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ASTRACT

Analysis of STRain Affected CharacTeristics of brittle SC cables

«Critical current evaluation of Nb3Sn samples from m(B) curves by SEM image processing in ANSYS APDL using Space Claim import tools»

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- Introducing the context
- The ASTRACT project
- Samples preparation
- Comparison of critical current and magnetization measurements
- ANSYS in our data analysis
- Results
- Summary

Introducing the contest

The next generation of Accelerator Magnets are targeting 16 T based on Nb3Sn cables. They need Jc up to 1500 A/mm² $@$ 16 T, 4.2 K

FALCOND strand: RRP Ti-doped (162/169)

EUROCIRCOL -> FALCOND model: the first $cos\theta$ dipole with bladder & key assembly technique

The TDR design featured a detachable titanium pole, a SS

 $_{c}^{*}I_{c} = 1355 A/mm^{2}$ (@16T, 4.2K)

Test laboratories

The ASTRACT project

A multidisciplinary approach to investigate the critical current degradation due to transverse strain in Nb3Sn strands.

The first part: pre-HT lamination from 0 to 25% then measuring Ic by V-I (VAMAS) and m(B) VSM

Samples preparation

We rolled 2 meter long strand up to 25% preparing VAMAS and VSM samples for the HT (CNR/SPIN)

VAMAS measured at LASA: $11 - 13$ T @ 4.2K; VSM samples measured at ENEA VSM 18T @ 4.2 K

E Ic: VAMAS vs m(B) \overline{d} ductor by measurements \overline{d} $\frac{1}{\sqrt{2}}$ \sqrt{S} M(B) \sqrt{S} is proposed. The similar general to assume a similar general to assume a similar general to a s $f(x) = f(x)$ $A \rightarrow V S \cap (B)$ \mathbf{r} , see [1], the current density inside the superconductor inside the superconductor inside the superconductor \mathbf{r} can be written as a local contract of the set $\cos(\nabla)$ $\mathsf{A}\mathsf{B}\mathsf{B}$ that the coupling between the sub-elements in a sub-element in a s

Evaluation of the Critical Current Density of
Multifilamentary $Nb₃Sn$ Wires From 1 - The Baumgartner's model Magnetization Measurements lam
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Multifilamentary Nb₂Sn Wires From Magnetization Measurements straight straight sample of the Critical Current Density of the Critical Current Density of the control of the Critical Current Density of the control of the Critical Current Density of the control of the Critical Current

factor, i.e.\ parallel hollow cylinders. This approximation should $\footnotesize \begin{array}{c} \textit{Baumgartner et al. 2012, DOI:10.1109/TASC.2011.2175350} \\ \textit{Baumgartner et al. 2012, DOI:10.1109/TASC.2011.2175350} \end{array}$

• Bean model, Jc constant over the sample ling between sub-elements in a strand: Baumgartner et al. 2012, DOI:10.1109/1ASC.2011.2175350
• Bean model, Jc constant over the sample and assuming no coupling between sub-elements in a strand: i model, Jc constant over the sample and assuming no coupling between sub-elements in a strand:
1 EXAMINED TYPES OF MULTIFILAMENTARY WIRES

1 – The Baumgartner's model

Ic: VAMAS vs m(B) *2 – Shape factor numerical calculation* \sim superconductor $\mathbf{A} \cap \mathbf{A}$ \Box and \Box if the sub-element cross section \Box sumarical calculation The momentum superconduction s

 dS

 $dS =$

We get the shape factor F by a first order numerical integration. For the i-th bundle:

Bean's model and current conservation $\Rightarrow \exists a \text{ line } \parallel to \vec{B} \text{ leading to domains } D_+ \text{ and } D_- \text{ such that }$ **Decute we deleted equation instead the superconductor** \vec{R} density integrating \vec{R} and \vec{R} and \vec{R} can be calculated by $dS = \int dS$ D icularly to domains D_+ and D_- such that \int_{D_+}

the
.. D_{+} and D_{-} > this defines the domains of integration for the magnetic moment We use the convention B || to the lateral side of the image. Starting from the bundle $\frac{B}{2}$ CM we look iteratively for the line parallel to B splitting the bundle in two equal areas $\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$

$$
\vec{m} = J_{\rm c} \cdot \frac{1}{2} \int \vec{r} \times \vec{e}_J(\vec{r}) d^3 r, \qquad \sum_{m_y^i = \frac{1}{2} J_{\rm c} L \left(\int_{D_{\perp}^i} y dx dy - \int_{D_{\perp}^i} y dx dy \right) = 0
$$
\n
$$
m_y^i = \frac{1}{2} J_{\rm c} L \left(\int_{D_{\perp}^i} x dx dy - \int_{D_{\perp}^i} x dx dy \right) = J_{\rm c} L F_y^i
$$

 $\left\{f_{w} \equiv \int \frac{xdxdy}{dx} \right\} = I L F^{i}$ we measure the intensity of \vec{m} We sum over the N bundles and

 $m_{irr} = J_c L \sum_{i=1}^{N} F_y^i = J_c L F$

For the i-th real bundle, meshing the domain D_+ and D_- with N_+ and N₋ elements we can approximate the shape factor as

Table 2. The shape factor evaluation for perfect bundles: theoretical value f vs numerical value F. $R_o = 29.6 \,\mu m$; $R_i = 18 \,\mu m$; $b_o = R_o; b_i = R_i; a_o = 26.2 \mu m; a_i = 13.8 \mu m$

Ic: VAMAS vs m(B) *3 – SEM images elaboration*

We averaged around 10 SEM images from different samples, in different position of the wire, per any kind of lamination.

 \sum ImageJ

We developed a elaboration protocol to avoid operator's dependencies

Ic: VAMAS vs m(B) *4 –* ANSYS in our data analysis

F and f are the numerical and analytical shape factor averaged over the SEM images; S_{SC} is the mean effective SC section

Ic: VAMAS vs m(B) *4 –* Methodology

- From m(B) -> $I_{c;VSM} = \frac{m_{irr}(B)}{I_F}$ $\frac{d_{irr}(B)}{d_{F}}S^{L}_{SC}$ with errors $\frac{\delta I_{c;VSM}}{I_{c;VSM}}$ $I_{c;VSM}$ $=\int \left(\frac{\delta m_{irr}}{m}\right)$ m_{irr} 2 + $\left(\frac{\delta l}{l}\right)$ ι 2 + $\left(\frac{\delta F}{r}\right)$ \overline{F} $+\frac{\delta S_{SC}}{s}$ $S_{\mathcal{S}\mathcal{C}}$ $\overline{2}$
- From (V,I) data -> Recovering I_c^{VAMAS} by fitting $V = V_c \left(\frac{I}{I_c} \right)$ I_c n
where $V_c = Voltag$ e taps lenght x $10 \frac{\mu V}{cm}$.
- I_c^{VAMAS} curves have to be corrected for self-field effects

• We scale original direct lc to VSM electrical criteria:
$$
I_c = I_c^{VAMAS} \left(\frac{E_{VSM}}{10 \,\mu V/cm}\right)^{\frac{1}{n}}
$$

where $E_{VSM} \cong \frac{D}{2} \vec{B} \cong 2 \div 4 \, 10^{-3} \frac{\mu V}{cm}$

• For both VAMAS and VSM measurements we claim degradation if normalized Ic $\left(\frac{I_C^{rolled}}{I_V^{V}}\right)$ I_c^{V} \overline{v} at same B and T) is less than 95%

VAMAS results

- The red curve is the one used to simulate the Falcon Dipole wire $(J_c(4.2 K, 16 T) = 1300 A);$
- The data are not self field corrected
- Virgin $1 = \text{vir}$ gin wire $14/10/2021$
- Virgin $2 = FD_V_A$
- Virgin $3 = FD$ V C
- **Rolled 20% was too much unstable to achieve Ic**

NORMALIZED CRITICAL CURRENT TO NO ROLLED SAMPLE VAMAS DATA

VAMAS vs VSM: 0% rolled

rel.diff = (VAMAS – VSM)/VAMAS 1

*t = $|VAMAS - VSM|/\delta VSM$

VAMAS vs VSM: 10 and 15 % rolled

 14

VAMAS vs VSM: 20 and 25 % rolled

WARNING: no L20 VAMAS data available -> compared to L15 VAMAS

Summary

- We presented a way to use ANSYS to extract critical current data from magnetic moments curves.
- Defining a database of strand SEM images is possible to achieve Ic values independently from transport data
- Testing the compatibility between VAMAS and VMS techniques requires to scale VAMAS data to the E_c^{VSM}
- Our method is compatible with transport data at 0, 10 and 15% of rolling.
- Both VAMAS and VSM show no critical current degradation due to rolling up to 25%. VSM normalized Ic seems to improve with a boost in performance at 25% of rolling (more than 115% Ic). This effect has been already observed in RRP rolled samples as a consequence of bundles merging and it could lead to a falling of our working hypothesis.

In the next future we will

- test the method respect to other $Nb₃Sn$ strand layouts or SC material as BSSCO
- improving the ANSYS code adding magnetic simulations
- find a way to determine the shape factor in case of merged bundles
- setup an experimental method to measure the n-value from m(B), which would allow to derive Jc from m(B) at any value of Ec

Thanks for your attention!