

# Muon Yukawa couplings at the high-energy muon collider

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Theory and Phenomenology  
of Fundamental Interactions  
UNIVERSITY AND INFN - BOLOGNA



University of  
**Pittsburgh**



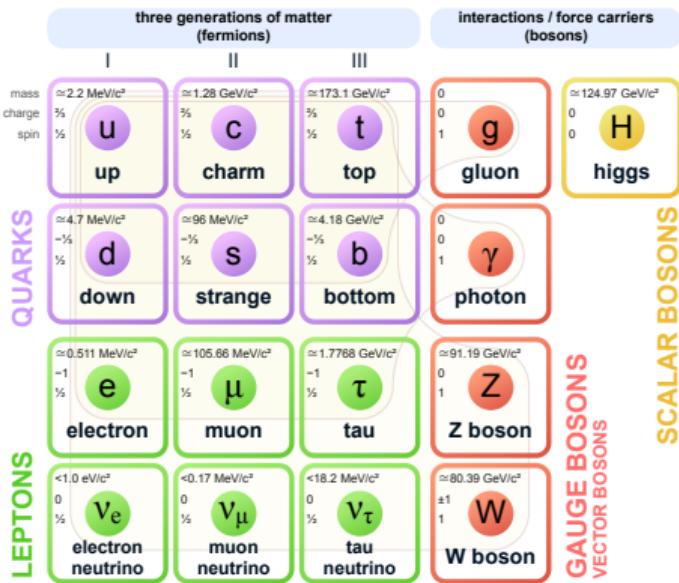
[T.Han, W.Kilian, N. Kreher, YM, J. Reuter, and K.Xie, JHEP 12 (2021) 162, 2108.05362 ]

[E. Celada, T.Han, W.Kilian, N. Kreher, YM, F. Maltoni, D. Pagani, J. Reuter, T. Striegl, and K.Xie, coming out soon ]

# Why Higgs?

A well understood and well tested model

## Standard Model of Elementary Particles

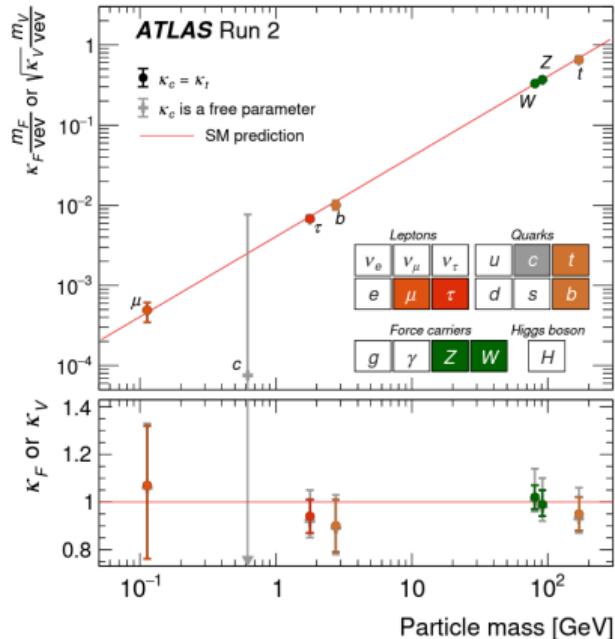


- Model doesn't make sense without Higgs or something like it
- The Higgs is a scalar particle whose interactions with other particles are predicted in terms of their masses
- It provides masses to all other elementary particles

Higgs physics: A portal to new physics

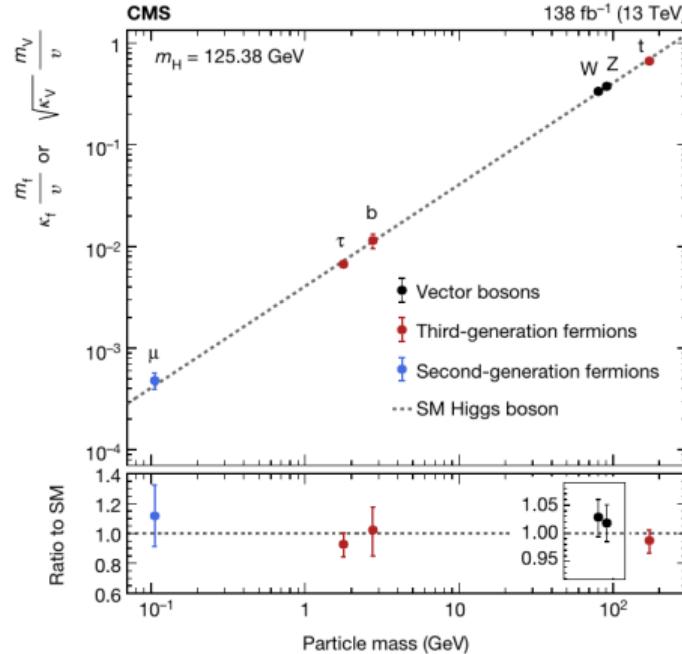
- LHC has gone **from discovery to precision**
- A telescope to high scale physics
- Interplay of theory and experiment is important

# Measure the Higgs couplings



[Nature 607 (2022) 52]

The next task is to complete the above plots



[Nature 607 (2022) 60]

# Why Yukawa couplings?

## Measuring the Yukawa couplings is important

- ▶ Directly test the SM Higgs mechanism that generates the masses.
- ▶ The Yukawa couplings for heavy fermions are well measured (to  $5\sigma$  level)

$$y_f = \sqrt{2}m_f^{\text{SM}}/v, \quad f = t, b, \tau$$

- ▶ The next target is the second generation fermions!

## The physics motivation:

- ▶ A naive question: We actually do not know whether the SM mass-generation mechanism applies just to the heavy particles, or also to the 1st/2nd generations.
- ▶ Logical possibility: The fermion mass is not (only) generated by SM Higgs.
  - ⇒ What if the possible BSM physics modifies the  $\mu\mu H$  coupling ( $y_\mu = \kappa_\mu y_\mu^{\text{SM}}$ )?
  - ⇒ Is there new vertex, i.e.  $\mu\mu H^n$  coupling?

## EFT parameterizations

- Nonlinear HEFT gives  $\kappa_\mu = \frac{v}{\sqrt{2}m_\mu} y_1$  [Coleman et al., PR1969, Weinberg, PLB1980, · · ·]

$$\begin{aligned}\mathcal{L}_{UH} = & \frac{v^2}{4} \text{Tr} \left[ D_\mu U^\dagger D^\mu U \right] F_U(H) + \frac{1}{2} \partial_\mu H \partial^\mu H - V(H) \\ & - \frac{v}{2\sqrt{2}} \left[ \bar{\ell}_L^i \tilde{Y}_\ell^{ij}(H) U (1 - \tau_3) \ell_R^j + \text{h.c.} \right]\end{aligned}$$

with  $F_U, V, \tilde{Y}$  expanded as

$$F_U(H) = 1 + \sum_{n \geq 1} f_{U,n} \left( \frac{H}{v} \right)^n, V(H) = v^4 \sum_{n \geq 2} f_{V,n} \left( \frac{H}{v} \right)^n, \tilde{Y}_\ell^{ij}(H) = \sum_{n \geq 0} \tilde{Y}_{\ell,n} \left( \frac{H}{v} \right)^n$$

- Linear SMEFT [Weinberg PRL1979, Abbott & Wise PRD1980, · · ·]

$$\mathcal{L} \supset - \sum_{n=1}^{\infty} \frac{c_\varphi^{(2n+4)}}{\Lambda^{2n}} \left( \varphi^\dagger \varphi - \frac{v^2}{2} \right)^{n+2} - \sum_{n=1}^{\infty} \frac{c_{\ell\varphi}^{(2n+4)}}{\Lambda^{2n}} \left( \varphi^\dagger \varphi - \frac{v^2}{2} \right)^n (\bar{\ell}_L \varphi e_R + \text{h.c.})$$

# HEFT in the unitary gauge: the extended $\kappa$ framework

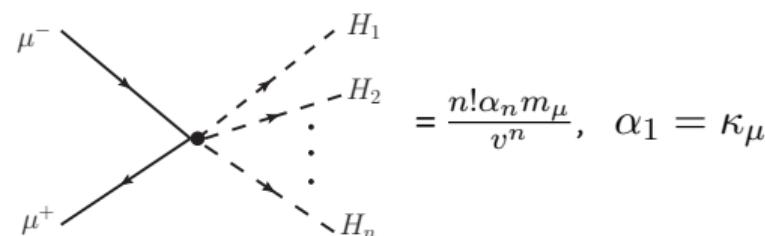
Introduce the form factors  $\alpha_n, \beta_n$

$$y_{\mu,n} = \frac{\sqrt{2}m_\mu}{v}\alpha_n, \quad f_{V,n} = \beta_n\lambda$$

In the unitary gauge, the HEFT formalism can be simplified to

$$\mathcal{L} \supset -\frac{m_H^2}{2}H^2 - m_\mu \bar{\mu}\mu - \sum_{n=3}^{\infty} \beta_n \frac{\lambda}{v^{n-4}} H^n - \sum_{n=1}^{\infty} \alpha_n \frac{m_\mu}{v^n} H^n \bar{\mu}\mu$$

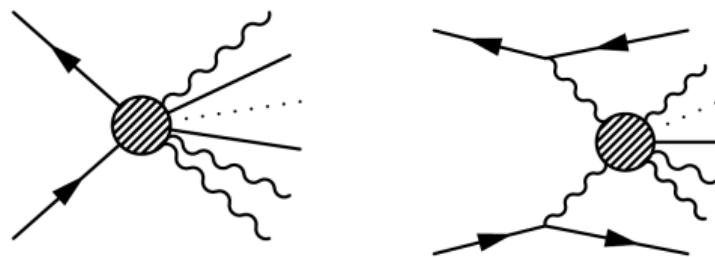
The regular “ $\kappa$  framework” is extended to include more vertices



# Physics at high-energy lepton colliders: electroweak Tevatron

## New phenomenology at a multi-TeV lepton collider:

1. Multi-boson production (annihilation)
2. ... and vector boson fusion (**VBF**) to multi-bosons,  
leading to multi-fermion final states with resonance structure.



[Barger, Cheung, Han, Phillips 1995] [Boos, He, Kilian, Pukhov, Yuan, Zerwas 1998]

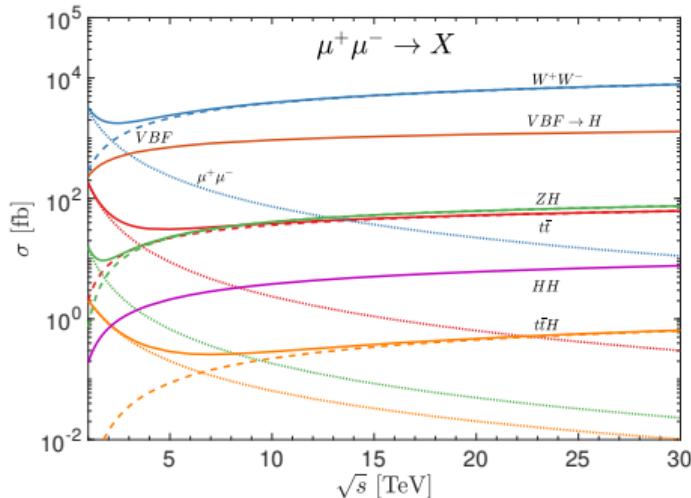
### Task:

Measure **all** interactions of multiple SM particles **exclusively** and with **precision**, from threshold to up to 2 orders of magnitude above EW scale.

# The full picture: Semi-inclusive processes

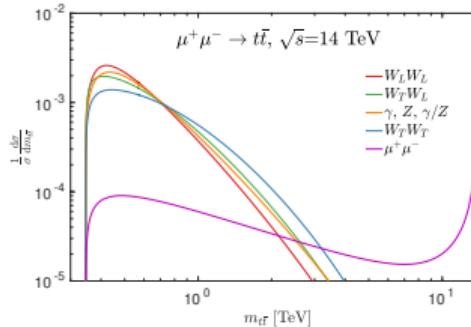
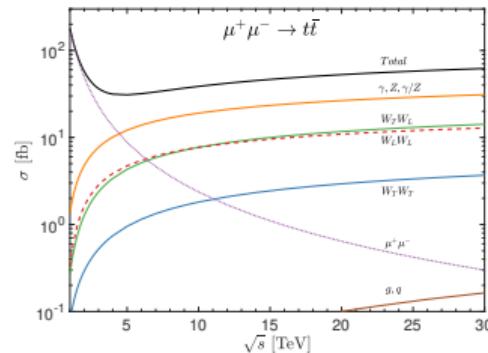
Just like in hadronic collisions:

$\mu^+ \mu^- \rightarrow$  exclusive particles + remnants



[T. Han, Y. Ma, K.Xie 2007.14300]

One example:  $\mu^+ \mu^- \rightarrow t\bar{t} + X$

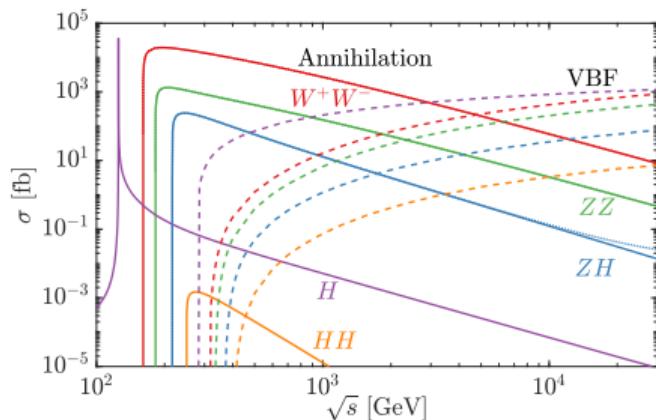


# Multi-boson production and the sensitivity to new physics

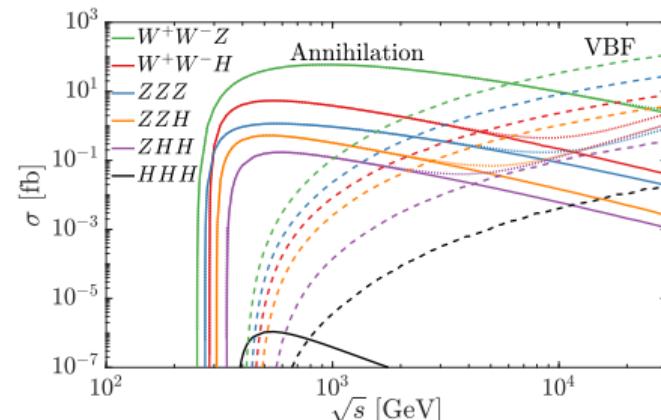
First glance:  $\alpha_1 = 1$  VS  $\alpha_1 = 0$

- **SM:**  $\lambda(\text{Muon} - \text{Higgs}) \sim y_\mu^{\text{SM}} = \sqrt{2}m_\mu^{\text{SM}}/v$
- **Possible BSM physics:**  $m_\mu = m_\mu^{\text{SM}}$ ,  $\lambda(\text{Muon} - \text{Higgs}) \sim \alpha_1 y_\mu^{\text{SM}}$ , e.g.  $\alpha_1 = 0$

## Two-boson final states

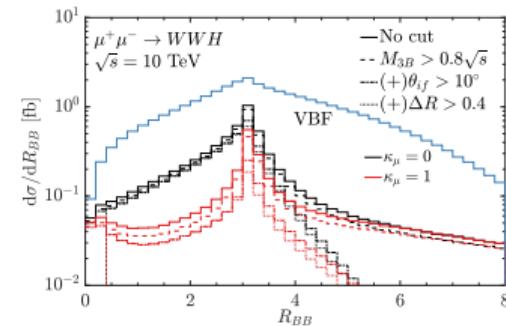
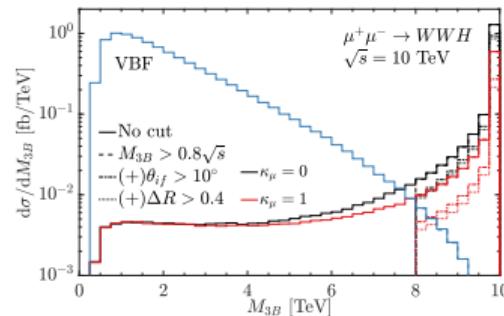
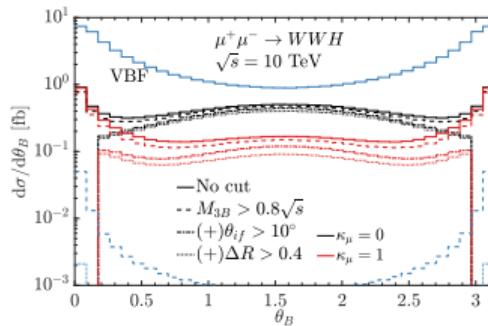


## Three-boson final states



New physics signal shows up in the high energy region

# WWH at a 10 TeV muon collider: Kinematics



- ▶ Background (VBF) is much larger than signal (annihilation)
- ▶ VBF events accumulate around threshold, and mostly forward
- ▶ Annihilation in the rest frame (central, and  $M \sim \sqrt{s}$  spread by ISR)
- ▶ Annihilation also has forward dominance, due to the gauge splitting  $W \rightarrow WH$

[T.Han, W.Kilian, N. Kreher, YM.J. Reuter, and K.Xie, JHEP 12 (2021) 162, 2108.05362]

# WWH at a 10 TeV muon collider: Cuts

Cut flow	$\kappa_\mu = 1$	w/o ISR	$\kappa_\mu = 0$ (2)	CVBF	NVBF
$\sigma$ [fb]	<i>WWH</i>				
No cut	0.24	0.21	0.47	2.3	7.2
$M_{3B} > 0.8\sqrt{s}$	0.20	0.21	0.42	$5.5 \cdot 10^{-3}$	$3.7 \cdot 10^{-2}$
$10^\circ < \theta_B < 170^\circ$	0.092	0.096	0.30	$2.5 \cdot 10^{-4}$	$2.7 \cdot 10^{-4}$
$\Delta R_{BB} > 0.4$	0.074	0.077	0.28	$2.1 \cdot 10^{-4}$	$2.4 \cdot 10^{-4}$
# of events	740	770	2800	2.1	2.4
$S/B$	2.8				

- ▶ Integrated luminosity  $\mathcal{L} = (\sqrt{s}/10 \text{ TeV})^2 \cdot 10 \text{ ab}^{-1}$  [[1901.06150](#)]
- ▶  $S = N_{\kappa_\mu} - N_{\kappa_\mu=1}$ ,  $B = N_{\kappa_\mu=1} + N_{\text{VBF}}$ .
- ▶ VBF and ISR are mostly excluded by invariant mass cut.
- ▶ Angular cut also weaken VBF further.

# Processes in consideration

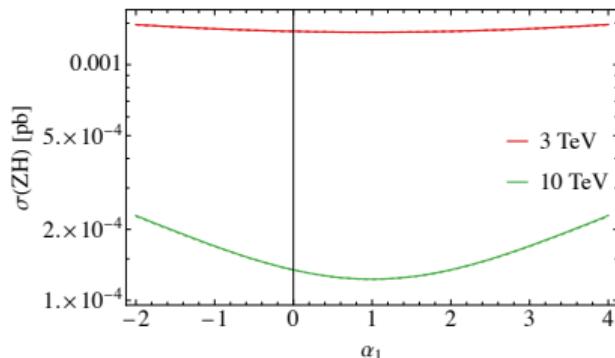
H \ V	0	1	2	3	4	5
0	-	Z	$Z^2, W^2$	$Z^3$ $W^2 Z$	$Z^4, W^4$ $W^2 Z^2$	$Z^5, W^2 Z^3$ $W^4 Z$
1	$H$	$ZH$	$W^2 H$ $Z^2 H$	$W^2 ZH$ $Z^3 H$	$W^4 H, Z^4 H$ $W^2 Z^2 H$	-
2	$H^2$	$ZH^2$	$W^2 H^2$ $Z^2 H^2$	$W^2 ZH^2$ $Z^3 H^2$	-	-
3	$H^3$	$ZH^3$	$W^2 H^3$ $Z^2 H^3$	-	-	-
4	$H^4$	$ZH^4$	-	-	-	-
5	$H^5$	-	-	-	-	-

[E. Celada, T.Han, W.Kilian, N. Kreher, YM, F. Maltoni, D. Pagani, J. Reuter, T. Striegl, and K.Xie, coming out soon ]

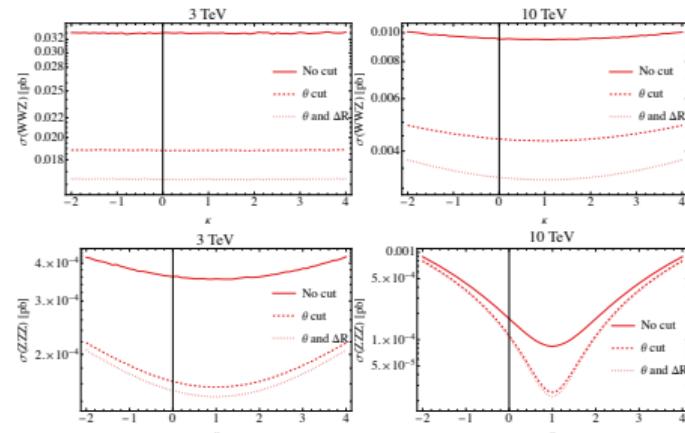
# Measure the $\mu\mu H$ vertex

There are processes that depend on only  $\alpha_1$ :  $\mu^+\mu^- \rightarrow ZH$  and  $\mu^+\mu^- \rightarrow V^3$

## ► $ZH$ production



## ► $V^3$ production



## ► Sign of Yukawa: $\alpha_1 = 1$ VS $\alpha_1 = -1$

$$\mathcal{S}_{3\text{ TeV}}^\pm = 1.23, \quad \mathcal{S}_{10\text{ TeV}}^\pm = 11.8$$

$$\mathcal{S}_{10\text{ TeV}} = 2.09, \quad |\Delta\alpha_1| \leq 0.8$$

## ► Measure $\alpha_1$ more precisely

$$\mathcal{S}_{10\text{ TeV}}^{WWZ} = 2.1, \quad |\Delta\alpha_1| \leq 0.8$$

$$\mathcal{S}_{10\text{ TeV}}^{ZZZ} = 2.2, \quad |\Delta\alpha_1| \leq 0.2.$$

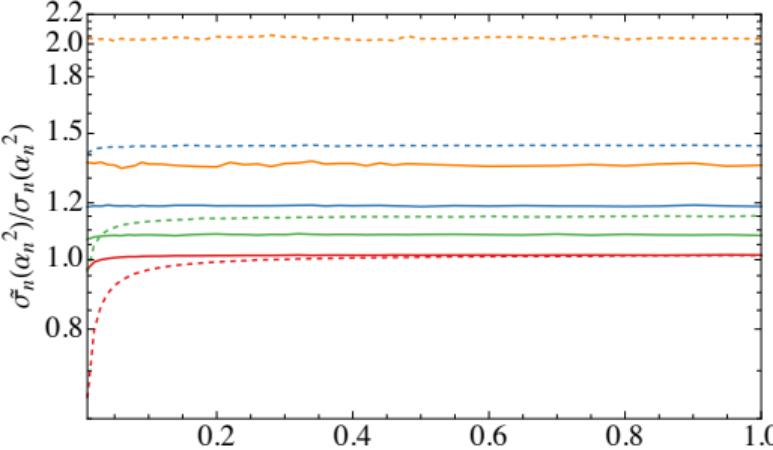
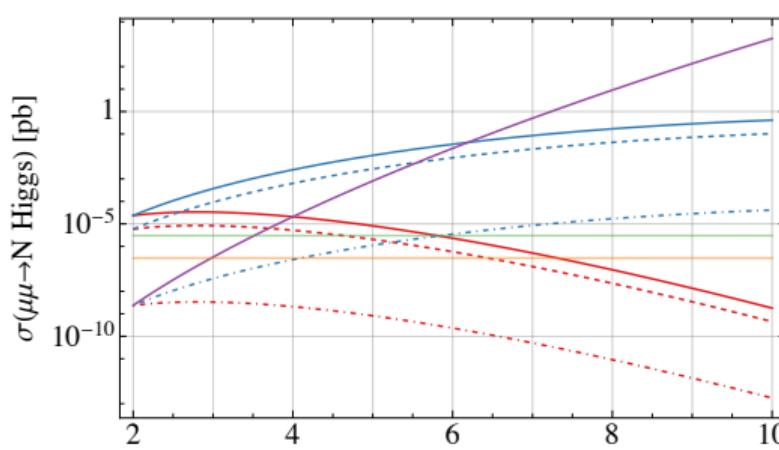
# Measure the $\mu\mu H^n$ vertices: Multi-Higgs production processes

Dominant contribution from the contact diagram:  $\mu\mu H^n$  vertex

$$\sigma_{\text{BSM}}(\mu^+ \mu^- \rightarrow H^n) = \sigma_{\text{SM}}^{(\text{loop})} + \sigma_n(\alpha_n^2) + \mathcal{O}(\alpha_m, \alpha_n), \quad m \leq n,$$

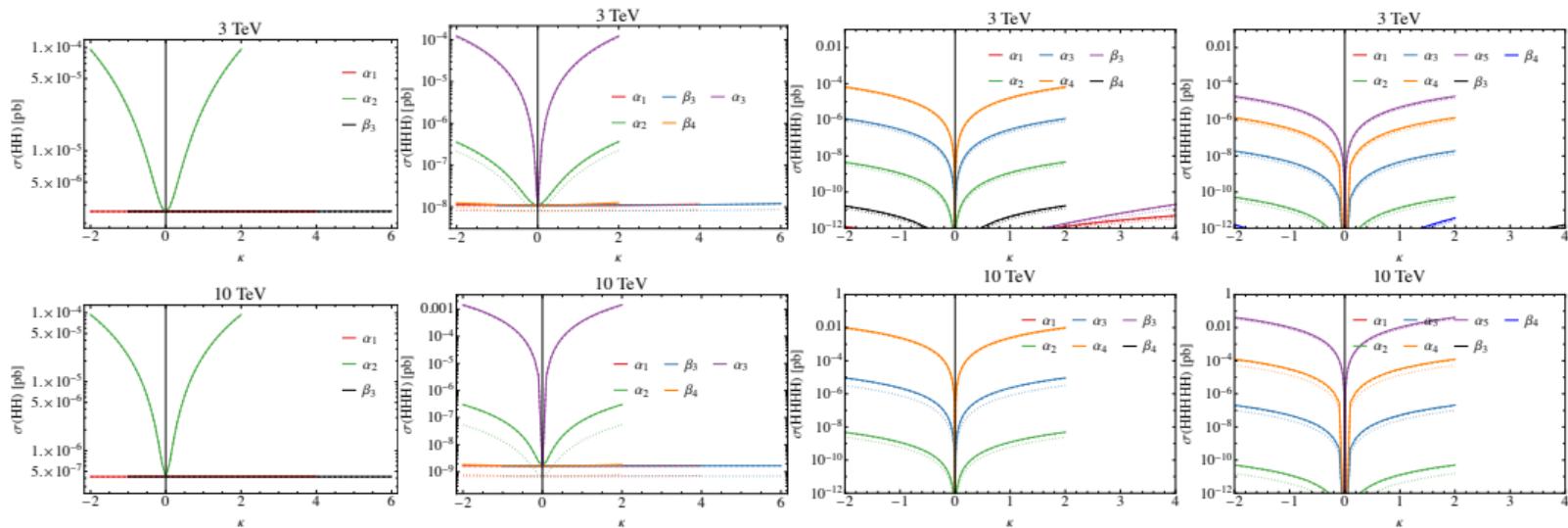
Approximate in the massless limit

$$\sigma_n(\alpha_n^2) \approx \tilde{\sigma}_n(\alpha_n^2) = \frac{n! m_\mu^2 s^{n-2} \alpha_n^2}{2^{4n-3} \pi^{2n-3} v^{2n} \Gamma(n) \Gamma(n-1)}.$$



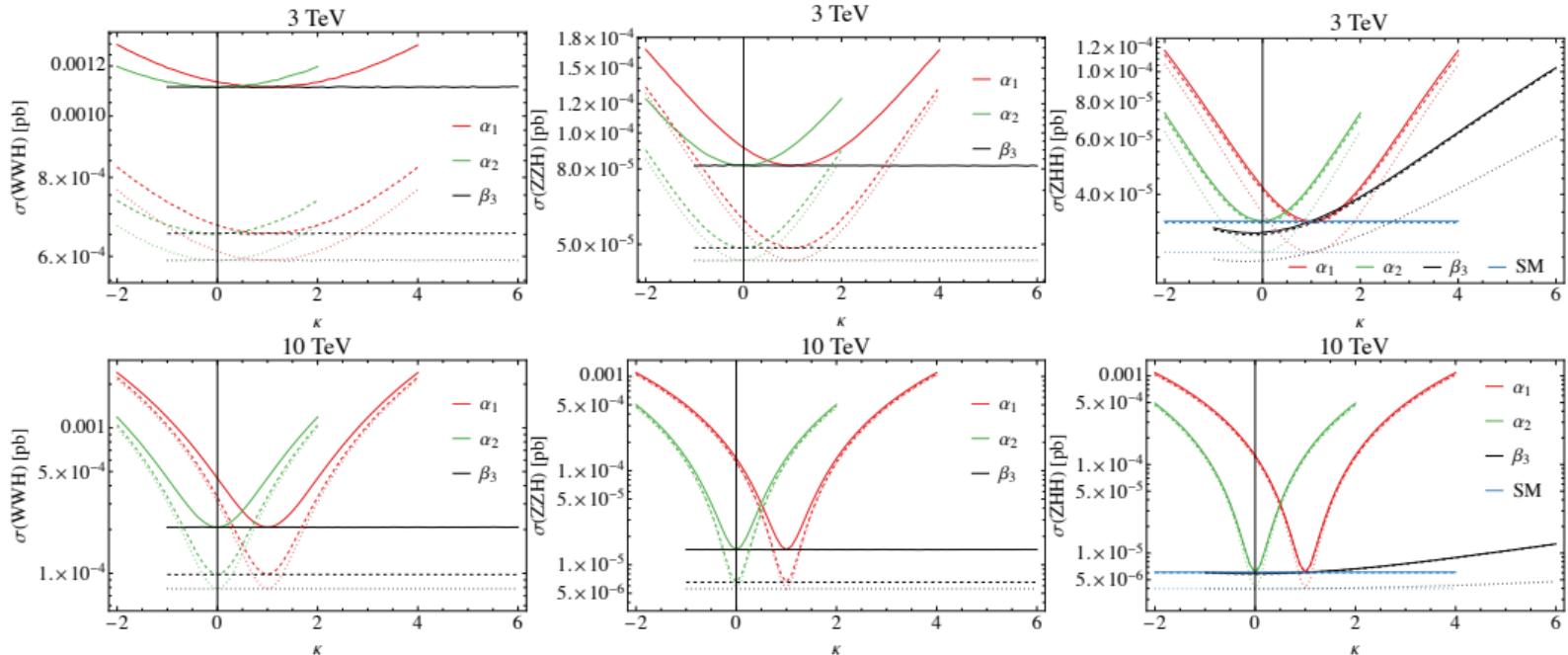
# Multi-Higgs production processes: $\mu^+ \mu^- \rightarrow H^n$

The cross section is solely dependent on  $\alpha_n \Rightarrow$  Measure  $\alpha_n$  directly

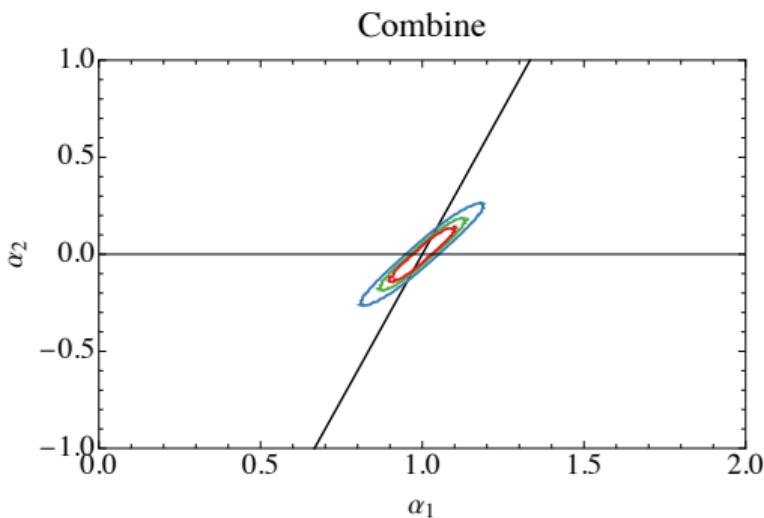
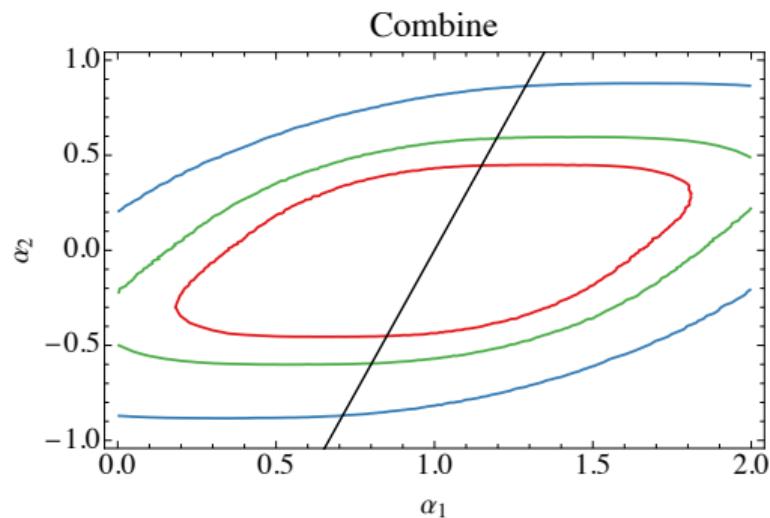


# Other processes: constrain $(\alpha_1, \alpha_2)$ simultaneously

Example:  $WWH, ZZH, ZHH$



# Constrains on $(\alpha_1, \alpha_2)$



# Summary and prospects

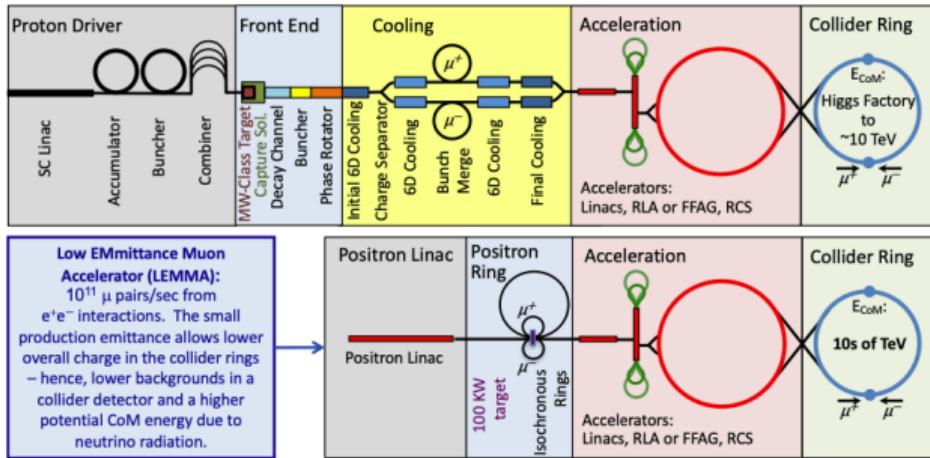
## Higgs is special and important

- ▶ The Higgs sector is the portal to new physics beyond SM.
  - ▶ Testing the SM mass generation mechanism indirectly helps BSM physics searches.
  - ▶ Measuring the vertices that do not exist in SM directly aim at BSM physics
  - ▶ The Yukawa couplings of the 3rd generation fermions are precisely measured  
⇒ The 2nd generation is the next target.

## Multi-TeV muon collider is cool

- ▶ A dream machine with many physics opportunities
- ▶ Two mechanisms with different kinematic features: lepton collision and VBF
- ▶ The main background is from VBF: Introduce kinematic cuts  $M_F > 80\% \sqrt{s}$
- ▶ The sign of  $\mu\mu H$  can be determined via  $\mu^+ \mu^- \rightarrow ZH$
- ▶ The  $\mu\mu H^n$  couplings can be determined via multi-boson production processes directly
- ▶ The  $\mu\mu H$  and  $\mu\mu H^2$  can be constrained simultaneously in the  $(\alpha_1, \alpha_2)$  contour plots

# Muon collider implementations



## Muon Accelerator Program [map.fnal.gov](http://map.fnal.gov)

[1901.06150,1907.08562]

- ▶ Protons → pions → muons
- ▶ 6D cooling is needed

## Low EMittance Muon Accelerator [web.infn.it/LEMMA](http://web.infn.it/LEMMA)

[1901.06150]

- ▶  $e^+e^- \rightarrow \mu^+\mu^-$ :  
45 GeV  $e^+$  to rest  $e^-$
- ▶ Cooling is not a problem
- ▶ High luminosity is challenging