

# High-pT v4 puzzle and importance of higher harmonics in QGP tomography

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# DREENA framework

- **Dynamical Radiative and Elastic ENergy loss Approach**
- fully optimized numerical procedure capable of generating high  $p_{\perp}$  predictions
- includes:
  - parton production
  - multi gluon-fluctuations
  - path-length fluctutations
  - fragmentation functions
- keeping all elements of the state-of-the art energy loss formalism, while introducing more complex temperature evolutions:
  - **DREENA-C: constant temperature medium**  
D. Z., I. Salom, J. Auvinen, M. Djordjevic and M. Djordjevic, J. Phys. G **46**, no. 8, 085101 (2019)
  - **DREENA-B: Bjorken expansion**  
D. Z., I. Salom, J. Auvinen, M. Djordjevic and M. Djordjevic, Phys. Lett. B **791**, 236 (2019)
  - **DREENA-A: smooth (2+1)D temperature evolution**  
D. Z., I. Salom, J. Auvinen, P. Huovinen and M. Djordjevic, Front. in Phys. **10**, 957019 (2022)
  - **ebe-DREENA: event-by-event fluctuating hydro background**  
D. Z., J. Auvinen, I. Salom, P. Huovinen and M. Djordjevic, Phys. Rev. C **106**, no.4, 044909 (2022)

- DREENA-C = constant temperature medium

D. Z., I. Salom, J. Auvinen, M. Djordjevic and M. Djordjevic, J. Phys. G **46**, no. 8, 085101 (2019)

- qualitatively good agreement with the data
- rough approximation - good for testing path-length dependence of energy loss

M. Djordjevic, D. Z, M. Djordjevic and J. Auvinen, Phys. Rev. C **99**, no.6, 061902 (2019)

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D. Z., I. Salom, J. Auvinen, M. Djordjevic and M. Djordjevic, J. Phys. G **46**, no. 8, 085101 (2019)

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- rough approximation - good for testing path-length dependence of energy loss

M. Djordjevic, D. Z, M. Djordjevic and J. Auvinen, Phys. Rev. C **99**, no.6, 061902 (2019)

- DREENA-B = 1D Bjorken evolution

D. Z., I. Salom, J. Auvinen, M. Djordjevic and M. Djordjevic, Phys. Lett. B **791**, 236 (2019)

- qualitatively and quantitatively good agreement with the data - only thermalization time as QGP property
- good for testing initial stages without impacting anything else

D. Z, B. Ilic, M. Djordjevic and M. Djordjevic, Phys. Rev. C **101**, no.6, 064909 (2020)

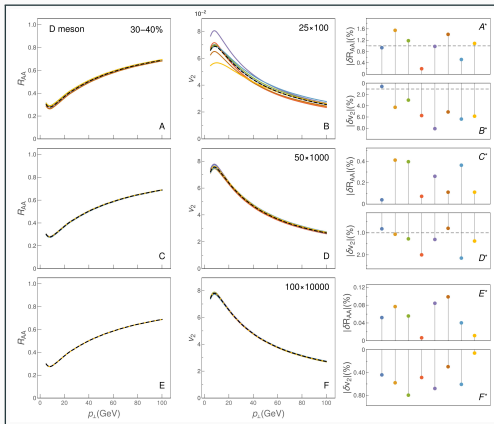


## A = adaptive

D. Z. I. Salom, J. Auvinen, P. Huovinen and M. Djordjevic, *Front. in Phys.* **10**, 957019 (2022)

- C & B good test for energy loss but no bulk medium properties to extract
- includes any, arbitrary, medium evolution as an input
- preserve all dynamical energy loss model properties
- generate a comprehensive set of light and heavy flavor suppression predictions
- needs to be an efficient (timewise) numerical procedure

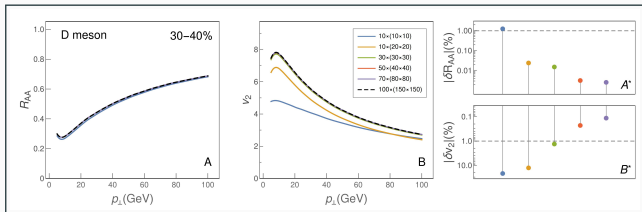
## Monte Carlo sampling



not very efficient

for  $v_2$ , one million trajectories needed to achieve a precision below 1%

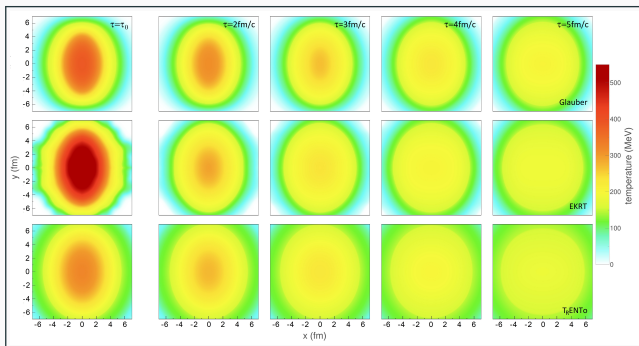
## equidistant sampling



two orders of magnitude increase in the efficiency  
for  $v_2$  , only 10k trajectories needed to achieve  $\sim 1\%$  precision

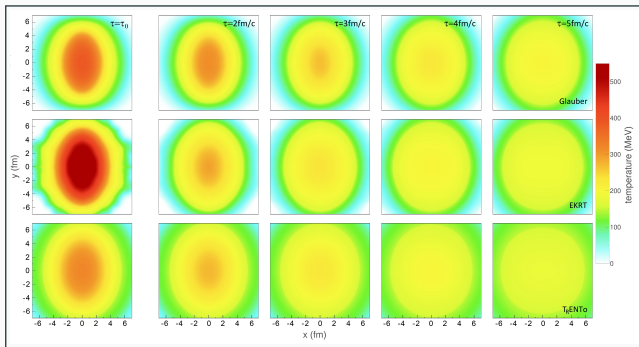
can efficiently generate predictions for all types of probes for  
arbitrary temperature profiles

Are high- $p_{\perp}$  observables indeed sensitive to different T evolutions?

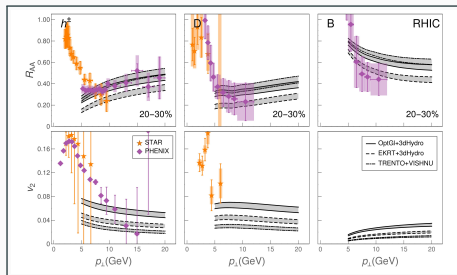
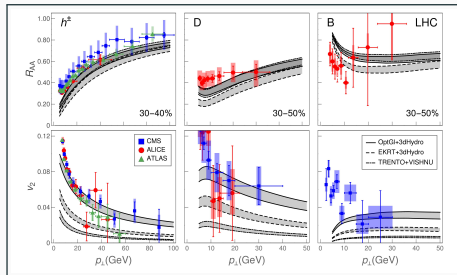


All three evolutions agree with low- $p_{\perp}$  data. Can high  $p_{\perp}$ -data provide further constraint?

## Qualitative differences



- Largest anisotropy for Glauber ( $\tau_0 = 1\text{fm}$ ) – expected differences in high- $p_{\perp}$   $v_2$
- EKRT shows larger temperature - smaller  $R_{AA}$  expected



- 'EKRT' initial conditions indeed lead to the smallest  $R_{AA}$
- anisotropy translates to  $v_2$  differences ('Glauber' largest, T<sub>R</sub>ENTo lowest)
- DREENA-A can differentiate between different T profiles
- heavy flavour even more sensitive to different T profiles
- additional (independent) constraint to low- $p_{\perp}$  data

## DREENA-A OUTLOOK

- DREENA-A is a fully optimized numerical implementation of the dynamical energy loss
- can include arbitrary *smooth* temperature profiles
- no additional fitting parameters within energy loss
- limitations: higher harmonics

## event-by-event DREENA

D. Z, J. Auvinen, I. Salom, P. Huovinen and M. Djordjevic, Phys. Rev. C **106**, no.4, 044909 (2022)

- generalization of DREENA-A
- high- $p_{\perp}$  energy loss on fluctuating hydro background
- different initial conditions and hydro models
- can produce high- $p_{\perp}$  higher harmonics
- 1st question: averaging over events



- cummulants:  $v_n\{2\}, v_n\{4\}$

A. Bilandzic, R. Snellings and S. Voloshin, Phys. Rev. C **83**, 044913 (2011).

- event plane:  $v_n\{EP\}$

Y. He, W. Chen, T. Luo, S. Cao, L. G. Pang and X. N. Wang, [arXiv:2201.08408 [hep-ph]].

- scalar product:  $v_n\{SP\}$

C. Andres, N. Armesto, H. Niemi, R. Paatelainen and C. A. Salgado, Phys. Lett. B **803**, 135318 (2020)

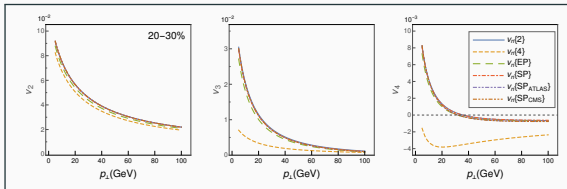
Y. He, W. Chen, T. Luo, S. Cao, L. G. Pang and X. N. Wang, [arXiv:2201.08408 [hep-ph]]

- scalar product - ATLAS:  $v_n\{SP_{ATLAS}\}$

M. Aaboud *et al.* [ATLAS], Eur. Phys. J. C **78**, no.12, 997 (2018)

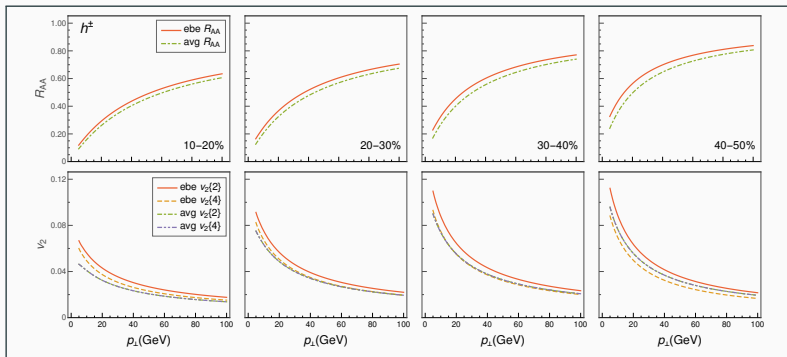
- scalar product - CMS:  $v_n\{SP_{CMS}\}$

A. M. Sirunyan *et al.* [CMS], Phys. Lett. B **776**, 195-216 (2018)



all methods, other than  $v_n\{4\}$  agree with each other  
no need for rapidity correlations!

- high- $p_T$  energy loss: ebe fluctuation vs smooth hydro background

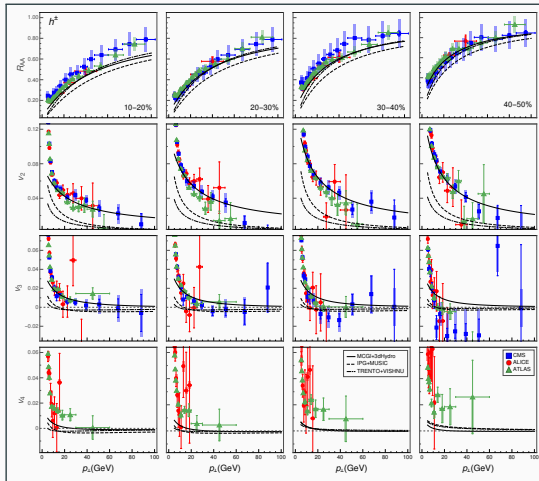


$R_{AA}$  differences small  $\sim 7\%$  and no centrality dependence

$v_2\{2\}$  differences from 14% in 40-50% up to 32% in 10-20%

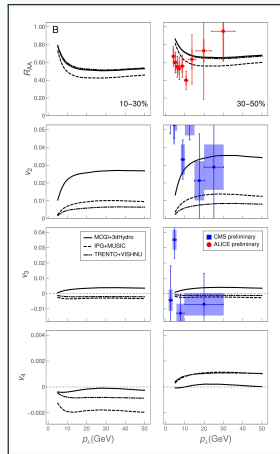
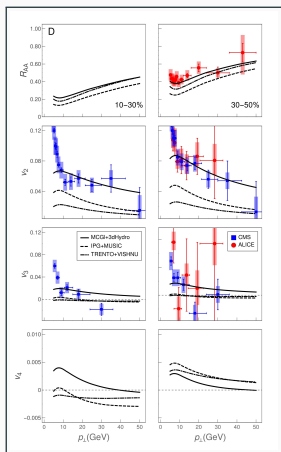
also  $p_\perp$  dependence of the differences

- charged hadrons,  $Pb + Pb$ ,  $\sqrt{s_{NN}} = 5.02 TeV$



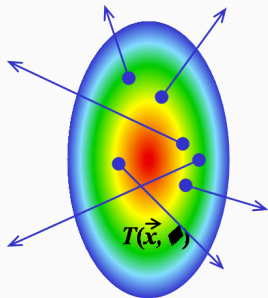
we can distinguish between different models with high- $p_T$  energy loss

- heavy flavour,  $Pb + Pb$ ,  $\sqrt{s_{NN}} = 5.02 TeV$



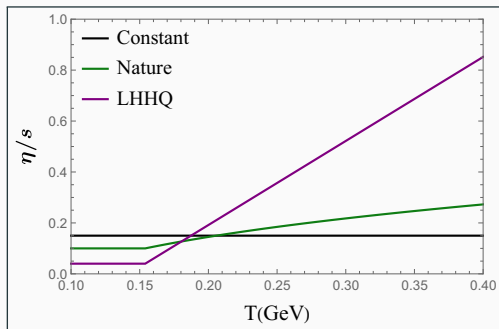
we can distinguish between different models with high- $p_T$  energy loss on heavy flavour as well - even more sensitive

- DREENA-A is a fully optimized numerical implementation of the dynamical energy loss formalism
- can include arbitrary temperature profiles, both smooth and event-by-event fluctuating
- no additional fitting parameters within energy loss
- high- $p_{\perp}$   $R_{AA}$  ,  $v_2$  , and higher harmonics show qualitative and quantitative sensitivity to details of T profile differences
- applicable to different types of flavor, collision systems, and energies
- OUTLOOK: an efficient QGP tomography tool for constraining the medium properties by both high- and low- $p_{\perp}$  data



- bulk GPQ properties traditionally explored by low- $p_{\perp}$  observables
- high energy particles lose energy
- energy loss sensitive to QGP properties
- predict the energy loss of high  $p_{\perp}$  probes
- use high  $p_{\perp}$  probes to infer QGP properties:
  - early evolution  
S. Stojku, J. Auvinen, M. Djordjevic, P. Huovinen and M. Djordjevic, Phys. Rev. C **105**, no.2, L021901 (2022)
  - jet anisotropy  
S. Stojku, J. Auvinen, L. Zivkovic, P. Huovinen and M. Djordjevic, Phys. Lett. B **835**, 137501 (2022)
- DREENA-A on github:  
<https://github.com/DusanZigic/DREENA-A>

## $\eta/s$ parametrization



- constant  $\eta/s$

- nature  $\eta/s$

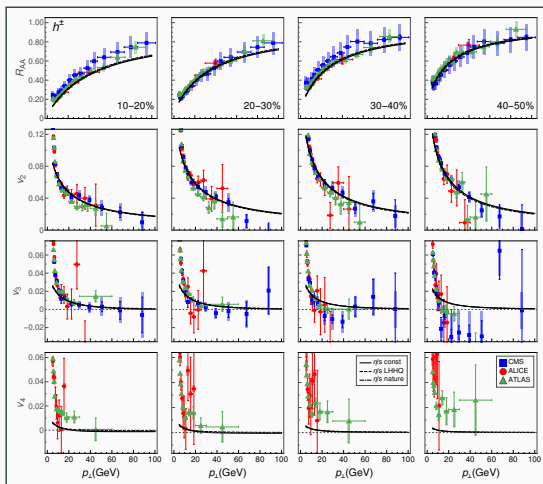
J. E. Bernhard, J. S. Moreland and S. A. Bass, Nature Phys. **15**, no.11, 1113-1117 (2019)

- LHHQ  $\eta/s$

H. Niemi, G. S. Denicol, P. Huovinen, E. Molnar and D. H. Rischke, Phys. Rev. Lett. **106**, 212302 (2011)

# $\eta/s$ parametrization

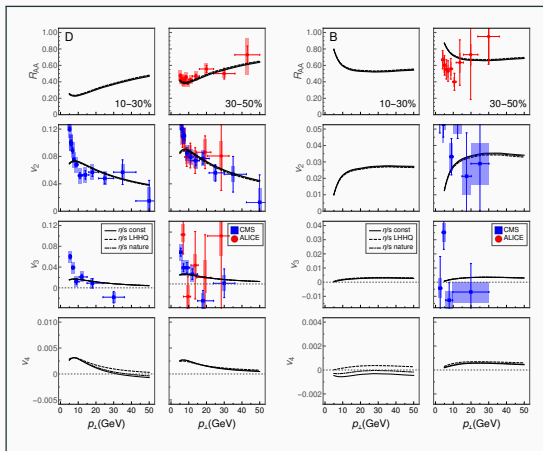
- charged hadrons,  $Pb + Pb$ ,  $\sqrt{s_{NN}} = 5.02 TeV$



high- $p_T$  not sensitive to different  $\eta/s$  as well

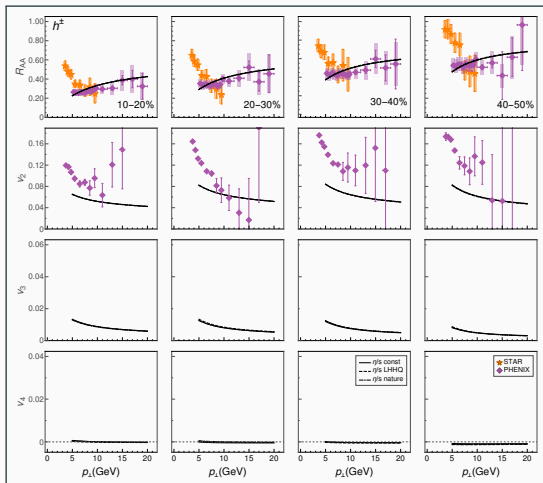


- heavy flavour,  $Pb + Pb$ ,  $\sqrt{s_{NN}} = 5.02 TeV$



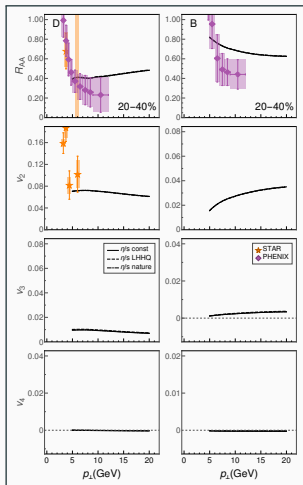
# $\eta/s$ parametrization

- charged hadrons,  $Au + Au$ ,  $\sqrt{s_{NN}} = 200\text{ GeV}$



# $\eta/s$ parametrization

- heavy flavour,  $Au + Au$ ,  $\sqrt{s_{NN}} = 200\text{GeV}$



# Acknowledgements



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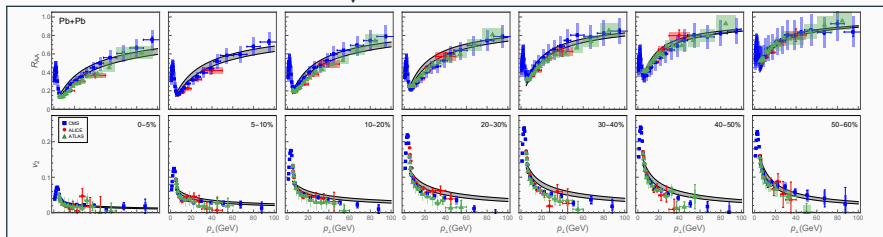


**МИНИСТАРСТВО ПРОСВЕТЕ,  
НАУКЕ И ТЕХНОЛОШКОГ РАЗВОЈА**

Thank you for your attention!

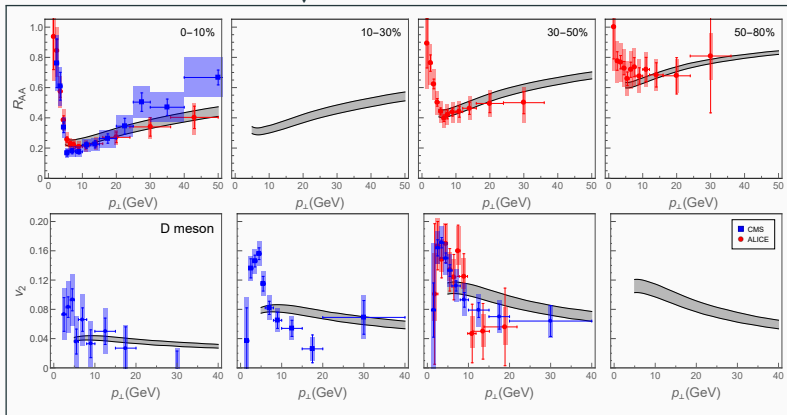
# Backup slides

DREENA-C,  $Pb + Pb$ ,  $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ ,  $h^\pm$



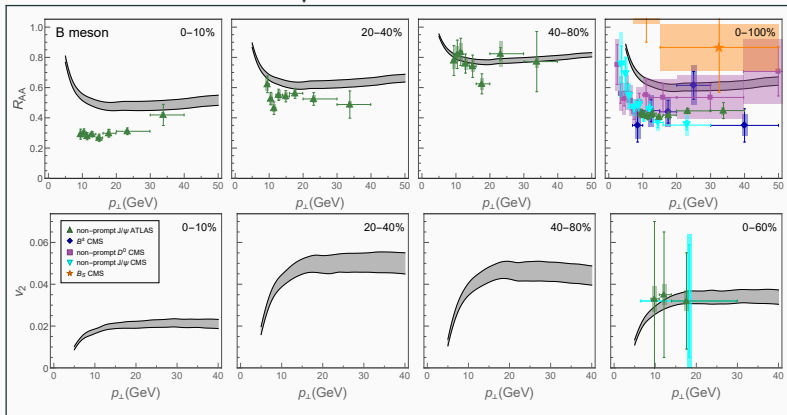
# Backup slides

DREENA-C,  $Pb + Pb$ ,  $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ ,  $D$

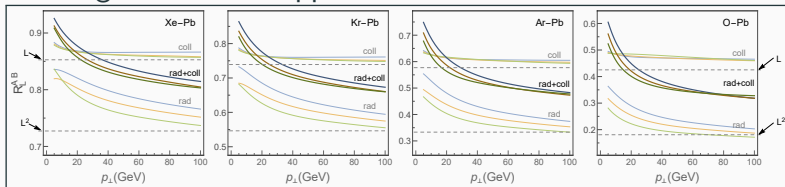


# Backup slides

DREENA-C,  $Pb + Pb$ ,  $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ ,  $B$



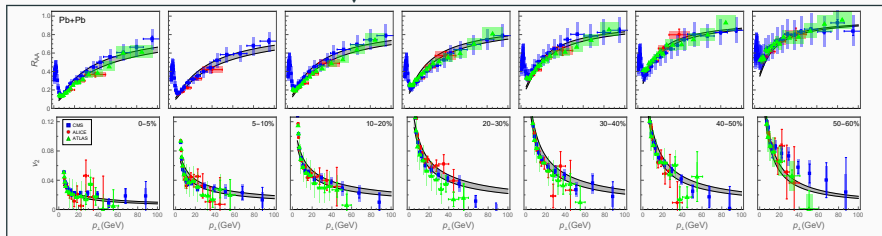
## path-length sensitive suppression ratio





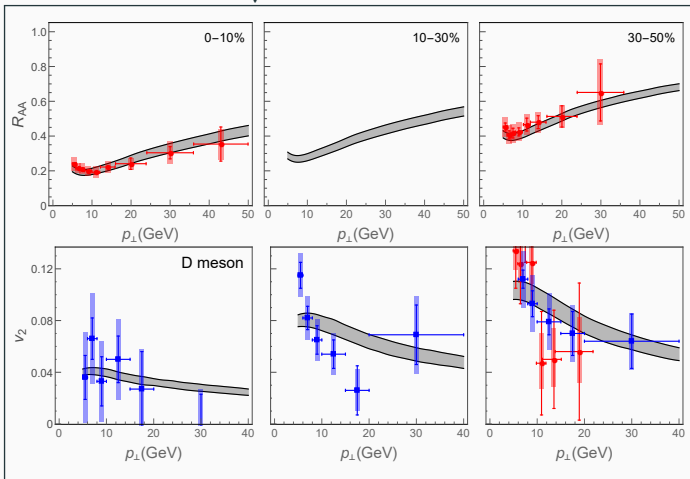
# Backup slides

DREENA-B,  $Pb + Pb$ ,  $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ ,  $h^\pm$



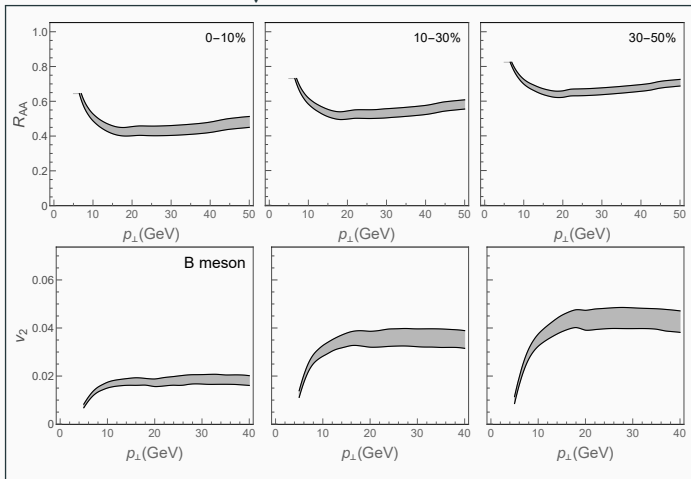
# Backup slides

DREENA-B,  $Pb + Pb$ ,  $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ ,  $D$



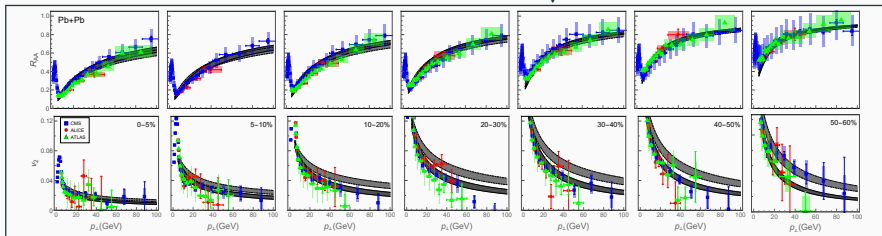
# Backup slides

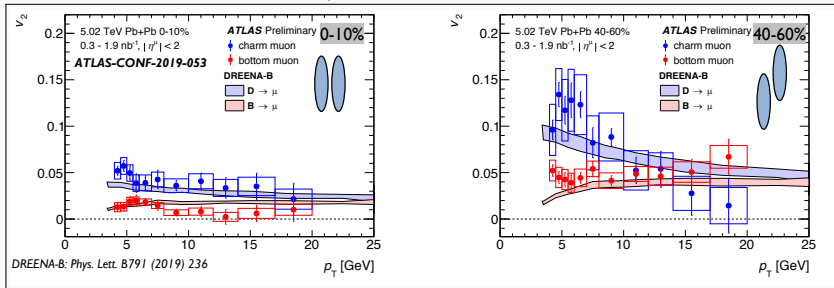
DREENA-B,  $Pb + Pb$ ,  $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ ,  $B$



# Backup slides

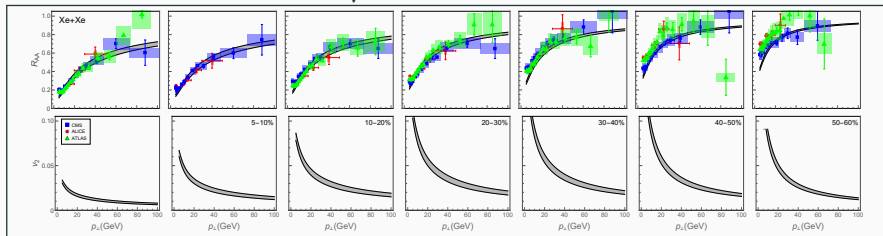
DREENA-B vs DREENA-C,  $Pb + Pb$ ,  $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ ,  $h^\pm$



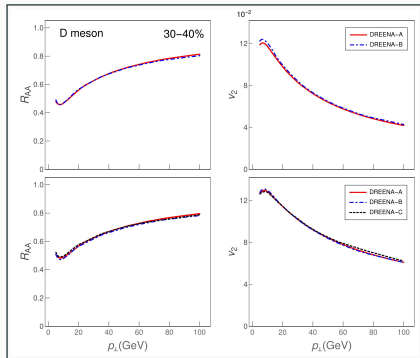
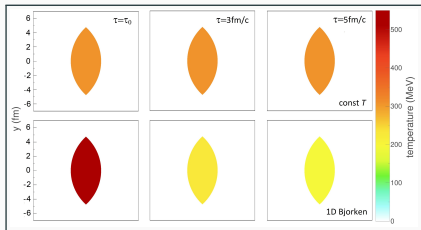
DREENA-B,  $Pb + Pb$ ,  $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ ,  $\mu$ 

# Backup slides

DREENA-B, Xe + Xe,  $\sqrt{s_{NN}} = 5.44 \text{ TeV}$ ,  $h^\pm$



## DREENA-A limits



ebe averaging methods:

$$Q_n = \frac{1}{M} \sum_{j=1}^M e^{in\phi_j} \equiv |v_n| e^{in\Psi_n}$$

$$R_{AA}(p_{\perp}) = \frac{1}{2\pi} \int_0^{2\pi} R_{AA}(p_{\perp}, \phi) d\phi$$

$$q_n^{\text{hard}} = \frac{\frac{1}{2\pi} \int_0^{2\pi} e^{in\phi} R_{AA}(p_{\perp}, \phi) d\phi}{R_{AA}(p_{\perp})}$$

$$v_n^{\text{hard}} = \frac{\frac{1}{2\pi} \int_0^{2\pi} \cos[n(\phi - \Psi_n^{\text{hard}}(p_{\perp}))] R_{AA}(p_{\perp}, \phi) d\phi}{R_{AA}(p_{\perp})}$$

$$\Psi_n^{\text{hard}}(p_{\perp}) = \frac{1}{n} \arctan \left( \frac{\int_0^{2\pi} \sin(n\phi) R_{AA}(p_{\perp}, \phi) d\phi}{\int_0^{2\pi} \cos(n\phi) R_{AA}(p_{\perp}, \phi) d\phi} \right)$$



# Backup slides

ebe averaging methods:

$$v_n^{\text{hard}}\{\text{SP}\} = \frac{\langle \text{Re}(q_n^{\text{hard}}(Q_n)^*) \rangle_{\text{ev}}}{\sqrt{\langle Q_n(Q_n)^* \rangle_{\text{ev}}}} = \frac{\langle |v_n^{\text{hard}}| |v_n| \cos[n(\Psi_n^{\text{hard}}(p_\perp) - \Psi_n)] \rangle_{\text{ev}}}{\sqrt{\langle |v_n|^2 \rangle_{\text{ev}}}}$$

$$v_n\{\text{EP}\} = \langle \langle \cos[n(\phi^{\text{hard}} - \Psi_n)] \rangle \rangle_{\text{ev}} = \langle v_n^{\text{hard}} \cos[n(\Psi_n^{\text{hard}} - \Psi_n)] \rangle_{\text{ev}}$$

$$v_n\{\text{SP}_{\text{ATLAS}}\} = \frac{\text{Re} \langle \langle e^{in\phi} (Q_n^{-|+})^* \rangle \rangle_{\text{ev}}}{\sqrt{\langle Q_n^- (Q_n^+)^* \rangle_{\text{ev}}}}$$

$$v_n\{\text{SP}_{\text{CMS}}\} = \frac{\text{Re} \langle Q_n Q_{nA}^* \rangle_{\text{ev}}}{\sqrt{\frac{\langle Q_{nA} Q_{nB}^* \rangle_{\text{ev}} \langle Q_{nA} Q_{nC}^* \rangle_{\text{ev}}}{\langle Q_{nB} Q_{nC}^* \rangle_{\text{ev}}}}}$$

ebe averaging methods:

- low- $p_{\perp}$

$$\tilde{Q}_n = \sum_{j=1}^M e^{in\phi_j}$$

$$v_n\{2\} = \sqrt{c_n\{2\}}, c_n\{2\} = \langle\langle 2 \rangle\rangle_{\text{ev}}, \langle 2 \rangle = \frac{|\tilde{Q}_n|^2 - M}{W_2}, W_2 = M(M-1)$$

$$v_n\{4\} = \sqrt[4]{-c_n\{4\}}, c_n\{4\} = \langle\langle 4 \rangle\rangle_{\text{ev}} - 2\langle\langle 2 \rangle\rangle_{\text{ev}}^2$$

$$\langle 4 \rangle = \frac{|\tilde{Q}_n|^4 + |\tilde{Q}_{2n}|^2 - 2\text{Re}|\tilde{Q}_{2n}\tilde{Q}_n^*\tilde{Q}_n^*|}{W_4} - 2\frac{2(M-2)|\tilde{Q}_n|^2 - M(M-3)}{W_4}$$

$$W_4 = M(M-1)(M-2)(M-3)$$

the averaging methods:

- high- $p_{\perp}$

$$q_n = \int_0^{2\pi} e^{in\phi} \frac{dN}{dp_{\perp} d\phi} d\phi, m_q = \int_0^{2\pi} \frac{dN}{dp_{\perp} d\phi} d\phi$$

$$W'_2 = m_q M, W'_4 = m_q M(M-1)(M-2)$$

$$\langle 2' \rangle = \frac{q_n \tilde{Q}_n^*}{W'_2}, \langle 4' \rangle = \frac{q_n \tilde{Q}_n \tilde{Q}_n^* \tilde{Q}_n^* - q_n \tilde{Q}_n \tilde{Q}_{2n}^* - 2Mq_n \tilde{Q}_n^* + 2q_n \tilde{Q}_n^*}{W'_4}$$

$$d_n\{2\} = \langle \langle 2' \rangle \rangle_{\text{ev}}, d_n\{4\} = \langle \langle 4' \rangle \rangle_{\text{ev}} - 2\langle \langle 2' \rangle \rangle_{\text{ev}} \langle \langle 2 \rangle \rangle_{\text{ev}}$$

$$v'_n\{2\} = \frac{d_n\{2\}}{\sqrt{c_n\{2\}}}, v'_n\{4\} = -\frac{d_n\{4\}}{(-c_n\{4\})^{3/4}}.$$

# Backup slides

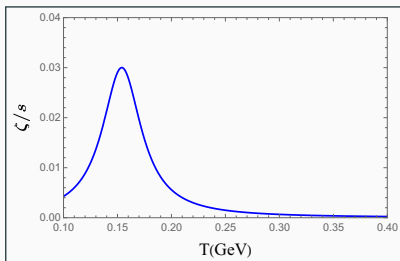
TRENTO parameters

Pb + Pb						
	p	x	n	k	w	d
const	0	7.0	90	1.19	0.5	0.5
LHHQ	0	7.0	80	1.19	0.5	0.5
nature	0	7.0	85	1.19	0.5	0.5

Au + Au						
	p	x	n	k	w	d
const	0	4.23	32	1.19	0.5	0.5
LHHQ	0	4.23	31	1.19	0.5	0.5
nature	0	4.23	32	1.19	0.5	0.5

VISHNU parameters

	$\tau_0$	Edec	VisT0	VisHRG	VisMin	VisSlope	VisCrv	VisBulkT0	VisBulkMax	VisBulkWidth
const	1.0	0.265	0.154	0.15 0.12	0.15 0.12	0.00	1.00	0.183	0.03	0.022
LHHQ	1.0	0.265	0.154	0.04	0.04	3.30	0.00	0.183	0.03	0.022
nature	1.0	0.265	0.154	0.10	0.10	1.11	-0.48	0.183	0.03	0.022



# Backup slides

energy loss:

$$\frac{dE_{col}}{d\tau} = \frac{2C_R}{\pi v^2} \alpha_S(E T) \alpha_S(\mu_E^2(T)) \times$$

$$\int_0^\infty n_{eq}(|\vec{k}|, T) d|\vec{k}| \left( \int_0^{|\vec{k}|/(1+v)} d|\vec{q}| \int_{-v|\vec{q}|}^{v|\vec{q}|} \omega d\omega + \int_{|\vec{k}|/(1+v)}^{|\vec{q}|_{max}} d|\vec{q}| \int_{|\vec{q}|-2|\vec{k}|}^{v|\vec{q}|} \omega d\omega \right) \times$$

$$\left( |\Delta_L(q, T)|^2 \frac{(2|\vec{k}| + \omega)^2 - |\vec{q}|^2}{2} + |\Delta_T(q, T)|^2 \frac{(|\vec{q}|^2 - \omega^2)((2|\vec{k}| + \omega)^2 + |\vec{q}|^2)}{4|\vec{q}|^4} (v^2|\vec{q}|^2 - \omega^2) \right)$$

$$\frac{d^2 N_{rad}}{dx d\tau} = \int \frac{d^2 k}{\pi} \frac{d^2 q}{\pi} \frac{2 C_R C_2(G) T}{x} \frac{\mu_E(T)^2 - \mu_M(T)^2}{(q^2 + \mu_M(T)^2)(q^2 + \mu_E(T)^2)} \frac{\alpha_S(E T) \alpha_S(\frac{\mathbf{k}^2 + \chi(T)}{x})}{\pi}$$

$$\times \frac{(\mathbf{k} + \mathbf{q})}{(\mathbf{k} + \mathbf{q})^2 + \chi(T)} \left( 1 - \cos \left( \frac{(\mathbf{k} + \mathbf{q})^2 + \chi(T)}{xE^+} \tau \right) \right) \left( \frac{(\mathbf{k} + \mathbf{q})}{(\mathbf{k} + \mathbf{q})^2 + \chi(T)} - \frac{\mathbf{k}}{\mathbf{k}^2 + \chi(T)} \right)$$