cases we study

CASE 2: reheaton \neq generaton

The generaton J_1 and the colored scalar ϕ need to be heavier than 500GeV, but now J_1 generates eV neutrino, so coupling $\frac{\rho_{\nu} \sqrt{J_1}}{f}$ is large enough for early out of equilibrium decay, leaving the reheating process undisturbed.

The reheaton J_2 can be arbitrary mass. For $\sqrt{J_2} > K C$, we get $\sqrt{J_2} > C$.

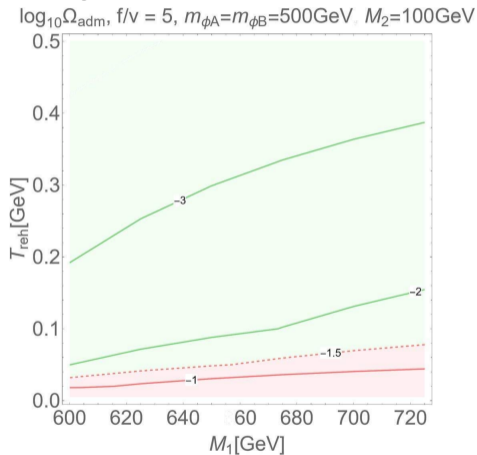
If $\sqrt{J_2} \sim 100 K C$, we can now have ~ 0.1 .

Further studies in progress for this case.

Preliminary Results

SM baryogenesis fixes coupling combinations.

We can predict the ADM density:



$\log_{10} \Omega_{\text{adm}} = ? [\text{eq} @ \text{S} \text{S} ^ \text{S} \text{z} \text{PC} < - \text{sC} \text{bH} \text{q} \text{PC} \text{z} \text{b} ^ \text{E} \text{LC} ^ \text{C} \text{q} \text{z} \text{b} ^ \text{i}$

