# Smoking-Gun Signatures for Indirect Detection from Bound State Formation of Electroweak Multiplets

Giovanni Armando

Collaborators: M. Aghaie, S. Bottaro, A. Dondarini, D. Gaggero, P. Panci

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Giovanni Armando

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### Electroweak Multiplets as Dark Matter

- Minimal Dark Matter (MDM): first proposed in 2005
   [M. Cirelli, N. Fornengo, A. Strumia, 2005, arXiv:hep-ph/0512090]
- Consider a generic Electroweak (EW) multiplet:

$$\chi \equiv \mathbf{1}_{C}, \begin{pmatrix} \chi_{1} \\ \chi_{2} \\ \dots \\ \chi_{n} \end{pmatrix} \bigg\} SU(2)_{EW} \text{ and } Y$$

The neutral component  $\chi_0$  is the Dark Matter (DM) candidate

- Consider real representations with Y = 0 and odd n
- DM mass fixed by requiring correct relic abundance  $\rightarrow$  no free parameters in the theory: **FULLY PREDICTIVE**

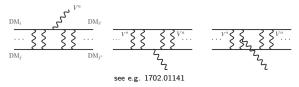
$$ightarrow$$
 For the fermion 5-plet,  $M_{\chi}=13.6$  TeV

## Non-Perturbative Effects and Phenomenology

• Sommerfeld Enhancement: In the non-relativistic limit, a long-range (attractive) potential between two DM particles distorts the two-body wavefunction, enhancing the annihilation cross section

$$\langle \sigma v \rangle \xrightarrow{\mathsf{SE}} R \times \langle \sigma v \rangle, \qquad R = \left| rac{\psi(\infty)}{\psi(0)} \right|^2$$

• The exchange of long range forces can lead to the Formation of Bound States through the emission of a gauge boson:  $\chi_i \ \chi_j \rightarrow BS \ V$ 

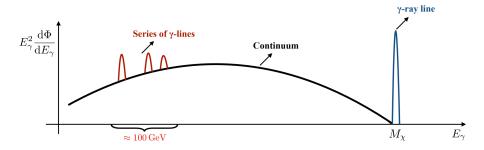


ightarrow Get an additional photon line with  $E_\gamma \sim 100$  GeV

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## Bound States Formation

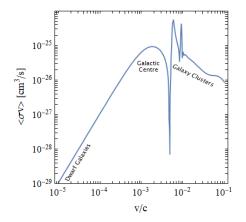
• Can get multiple lines for bound states with different binding energies



• Multiple line search can be extremely promising to discover/rule-out the candidate! But where to look?

#### Bound States Formation

 $\chi_0 \ \chi_0 \rightarrow B_{1s3} \ \gamma, M_{\chi} = 13.6 \text{ TeV}$ 



 $\bullet$  Larger relative velocity gives a more distinct signature  $\rightarrow$  galaxy clusters

# Summary & Outlook

- Electroweak Multiplets as WIMP candidates have no free parameters: extremely predictive!
- Phenomenology affected by non-perturbative effects: Sommerfeld Enhancement and the formation of Bound States
- Bound States give rise to distinctive features in the photon spectrum: smoking-gun signatures
- Consider available data e.g. Fermi-LAT to perform a correlated statistical analysis

Thank you for listening!

# Backup Slides

#### Electroweak Multiplets as Dark Matter

$$\mathcal{L}_{f} = \frac{1}{2}\bar{\chi}(i\not{D} - M_{\chi})\chi, \quad \chi = \chi^{c} \quad \text{(fermion)}$$

$$\mathcal{L}_{s} = \frac{1}{2}(D_{\mu}\chi)^{2} - \frac{1}{2}M_{\chi}^{2}\chi^{2} - \frac{\lambda_{H}}{2}\chi^{2}|H|^{2} - \frac{\lambda_{\chi}}{4}\chi^{4} \quad \text{(scalar)}$$

$$D_{\mu} = \partial_{\mu} + ig_{2}W_{\mu}^{a}T^{a}$$

- The DM mass is fixed by requiring that the WIMP make up the whole DM content of the Universe  $\rightarrow$  TeV scale DM
- DM stability is obtained accidentally: no renormalisable operators break  $\mathcal{Z}_2$  symmetry

## DM Stability

- In an Effective Field Theory (EFT), can have operators of arbitrarily high dimensions suppressed by powers of a cut-off energy,  $\Lambda_{UV}$
- For n=3, have fast DM decay at the renormalisable level

$$\mathcal{L} \supset \lambda_3 \chi HL$$

• For n > 3, have an accidental  $\mathcal{Z}_2$  symmetry explicitly broken by operators with  $d \ge 5$ 

$$\mathcal{L}_{EFT} \supset \frac{\mathcal{C}_{1}}{\Lambda_{UV}^{n-3}} (\chi HL) (H^{\dagger}H)^{\frac{n-3}{2}} + ... + \frac{\mathcal{C}_{3\chi}}{\Lambda_{UV}^{3}} \chi^{3} HL$$

A sufficiently large  $\Lambda_{UV}$  makes sure these operators are sufficiently suppressed  $\rightarrow$  DM is stable

• These non-perturbative effects are very important both for the computation of the DM mass...

DM spin	EW n-plet	$M_{\chi}$ (TeV)	$(\sigma v)_{\rm tot}^{J=0}/(\sigma v)_{\rm max}^{J=0}$	$\Lambda_{\rm Landau}/M_{\rm DM}$	$\Lambda_{\rm UV}/M_{\rm DM}$
Real scalar	3	$2.53\pm0.01$	_	$2.4 \times 10^{37}$	$4 \times 10^{24}$
	5	$15.4\pm0.7$	0.002	$7  imes 10^{36}$	$3  imes 10^{24}$
	7	$54.2\pm3.1$	0.022	$7.8  imes 10^{16}$	$2 \times 10^{24}$
	9	$117.8 \pm 15.4$	0.088	$3 \times 10^4$	$2 \times 10^{24}$
	11	$199\pm42$	0.25	62	$1 \times 10^{24}$
	13	$338 \pm 102$	0.6	7.2	$2 \times 10^{24}$
Majorana fermion	3	$2.86\pm0.01$	-	$2.4 \times 10^{37}$	$2 \times 10^{12*}$
	5	$13.6\pm0.8$	0.003	$5.5  imes 10^{17}$	$3 \times 10^{12}$
	7	$48.8\pm3.3$	0.019	$1.2 \times 10^4$	$1 \times 10^8$
	9	$113 \pm 15$	0.07	41	$1 \times 10^8$
	11	$202\pm43$	0.2	6	$1 \times 10^8$
	13	$324.6\pm94$	0.5	2.6	$1 \times 10^8$

- ...and for the observational tests:
  - $\rightarrow$  Collider Searches out of reach
  - $\rightarrow$  Direct Detection only at loop-level:  $\sigma_{\rm DD} \lesssim 10^{-45} {\rm cm}^2$
  - $\rightarrow$  Turn to Indirect Detection

## Indirect Detection

#### • Procedure:

- 1) Calculate DM annihilation cross section
- 2) Calculate flux of stable SM particles: photons and neutrinos
- 3) Compare with the observed flux
- For each DM annihilation channel, the final state (primary channel) can then decay into other SM particles (secondary channel)

$$\frac{d\Phi_{\gamma}}{d\Omega \ dE} = \underbrace{\frac{1}{2} \frac{r_{\odot}}{4\pi} \left(\frac{\rho_{\odot}}{M_{\chi}}\right)^2 \int_{\text{l.o.s.}} \frac{ds}{r_{\odot}} \left(\frac{\rho_{\text{DM}}(s)}{\rho_{\odot}}\right)^2}_{\text{Astrophysics}} \underbrace{\sum_{f} \langle \sigma v \rangle_f \frac{dN_{\gamma}^f}{dE}}_{\text{Particle Physics}}$$

#### Expected Photon Flux

