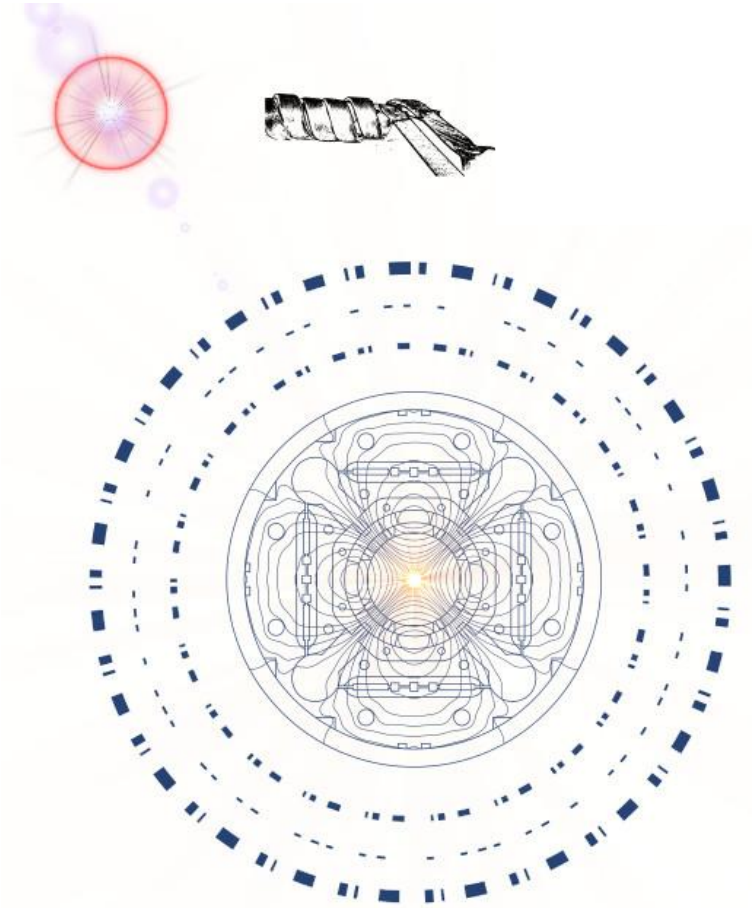




EXPERIMENTAL AND NUMERICAL ASSESSMENT OF THE ELECTRO-MECHANICAL LIMITS OF Nb_3Sn ACCELERATOR MAGNETS

*Presented at CHATS on Applied Superconductivity.
5th May 2023*

J. Ferradas Troitino, B. Bordini, A. Devred, P. Ferracin, S. Izquierdo Bermudez,
F.J. Mangiarotti, A. Milanese, C. Senatore, E. Todesco and G. Vallone



Outline



3D rendered view of a MQXF short model magnet, whose analysis is treated in the slides.

- Introduction / motivation
- Methodology
- Results at the strand level
- Results at the magnet level
- Comparison
- Conclusions

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Nb₃Sn accelerator magnets

On the search for higher magnetic fields in accelerator magnets

Presently, Nb₃Sn superconductors are the preferred solution for the windings of high-field accelerator magnets:

- Performance above Nb-Ti limits.
- Proven large-scale industrial production.

State-of-the-art Nb₃Sn wires surpass the values of $J_c > 1000 \text{ A/mm}^2$ (SC area) at 16 T and 4.2 K.

However, these appealing properties are impaired by:

The strain sensitivity and **brittle** nature of the superconductor.

The determination of sound electro-mechanical limits for Nb₃Sn technology becomes of paramount importance!

High-field accelerator units

$$\mathbf{j} = \mathbf{j} \times \mathbf{B}$$

Magnets are intrinsically characterized by the large electro-magnetic forces arising from the combination of the required high current densities, J , and produced magnetic fields, B .



Nb₃Sn Rutherford cable. Source: [CERN]

A bit of history and motivation for our work

During the past years, our community has tried to shed some light on the electro-mechanical properties of Nb₃Sn superconductors.

- To establish safe strain/stress limits for operation.

- Developing the necessary mathematical background to describe the strain effects on the superconducting properties.

Laboratory-scale experiments

Exploring the electro-mechanical response of single wires, up to full cables, when subjected to mechanical loads.



- Back in the 1980's by Ekin and co-workers
[1] – [2]
- **University of Geneva**
[3] – [11]
- **CERN**
[12] – [16]
- University of Twente
[17] – [18]
- NHMFL
[19] – [21]
- Fermilab
[22] – [24]

Direct assessment on magnet configuration

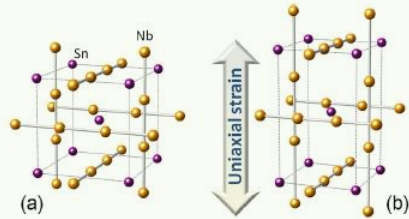
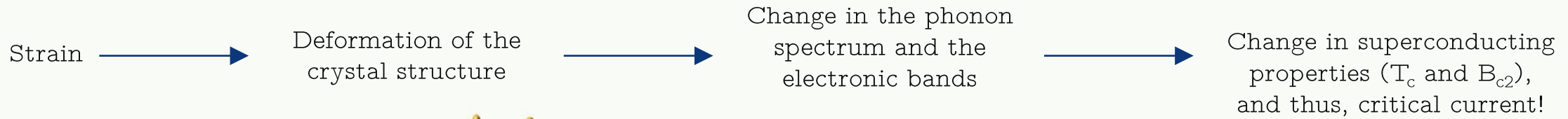
Magnet performance studied as a function of different conductor stress levels.



- Lawrence Berkeley National Lab. (LBNL)
TQ magnet
[25]
- CERN
SMC magnet
[26 - 27]
MQXF magnet
[28]

Overview on Nb₃Sn

- Strain sensitive and **brittle** superconductor



G. De Marzi et al. arXiv:1210.3705, 2012.

1

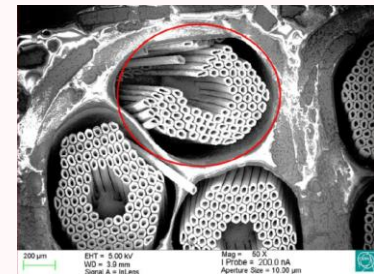
Brittle fracture of the superconducting phase



Reduction of the current-carrying surface



Reduction in conductor performance / critical current!



A. Moros et al., IEE TAS 2023

2

Overview on Nb₃Sn

Critical surface → *Exponential scaling law proposed by Bordini, Bottura et al. [29] – [30]*

Initially postulated to predict the effects of mechanical strain in axially loaded conductor → strain region!

I_1 - first invariant

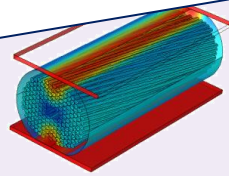
For magnet design, instrumental to apply our understanding acquired at the conductor level.

- Are conductor tests applicable to magnet configuration?
- Can we provide reliable limits for the technology?

Based on experimental results and numerical modelling, we will try to provide answers to these fundamental questions **at the magnet scale.**

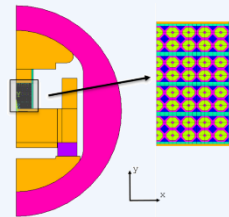
- Transversally

- 2D and 3D detailed models [31]-[32] Bordini, Cattabiani, Baffari et al. [5] Senatore, Bagni, Calzolaio et al. [33] Chiesa, Wang et al.



Cattabiani, Baffari et al.

Transversally loaded cable stacks (macro-scale model)



Vallone, Ferracin, Bordini et al.

[34] – [35] G. Vallone et al.

It can be applied to a 2D magnet cross-section (already published).

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Methodology

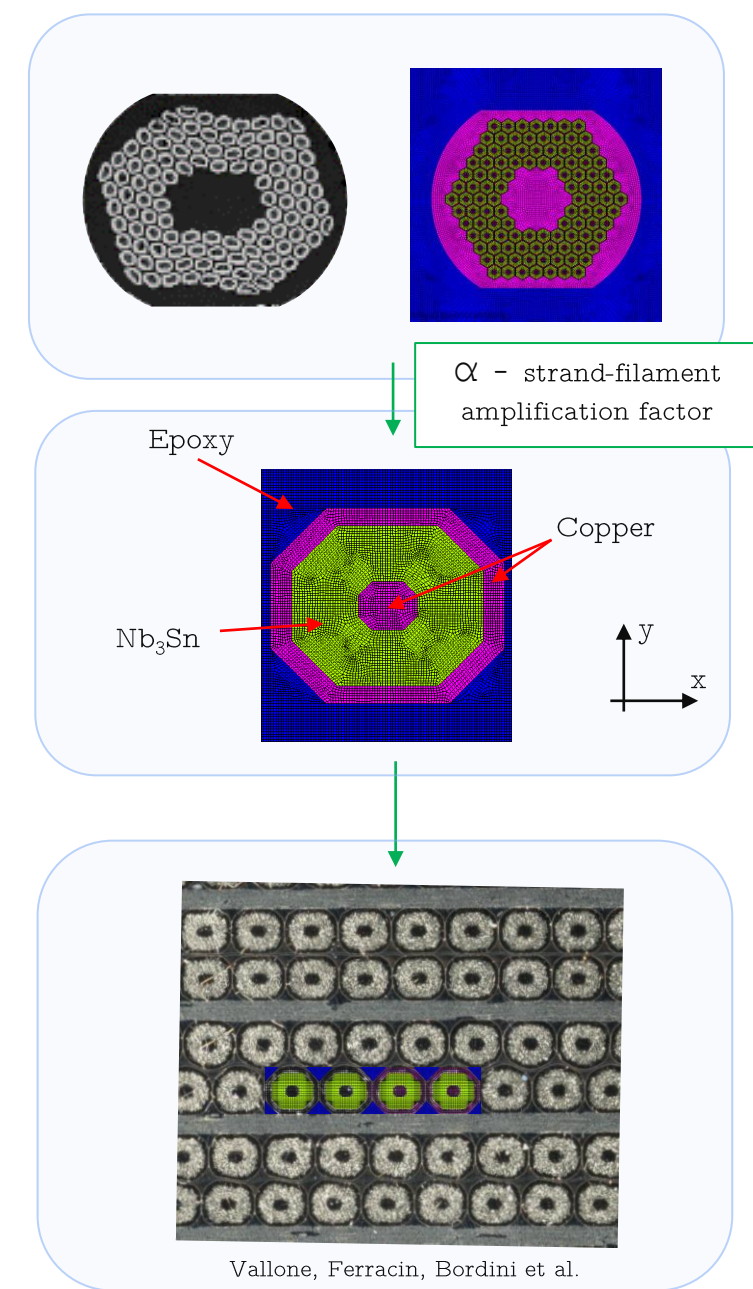
- Domain homogenized down to an octagonal strand representation (Cu, Nb₃Sn).
- 2D static structural analysis. Material properties from literature. Described in details in [34] – [35].
- Boundary conditions play a critical role. Plane stress used in the simulations shown in these slides.
- The strand critical current is computed averaging the critical current over the strand area (perfect current redistribution within the strand).
- The horizontal and vertical strain in the Nb₃Sn area are amplified with a constant factor:

$$\varepsilon_i \rightarrow \alpha \varepsilon_i$$

Amplification when scaling the model to the filament level (strand-to-filament amplification factor).

- Furthermore, the approach considers the effect of the thermal contraction mismatch on strand constituents (from the reaction to cryogenic temperatures). This “pre-compression strain tensor” is added to the one computed by the mechanical model:

$$\begin{aligned}\varepsilon_{zz} &= \varepsilon'_{zz} + \varepsilon_{l0} \\ \varepsilon_{xx} &= \varepsilon'_{xx} - \nu \varepsilon_{l0} \\ \varepsilon_{yy} &= \varepsilon'_{yy} - \nu \varepsilon_{l0}\end{aligned}$$

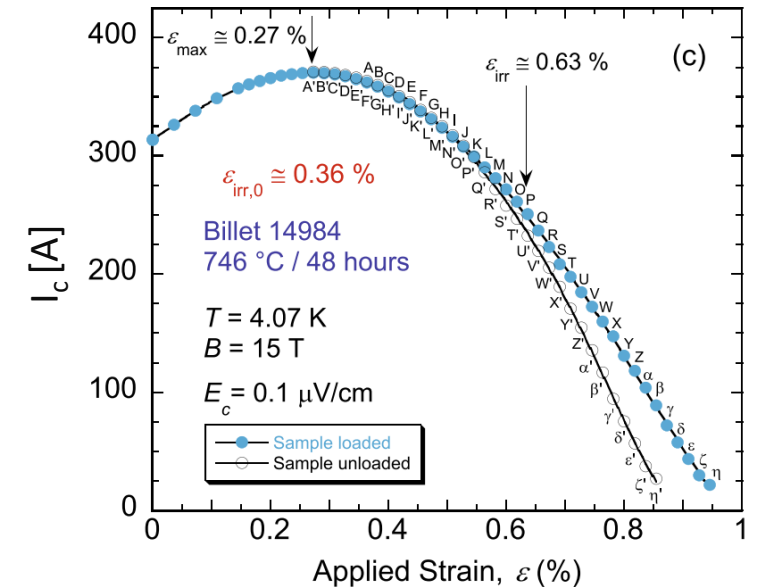


Methodology

- The following fitting parameters are needed to feed the scaling law:
 - Non-strain related: C_0 , $B_{c2}(0,0)$, $T_c(0)$
 - Strain related: C_1 , ϵ_{l0}

Requirement: Characterization of the desired strand using a standard WASP measurement.

- **Recall:** This approach is valid when the mechanism governing the I_c change is a modification on the strain state of the superconducting phase:
 - Elastic (recovered when the load is removed)
 - Plastic (mainly coming from matrix plasticization)
- A failure criteria is needed (separately) to deal with the brittle fracture of the Nb_3Sn phase.
 - Proposal from Vallone et al. in <https://indico.cern.ch/event/1177999/>
 - Indications of clear differences depending on the test configuration (4-wall fully supported strand, 2-wall)
 - Not covered in these slides!



Plot from N. Cheggour et al. Scientific Reports (2019)

Outline



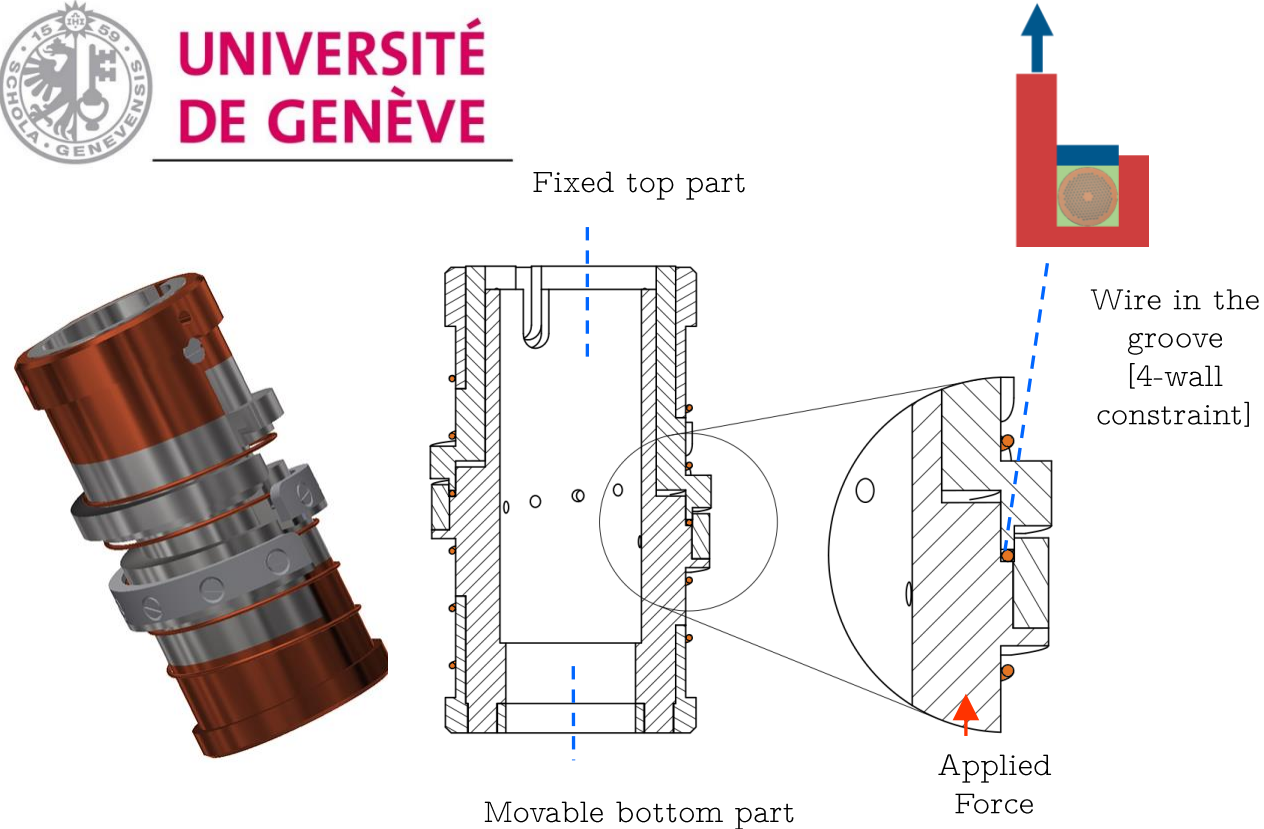
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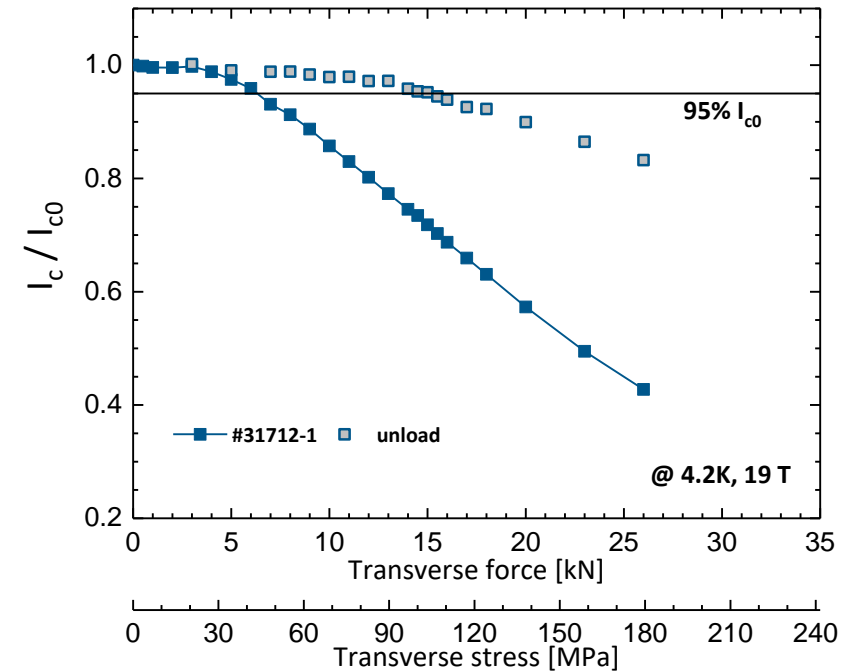
Ic vs. transverse stress on single wires



UNIVERSITÉ
DE GENÈVE



Example on measurement results



C-WASP – I_c vs transverse stress measurements

Trying to reproduce the operating conditions of accelerator magnets.

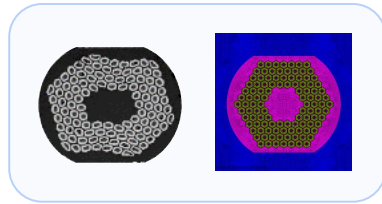
To compare results among wires, a conventional irreversible threshold is generally defined as the stress level leading to a permanent 5% loss of I_c .

Plot shows results at 19 T.

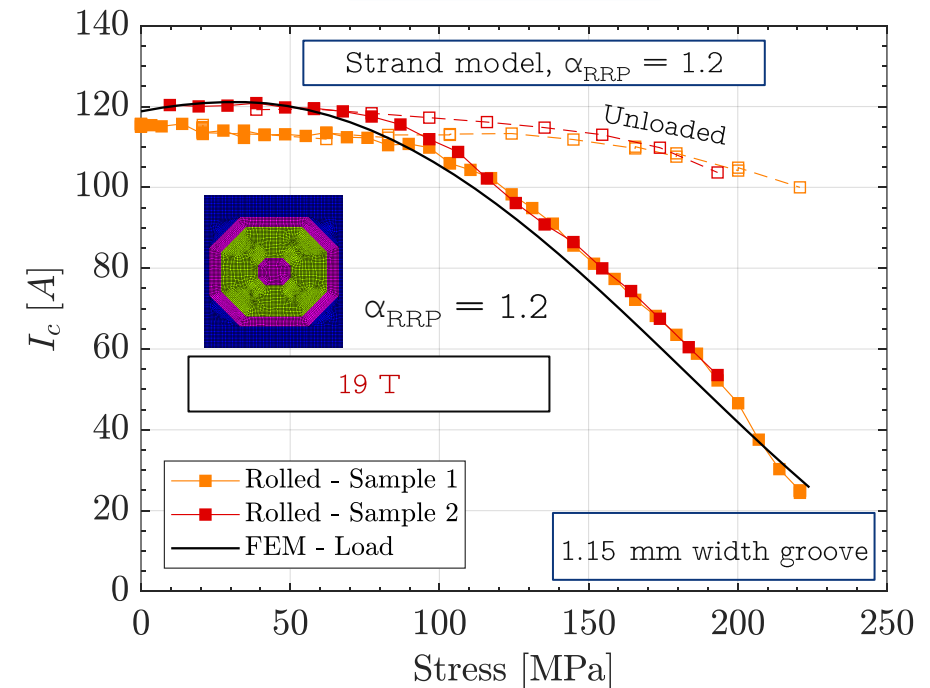
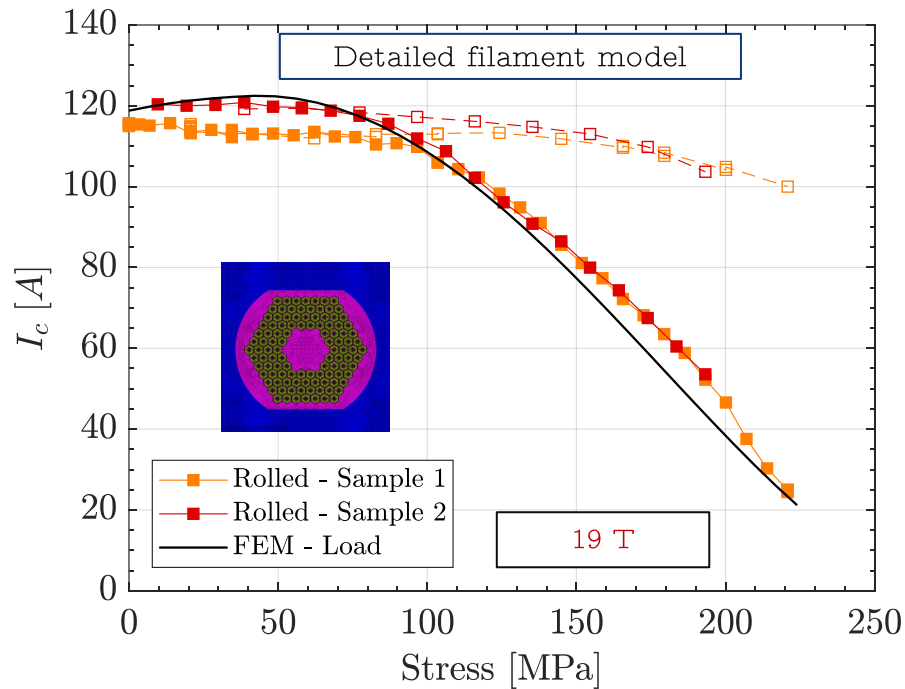
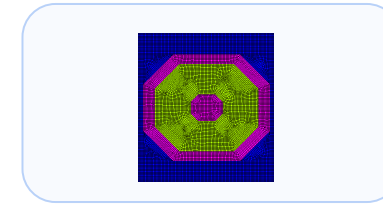
Ic vs. transverse stress on single wires



Amplification factor assessment:



$$\alpha_{RRP} = 1.2$$



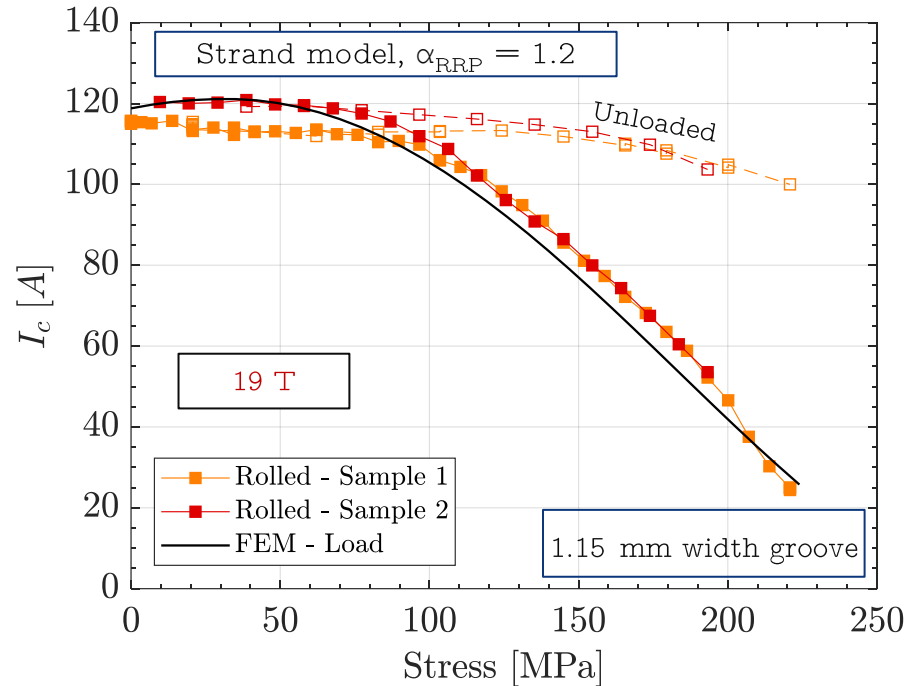
0.85 mm RRP 108/127 15% rolled (HL-LHC MQXF quadrupole)

Ic vs. transverse stress on single wires

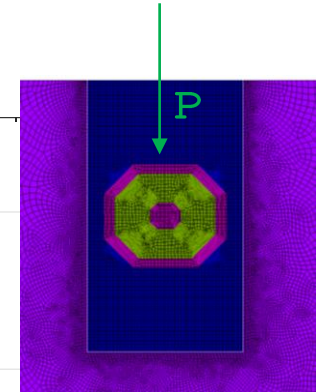
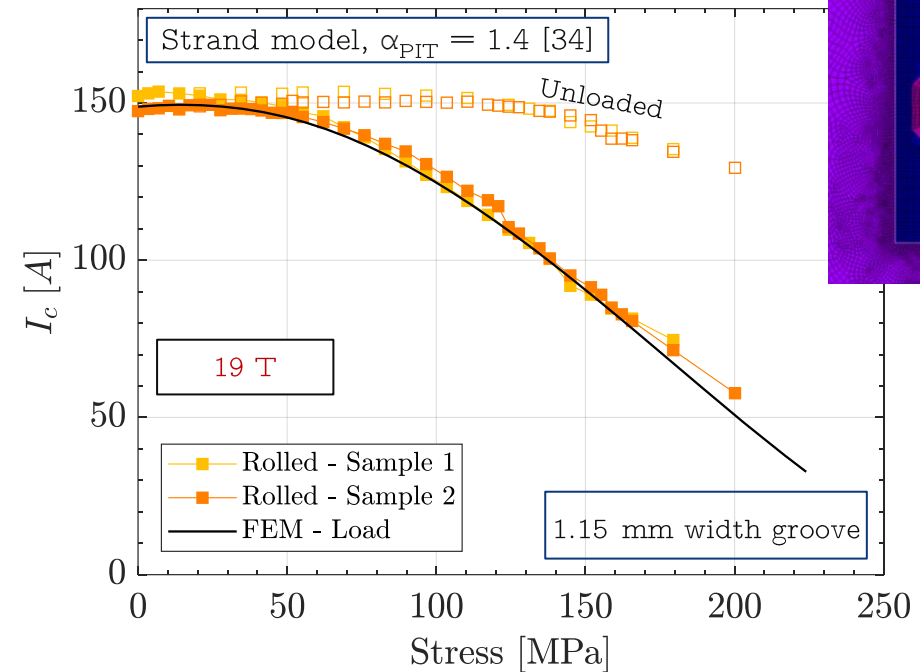


Simulation and experimental results:

0.85 mm RRP 108/127 15% rolled
(HL-LHC MQXF quadrupole)

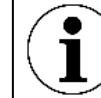


1 mm PIT 192 15% rolled
(FRESCA2 dipole)



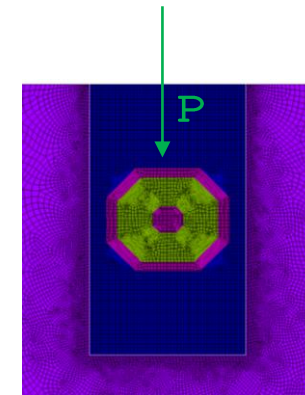
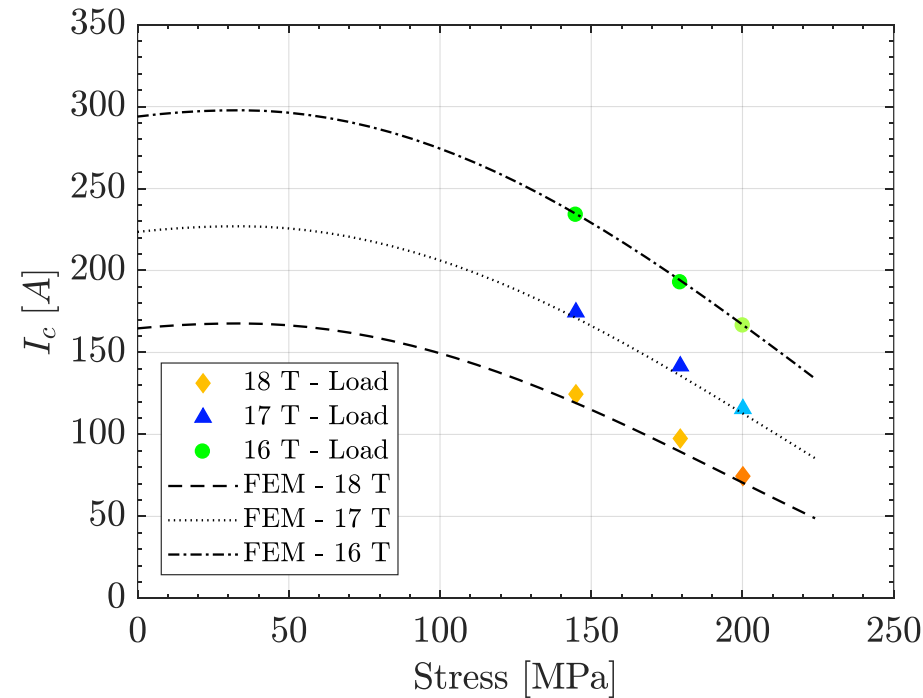
- Model catches the Ic reduction when the transverse load is applied for both wire technologies.
 $\alpha_{PIT} = 1.4$, consistent with cable model [34] → calculated by direct fitting of experimental data. Larger strain sensitivity than RRP (core, etc [32]).
- Simulation results after force unload are inferred using the experimental data. Model is not predictive, but it can still provide I_{c_unload} when coupled to a C-WASP measurement (additional slides). Development on-going.

I_c vs. transverse stress on single wires



Simulation and experimental results:

0.85 mm RRP 108/127 15% rolled
(HL-LHC MQXF quadrupole)



Model validated using the complete magnetic field range available experimentally.

Wrap up on tests on single wires

Important information for magnet application:

- As already shown by previously published models, the exponential scaling law is able to correctly describe the reduction in I_c due to applied transverse stresses as high as 200 MPa.

Indication that the dominating mechanism is the strain effect on B_{c2} and T_c

For this test configuration (4-wall)!

- Independent confirmation on three different wire types has been recently published by Senatore et al. [11] - detailed analysis of the experimental $B_{c2}(\epsilon)$ evolution.

RRP up to ~ 250 MPa (anvil applied normal stress)

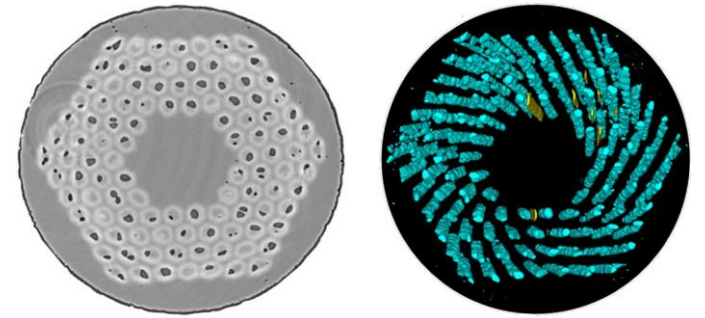
PIT up to ~ 200 MPa (anvil applied normal stress)

- Post-mortem inspection of C-WASP test samples revealed the absence of transversal cracks sectioning the Nb_3Sn filaments [8] – [10].
- Finally, consistent results found on PIT **cable** tests, where the onset of cracks was identified at around 180 – 200 MPa [12] – [14].

Bladder and key set-up.

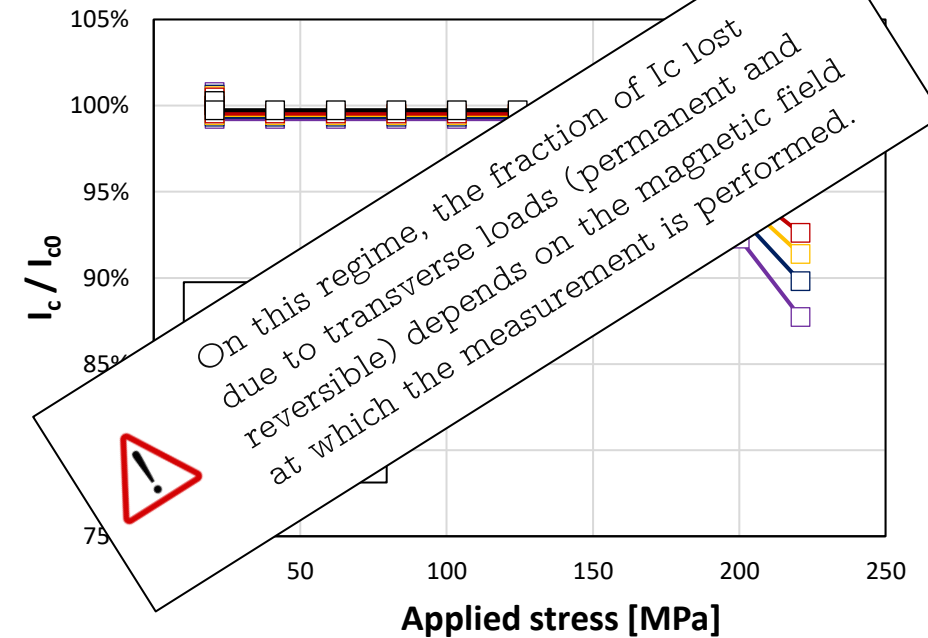


The behavior can be successfully reproduced using the scaling law (already shown by previous studies). This time, also with a simplified strand representation!



Bagni, et al. SUST - 2022

0.85 mm RRP 108/127 15% rolled after force unload (HL-LHC MQXF quadrupole)



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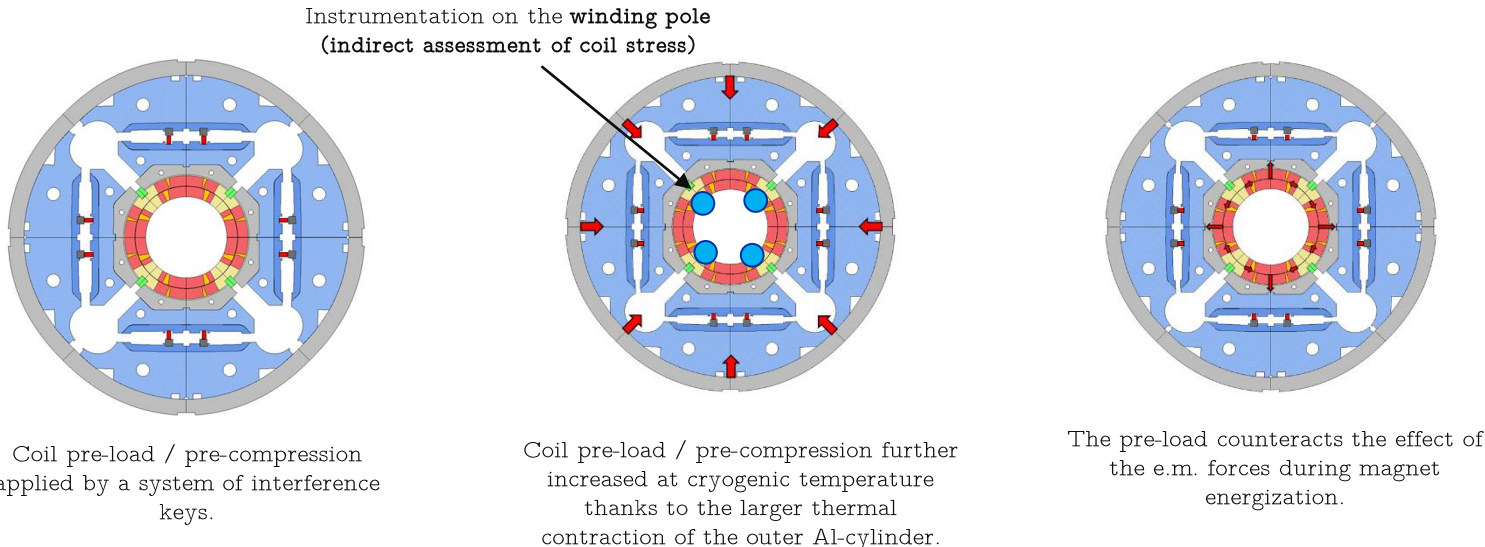
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Electro-mechanical limits on magnet configuration

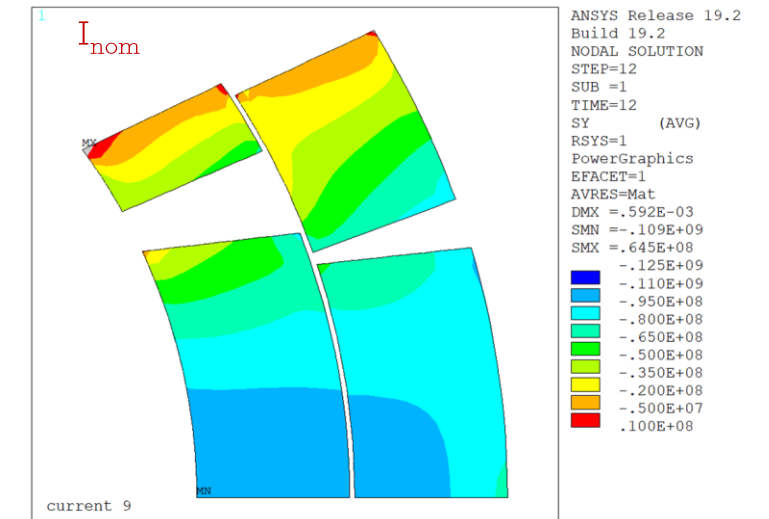
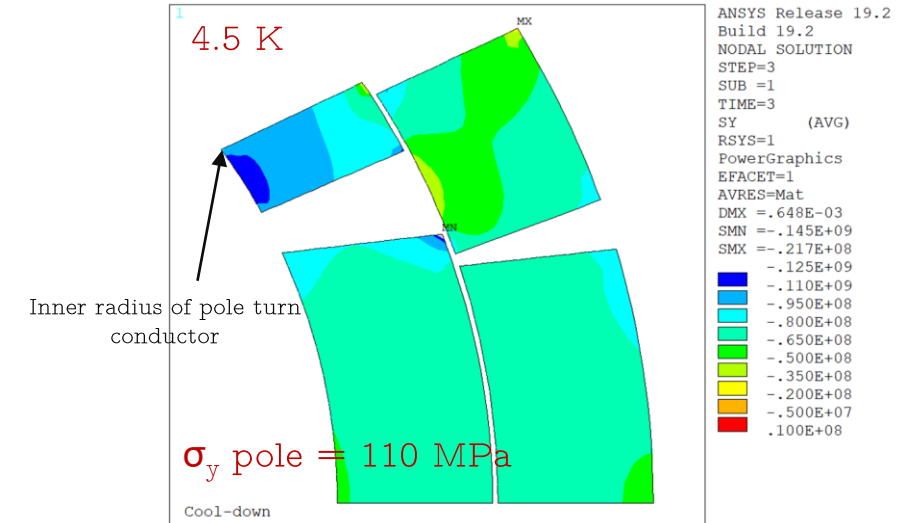
HL-LHC MQXF project:

One short model magnets is being used to explore the electro-mechanical limits of the conductor in magnet configuration.

- MQXFS7 – 2 coils with PIT and 2 coils RRP conductor.
- Magnet pre-load gradually increased searching for signs of performance reduction. Similar test as the one reported by H. Felice et al. [25].
- **Pre-load** level at cold **estimated** using mechanical **instrumentation** placed on the **winding pole** (representative of stress in inner radius pole turn conductor, standard FE model).



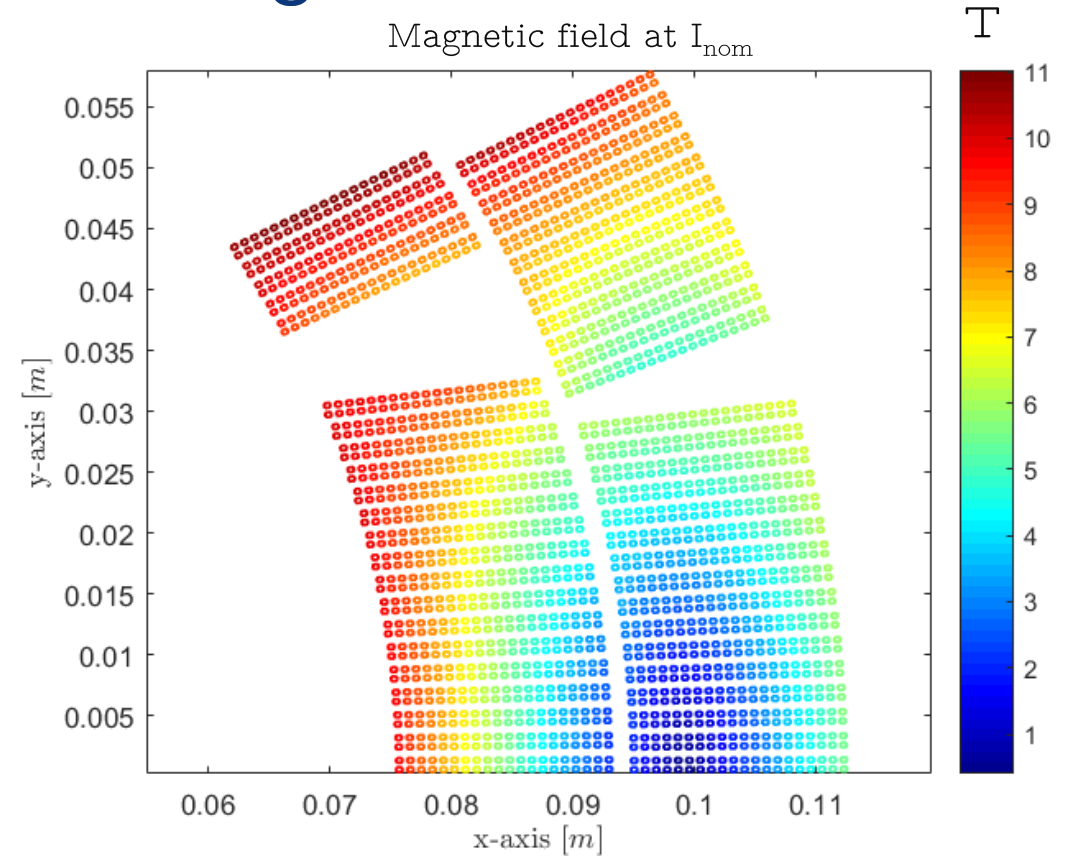
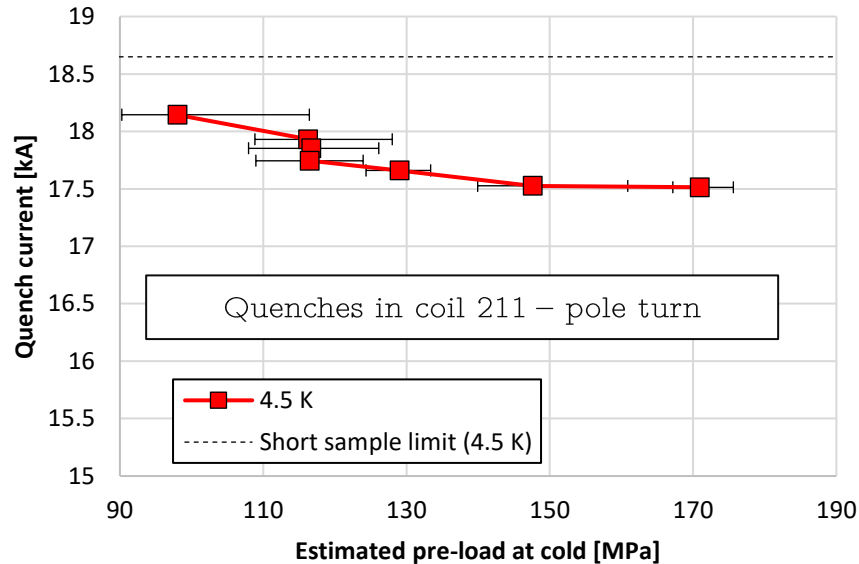
Reference pre-load. Standard FE model Coil azimuthal / transverse stress



Electro-mechanical limits on magnet configuration

Results on MQXFS7 experiment so-far:

- Critical current characterization on extracted strands for each coil. Coil 211 (PIT) being the limiting coil according to these measurements.
- At 4.5 K, the magnet quenches consistently in the inner layer pole turn of coil 211.
 - Peak field, low stress during powering
 - Close to the peak stress region at cryogenic temperature

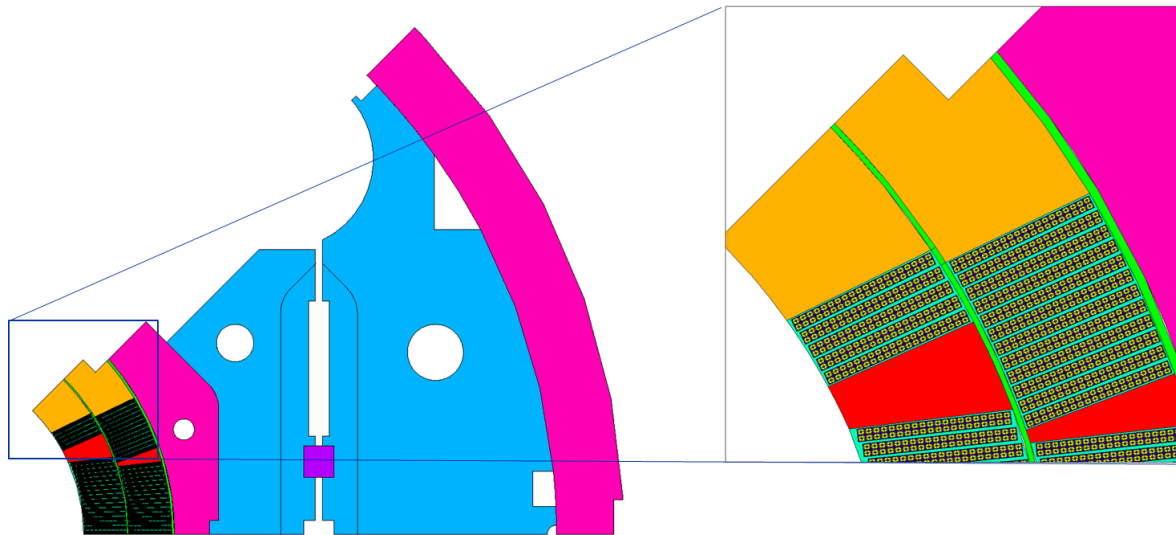


Small decrease in quench current from 100 MPa (97% of SS) to 150 MPa (94% of SS), more assembly iterations to come!

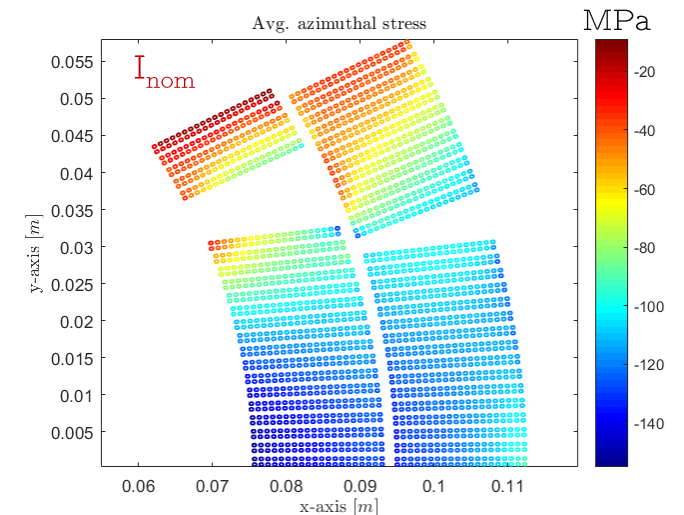
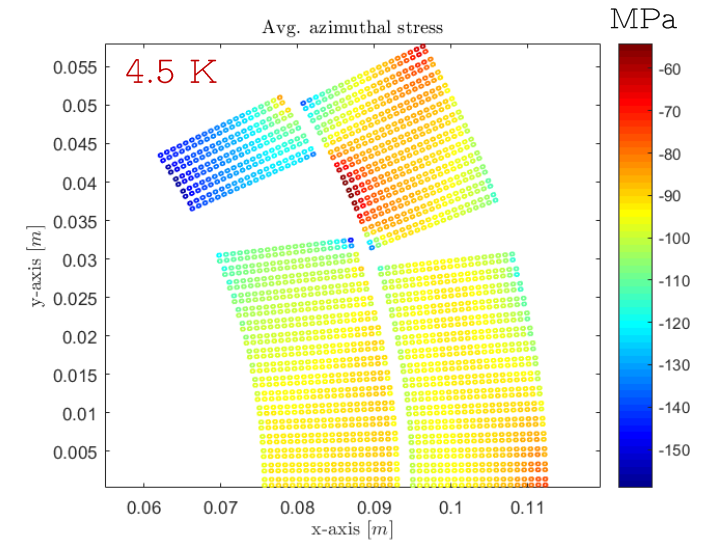
Electro-mechanical limits on magnet configuration

Simulation set-up:

- MQXFS7 experiment modelled using our proposed methodology. Model validated mechanically in [35]. Recall: material properties from literature.
- Strain related fitting parameters (including expected I_c decrease due to permanent strain effects): extracted from available measurements at UNIGE, 1 mm PIT192 strand (**not exactly the same**). Non-strain related: coil extracted strands.
- For each assembly iteration, the pre-load is adjusted to match the experimental results (**we match the indirect measurements at the winding pole**).
- Quench current defined as the point where the critical surface is crossed.



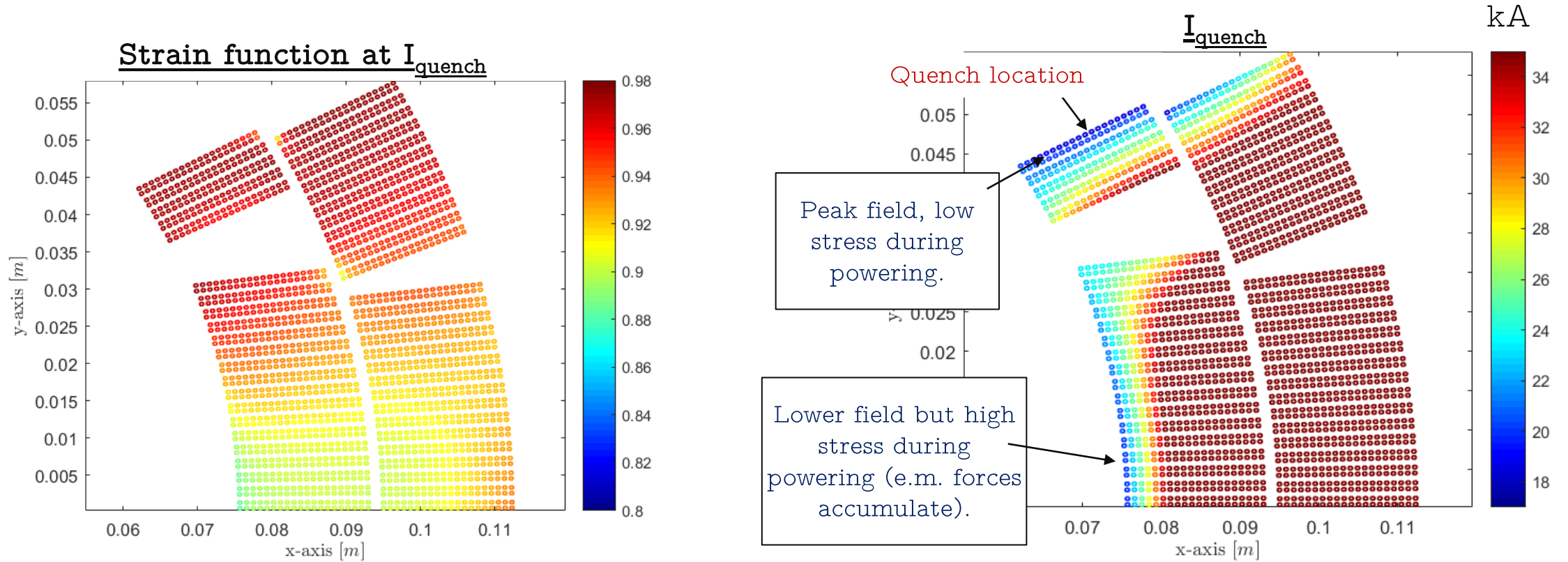
Reference pre-load (Nb_3Sn region)



Electro-mechanical limits on magnet configuration

Results:

Reference pre-load (4.5 K / 110 MPa)



Following the magnet design, the peak field region during powering is found in the pole turn block, which unloads under the action of the electro-magnetic forces. Peak field and highest stress area during powering disentangled!

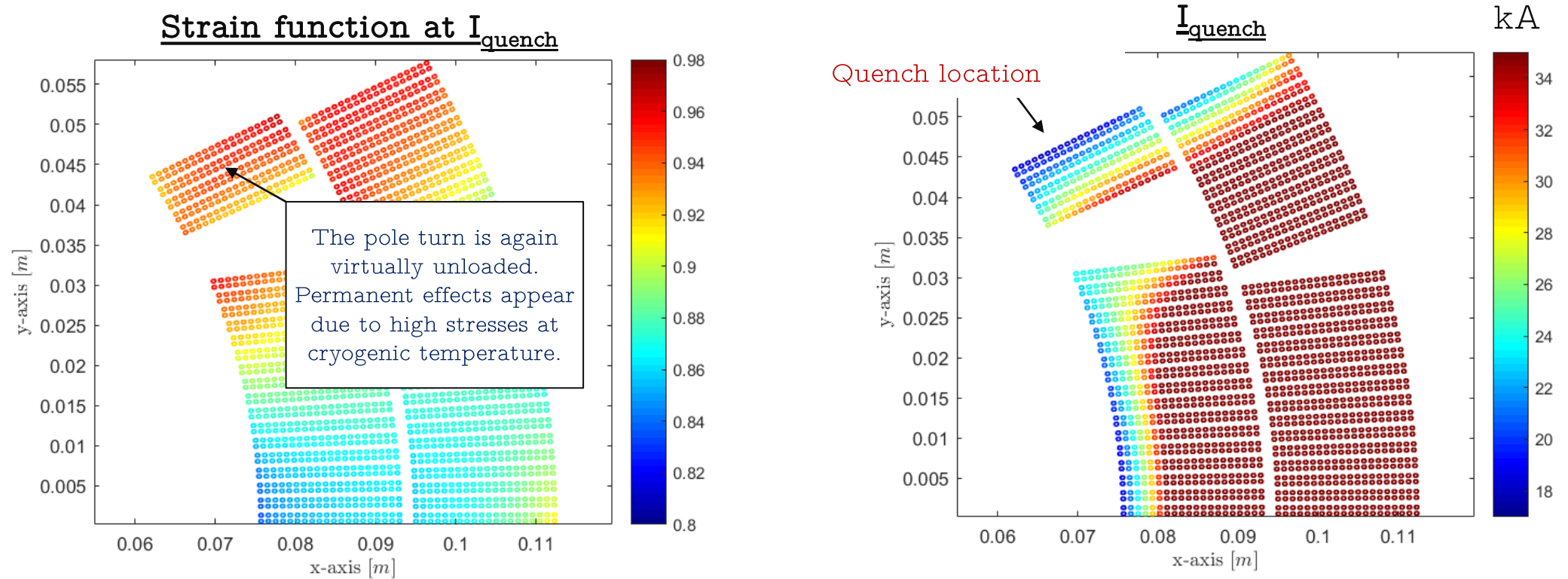
$$B_{p \text{ mid-plane (high stress)}} / B_{p \text{ pole turn (low stress)}} = 0.86$$

$$\text{At } I_{\text{nominal}} \rightarrow B_{p \text{ mid-plane (high stress)}} - B_{p \text{ pole turn (low stress)}} \sim 1.5 \text{ T}$$

Electro-mechanical limits on magnet configuration

Results:

200 MPa pre-load (4.5 K)

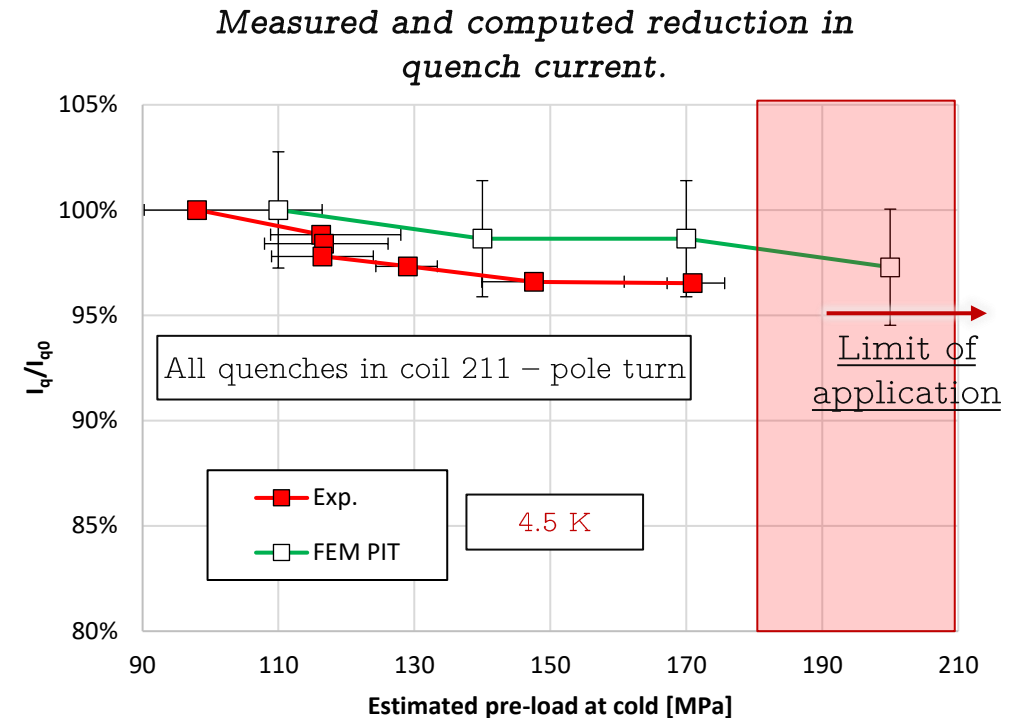


In this particular case for MQXF (4.5 K and strand characteristics), when the magnet is “pre-loaded to the extreme” (200 MPa), a reduction of the short sample limit still appears in the pole turn block due to permanent effects originated by large stresses after cool-down. We do not “jump” to the high stress region (lower field).

Electro-mechanical limits on magnet configuration

Results:

- The magnet strand model (fed with the data from transverse compressive strand tests) shows a similar tendency than the (so-far) measured reduction in quench current.
 - Note that our input data is not strictly the one for the MQXFS7 strand.
 - Experimental and simulated reduction in the same order of magnitude (< 5 %).
- Disregarding the absolute value of $I_{\text{quench}}/I_0 \rightarrow$ Quench happens in the peak field pole turn conductor, which is unloaded during magnet ramping. Same **result predicted** by the **model**.
- The small short sample reduction with increasing magnet pre-load in this case is thus due to permanent effects arising from the large stress levels at cold in that region.
 - From conductor tests: reduction in I_c under “unloaded conditions” is significantly smaller than in the loaded regime \rightarrow decrease in quench current contained within 5 % even at high pre-load levels.
- Note that in the real system, geometrical tolerances/imperfections, can modify the conductor strain state. These cases can be simulated as well, but presented results correspond to a perfect geometry.



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Correlation between strand and magnet results

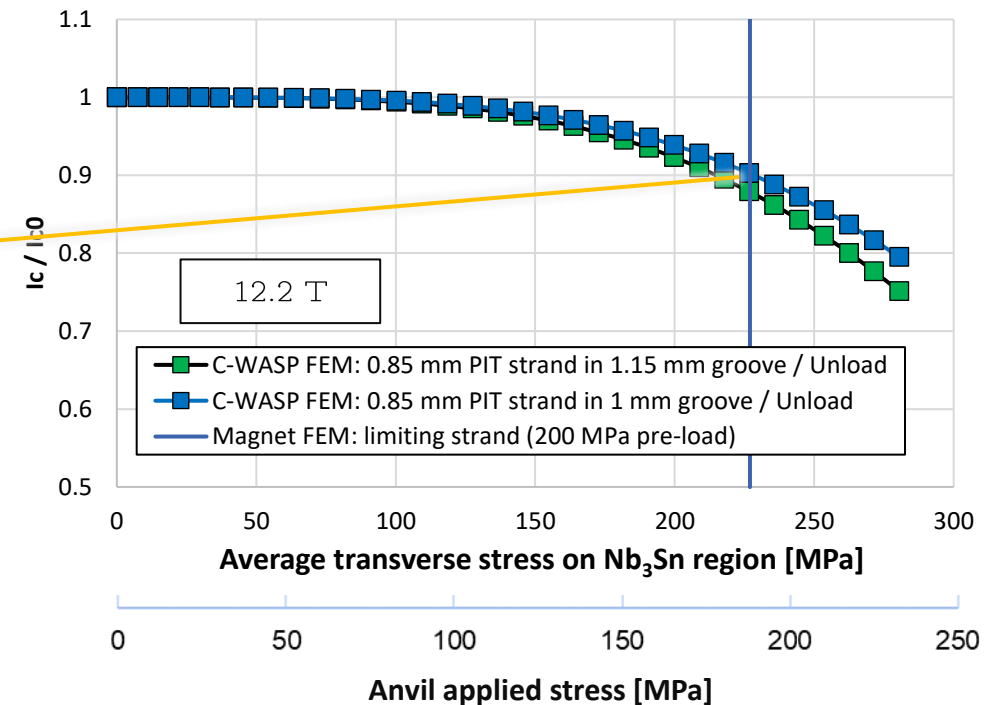
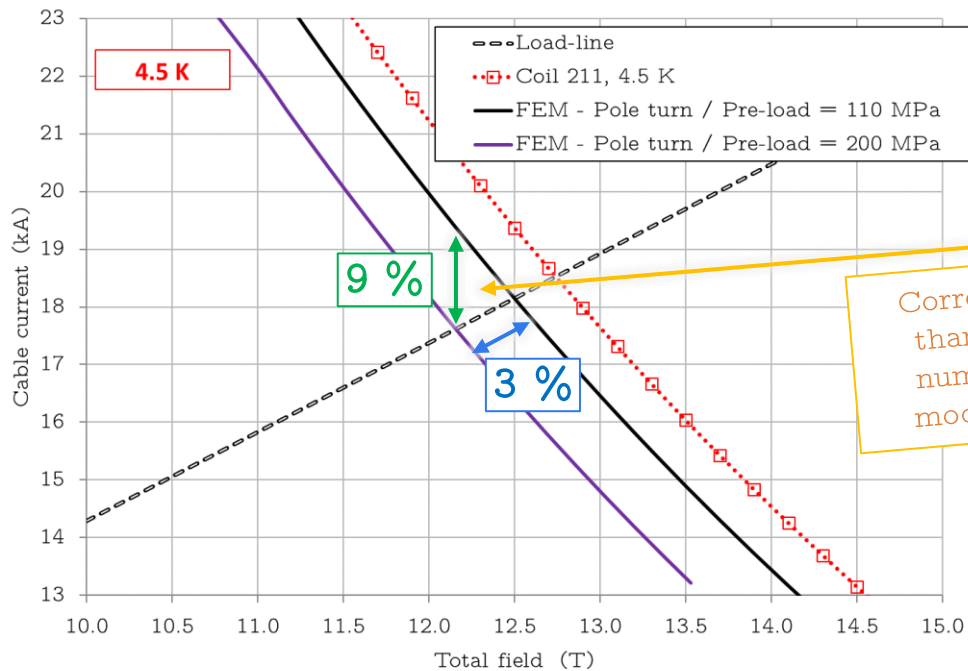
Correlation between both experiments:

- Methodology allows to compute the magnet limiting current thanks to the inclusion of strain effects on the conductor critical surface (assuming no cracks on the Nb₃Sn region).

Let's focus on the simulated case at 200 MPa (in the limit of validity, based on our current conductor understanding)

Note that: Load line margin is usually 1/3 of the corresponding critical surface one.

- No C-WASP tests are available for the 0.85 mm MQXF PIT strand, but this particular case can be simulated.



Outline



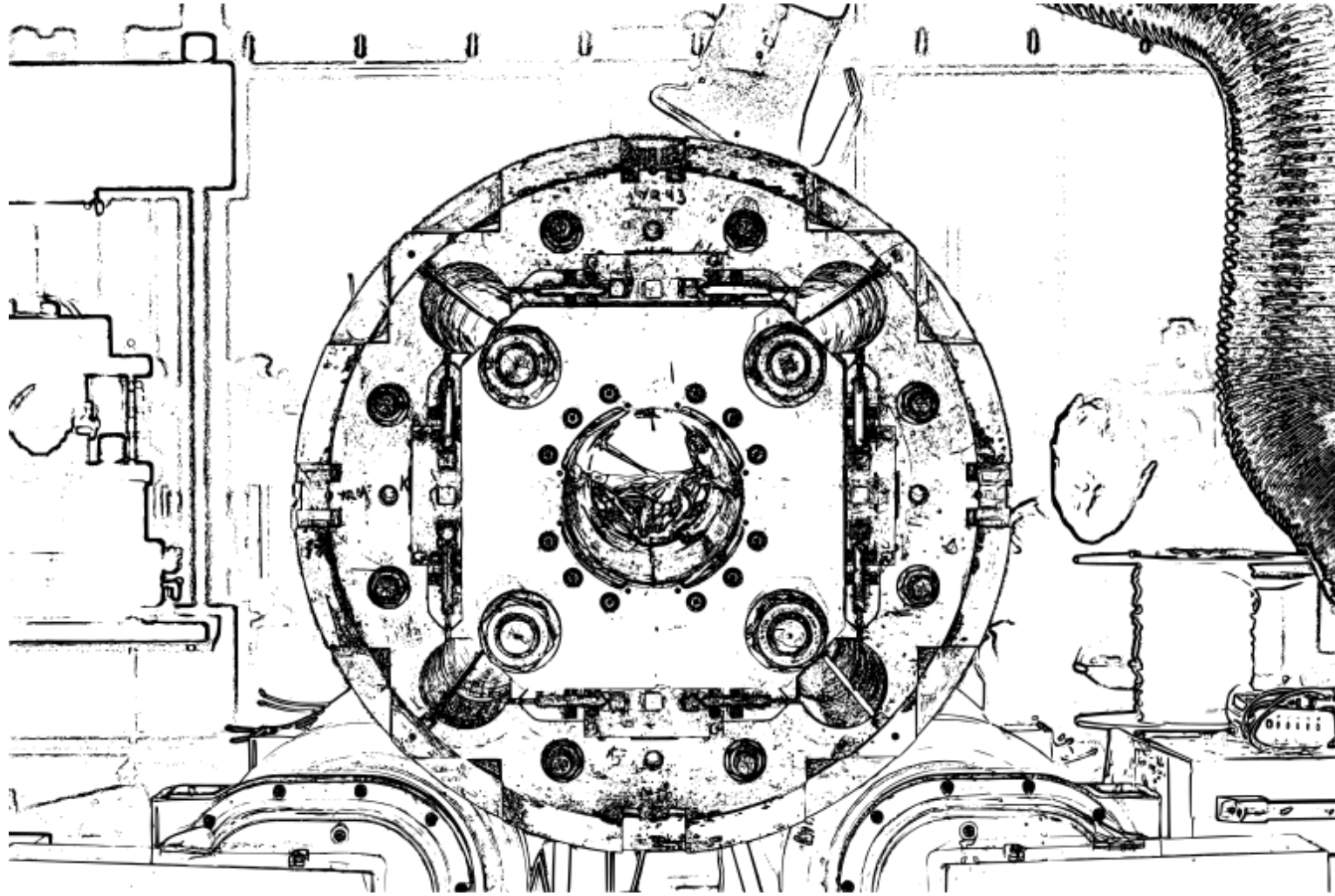
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Conclusions

- The accurate reproduction of the strain effects in the Nb₃Sn critical surface (pure strain effects, no cracks) is today a well-known domain of study. A three-dimensional mathematical description is available in the form of the exponential scaling law.
 - The latter has been successfully applied to detailed 2D and 3D FE models, matching experimental data at the conductor level.
- Meaningful results at the conductor scale (cable and strand/wire) can be obtained as well with a simplified geometry of the domain. This reduces significantly the computational cost and allows the application of the scaling law at the magnet scale.
 - Amplification factor needed to go from the strand to the filament level. This factor has been obtained by combining detailed and simplified geometry models. Consistent results with the one obtained by direct fitting of experimental data.
- Under certain support / boundary conditions Nb₃Sn conductors can withstand large stress levels (up to / above 200 MPa), in the absence of cracks sectioning the superconductor. Proven experimentally (B_{c2} evolution, post-mortem inspection) and numerically in wire and cable configuration.
 - In such a regime, the strain sensitivity is strongly magnetic field dependent. It is crucial to consider this dependency for the correct interpretation of test results.
- The direct application of the strand model to the magnet allows to include the strain effects modifying the short sample limit (again, no cracks regime). Good agreement between the obtained simulation results and recent experiment in MQXFS7.
 - Small loss (< 5 %) in magnet quench current even up to high pre-load levels (170 MPa measured in the winding pole).
 - Quench location corresponds to the pole turn block (peak field, low stress during magnet ramping). Caught by the model.
 - Since results are caught by the scaling law (up to this level) → indirect indication that the magnet support conditions are coherent with the conductor tests (C-WASP, FRESCA sample holder). More work on-going to substantiate/support this statement!
- The methodology provides important information to magnet designers. Always interesting to disentangle the peak field and high stress regions. Strand model provides the essential information on where the magnet is limited (peak field-low stress vs. lower field-high stress).
 - From conductor tests, the loss in I_c is significantly smaller when the sample is unloaded (permanent effects: copper plasticization).
 - Design / assembly parameters can be optimized to stay as much as possible in this domain.

Thank you very much for your attention!



**Special thanks (and credits!)
to everybody involved:**

Susana Izquierdo
Giorgio Vallone
Bernardo Bordini
Tommaso Bagni
Carmine Senatore
Paolo Ferracin

ANNEX

The scaling law

Important to recognize the three-dimensional nature of the scaling law.

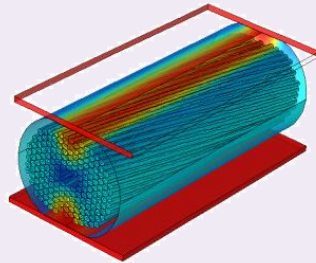
It has been proven that for (1) the scaling law can be used under more complex loading conditions.

Focus on transverse loads, those usually of largest magnitude on accelerator magnets

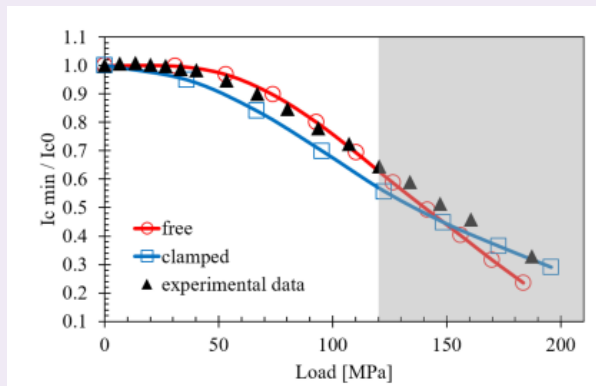
Studies employ Finite Element models to extract the strain tensor. Validated in:

- Transversally loaded wires (micro-scale models)

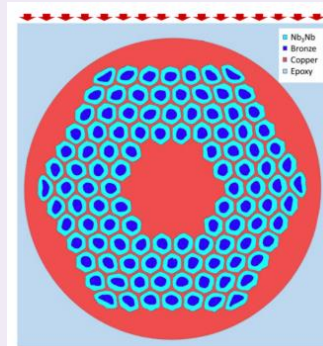
- 2D and 3D models computing I_c
- Detailed sub-element geometry
 - [31]-[32] Bordini, Cattabiani, Baffari et al.
 - [5] Senatore, Bagni, Calzolaio et al.
 - [33] Chiesa, Wang et al.



Cattabiani, Baffari et al.



Cattabiani, Baffari et al.

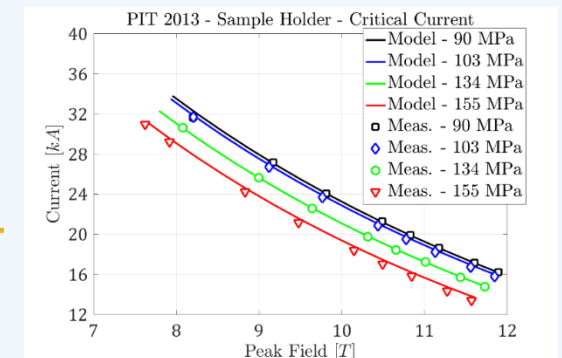
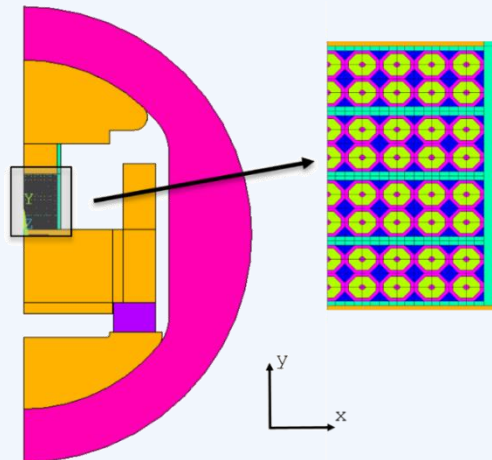


Bagni, Calzolaio et al.

- Transversally loaded cable stacks (macro-scale model)

[34] – [35] G. Vallone et al.

It can be applied to a 2D magnet cross-section (already published).

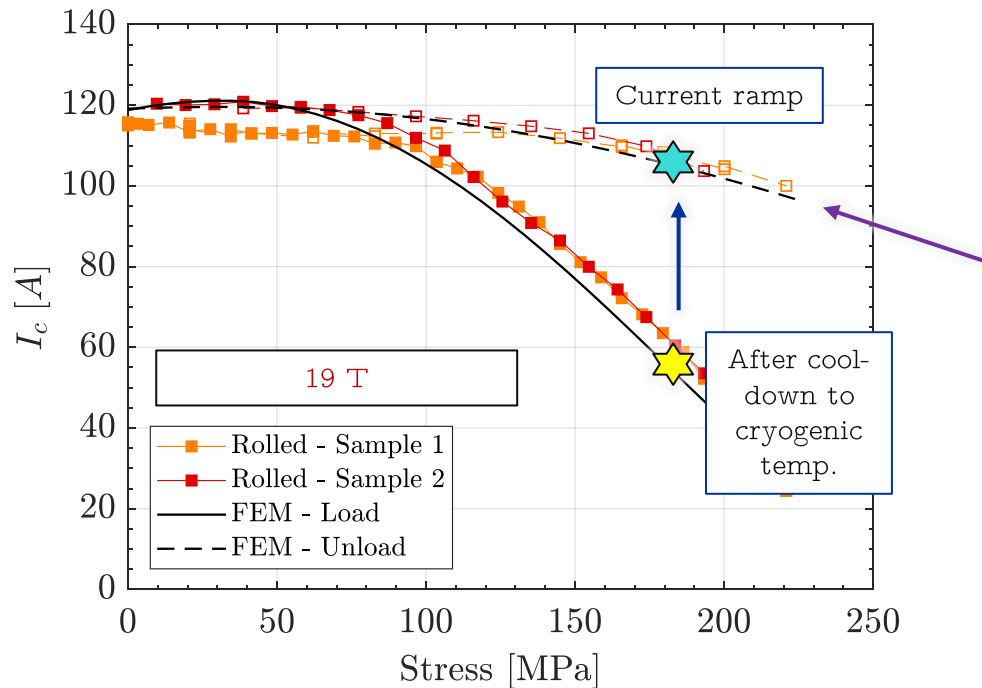


Vallone, Ferracin, Bordini et al.

Electro-mechanical limits on magnet configuration

A quick analogy, why is this important?

**Example! C-WASP measurement for
0.85 mm RRP 108/127 15% rolled**



For the simulated magnet case - 4.5 K

In our peak field region, we follow the unload curve, whose slope is significantly lower than the load one!

This is an interesting feature of a magnet design, allowing to contain the decrease in quench current under reasonable limits for high stress levels (without decreasing the overall current density).

Correlation between strand and magnet results

- In the domain of application for the scaling law (no cracks), the conductor results could be correlated with the magnet ones when:
 - 1) The field dependence is considered
 - 2) The test configuration is such that is representative of the magnet one:
 - ❖ Strand / insulation / groove ratios consistent with the final cable ones
Defines the fraction of force reaching the strand in the C-WASP experiment.
 - ❖ Support conditions are adequate (4-wall configuration, strain conditions)

Essential to consider all these factors in order to obtain meaningful information from strand / cable tests !

Magnet strand model:

Location of the quenching strand
Strain / stress state
Support conditions

The “measurable” (experimentally) is the transverse pre-load value.

C-WASP strand model:

Strain / stress state
Support conditions

The “measurable” (experimentally) is the normal applied stress value.

→
Consistent?
Domain of application of the scaling law?

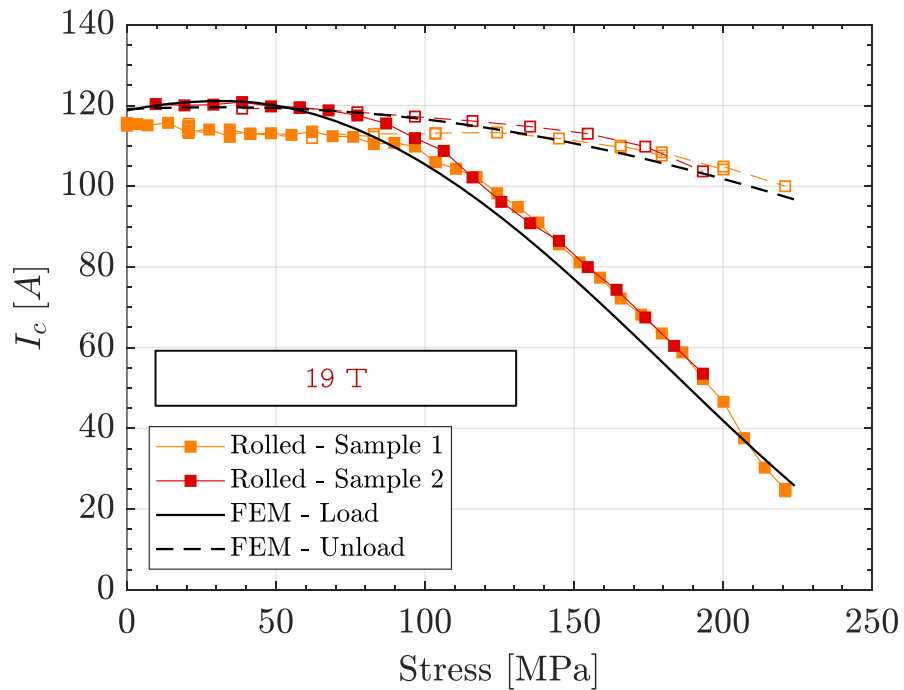
Correlation between both results.
Conductor data to magnet application!

I_c vs. transverse stress on single wires



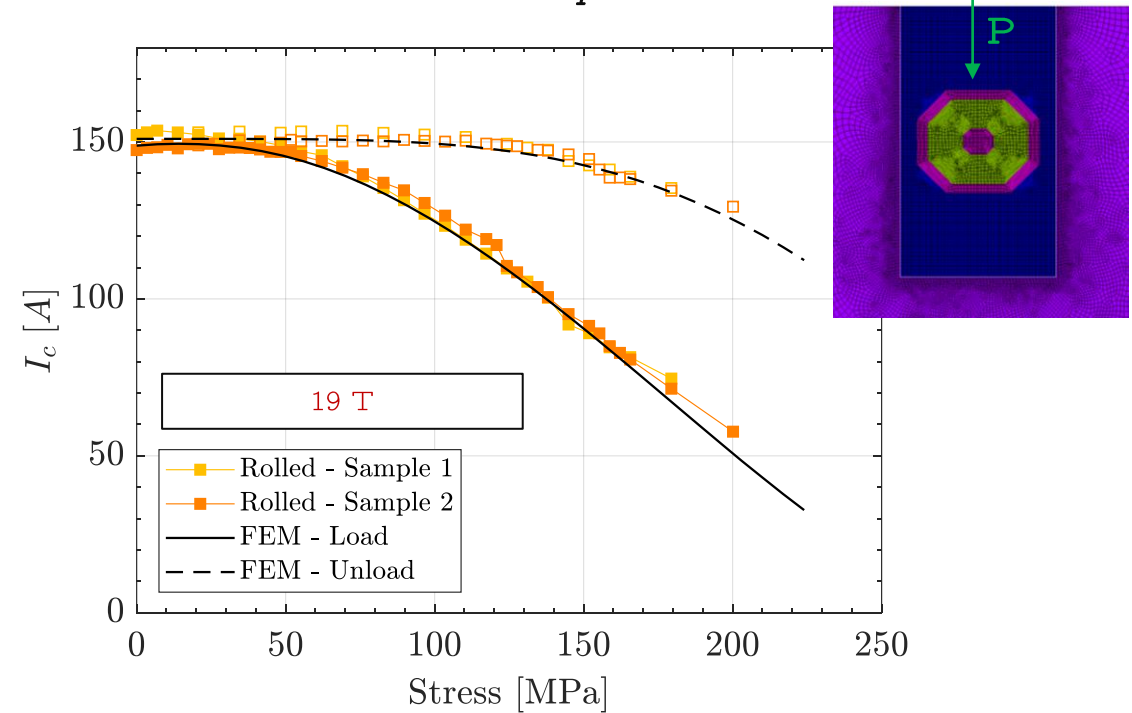
Simulation and experimental results:

0.85 mm RRP 108/127 15% rolled
(HL-LHC MQXF quadrupole)



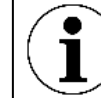
1.15 mm width groove

1 mm PIT 192 15% rolled
(FRESCA2 dipole)

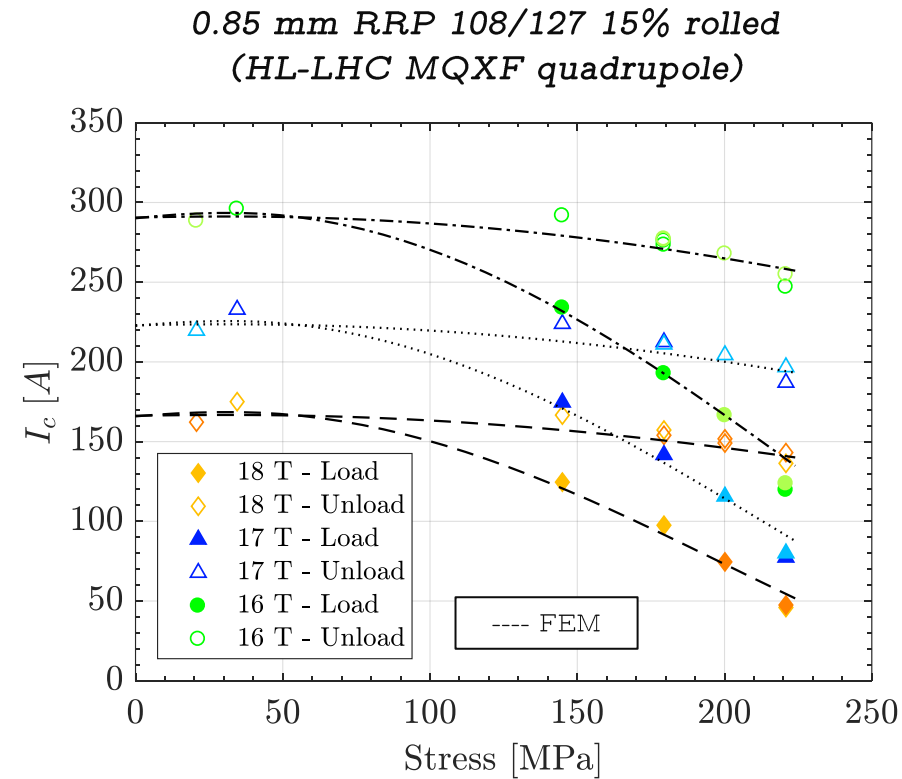


1.15 mm width groove

I_c vs. transverse stress on single wires

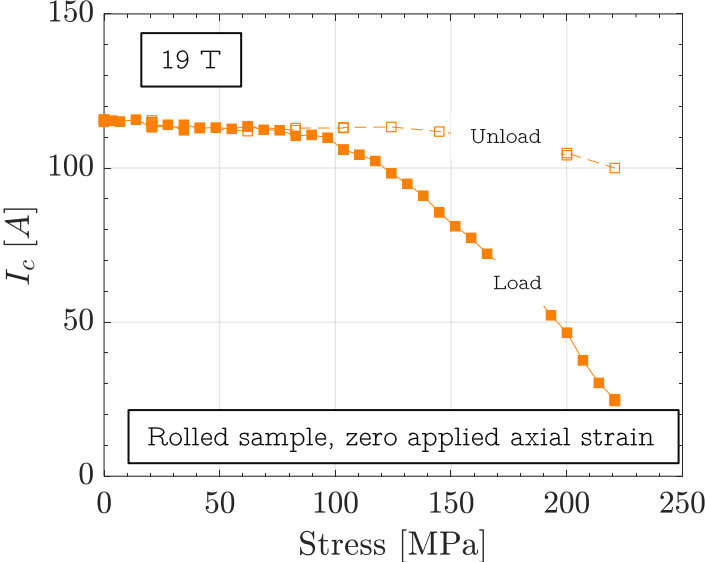


Simulation and experimental results:



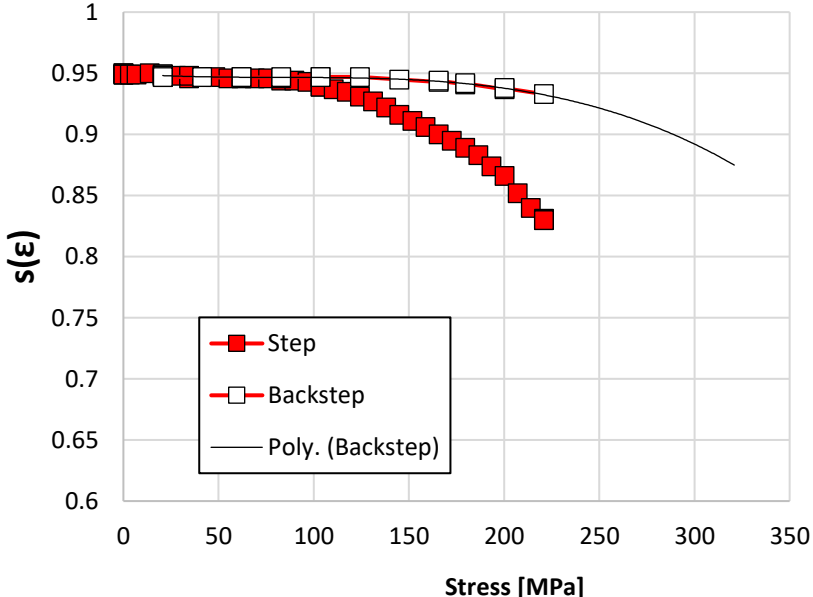
Permanent effect based on experimental data

Under the assumption of no cracks, we can use our T_{c0} , B_{c20} , C_0 obtained from the axial test and map a “experimental” strain function from the measurements at 19 T



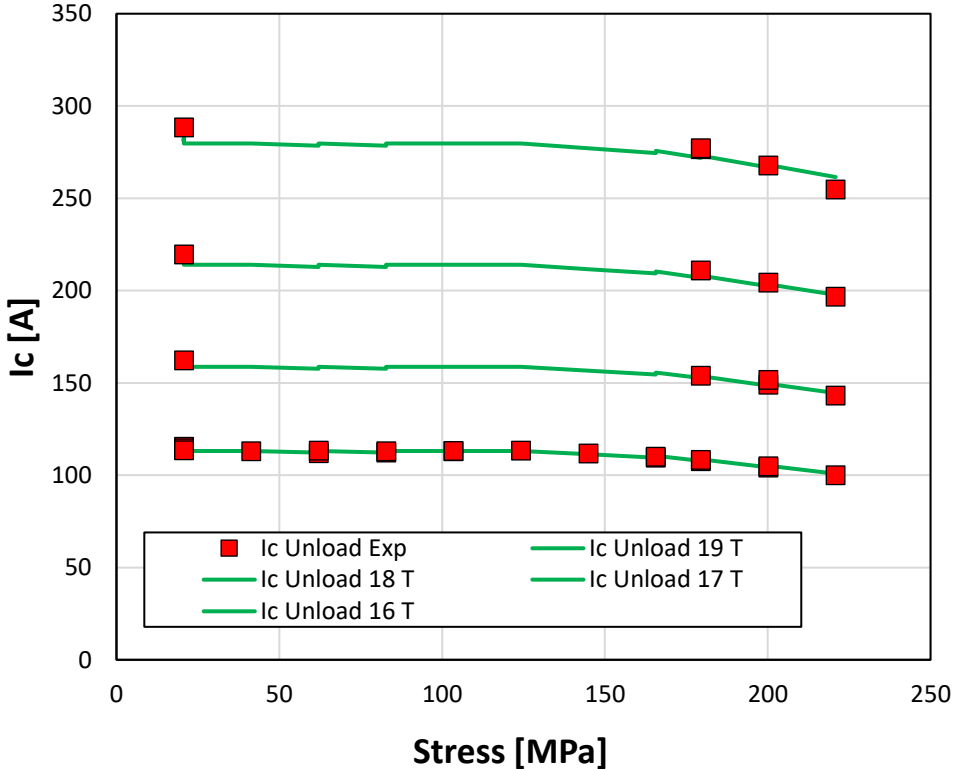
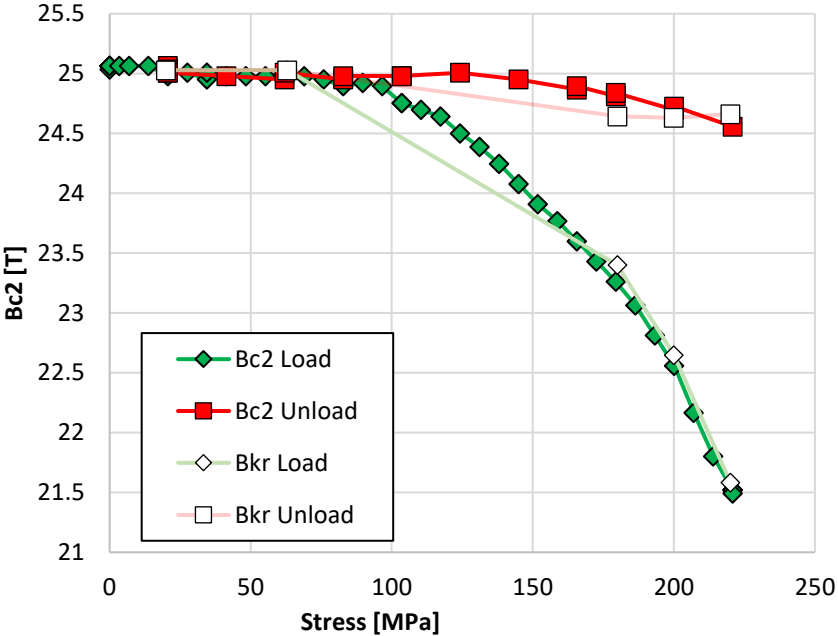
Assuming T_{c0} , B_{c20} , C_0 from axial WASP

→
Optimization algorithm



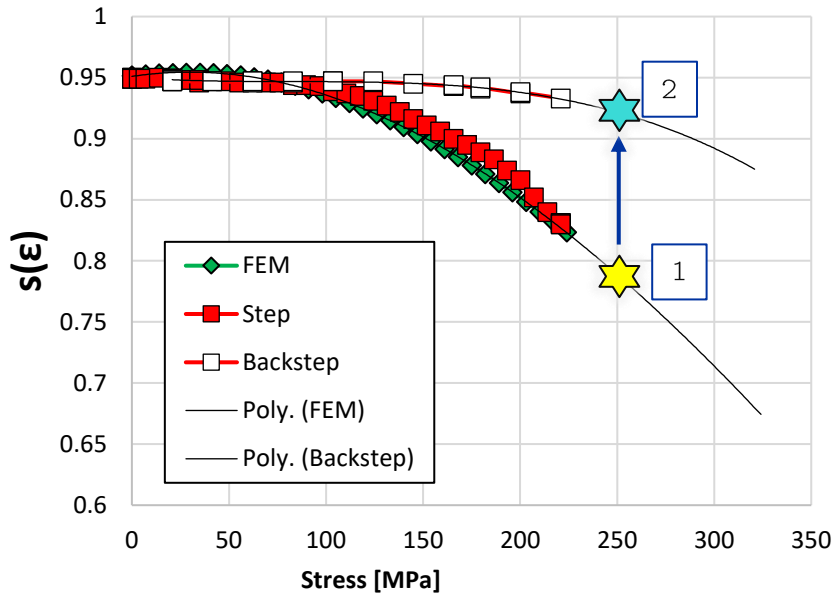
Permanent effect based on experimental data

Knowing the strain function (from measurements at 19T) and using the fitting parameters from the axial WASP we can compute I_c at any field, or B_{c2} at any stress.



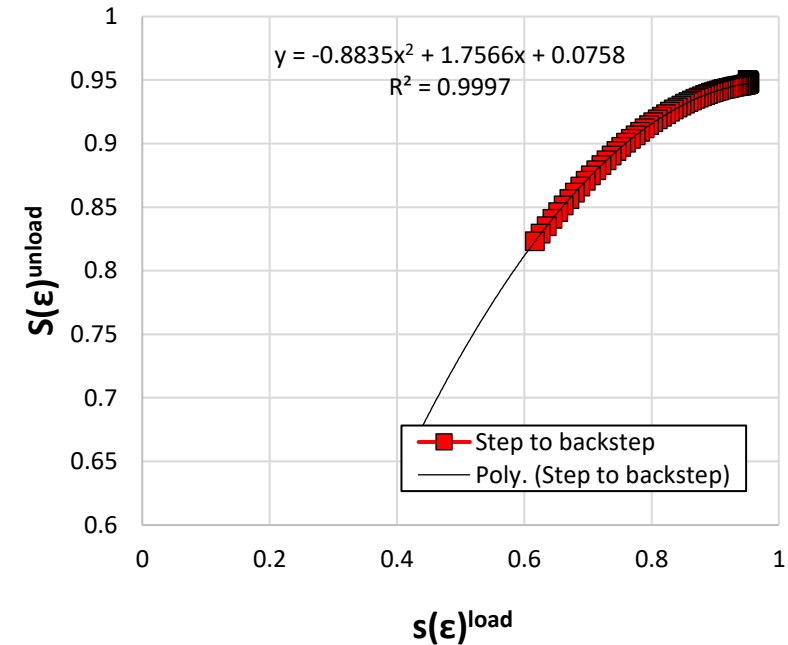
Permanent effect based on experimental data

Knowing the strain function (from measurements at 19T) and using the fitting parameters from the axial WASP we can compute I_c at any field, or $Bc2$ at any stress.



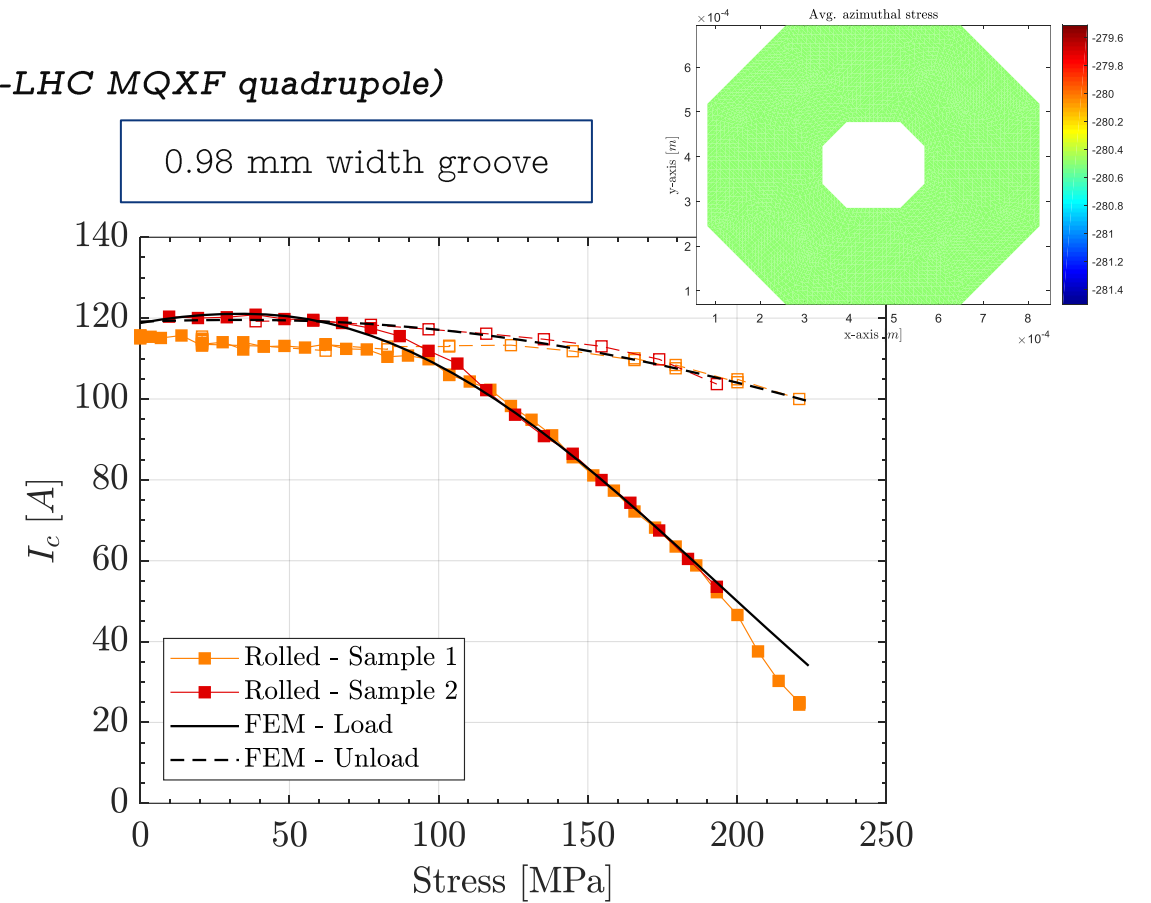
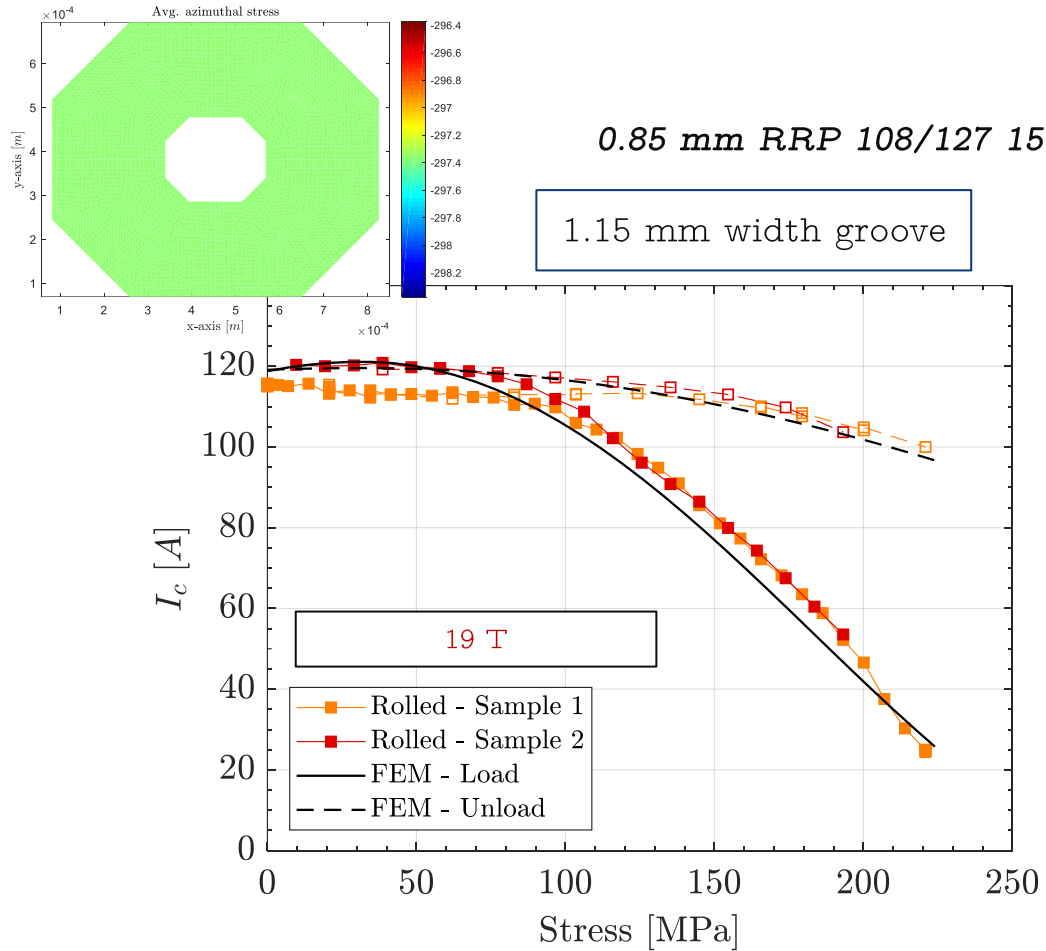
Furthermore, we know that our FEM matches well the part under load. To infer the I_c after unload, a simplistic approach could be to use the experimental strain function plot to determine $s(\epsilon)$ for the unloaded sample.

Note.- If our model would catch better the residual strain, this would not be needed.



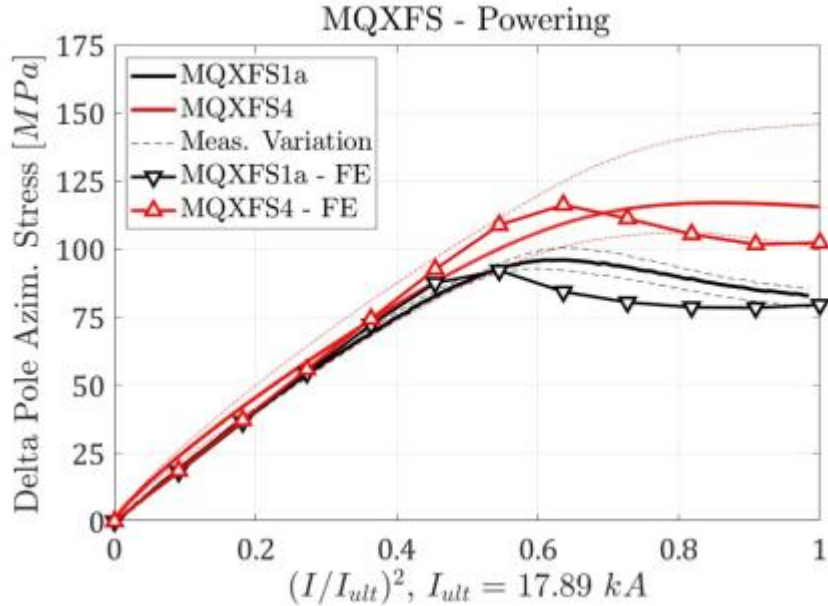
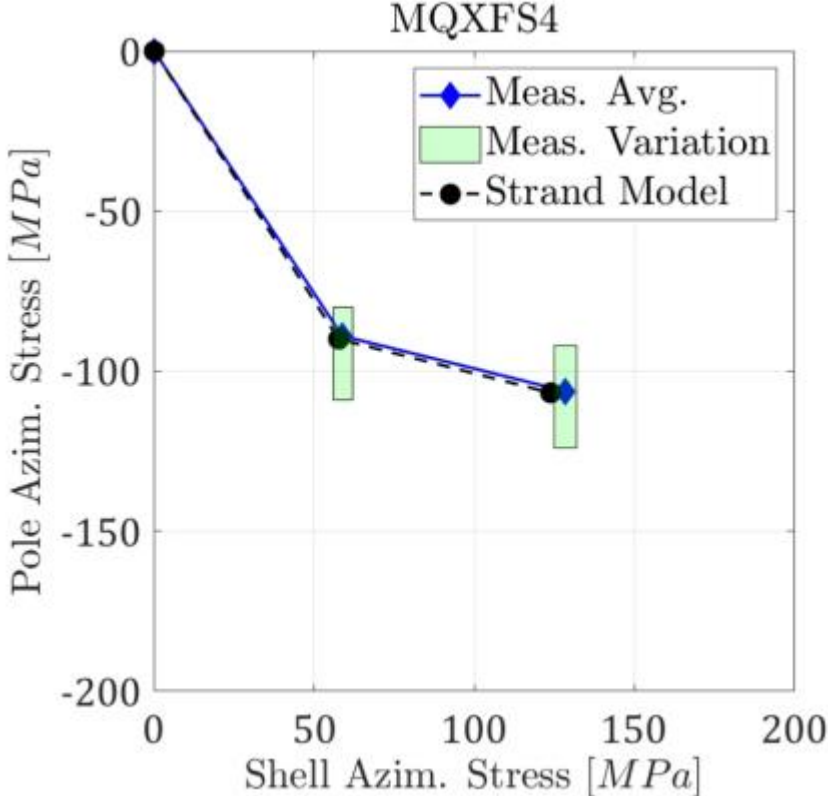
Thus, we can set our transfer function between $s(\epsilon)_{\text{load}}$ and $s(\epsilon)_{\text{unload}}$ based on UNIGE measurements.

Electro-mechanical limits on magnet configuration



Setting the pressure, the total force in the system is larger for larger grooves (larger total force).
 For the same groove and pressure (same total force), the force reaching the strand is larger for larger diameters.
 Force fraction between UNIGE tests varies within 10%. See [5].

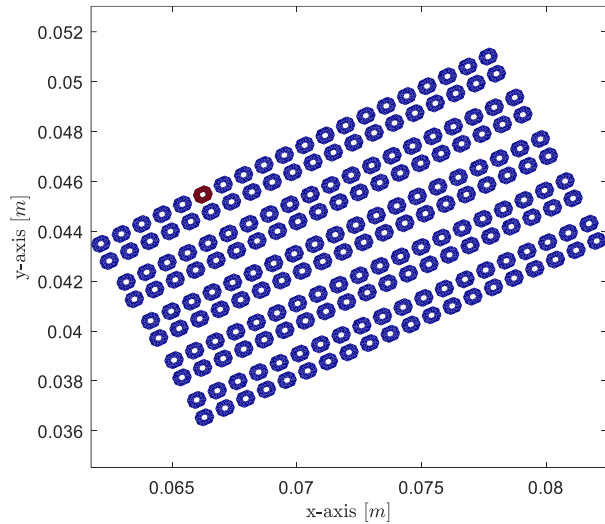
Magnet strand model



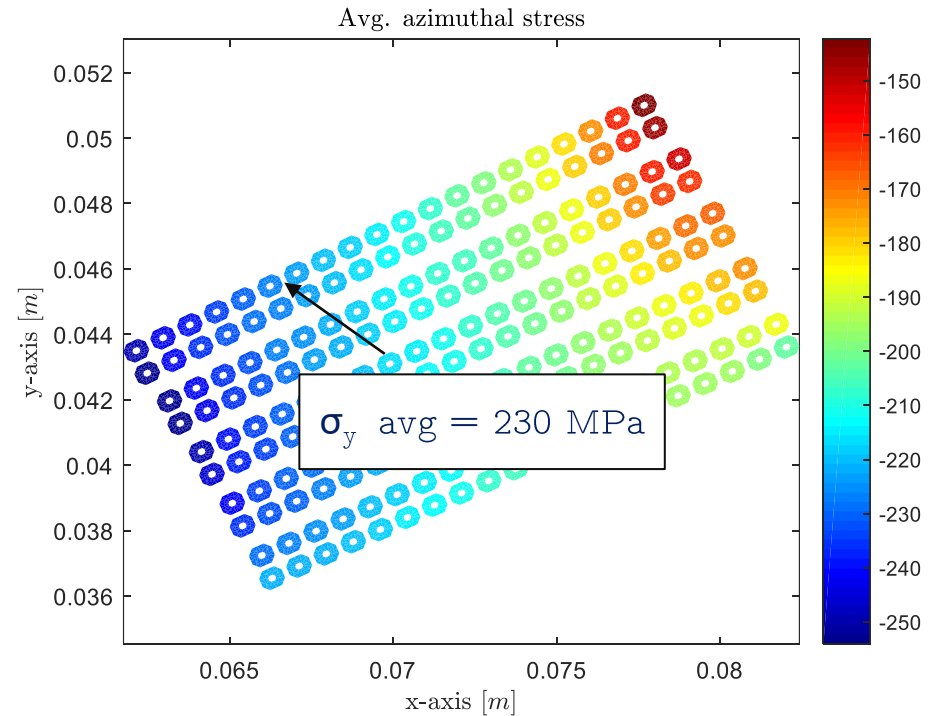
Vallone et al.

Correlation between strand and magnet results

- Average normal stress in the Nb₃Sn region (magnet model) after cool-down to cryogenic temperature.



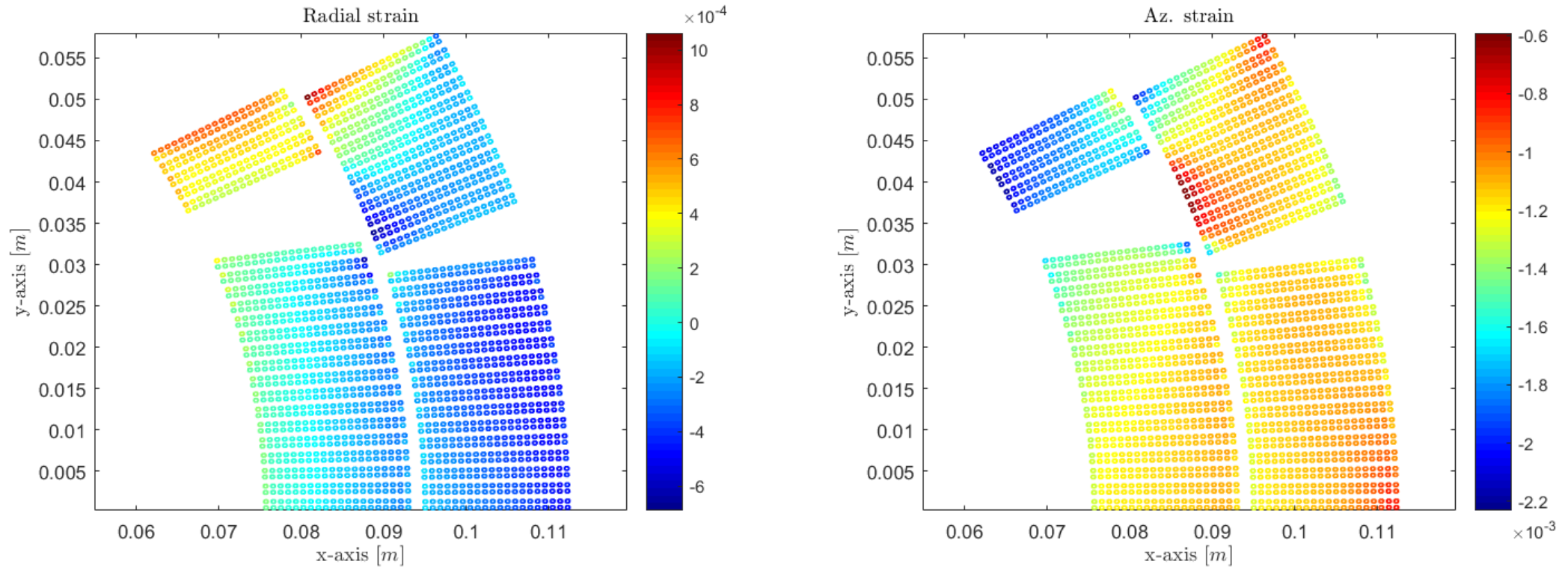
200 MPa pre-load – magnet model (1.9 K)



Electro-mechanical limits on magnet configuration

Magnet stress state:

Reference pre-load (4.5 K - Nb₃Sn region)



Azimuthal, radial and axial directions ~ principal.
Dominated by the azimuthal component.

Permanent effect based on experimental data.

Fixed temperature simplification:

$$J_c(B, \varepsilon) = C^* \left(\frac{B}{B_{c2}^*(\varepsilon)} \right)^{p-1} \left[1 - \left(\frac{B}{B_{c2}^*(\varepsilon)} \right) \right]^q$$

This is valid when the I_c reduction is governed by the lattice deformation (strain, no cracks) and far from T_c .

Since in our experiment at UNIGE we have found that no cracks are present in the wires (by inspection and by consistency in B_{c2} decrease), we can use the expression above.

- $B_{c2}(\varepsilon)$ is obtained by Kramer extrapolation* from the different field sweeps.
- One can then compute I_c at any magnetic field and any strain, if we know $B_{c2}(\varepsilon)$.
- The same applies when the load is removed, there the permanent effect comes from $B_{c2}^{\text{unload}}(\varepsilon)$.
- Using the expression we can get the permanent reduction by using: $I_c^{\text{unload}}(\varepsilon)/I_{c0}$ [as a function of $B_{c2}^{\text{unload}}(\varepsilon)$ and $B_{c2}^{\text{unload}}(0)$]
- Very interestingly for us, we can also get the ratio $I_c^{\text{unload}}(\varepsilon)/I_c^{\text{load}}(\varepsilon)$. In this case as a function of $B_{c2}^{\text{unload}}(\varepsilon)/B_{c2}^{\text{load}}(\varepsilon)$.

With $p=0.5$ and $q=2$:

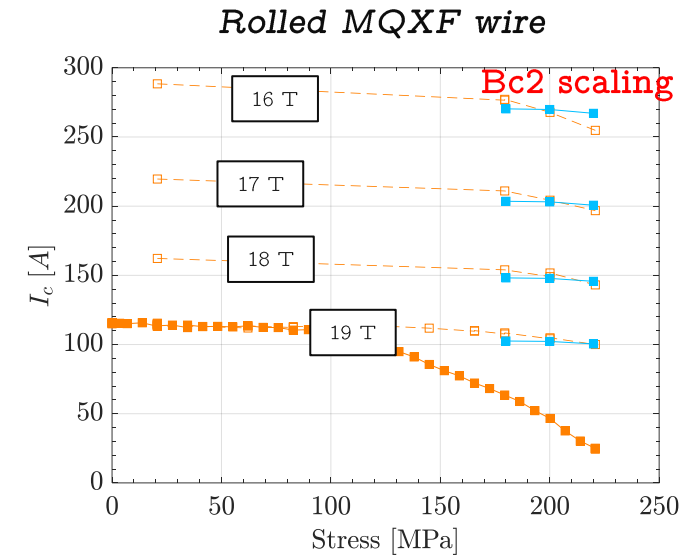
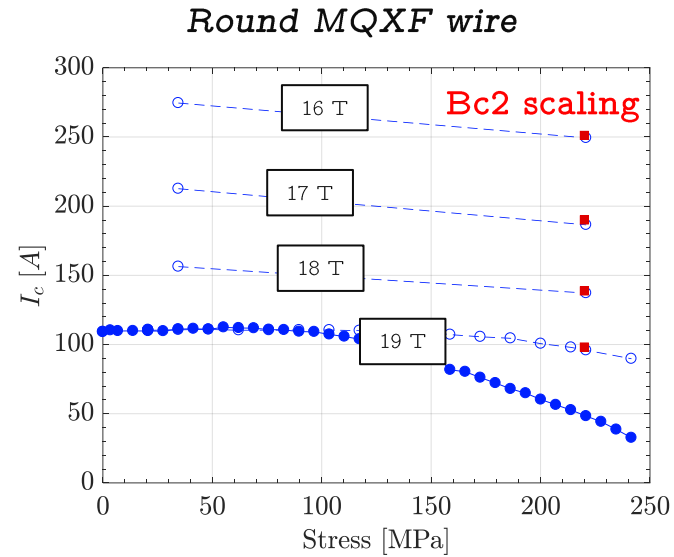
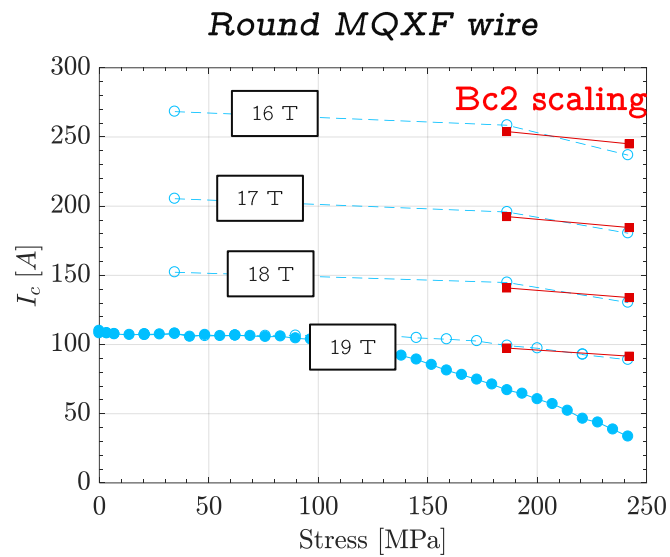
$$\frac{I_c^{\text{unload}}}{I_c^{\text{load}}}(B, \sigma) = \left[\frac{B_{c2}^{\text{load}}(\sigma)}{B_{c2}^{\text{unload}}(\sigma)} \right]^{3/2} - \left[\frac{B_{c2}^{\text{unload}}(\sigma) - B}{B_{c2}^{\text{load}}(\sigma) - B} \right]^2$$

$$\frac{I_c^{\text{unload}}(B, \sigma)}{I_{c0}(B, \sigma = 0)} = \left[\frac{B_{c2}(0)}{B_{c2}^{\text{unload}}(\sigma)} \right]^{3/2} - \left[\frac{B_{c2}^{\text{unload}}(\sigma) - B}{B_{c2}(0) - B} \right]^2$$

Permanent effect based on experimental data.

Results

Measured I_c and computed using $Bc2(\epsilon)$ after force unload agrees well.



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