

SUPPORTED BY:

NKFIH K-133046, ÚNKP-21-3 & ÚNKP-22-3 New National Excellence Program of the Ministry for Innovation and Technology from the source of the National Research, Development and Innovation Fund

Odderon and Pomeron from the ReBB model

based on

[Eur. Phys. J. C **82**, 827 \(19 September 2022\)](#) & [Eur. Phys. J. C **81**, 611 \(13 July 2021\)](#)

and other recent results

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Day of Femtoscopy 2022
15 November 2022, Gyöngyös, Hungary

Bialas-Bzdak p=(q,d) model

$$\tilde{\sigma}_{in}(\vec{b}) = \int_{-\infty}^{+\infty} \dots \int_{-\infty}^{+\infty} d^2\vec{s}_q d^2\vec{s}'_q d^2\vec{s}_d d^2\vec{s}'_d D(\vec{s}_q, \vec{s}_d) D(\vec{s}'_q, \vec{s}'_d) \sigma(\vec{s}_q, \vec{s}_d; \vec{s}'_q, \vec{s}'_d; \vec{b})$$

- quark-diquark distribution inside the proton:

$$D(\vec{s}_q, \vec{s}_d) = \frac{1 + \lambda^2}{R_{qd}^2 \pi} e^{-\frac{s_q^2 + s_d^2}{R_{qd}^2}} \delta^2(\vec{s}_q + \lambda \vec{s}_d)$$

$$\lambda = \frac{m_q}{m_d}$$

$$\vec{s}_d = -\lambda \vec{s}_q$$

$$\vec{s}'_d = -\lambda \vec{s}'_q$$

[A. Bialas, A. Bzdak, Acta Phys.Polon. B 38, 159-168 \(2007\), Phys.Lett.B 649: 263-268 \(2007\)](#)

- sum of inelastic interaction probabilities of the constituents:

$$\sigma(\vec{s}_q, \vec{s}_d; \vec{s}'_q, \vec{s}'_d; \vec{b}) = 1 - \prod_a \prod_b [1 - \sigma_{ab}(\vec{b} + \vec{s}'_a - \vec{s}_b)]$$

$$\sigma_{ab}(\vec{s}) = A_{ab} e^{-|\vec{s}|^2 / S_{ab}^2}$$

$$S_{ab}^2 = R_a^2 + R_b^2$$

$$a, b \in \{q, d\}$$

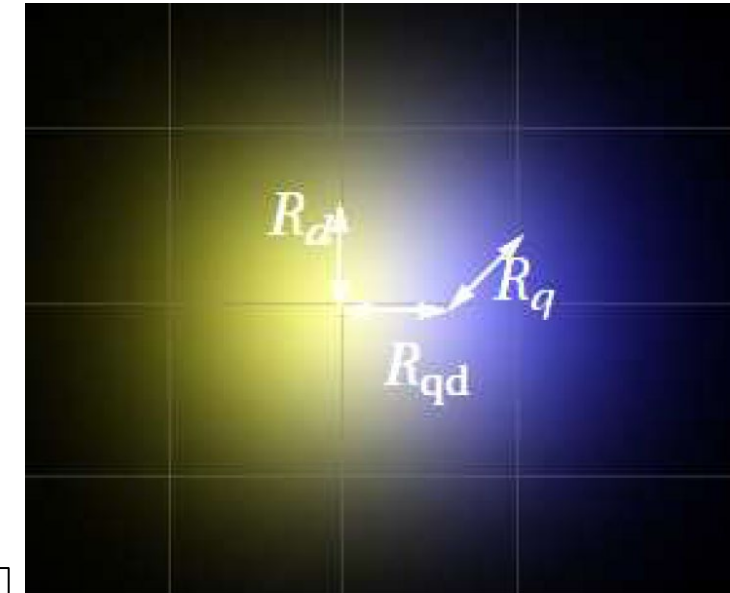
- inelastic cross-sections of quark, diquark scatterings :

$$\sigma_{ab,in} = \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \sigma_{ab}(\vec{s}) d^2\vec{s}$$

$$\sigma_{qq,in} : \sigma_{qd,in} : \sigma_{dd,in} = 1 : 2 : 4$$

- free parameters:

$$A_{qq}, \lambda, R_q, R_d, R_{qd}, \quad (A_{qq} = 1 \text{ and } \lambda = 0.5 \text{ can be fixed})$$



Proton-(anti)proton scattering in the quark-diquark model (Glauber style calculation).

Unitarily Real Extended Bialas-Bzdak (REBB) model

- elastic scattering amplitude in the impact parameter space:

$$t_{el}(s, \vec{b}) = i \left[1 - e^{-\Omega(s, \vec{b})} \right]$$

arXiv:1505.01415

F. Nemes, T. Csörgő, M. Csanád, Int. J. Mod. Phys. A Vol. 30 (2015) 1550076

- the opacity function:

$$\Omega(s, \vec{b}) = \text{Re}\Omega(s, \vec{b}) + i \text{Im}\Omega(s, \vec{b})$$

$\text{Im}\Omega \neq 0$ as the real part of the amplitude is not negligibly small

$$\text{Re}\Omega(s, \vec{b}) = -\frac{1}{2} \ln[1 - \tilde{\sigma}_{in}(s, \vec{b})]$$

$$\text{Im}\Omega(s, \vec{b}) = -\alpha \tilde{\sigma}_{in}(s, \vec{b})$$

NEW FREE PARAMETER,
has different values for pp and $\bar{p}p$

- elastic scattering amplitude in momentum space:

$$T(s, t) = 2\pi \int_0^\infty t_{el}(s, |\vec{b}|) J_0(|\vec{\Delta}| |\vec{b}|) |\vec{b}| d|\vec{b}|$$

$$|\vec{\Delta}| \equiv \sqrt{-t} \text{ as } \sqrt{s} \rightarrow \infty$$

(t is the squared momentum transfer)

ReBB model analysis of pp and p \bar{p} data

→ fits for pp d σ /dt data at 2.76 TeV and 7 TeV and for p \bar{p} d σ /dt data at 0.546 TeV and 1.96 TeV

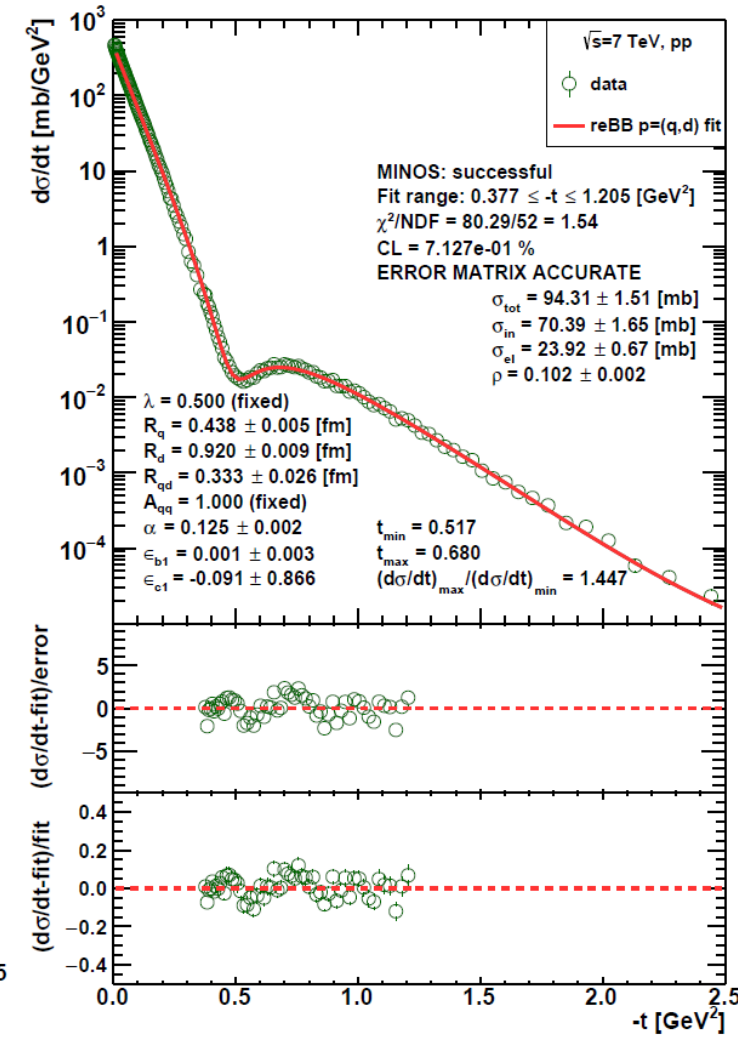
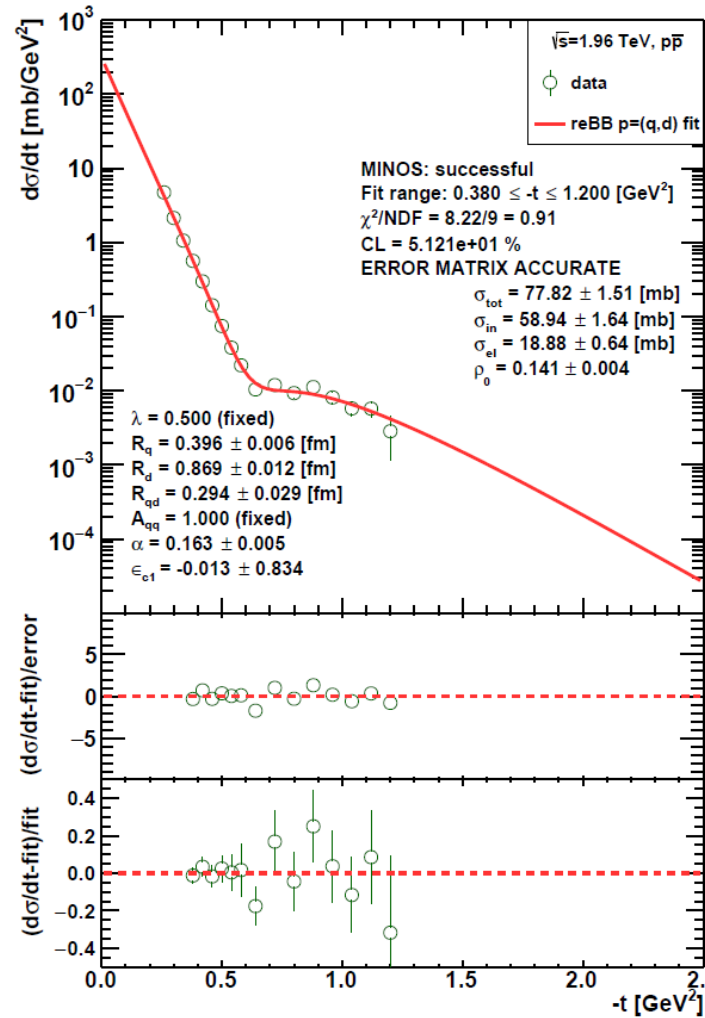
→ use of the χ^2 definition developed by PHENIX

→ determination of the energy dependences of the model parameters

→ satisfactory description in the kinematical range: $0.546 \leq \sqrt{s} \leq 8$ TeV & $0.37 \leq -t \leq 1.2$ GeV²

I. Szanyi, T. Csörgő, *Eur. Phys. J. C* 82, 827 (2022)

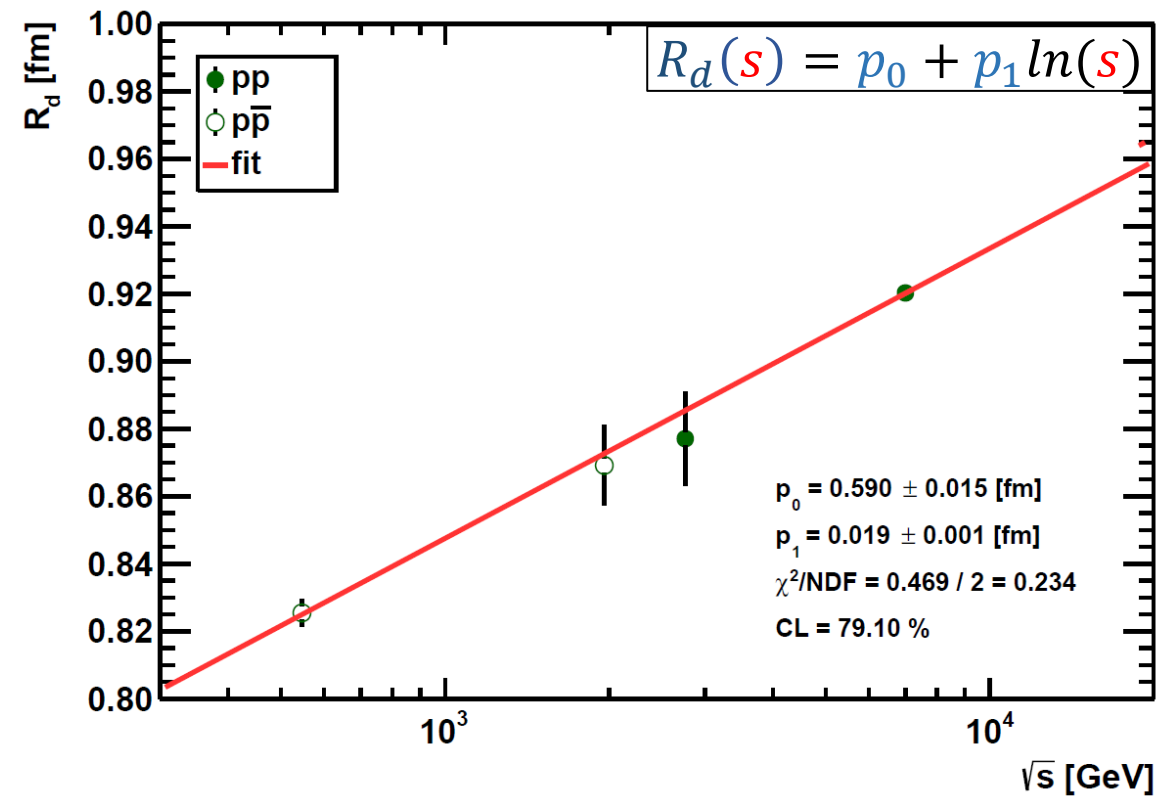
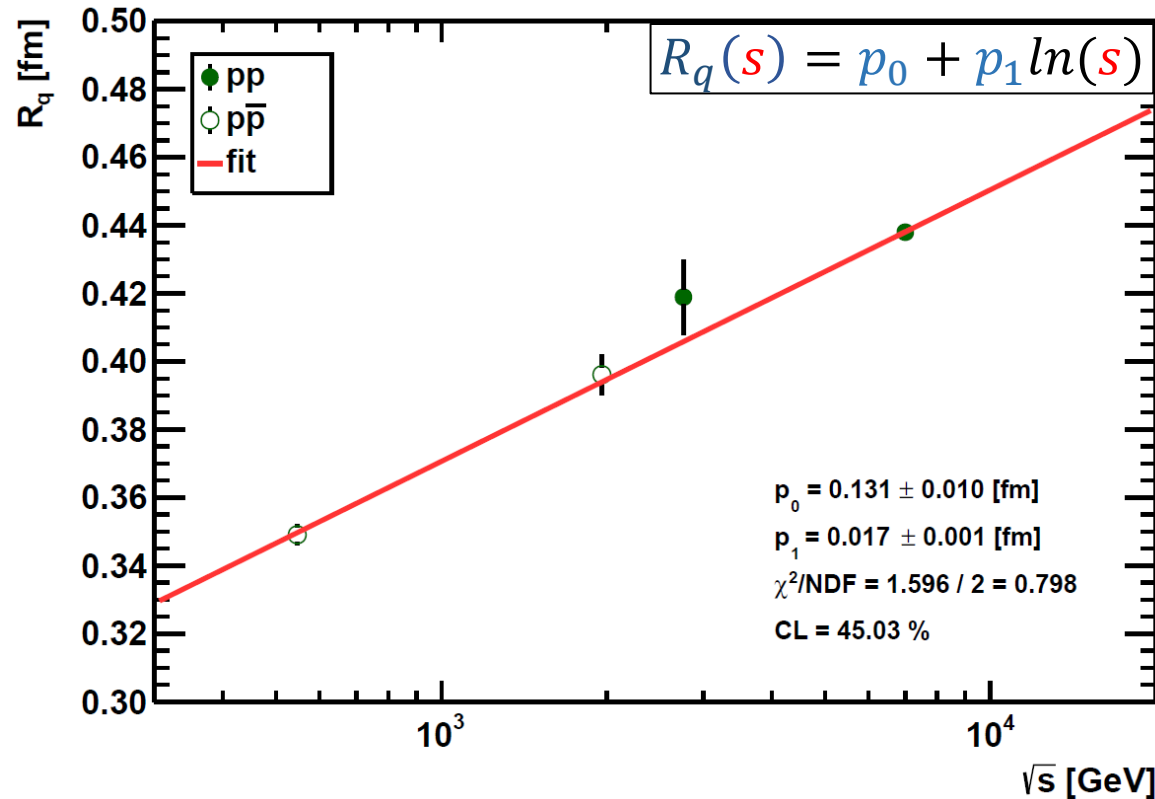
T. Csörgő, I. Szanyi, *Eur. Phys. J. C* 81, 611 (2021)



Examples of ReBB model fits for pp and p \bar{p} differential cross section data.

Energy dependences of the ReBB model parameters

T. Csörgő, I. Szanyi, *Eur. Phys. J. C* 81, 611 (2021)

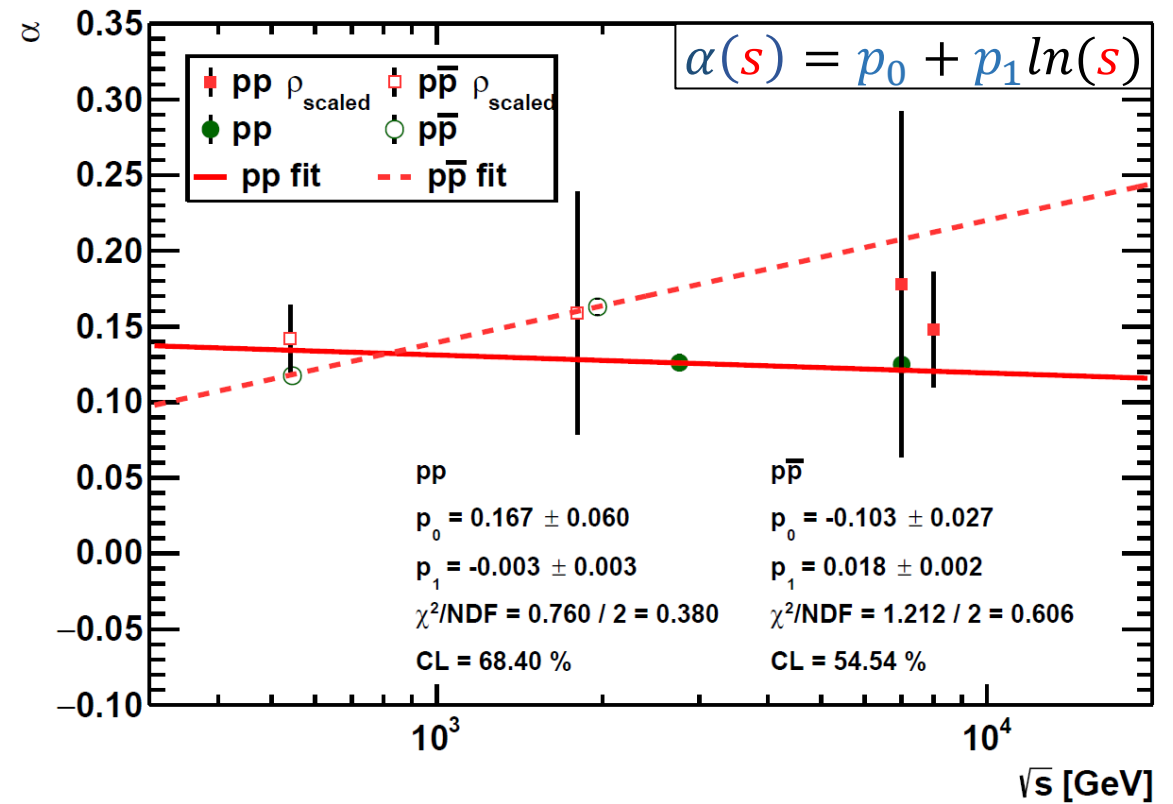
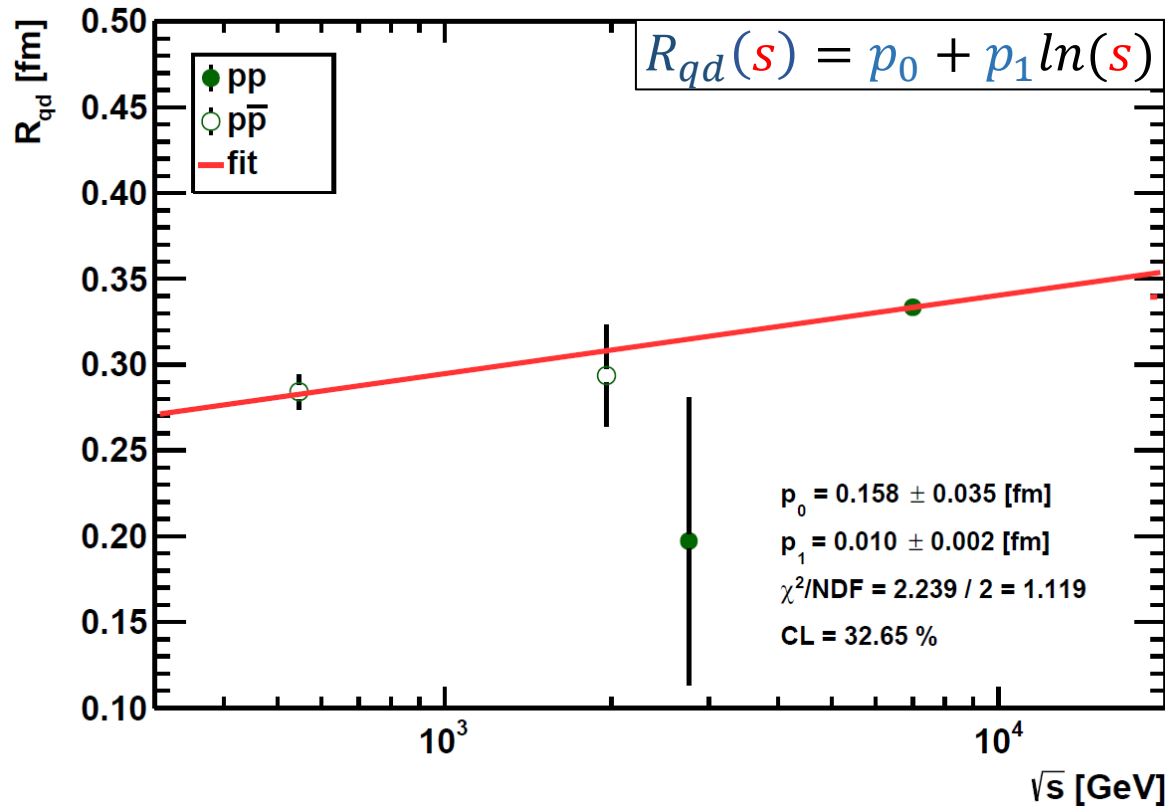


The energy dependences of the scale parameters, $R_q(s)$, $R_d(s)$, and $R_{qd}(s)$ are **linear logarithmic** and the **same** for pp and $p\bar{p}$ processes!

The energy dependence of the α parameter, $\alpha(s)$ is **linear logarithmic** too, but **not** the same for pp and $p\bar{p}$ processes!

Energy dependences of the ReBB model parameters

T. Csörgő, I. Szanyi, *Eur. Phys. J. C* 81, 611 (2021)



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Pomeron & odderon from pp & $p\bar{p}$ elastic scattering amplitudes

- according to the Regge formalism the strong scattering amplitude for pp and $p\bar{p}$ scattering is written in terms of $C = +1$ and $C = -1$ exchange components

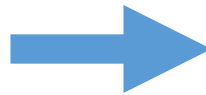
$$T^{pp}(s, t) = T^+(s, t) - T^-(s, t)$$

$$T^{p\bar{p}}(s, t) = T^+(s, t) + T^-(s, t)$$

- for $\sqrt{s} \gtrsim 1$ TeV the mesonic reggeon exchanges are negligible and essentially only the gluonic Pomeron and Odderon exchanges are present implying that

$$T^+(s, t) \equiv T^P(s, t)$$

$$T^-(s, t) \equiv T^O(s, t)$$



$$T^P(s, t) = \frac{1}{2} (T^{pp}(s, t) + T^{p\bar{p}}(s, t))$$

$$T^O(s, t) = \frac{1}{2} (T^{p\bar{p}}(s, t) - T^{pp}(s, t))$$

- In the ReBB model the pp and $p\bar{p}$ scattering amplitudes are defined and fitted to the data; the odderon and pomeron amplitudes are expressed in terms of those

Pomeron & odderon from the ReBB model

- In the ReBB model the pp and $p\bar{p}$ scattering amplitudes depend on energy through four energy dependent parameters:

$$T_{el}^{pp}(s, t) = F(R_q^{pp}(s), R_d^{pp}(s), R_{qd}^{pp}(s), \alpha^{pp}(s); t)$$

$$T_{el}^{p\bar{p}}(s, t) = F(R_q^{p\bar{p}}(s), R_d^{p\bar{p}}(s), R_{qd}^{p\bar{p}}(s), \alpha^{p\bar{p}}(s); t)$$

- It is found that:

$$\begin{aligned} R_q(s) &\equiv R_q^{pp}(s) = R_q^{p\bar{p}}(s) \\ R_d(s) &\equiv R_d^{pp}(s) = R_d^{p\bar{p}}(s) \\ R_{qd}(s) &\equiv R_{qd}^{pp}(s) = R_{qd}^{p\bar{p}}(s) \\ \alpha^{pp}(s) &\neq \alpha^{p\bar{p}}(s) \end{aligned}$$

thus

$$T_{el}^{\mathbb{P}}(s, t) = G(R_q^{pp}(s), R_d^{pp}(s), R_{qd}^{pp}(s), \alpha^{pp}(s), \alpha^{p\bar{p}}(s); t)$$

$$T_{el}^{\mathbb{O}}(s, t) = H(R_q^{pp}(s), R_d^{pp}(s), R_{qd}^{pp}(s), \alpha^{pp}(s), \alpha^{p\bar{p}}(s); t)$$

Pomeron & odderon from the ReBB model in b space

- the ReBB model b-dependent pp and $p\bar{p}$ scattering amplitudes are:

$$t_{el}^{pp}(s, b) = i \left(1 - e^{i \alpha^{pp}(s) \tilde{\sigma}_{in}(s, b)} \sqrt{1 - \tilde{\sigma}_{in}(s, b)} \right)$$

$$t_{el}^{p\bar{p}}(s, b) = i \left(1 - e^{i \alpha^{p\bar{p}}(s) \tilde{\sigma}_{in}(s, b)} \sqrt{1 - \tilde{\sigma}_{in}(s, b)} \right)$$

- the C-even part is:

$$\text{Re} t_{el}^{\mathbb{P}}(s, b) = \sqrt{1 - \tilde{\sigma}_{in}} \sin \left(\frac{\alpha^{pp} + \alpha^{p\bar{p}}}{2} \tilde{\sigma}_{in} \right) \cos \left(\frac{\alpha^{p\bar{p}} - \alpha^{pp}}{2} \tilde{\sigma}_{in} \right)$$

$$\text{Im} t_{el}^{\mathbb{P}}(s, b) = 1 - \sqrt{1 - \tilde{\sigma}_{in}} \cos \left(\frac{\alpha^{pp} + \alpha^{p\bar{p}}}{2} \tilde{\sigma}_{in} \right) \cos \left(\frac{\alpha^{p\bar{p}} - \alpha^{pp}}{2} \tilde{\sigma}_{in} \right)$$

- the C-odd part is:

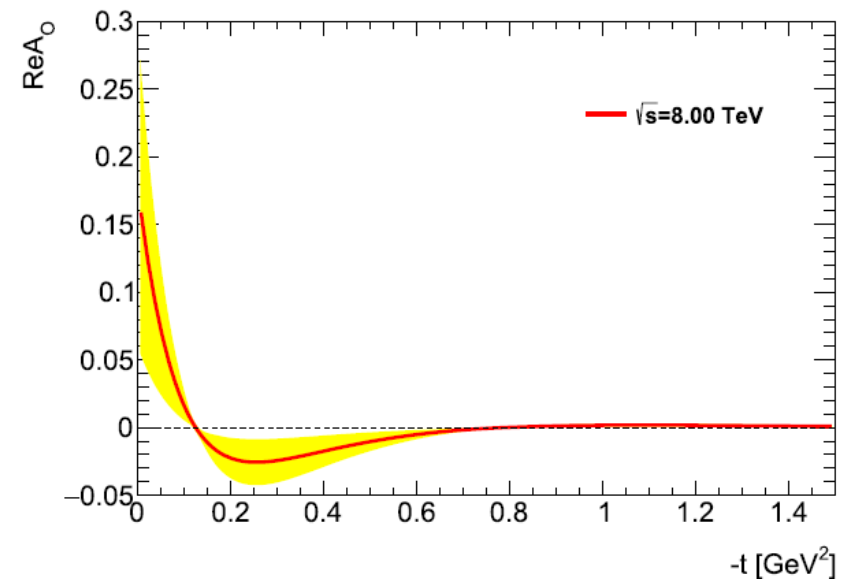
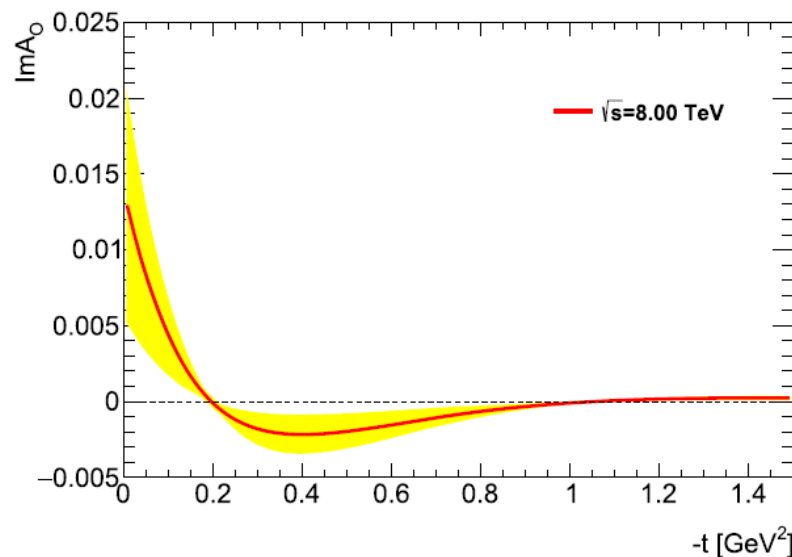
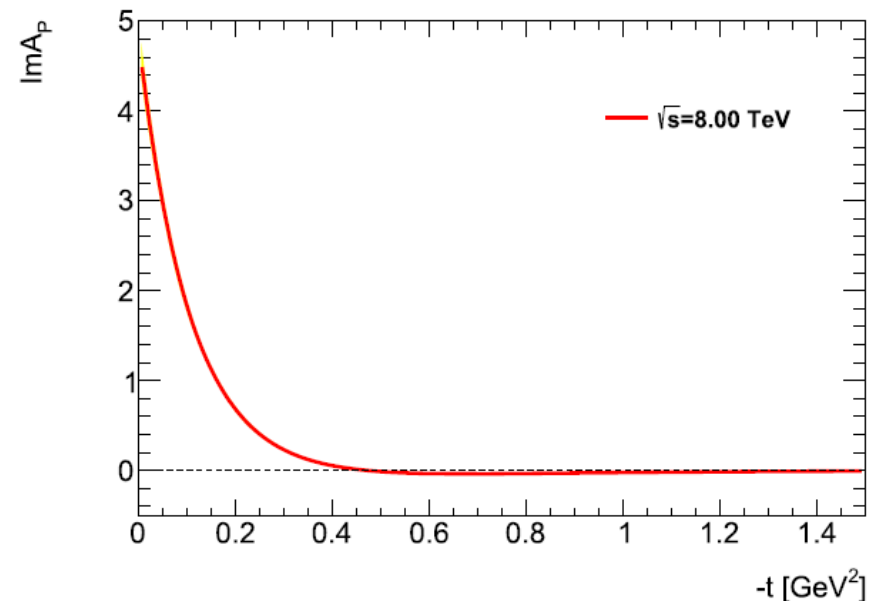
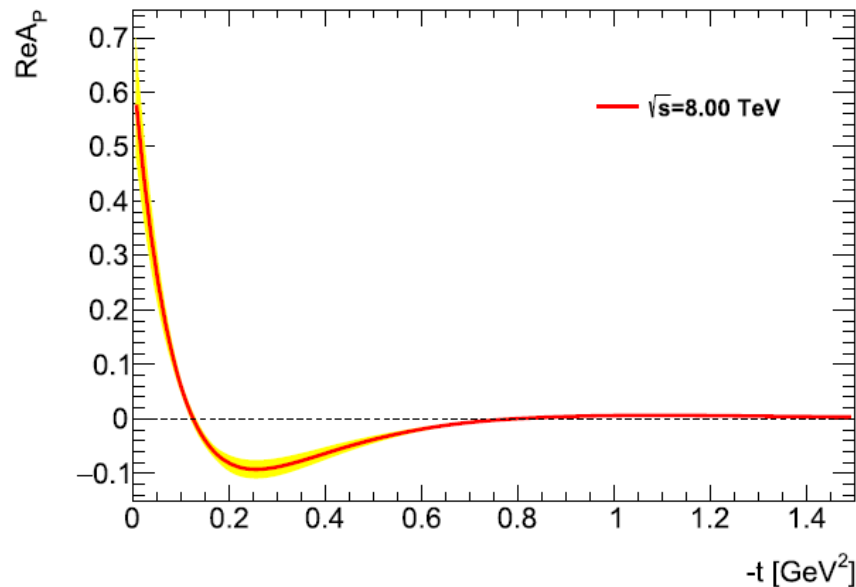
$$\text{Re} t_{el}^{\mathbb{O}}(s, b) = \sqrt{1 - \tilde{\sigma}_{in}} \sin \left(\frac{\alpha^{p\bar{p}} - \alpha^{pp}}{2} \tilde{\sigma}_{in} \right) \cos \left(\frac{\alpha^{p\bar{p}} + \alpha^{pp}}{2} \tilde{\sigma}_{in} \right)$$

$$\text{Im} t_{el}^{\mathbb{O}}(s, b) = \sqrt{1 - \tilde{\sigma}_{in}} \sin \left(\frac{\alpha^{p\bar{p}} - \alpha^{pp}}{2} \tilde{\sigma}_{in} \right) \sin \left(\frac{\alpha^{pp} + \alpha^{p\bar{p}}}{2} \tilde{\sigma}_{in} \right)$$

- the C-odd part vanishes if, and only if $\alpha^{pp}(s) \equiv \alpha^{p\bar{p}}(s)$

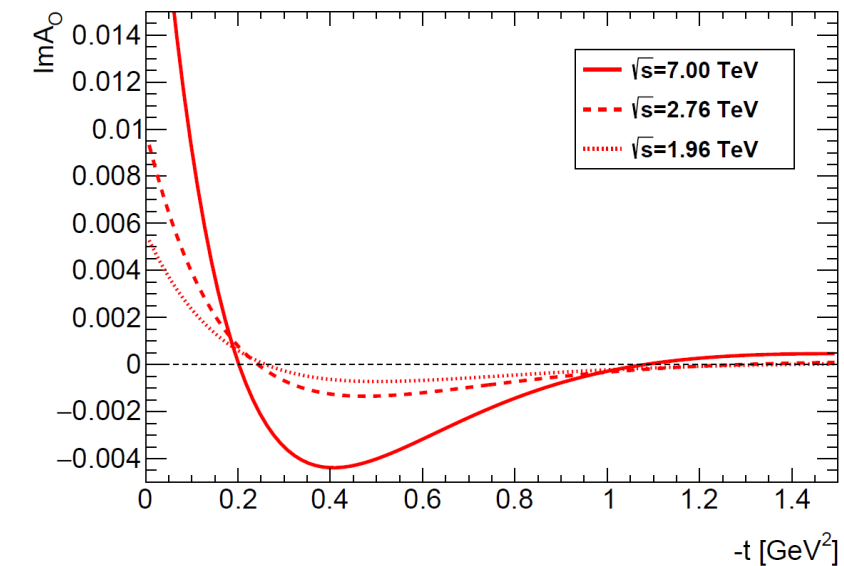
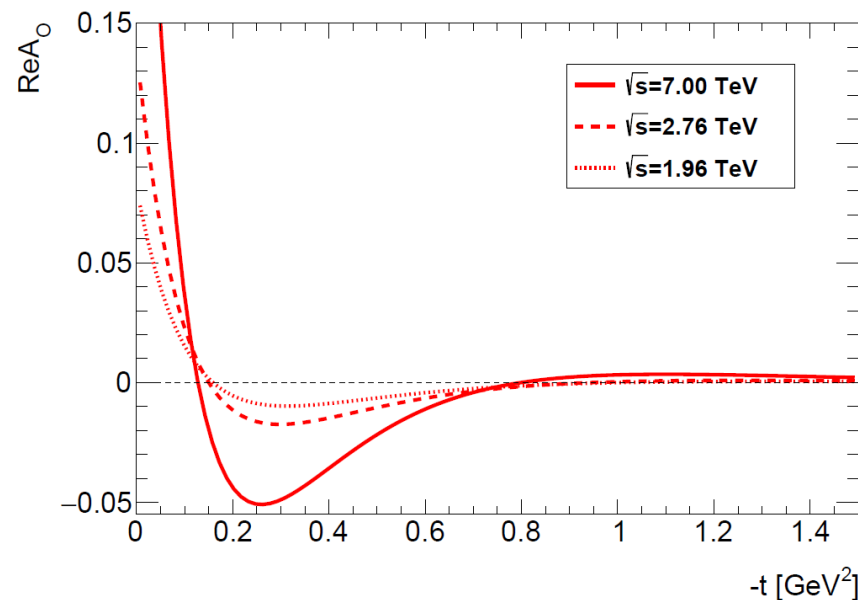
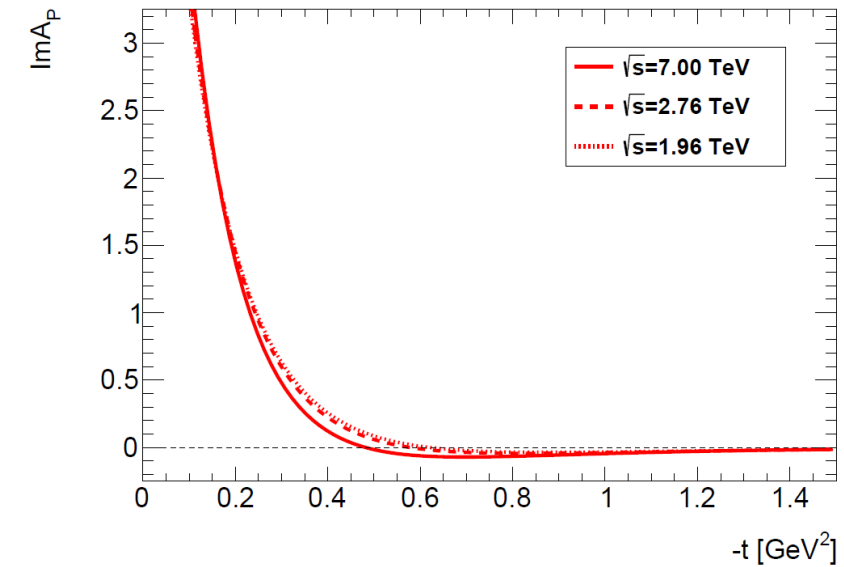
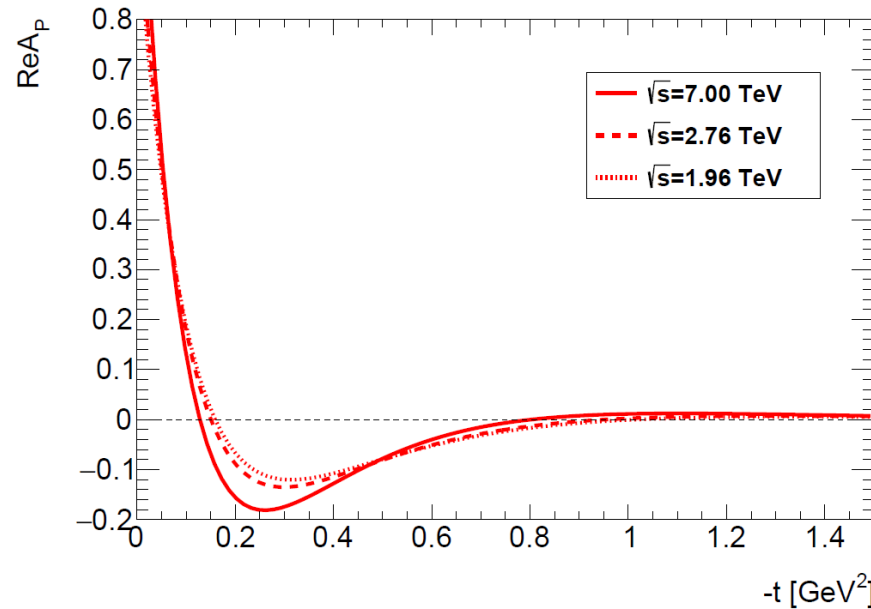
ReBB pomeron & odderon amplitudes @ 8 TeV

- both the real and the imaginary parts are dominated by the Pomeron at low- t
- the real part of the Odderon is about an order of magnitude larger, than its imaginary part; the opposite is true for the Pomeron
- however, the ReBB model is not calibrated at low- t (Levy generalization & Coulomb corrections are needed)



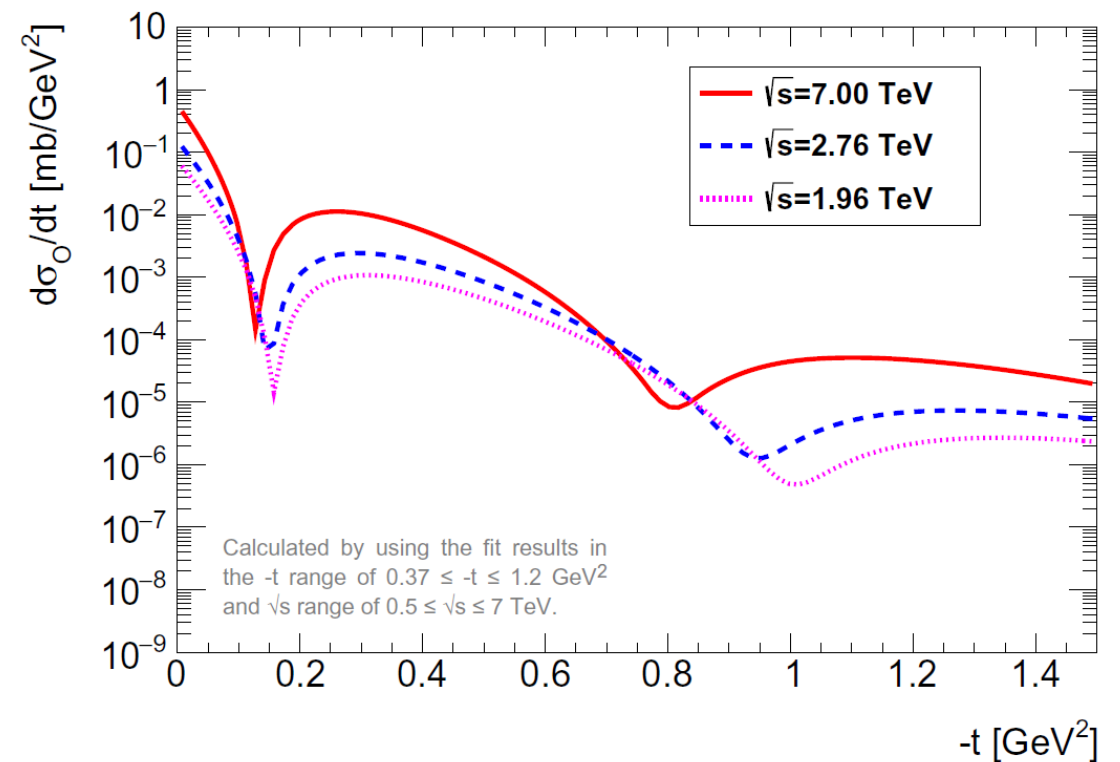
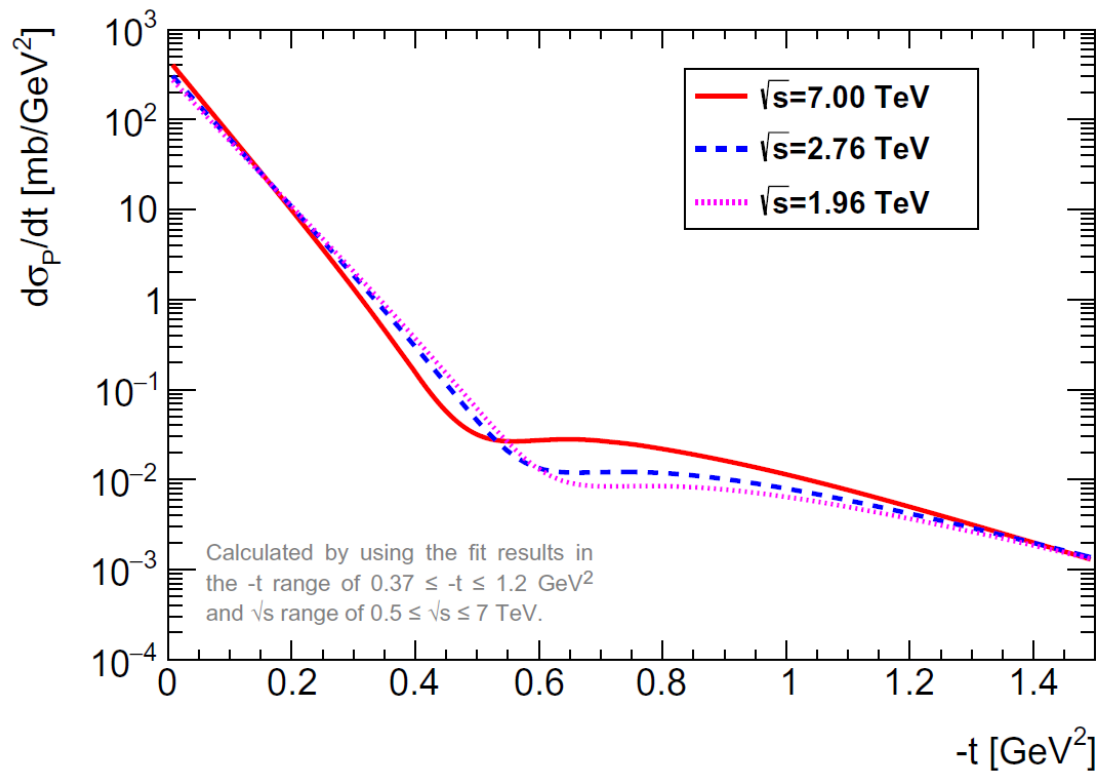
ReBB pomeron & odderon amplitudes @ TeV energies

- the real and imaginary parts of both pomeron and odderon amplitudes have fixed points
- The real part of the pomeron and the real and imaginary parts of the odderon have two zeros
- however, the ReBB model is not calibrated at low- t (Levy generalization & Coulomb corrections are needed)



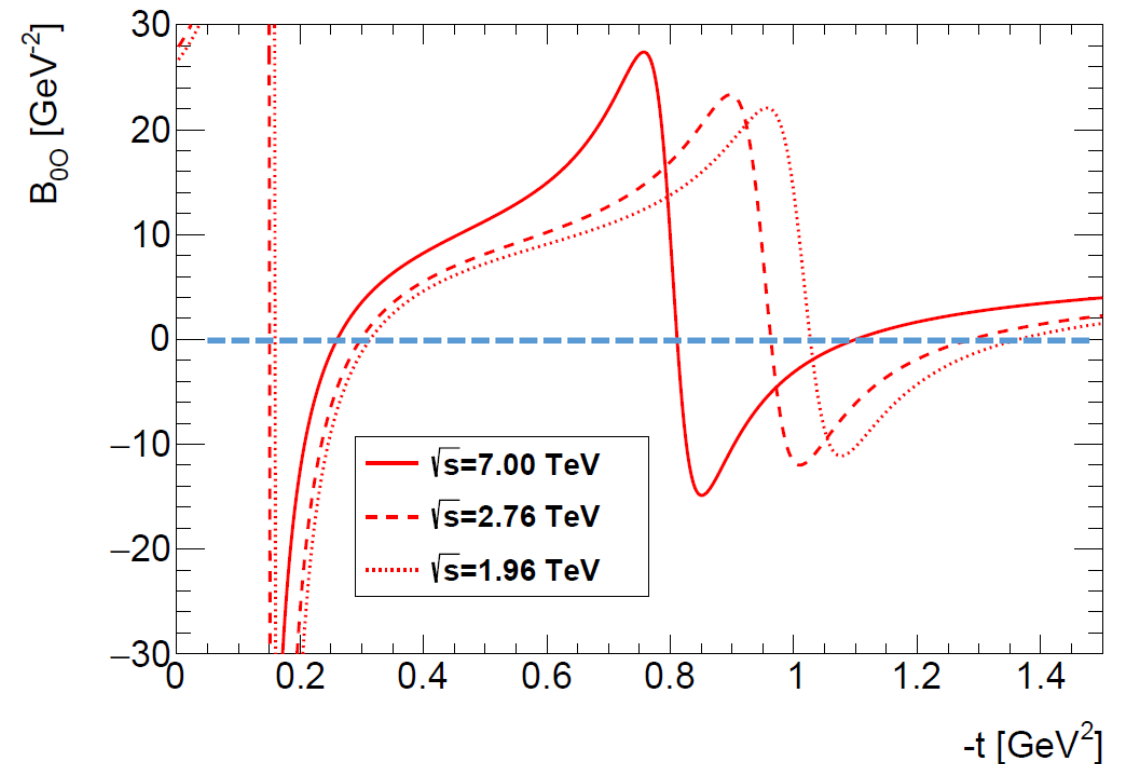
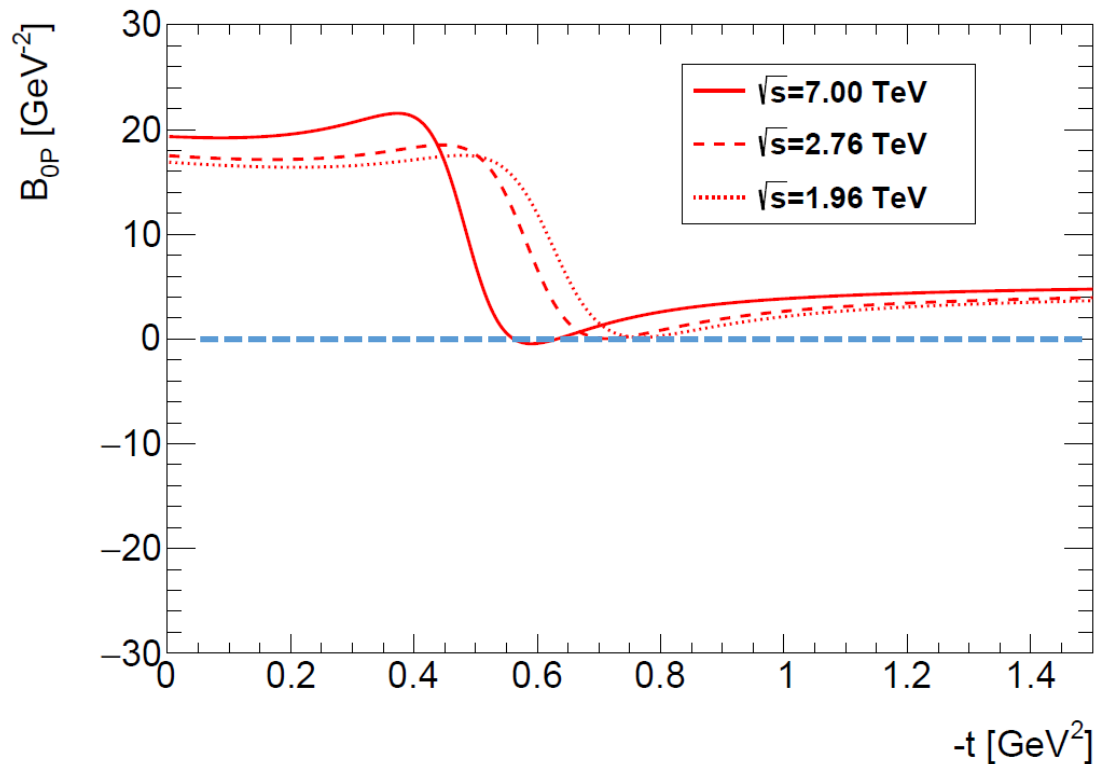
ReBB pomeron & odderon $d\sigma/dt$ @ TeV energies

- the pomeron exchange does not lead to a pronounced diffractive minimum structure
- the odderon exchange may lead even to two pronounced diffractive minima
- however, the interference between the pomeron and the odderon exchange leads to a single diffractive minimum in elastic pp collisions at the TeV scale



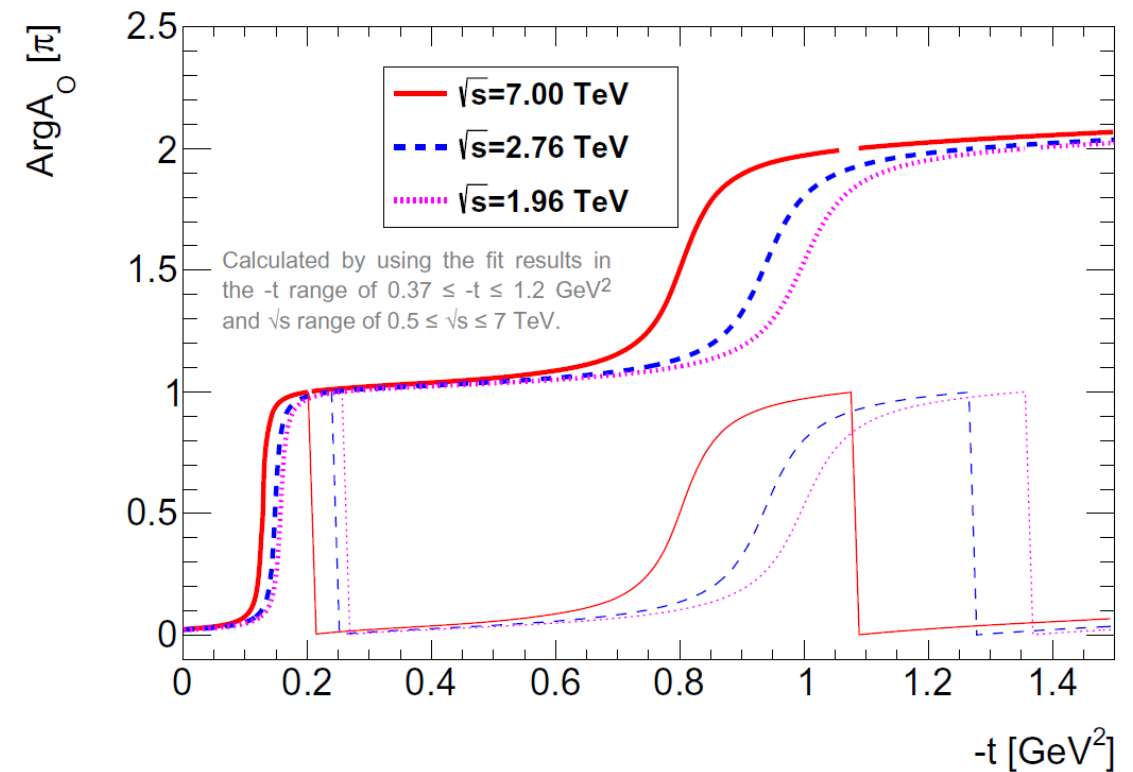
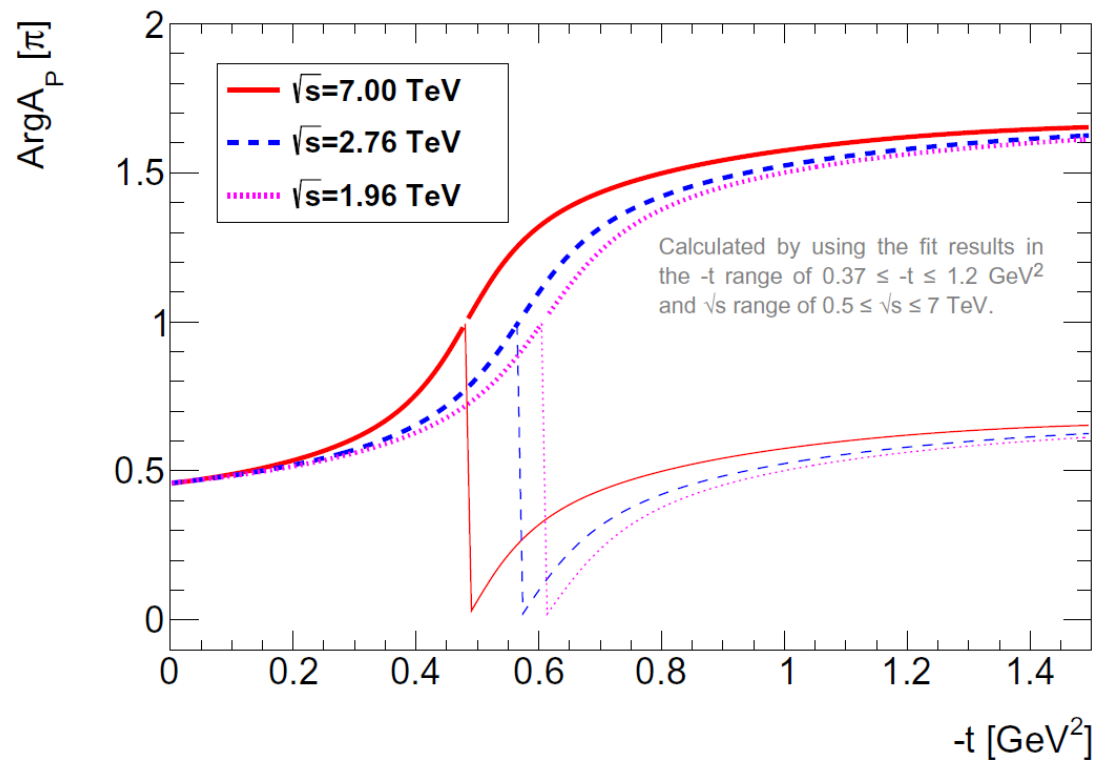
ReBB pomeron & odderon local slope @ TeV energies

- Near the diffractive minimum of the pomeron and odderon components of the differential cross section the local slopes reach or even cross zero



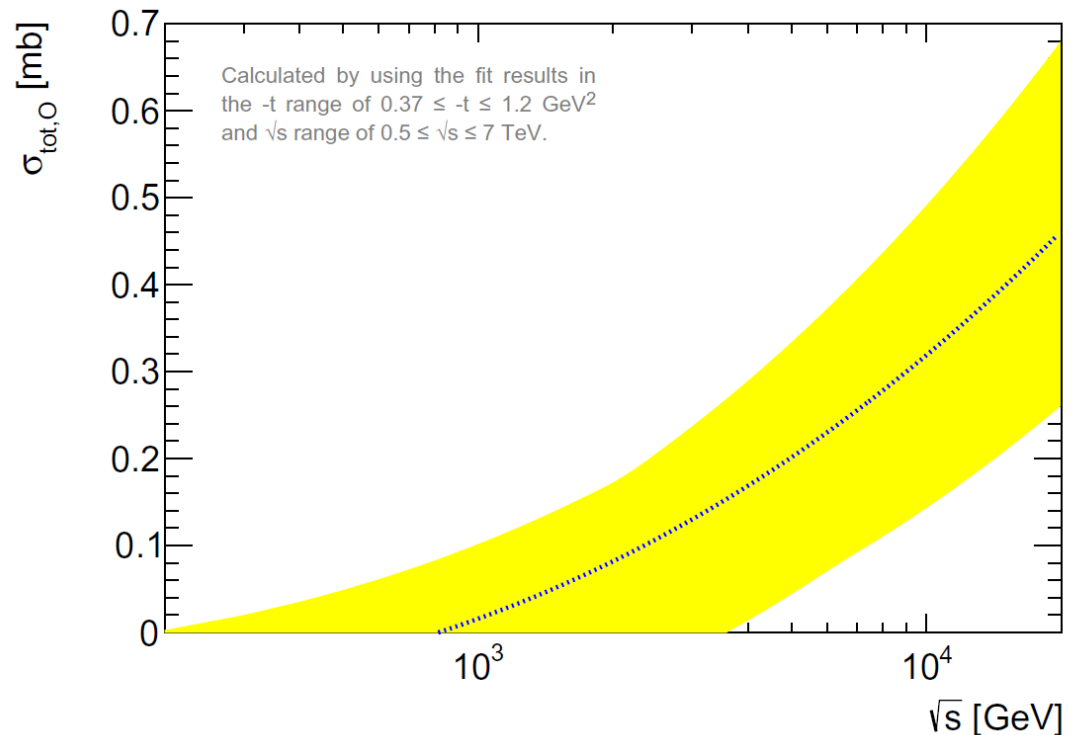
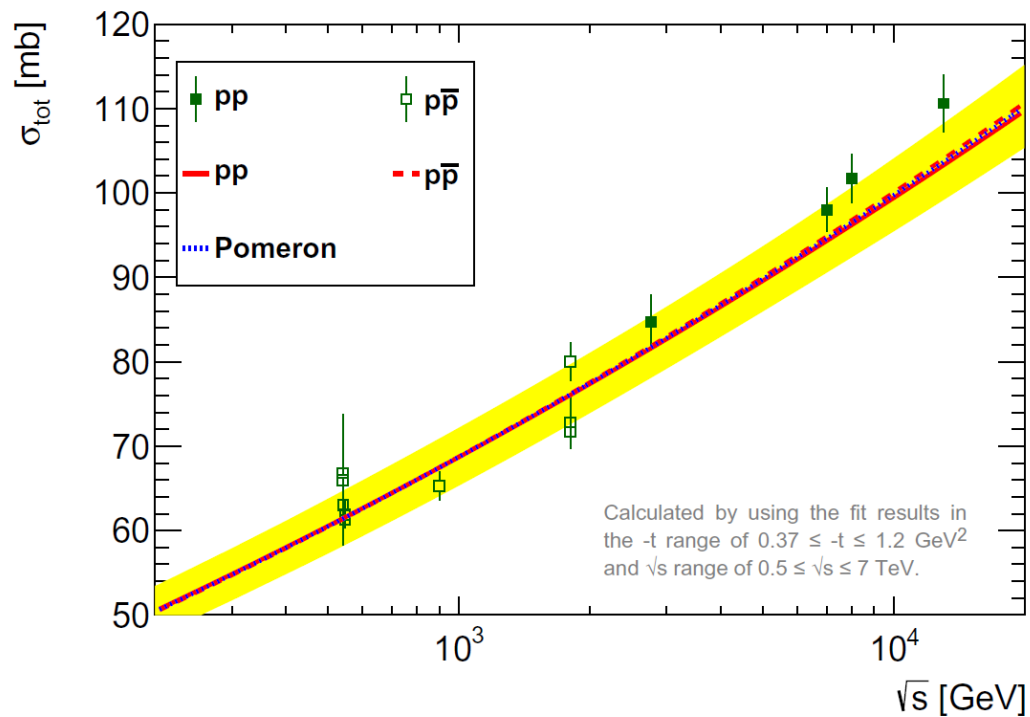
ReBB pomeron & odderon phase @ TeV energies

- at low- t , the Pomeron contribution is predominantly imaginary; near the diffractive minimum the real part is starting to be important
- at very low- t and near the diffractive minimum, the Odderon contribution is predominantly real

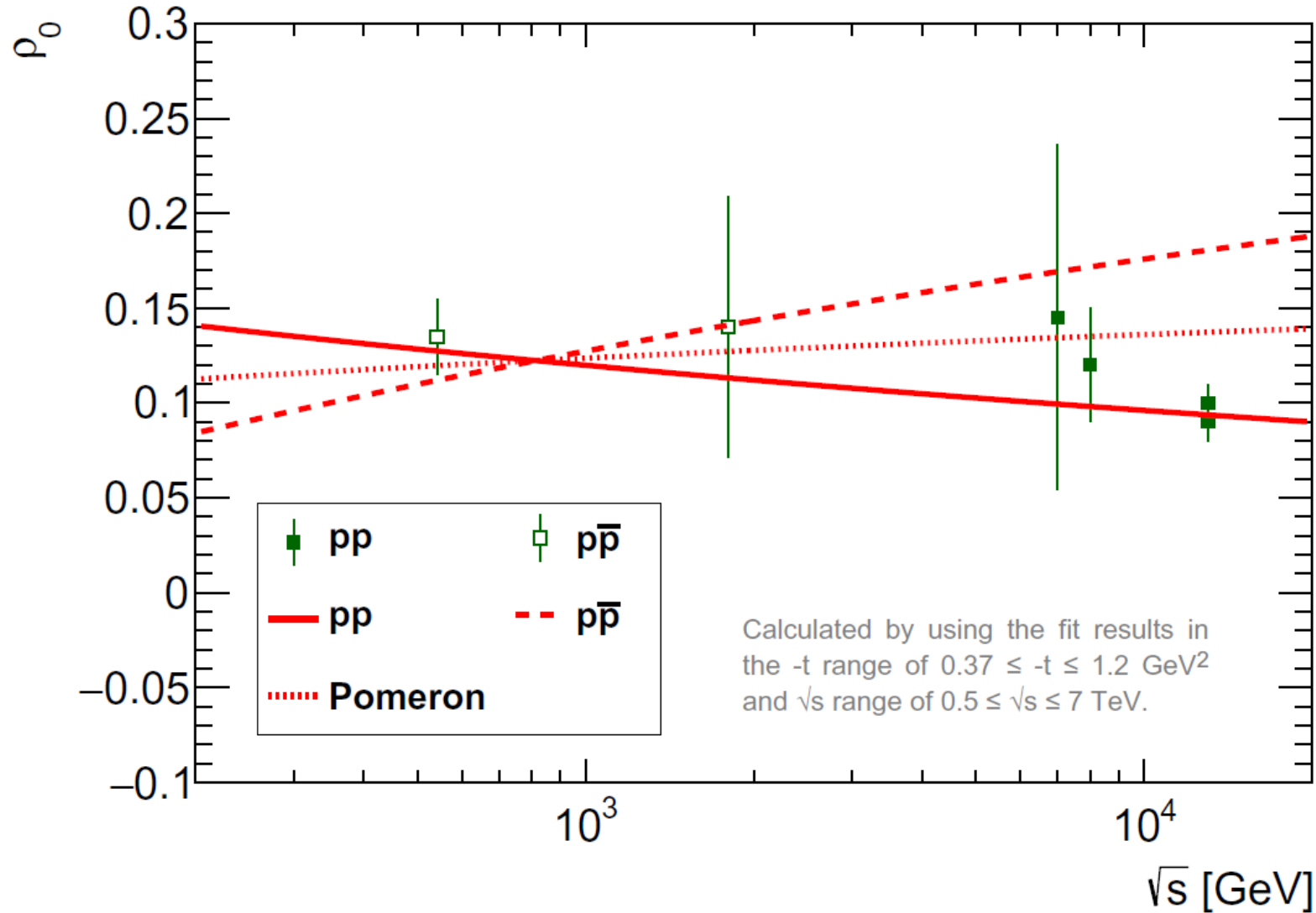


ReBB pomeron & odderon σ_{tot}

- within experimental uncertainties the effect of the odderon is negligible if one considers only σ_{tot} data
- the odderon starts to be present from about 1 TeV
- the pomeron is about two orders of magnitude greater than the odderon at LHC energies



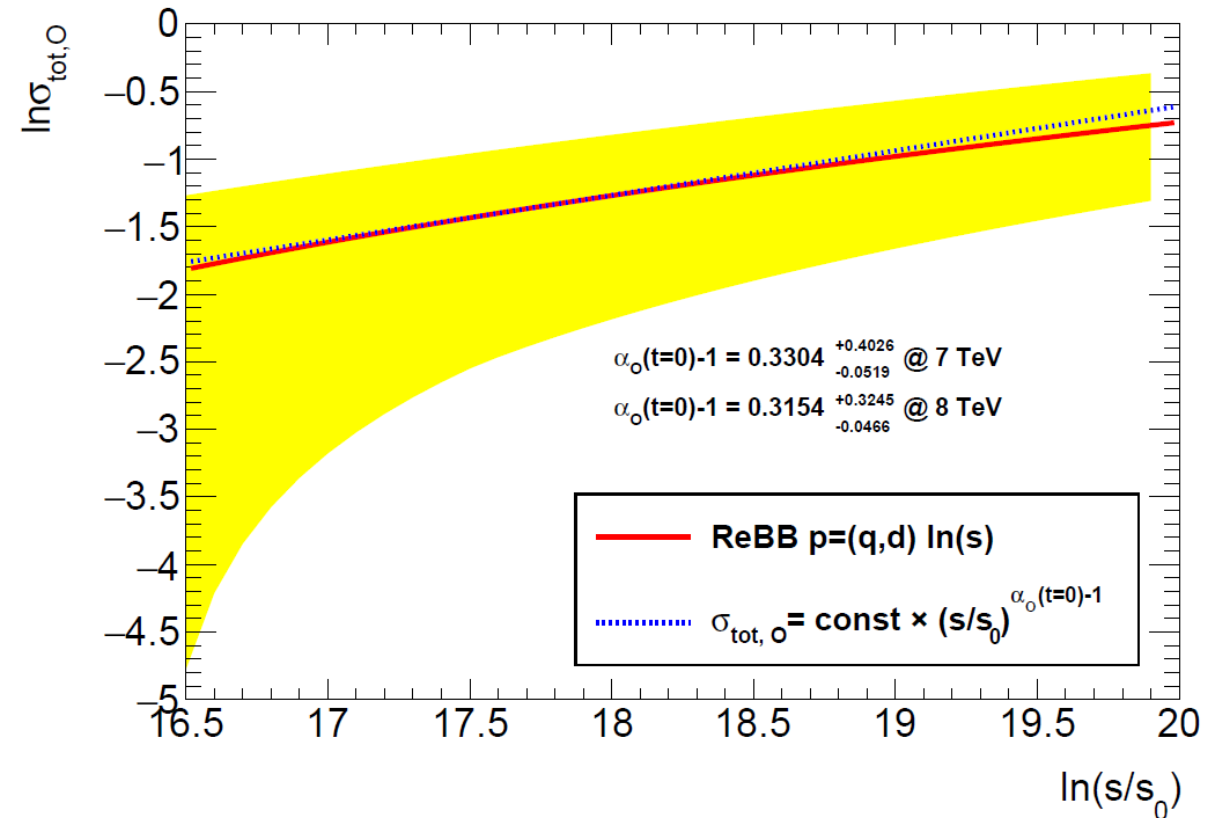
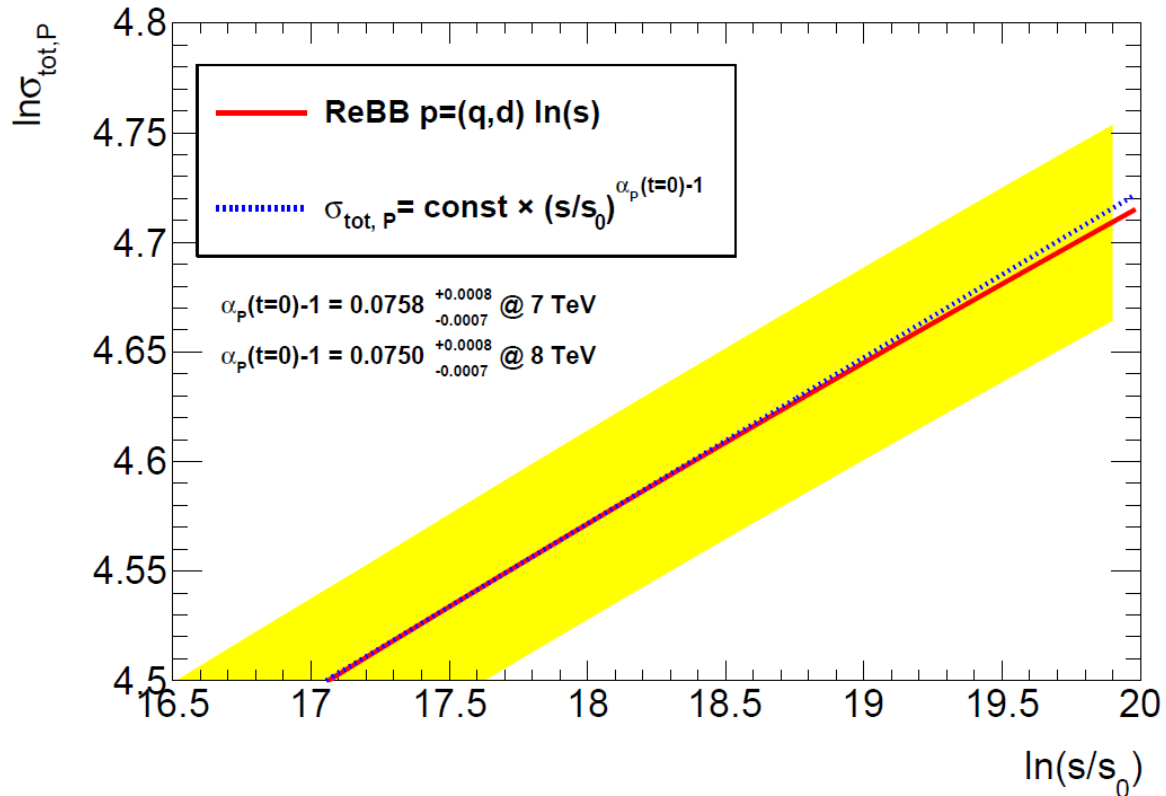
ReBB rho parameter with and without odderon



- Lévy generalization and Coulomb corrections are needed to get a decisive result
- A. Donnachie, P. V. Landshoff, Phys. Lett. B 831 (2022): weak odderon signal from ρ at 13 TeV

Pomeron & odderon intercepts

- the pomeron intercept is normal
- the odderon intercept is quite large



Summary & conclusions

- the ReBB model represents all available pp and $p\bar{p}$ $d\sigma/dt$ data in the kinematical ranges $0.546 \leq \sqrt{s} \leq 8$ TeV and $0.37 \leq -t \leq 1.2$ GeV² in a statistically acceptable manner
- from the pp and $p\bar{p}$ ReBB model amplitudes the ReBB model Pomeron and Odderon amplitudes are calculated
- from the ReBB model Pomeron and Odderon amplitudes the corresponding components of the measurable quantities are calculated
- there is a need for Lévy generalization and Coulomb corrections in order to get statistically acceptable and hence decisive results

Thank you for your attention!

