

Neutron stars in 1+4D with interacting Fermi gas

Anna Horváth

*Wigner Research Centre for Physics
Eötvös Loránd University*

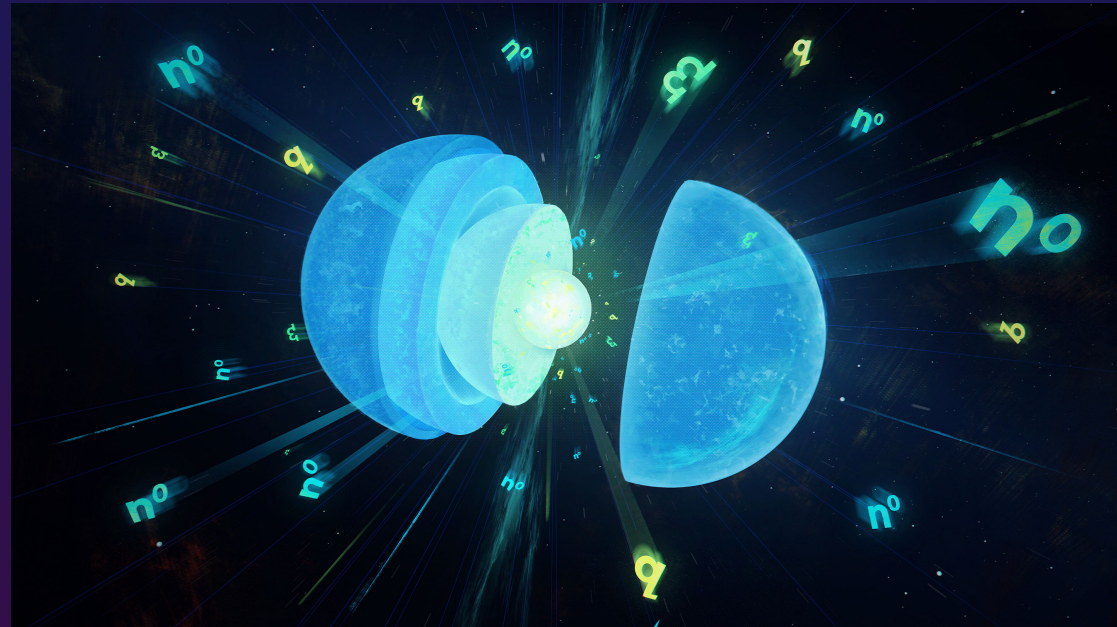
In collaboration with:

Gergely Gábor Barnaföldi

Wigner Research Centre for Physics

Emese Forgács-Dajka

Eötvös Loránd University

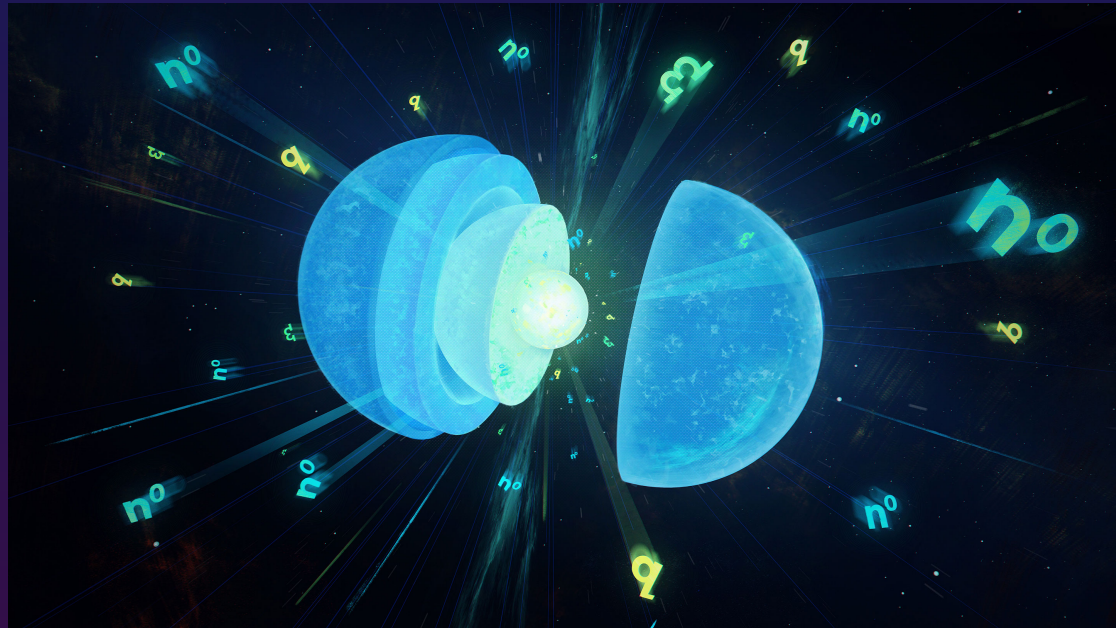


Neutron stars in 1+4D with interacting Fermi gas

- Remnants of supernovae
- Supported by baryon degeneracy
- Compact objects:
 $R \sim 10\text{km}$ $M \sim 1.4M_{\odot}$
- High density, low temperature



Probe physics in environments
not available on Earth

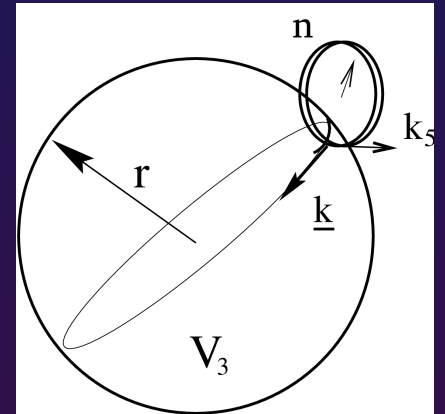


Treatment in $1+3+1_c$ dimensions

- Assume one extra **compactified spatial dimension** with size R_C
- At **each point** in ordinary 3D space **particles** with enough energy **can move into it**
 - **3D**: particles with **different masses**
 - **3+1_cD**: **one particle** but with **different quantized momenta** in the extra dimension

$$E_5 = \sqrt{\underline{k}^2 + \left(\frac{n}{R_C}\right)^2} + m^2 = \sqrt{\underline{k}^2 + \bar{m}^2}$$

$$\bar{m}^2 = \left(\frac{n}{R_C}\right)^2 + m^2$$



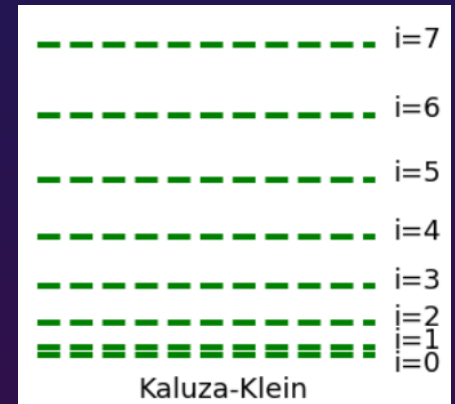
- With the **right** choice of R_C the **mass spectrum** of particles could be **reproduced**

Treatment in 1+3+1_c dimensions

- Assume one extra **compactified spatial dimension** with size R_C
- At **each point** in ordinary 3D space **particles** with enough energy **can move into it**
 - **3D**: particles with **different masses**
 - **3+1_cD**: **one particle** but with **different quantized momenta** in the extra dimension

$$E_5 = \sqrt{\underline{k}^2 + \left(\frac{n}{R_C}\right)^2} + m^2 = \sqrt{\underline{k}^2 + \bar{m}^2}$$

$$\bar{m}^2 = \left(\frac{n}{R_C}\right)^2 + m^2$$



- With the **right** choice of R_C the **mass spectrum** of particles could be **reproduced**

Building up stars

- Two equations needed:
 - Tolmann-Oppenheimer-Volkoff (TOV)

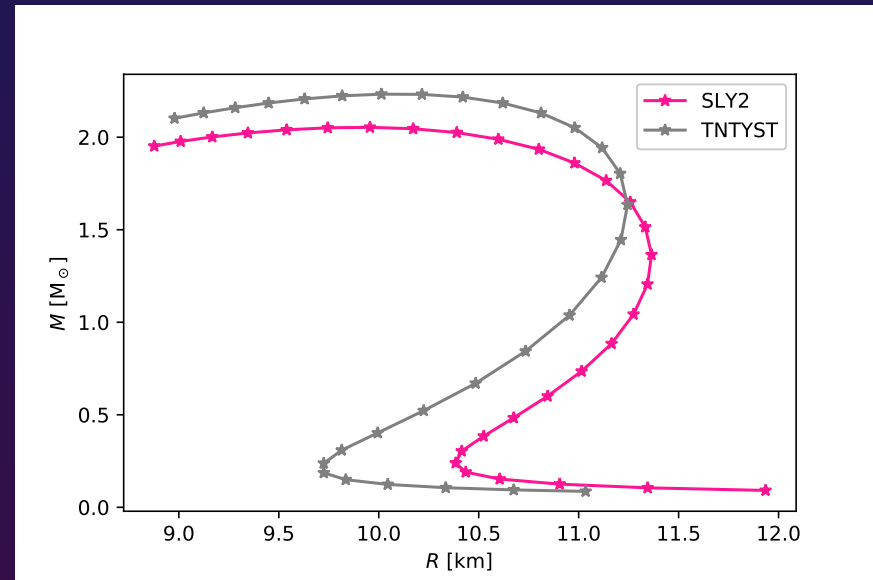
$$\frac{dp(r)}{dr} = -\frac{GM(r)\varepsilon(r)}{r^2} \times \left[1 + \frac{p(r)}{\varepsilon(r)}\right] \left[1 + \frac{4\pi r^3 p(r)}{M(r)}\right] \left[1 - \frac{GM(r)}{r}\right]^{-1}$$

$$M(r) = \int_0^r dr' 4\pi r'^2 \varepsilon(r')$$

- Equation of state (EoS) $\varepsilon(p)$
- Boundary conditions:
 - Pressure at the surface $p(R) = 0$
(in practice $p(R) = p_{\min} = 10^{-5} \text{ km}^{-2}$)
 - Central energy density ε_C



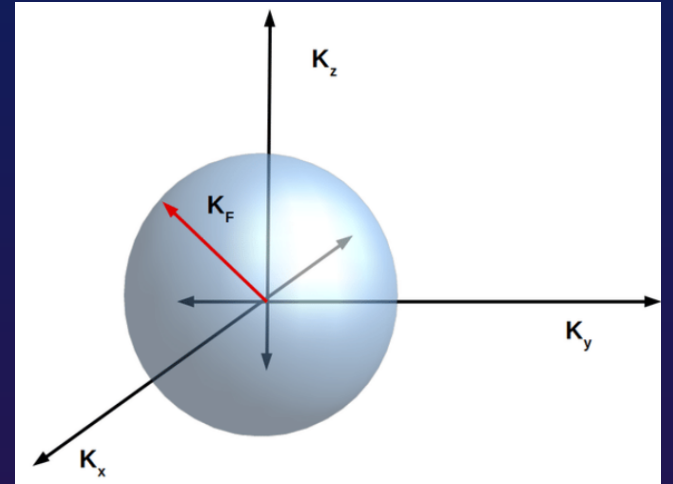
M-R diagrams



EoS in 1+4D

- **Interacting** degenerate **Fermi gas**
- **Potential** is a **linear** function of density:
$$U(n) = \xi n \quad \xi = \text{const}$$
- **Thermodynamic potential** on $T=0$ MeV

$$\tilde{\Omega} = \frac{-2k_B T V_{(d)}}{h^d} \int \ln \left(1 + e^{\frac{\mu - E(\mathbf{p})}{k_B T}} \right) d^d \mathbf{p}$$



- **Extra dimension** \longrightarrow calculate with **excited mass**
- **Interaction** \longrightarrow **chemical potential shifted by $-U(n)$**

$$\epsilon(\mu) = \epsilon_0(\mu - U(n)) + \epsilon_{int}$$

$$p(\mu) = p_0(\mu - U(n)) + p_{int}$$

$$n(\mu) = n_0(\mu - U(n))$$

$$\epsilon_{int} = p_{int} = \int U(n) dn = \int \xi n dn = \frac{1}{2} \xi n^2$$

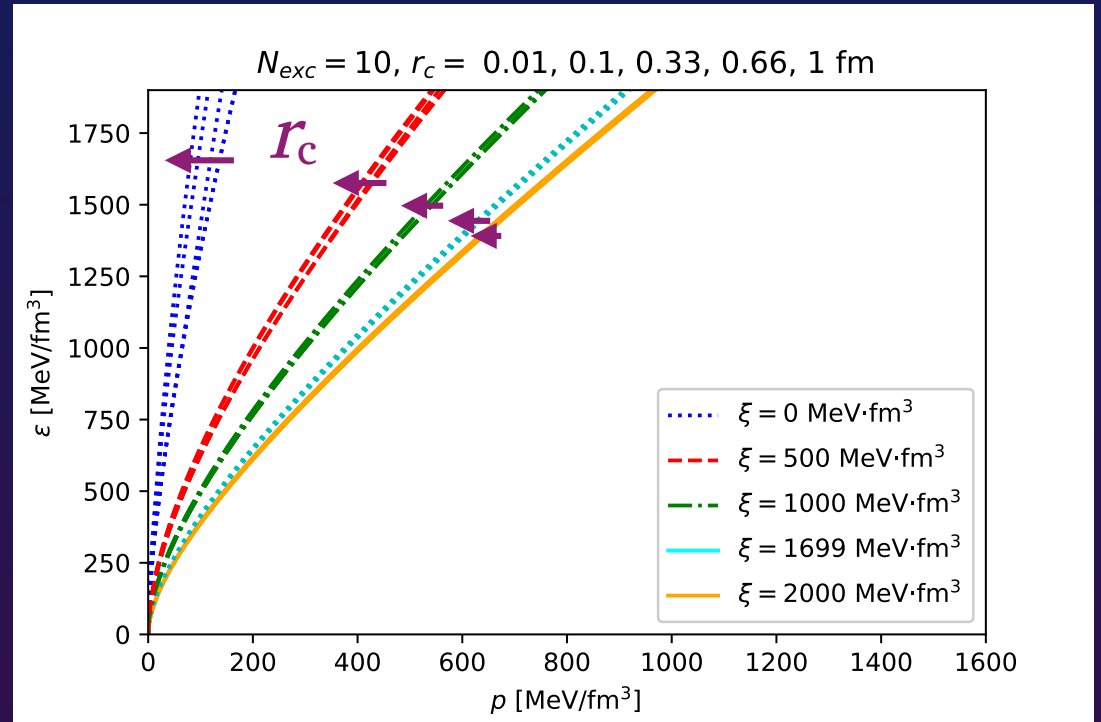
Relativity in 1+4D

- Assume (for **TOV**):
 - ◊ Spherical symmetry
 - ◊ Time-independence
 - ◊ Isotropic relativistic ideal fluid
- Assume (for **extra dimension**):
 - Microscopic
 - ◊ 4D metric does not depend on g_{55}
 - ◊ Causality postulates hold
 - ◊ Full Killing symmetry

$$g_{\mu\nu} = \begin{bmatrix} g_{00} & \cancel{g_{01}} & 0 & 0 & \cancel{g_{05}} \\ \cancel{g_{01}} & g_{11} & 0 & 0 & \cancel{g_{15}} \\ 0 & 0 & g_{22} & 0 & 0 \\ 0 & 0 & 0 & g_{22} \sin^2 \vartheta & 0 \\ \cancel{g_{05}} & \cancel{g_{15}} & 0 & 0 & \boxed{g_{55}} \end{bmatrix}$$

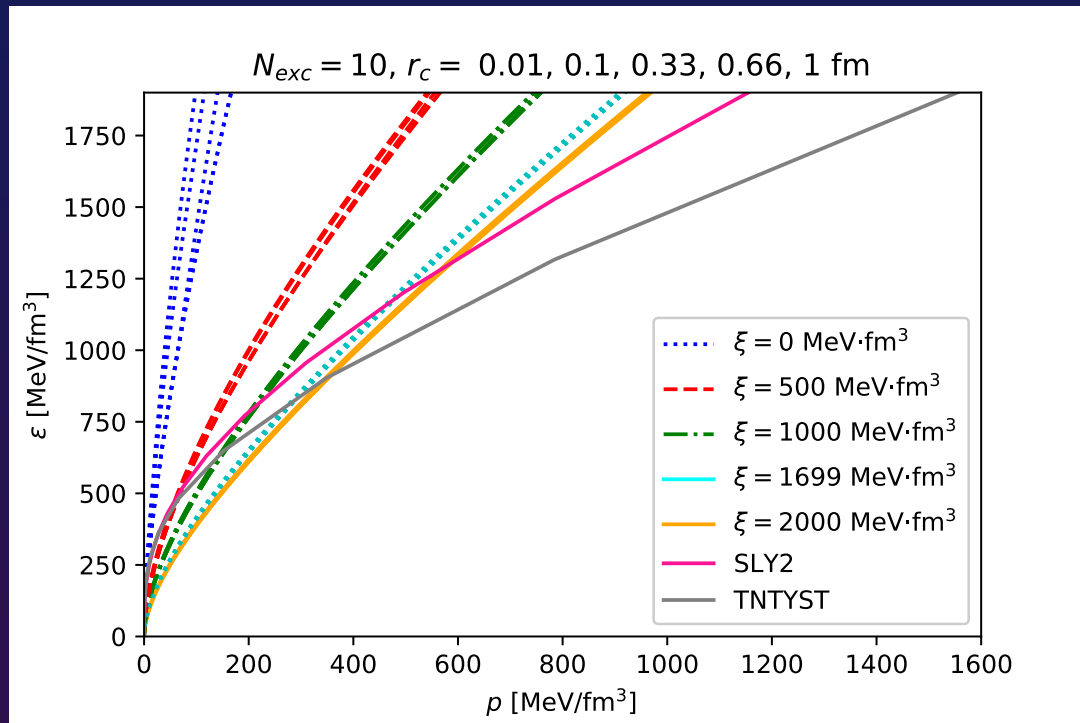
Equation of state

- ξ dependence is much more dominant than r_c
- The bigger ξ , the less important r_c becomes



Equation of state

- ξ dependence is much more dominant than r_c
- The bigger ξ , the less important r_c becomes
- For lower energies small ξ approximates more realistic EoSs
- For high energies a large ξ is a better approximation

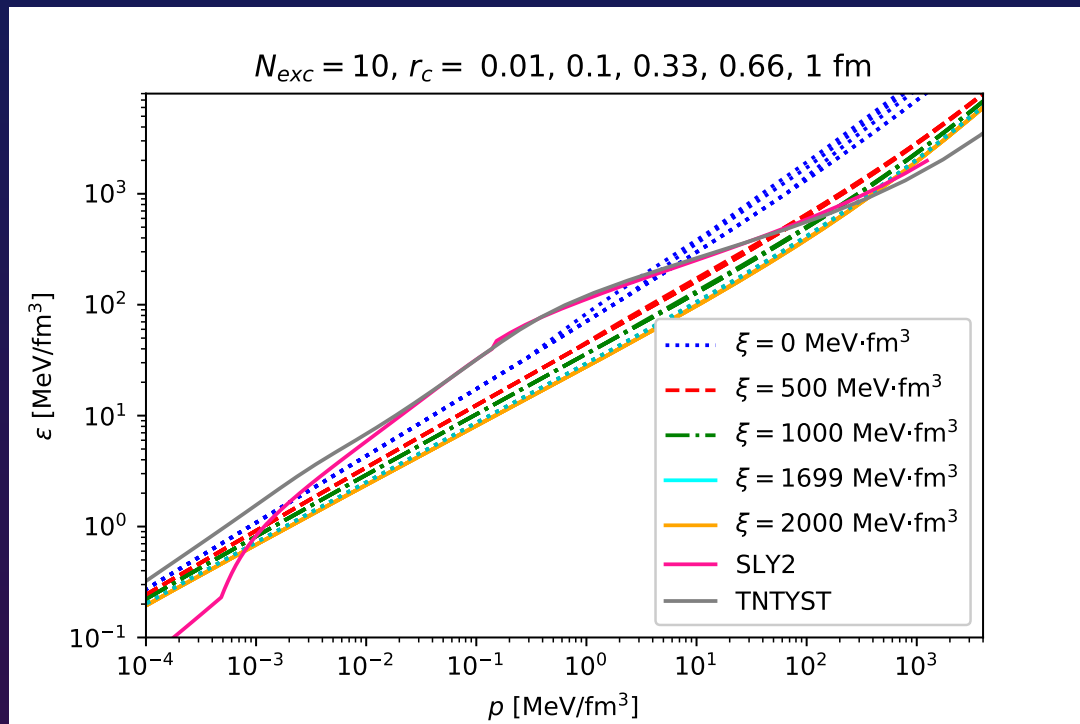


<https://compose.obspm.fr/>

1. H. Togashi, K. Nakazato, Y. Takehara, S. Yamamuro, H. Suzuki and M. Takano, Nucl. Phys. A 961 (2017) 78
1. F. Gulminelli and Ad. R. Raduta, arXiv:1504.04493.
2. E. Chabanat, Ph.D. thesis, University Claude Bernard Lyon-1, Lyon, France, 1995.
3. P. Danielewicz et J. Lee, Nucl. Phys. A818, 36 (2009).

Equation of state

- ξ dependence is much more dominant than r_c
- The bigger ξ , the less important r_c becomes
- For lower energies small ξ approximates more realistic EoSs
- For high energies a large ξ is a better approximation

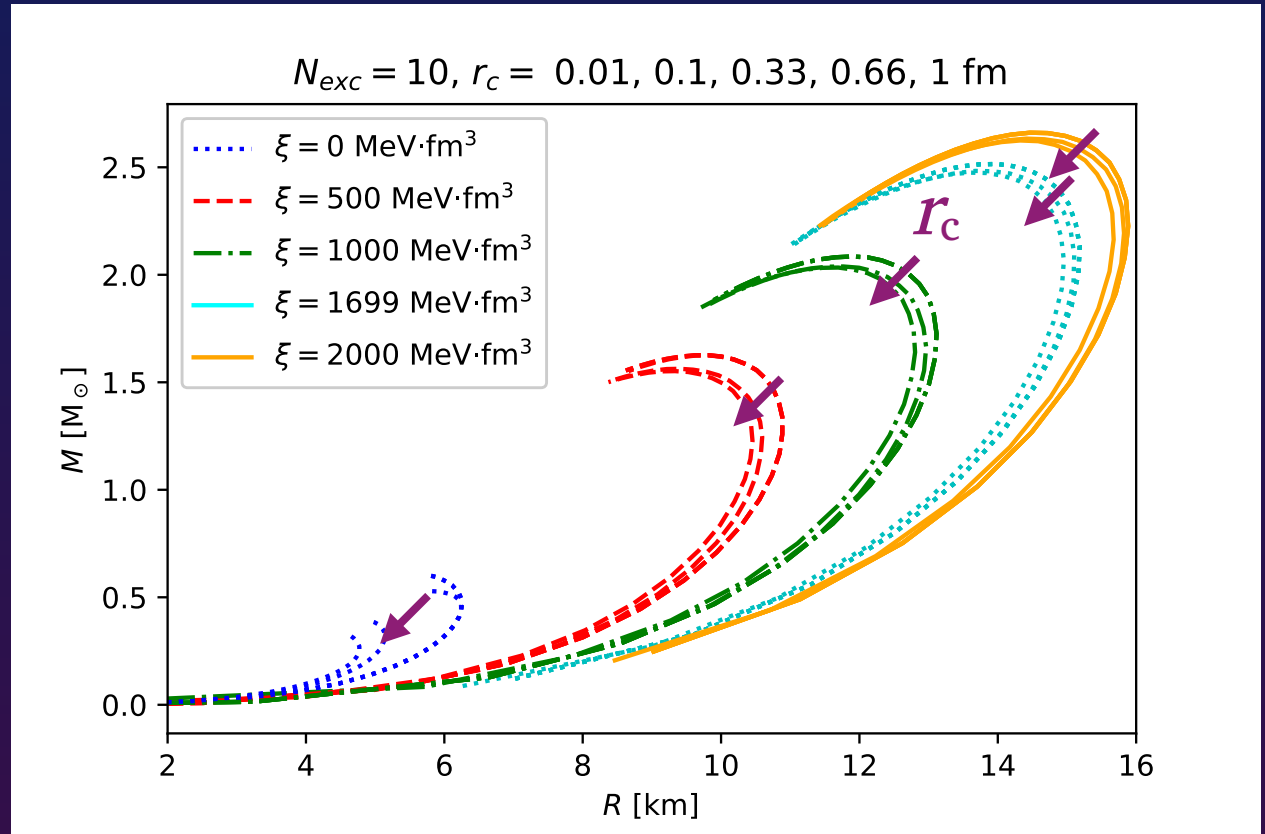


<https://compose.obspm.fr/>

1. H. Togashi, K. Nakazato, Y. Takehara, S. Yamamuro, H. Suzuki and M. Takano, Nucl. Phys. A 961 (2017) 78
1. F. Gulminelli and Ad. R. Raduta, arXiv:1504.04493.
2. E. Chabanat, Ph.D. thesis, University Claude Bernard Lyon-1, Lyon, France, 1995.
3. P. Danielewicz et J. Lee, Nucl. Phys. A818, 36 (2009).

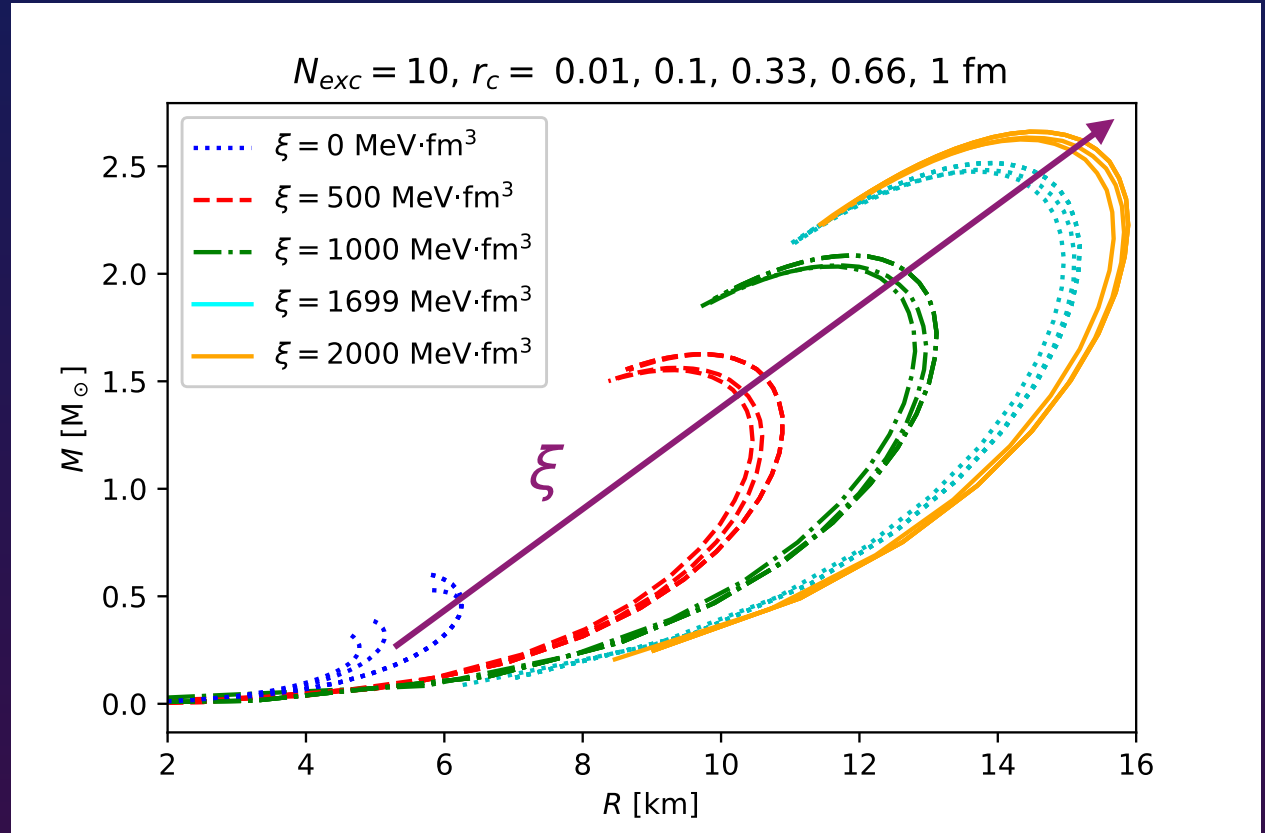
M - R diagrams of the EoS

- ξ dependence is much more dominant than r_c (latter only $\sim 5\%$)
- The bigger ξ , the less important r_c becomes



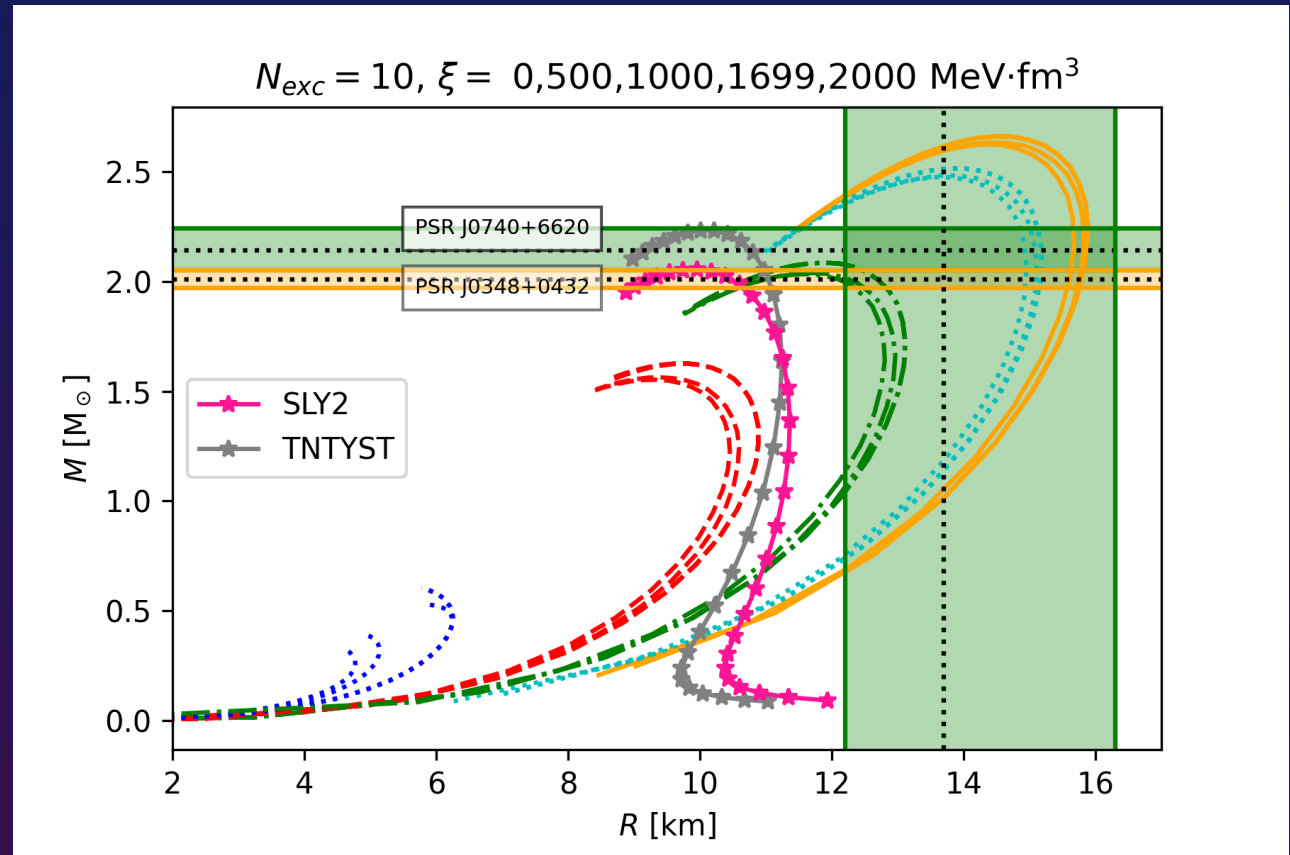
M - R diagrams of the EoS

- ξ dependence is much more dominant than r_c (latter only $\sim 5\%$)
- The bigger ξ , the less important r_c becomes



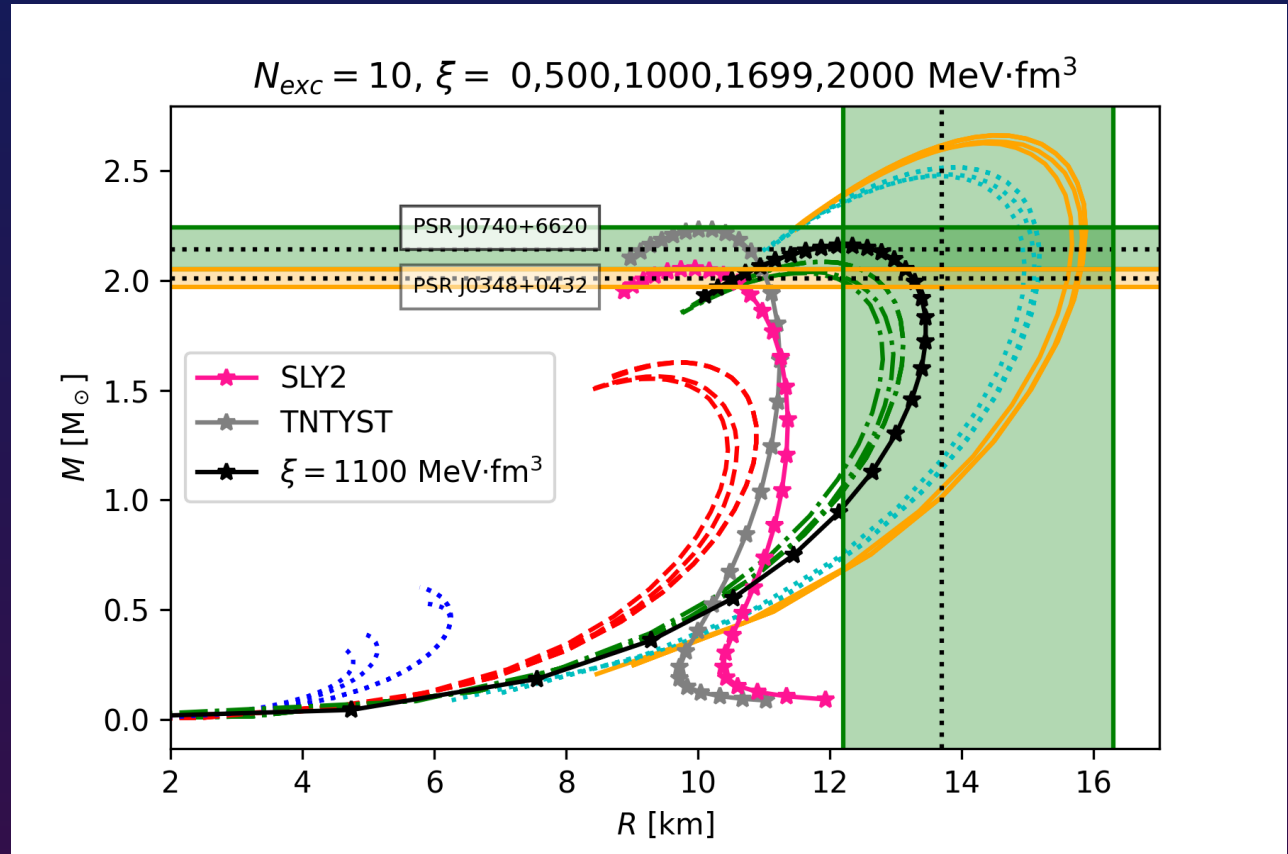
M - R diagrams

- + measurement data
- + 2 more realistic EoSs



M - R diagrams

- + measurement data
- + 2 more realistic EoSs
- + approximation for ξ

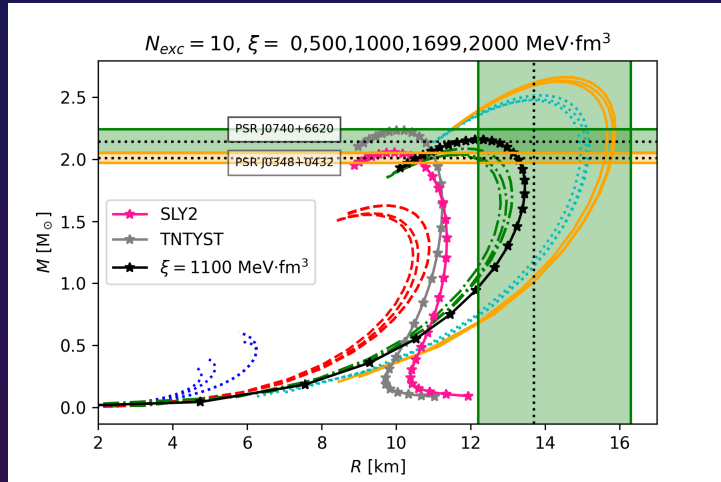


Summary

- Model with the possibility of probing **beyond standard model physics**
- One **extra spatial compactified dimension**



Ordinary **mass** can be described as **quantized 5thD momenta**



- Effective nuclear field theory with **linear repulsive potential**
- It is **possible** to **build** compact stars with **realistic** properties
- **Refinement** of the EoS is **needed** to get more accurate results