

Bayesian inference of the EoS of dense nuclear matter from heavy-ion collision data

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Density dependent EoS in UrQMD

[EPJ C 82.5 \(2022\): 1-12](#),
[EPJ C 82.10 \(2022\): 1-12](#)

- ❖ A density dependent potential enters the QMD equations: $\dot{\mathbf{r}}_i = \frac{\partial \mathbf{H}}{\partial \mathbf{p}_i}$, $\dot{\mathbf{p}}_i = -\frac{\partial \mathbf{H}}{\partial \mathbf{r}_i}$.
 - The potential energy term of the Hamiltonian \mathbf{H} is density dependent.
- ❖ The potential energy $V(n_B)$ is related to the pressure as:

$$P(n_B) = P_{id}(n_B) + \int_0^{n_B} n' \frac{\partial U(n')}{\partial n'} dn', \quad U(n_B) = \frac{\partial (n_B \cdot V(n_B))}{\partial n_B}$$

$U(n_B) \Rightarrow$ single particle potential, $P_{id}(n_B) \Rightarrow$ pressure of an ideal Fermi gas of baryons

constraining the potential energy $V(n_B) \Rightarrow$ constraining the EoS

Parameterisation of the potential energy

- ❖ Upto $2n_0$, EoS reasonably constrained by
 - flow data at moderate energies *P. Danielewicz, Et al Science 298, 1592 (2002), H. Kruse, Et al. Phys. Rev. Lett. 54, 289 (1985)*
 - nuclear incompressibility data *Y. Wang, Et al. PLB 778, 207 (2018)*
 - bayesian analysis *S. Huth et al., Nature 606, 276 (2022)*
- ❖ Upto $2n_0$, CMF model-fit is used *A. Motornenko et al., PRC 103.5 (2021) , A. Motornenko et al., PRC 101.3 (2020)*

- ❖ We constrain the $V(n_B) > 2n_0$
 - 7th degree polynomial is used
 - $h = -22.07$ MeV to match CMF at $2n_0$

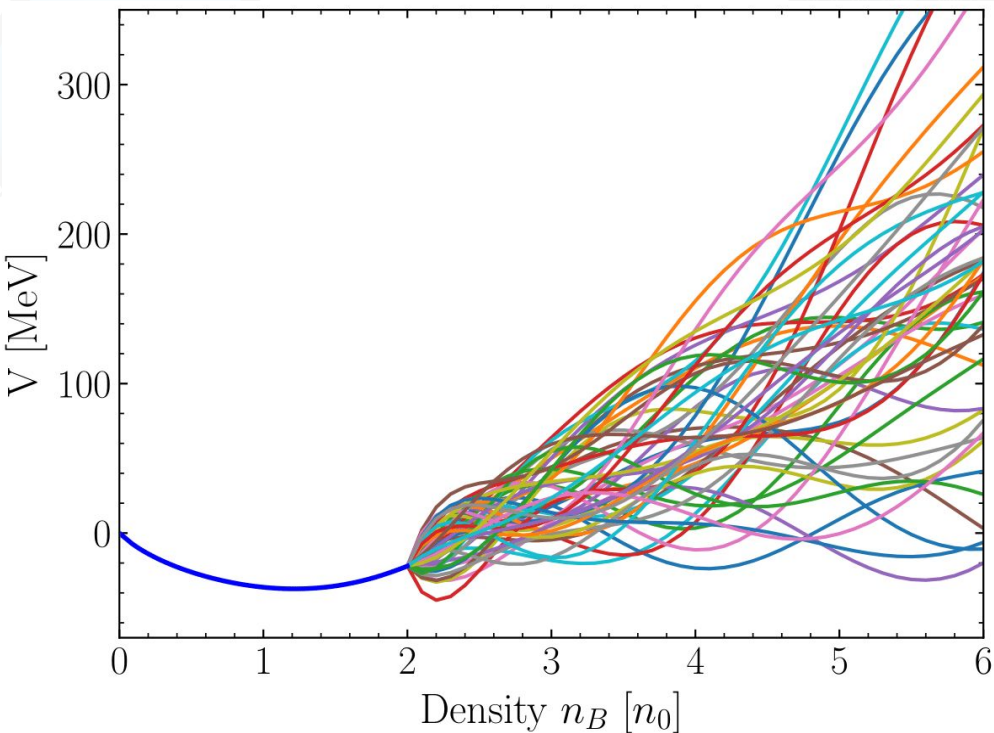
CMF

$$V(n_B) = \sum_{i=1}^7 \theta_i \left(\frac{n_B}{n_0} - 2 \right)^i + h$$

0- $2n_0$

> $2n_0$

Few examples...



$$V(n_B) = \sum_{i=1}^7 \theta_i \left(\frac{n_B}{n_0} - 2 \right)^i + h$$

We constrain the polynomial coefficients

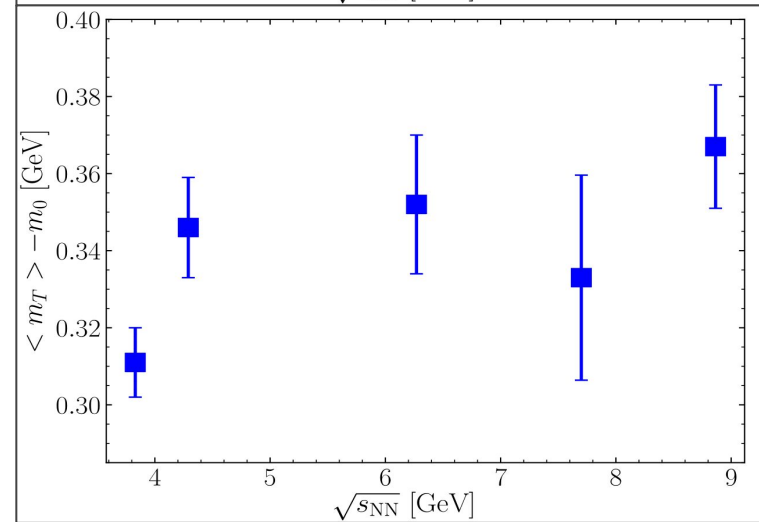
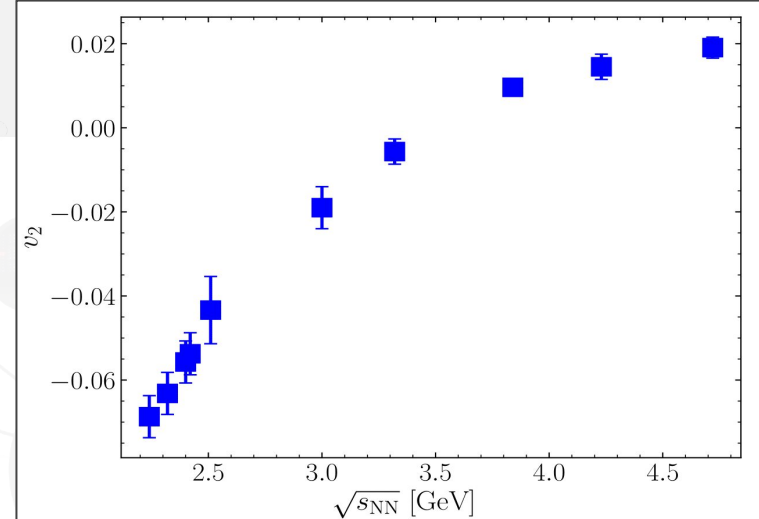
$$\theta = \{\theta_1, \theta_2, \dots, \theta_7\}$$

The experimental data

❖ Proton observables (mid rapidity)

- Elliptic flow : 10 data points
 - *E895, CERES, FOPI, STAR, HADES*
 - Mid-central collisions
- Transverse kinetic energy: 5 data points
 - *E802, NA49, STAR*
 - Central collisions

The data, $\mathbf{D} = \{v_2^{exp}, \langle m_T \rangle^{exp} - m_0\}$ is used to constrain the parameters θ .



Fast emulators are necessary

- ❖ Bayesian inference involves numerous UrQMD simulations
 - UrQMD ~ 100 s/event

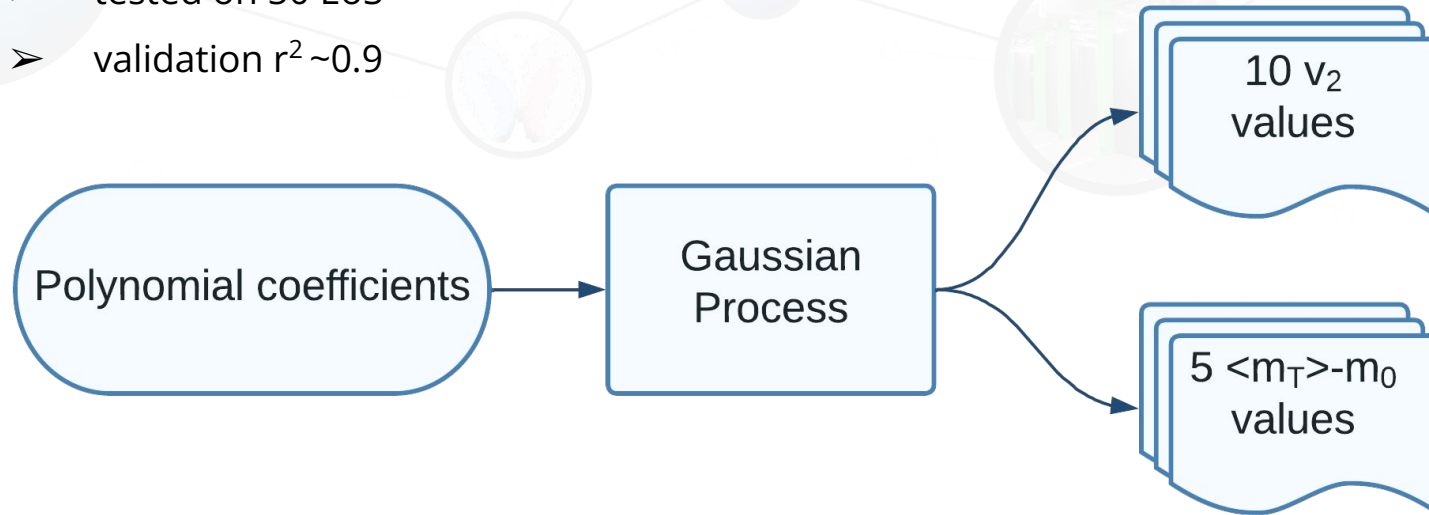
<i>observable</i>	<i>No of beam energies</i>	<i>No of events per beam energy</i>
v_2	10	12000
$\langle m_T \rangle - m_0$	5	1000

- **For one parameter set** ~ 125000 events* 100 s = **~ 3500 hrs**
- ❖ MCMC sampling requires random walk through 1000s of parameter sets
 - **not feasible to run UrQMD for MCMC !**

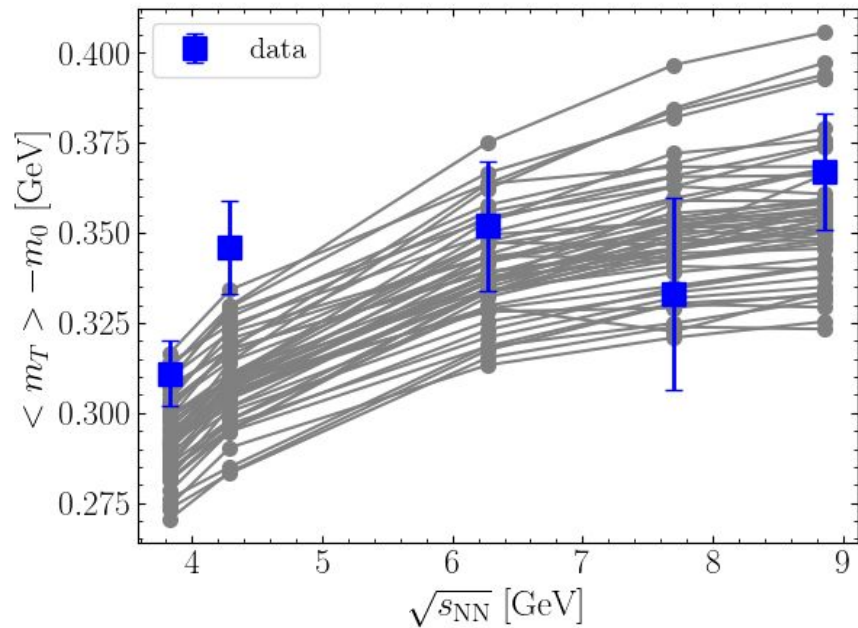
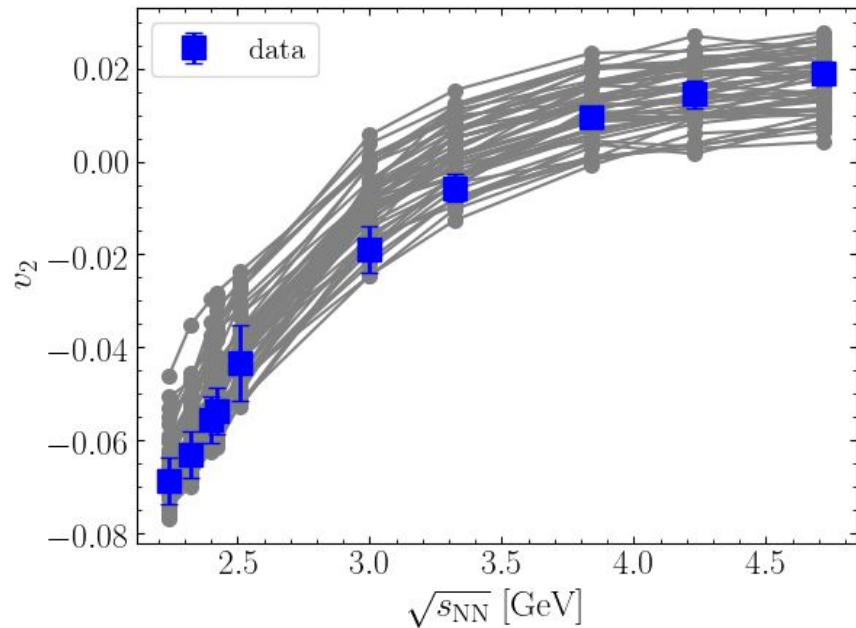
Training the GP models

❖ Gaussian Process models are trained as fast emulators

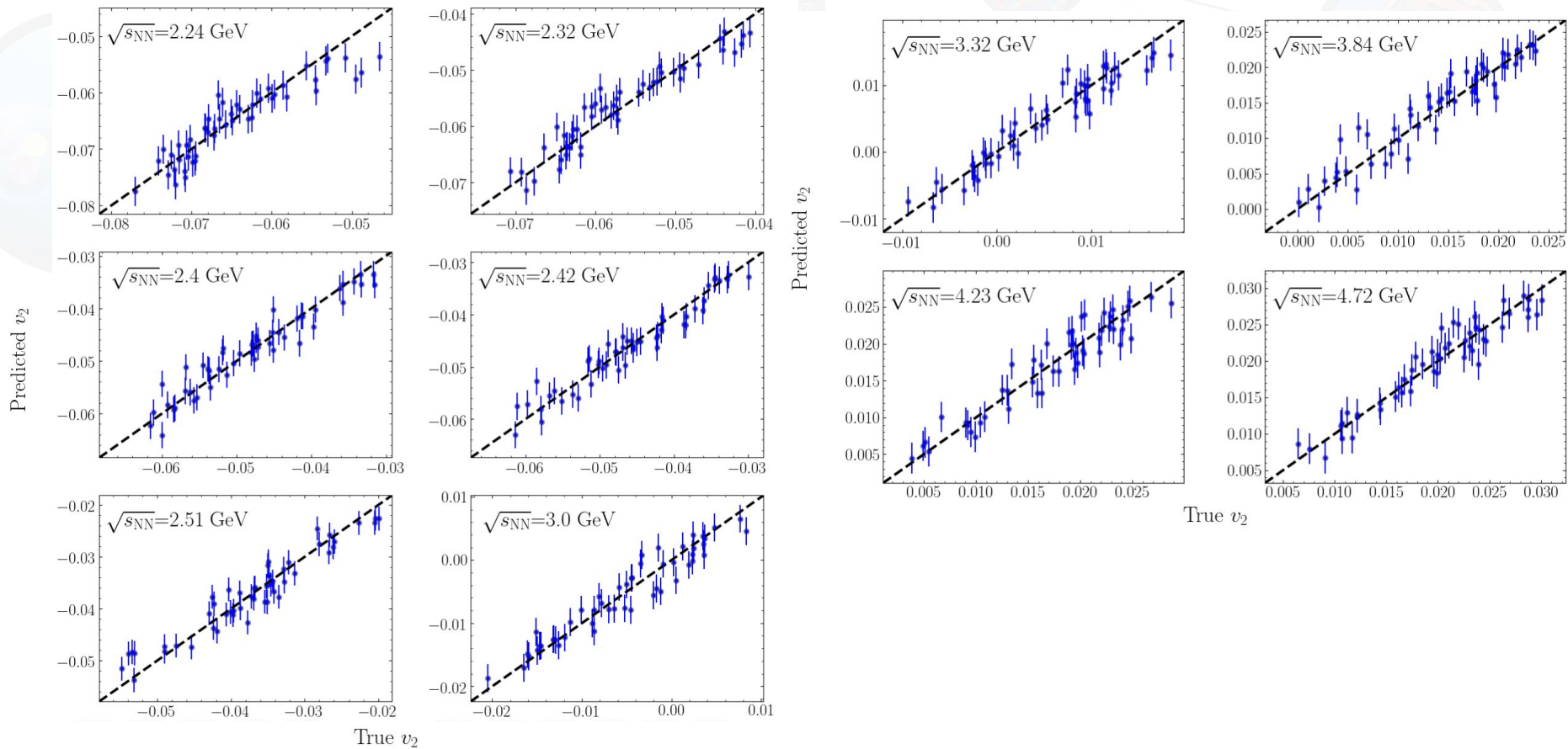
- trained on 200 randomly generated EoSs
- tested on 50 EoS
- validation $r^2 \sim 0.9$



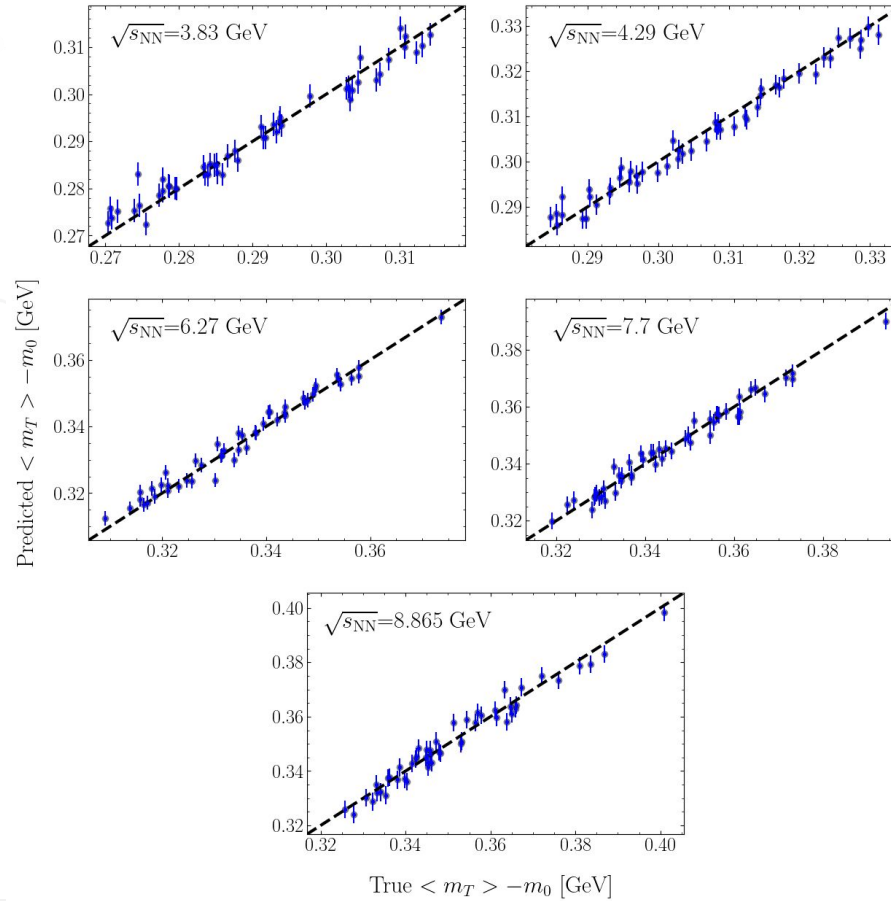
Potentials for training the GP models



GP models: performance



GP models: performance



Bayesian inference

Posterior

- ❖ probability that the parameters θ explains the data \mathbf{D}

$$P(\theta|\mathbf{D}) \propto P(\mathbf{D}|\theta)P(\theta)$$

Log-likelihood

$$\ln P(\mathbf{D}|\theta) = -\frac{1}{2} \sum_i \left[\frac{(x_i^\theta - d_i)^2}{\sigma_i^2} + (\ln(2\pi\sigma_i^2)) \right]$$

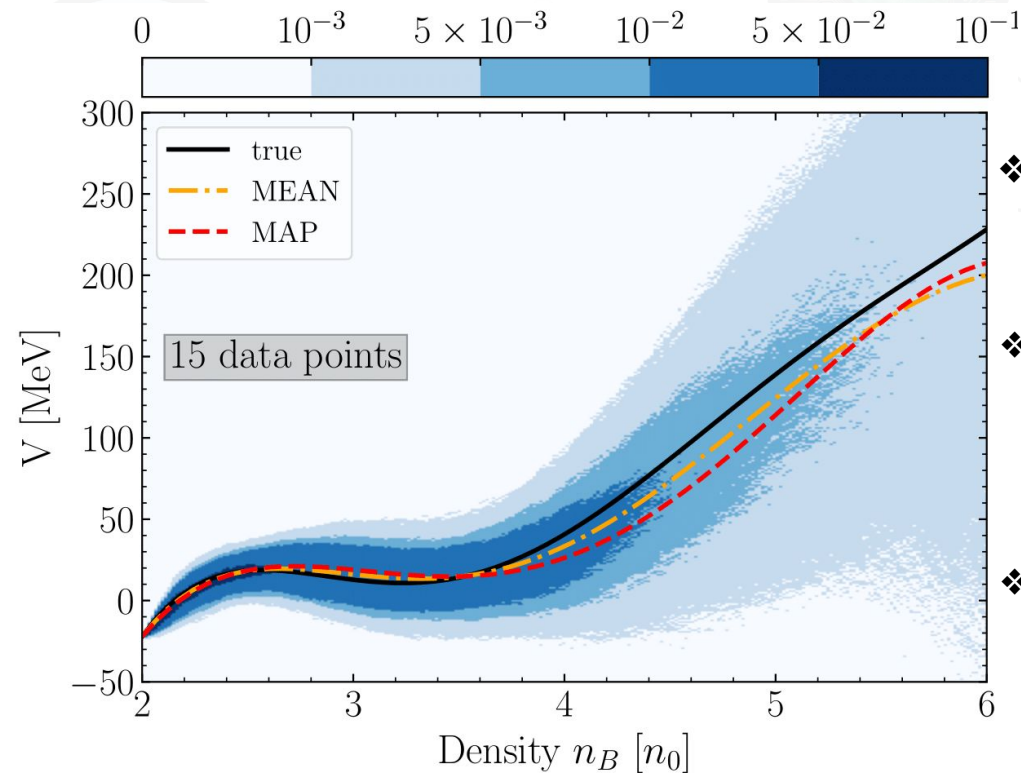
Includes uncertainty from experiment & GP model

- ❖ gaussian priors are used
- ❖ μ, σ from GP training data

Closure test

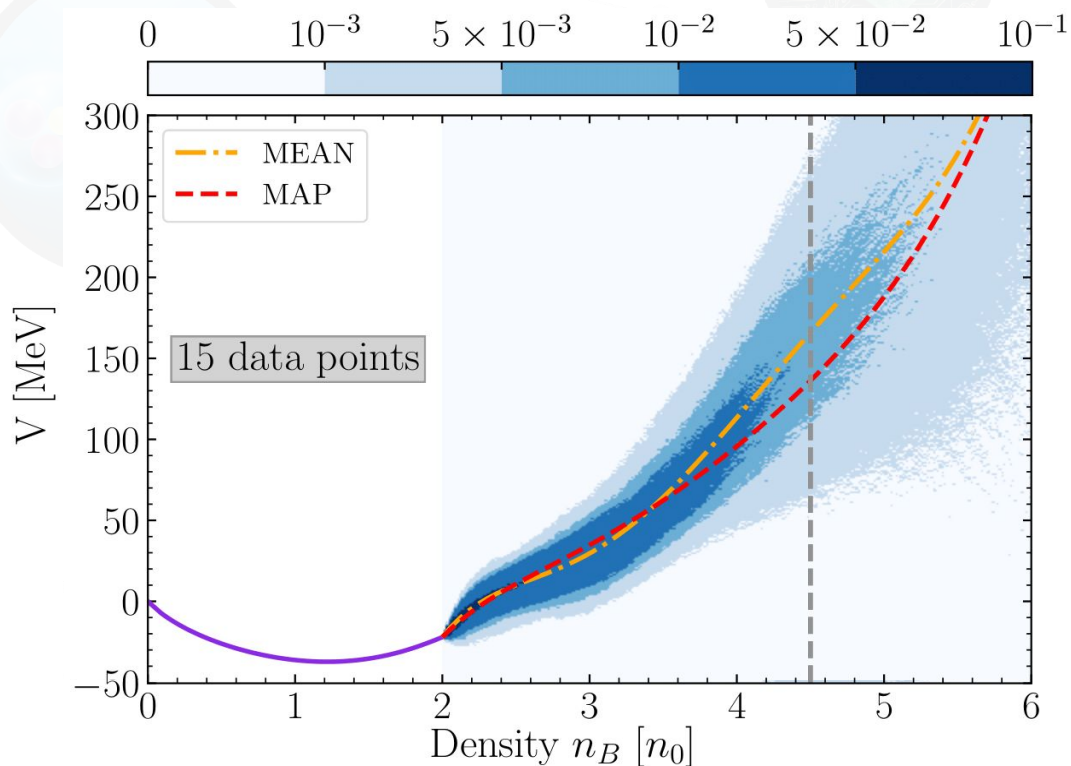
1. Consider a random EoS as “ground-truth”
2. Calculate observables from UrQMD
 - a. 10 values of v_2 and 5 values of $\langle m_T \rangle - m_0$
3. Assume the UrQMD observables for this EoS to be the “data”
 - a. assume uncertainty similar to experimental data
4. Construct the posterior
5. Compare with the “ground-truth”

Closure test: results



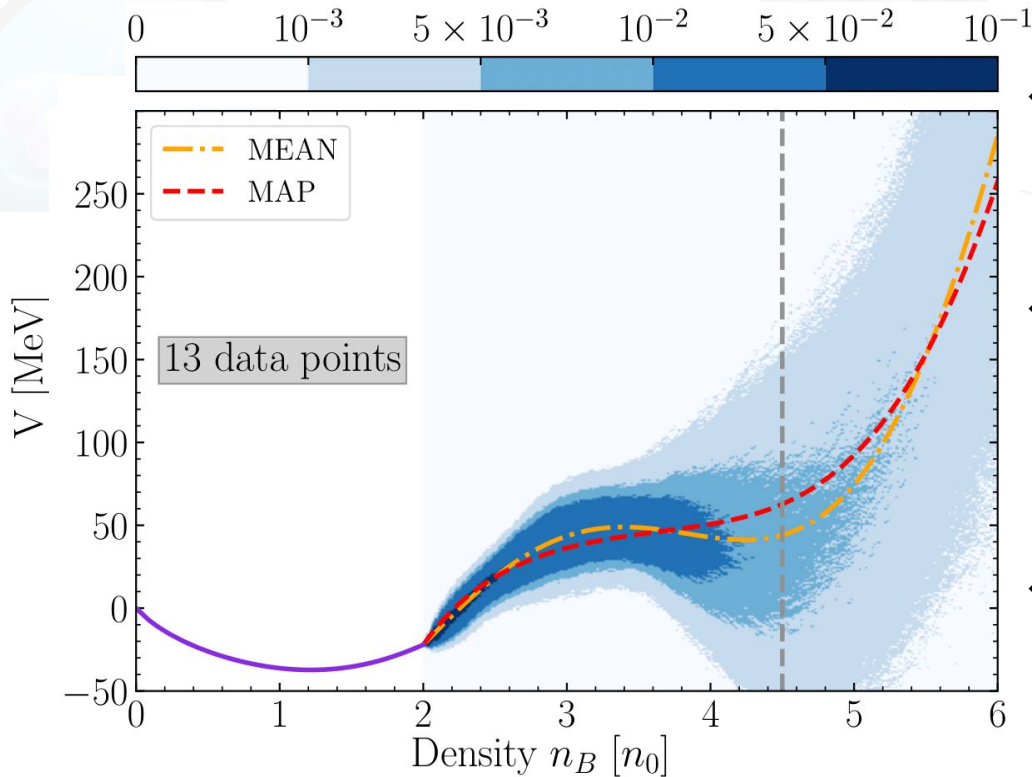
- ❖ Tight constraints up to $4n_0$
 - large uncertainty above $4n_0$
 - yet mean closely follows “ground-truth”
- ❖ Two curves extracted:
 - “MAP”: mode of posterior
 - “MEAN”: mean of posterior
- ❖ MEAN and MAP closely follows “ground-truth” upto $6n_0$

Result from experimental data



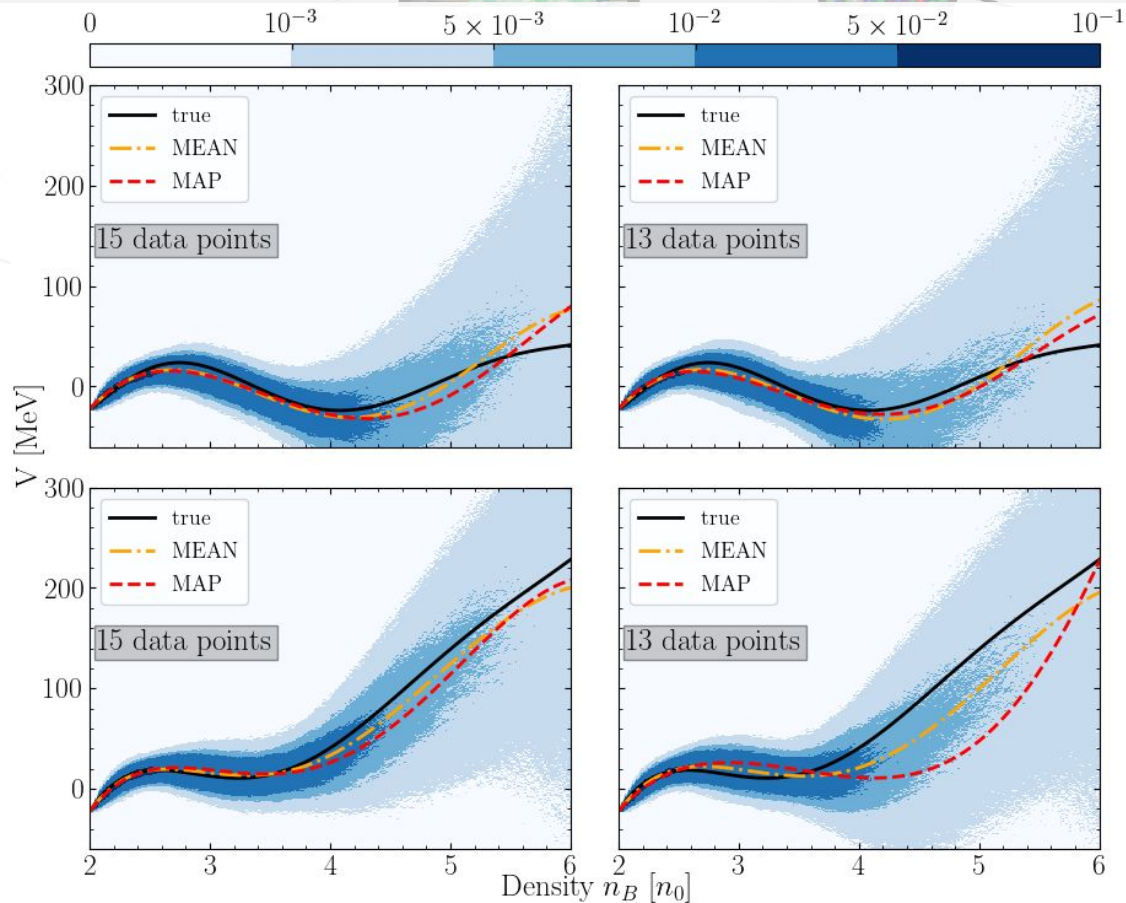
- ❖ Posterior from experimental data
 - 10 measurements of v_2
 - 5 measurements of $\langle m_T \rangle - m_0$
- ❖ Tight constraints upto $4n_0$
 - MEAN, MAP suggests stiff EoS
 - No phase transition

Sensitivity to choice of observables

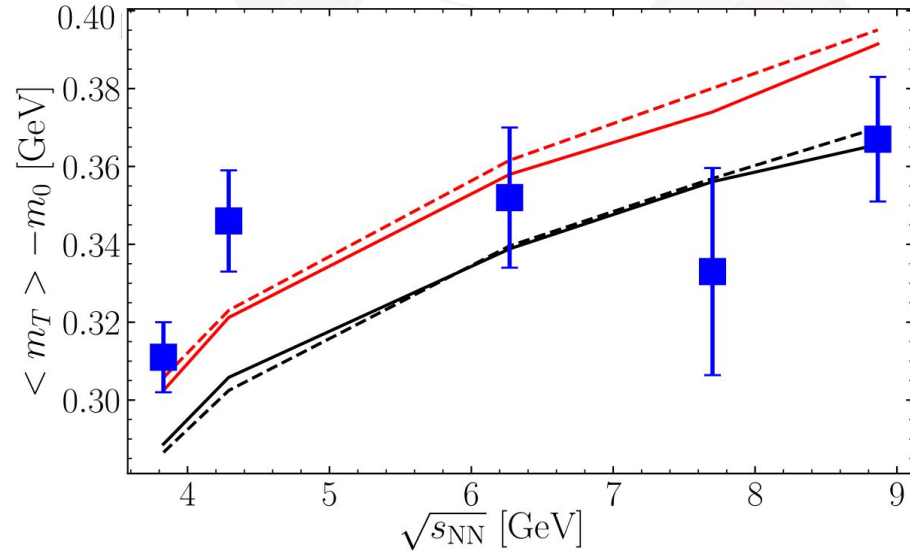
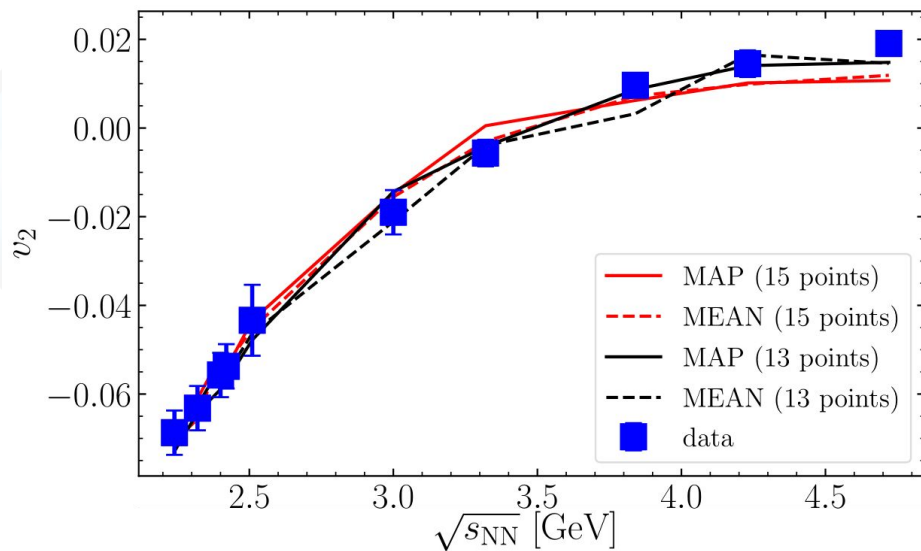


- ❖ Only 13 data points are used
 - $\langle m_T \rangle - m_0$ at 3.83, 4.29 GeV not used
- ❖ Significant differences in posterior
 - softening at $3-5n_0$
 - phase transition
- ❖ Beyond $3n_0$ strong dependence to choice of observables

More closure test results

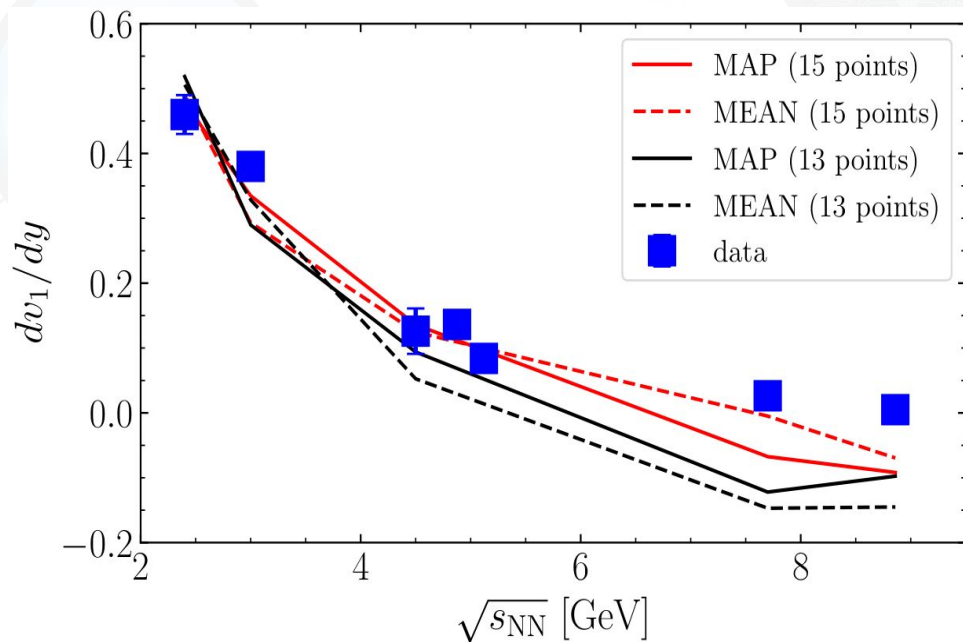


Reconstructed EoSs: v_2 , $\langle m_T \rangle - m_0$



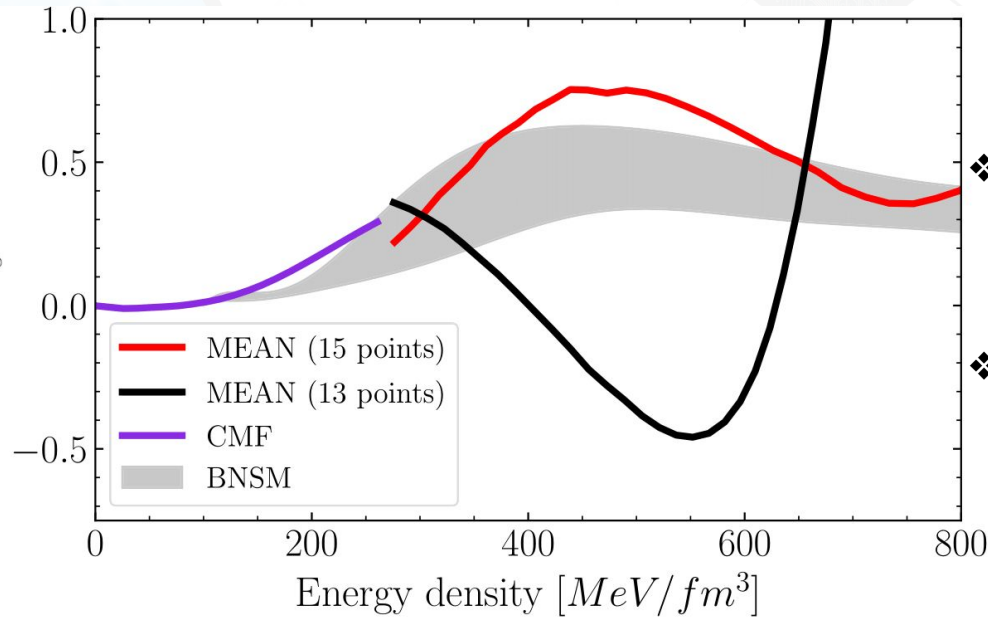
- ❖ better v_2 predictions at high energies when 2 data points are removed
 - but also results in lower $\langle m_T \rangle - m_0$
- ❖ large $\langle m_T \rangle - m_0$ values for the stiff EoS (extracted using all data points)
- ❖ possible tension in data at ~ 4 GeV!
 - Measurement uncertainty? or limitation of the model?

Reconstructed EoSs: dv_1/dy



- ❖ dv_1/dy data was **not used for inference**
 - yet consistent with reconstructed EoSs
 - especially with all 15 data points

Reconstructed EoSs: c_s^2



❖ 15 points, predicts a stiff EoS

➤ consistent with astrophysical constraints

[arXiv:2203.14974 \(2022\)](https://arxiv.org/abs/2203.14974)

■ broad peak structure

❖ 13 points, drastic drop in c_s^2

➤ first order phase transition

Summary

- ❖ Bayesian inference using polynomial parameterization of the density dependence of EoS
 - v_2 , $\langle m_t \rangle - m_0$ of protons are used for inference
- ❖ inference using all 15 data points:
 - constraints the EoS upto $4n_0$
 - stiff EoS upto $4n_0$, no phase transition
 - consistent with BNSM constraint, dv_1/dy data
- ❖ strong dependence on choice of observables for $> 3n_0$
- ❖ tension in data at ~ 4 GeV
 - measurement uncertainty or model limitation?

<https://arxiv.org/abs/2211.11670>

For stricter, robust constraints on the EoS below $4n_0$, significant improvements and consistency in flow measurements are necessary for $E_{lab} = 2-10$ A GeV



Backup slides

Microscopic transport with density dependent potential

- ❖ Non-equilibrium MD part of UrQMD is used
- ❖ UrQMD:
 - Propagation of hadrons on classic trajectories
 - stochastic binary scattering , color string formation, resonance excitation and decays
 - Imaginary part of interactions:
 - geometric interpretation of cross section
 - Experiment, detailed balance
 - Hadronic cascade
 - effective EoS of HRG with respective dof
- ❖ Real part of interactions in UrQMD
 - QMD + density dependent potential
 - Unlike other mean field models, QMD is an n-body theory of interactions between n nucleons

Microscopic transport with density dependent potential

A density dependent potential enters QMD equations

$$\dot{\mathbf{r}}_i = \frac{\partial \mathbf{H}}{\partial \mathbf{p}_i}, \quad \dot{\mathbf{p}}_i = -\frac{\partial \mathbf{H}}{\partial \mathbf{r}_i}.$$

The total hamiltonian function is sum over all hamiltonians of the i baryons

$$\mathbf{H} = \sum_i H_i, \quad H_i = E_i^{kin} + V_i$$

This include KE and total potential energy \mathbf{V}

$$\mathbf{V} = \sum_i V_i \equiv \sum_i V(n_B(r_i))$$

The change in momentum for baryon 'i' is then

$$\begin{aligned} \dot{\mathbf{p}}_i &= -\frac{\partial \mathbf{H}}{\partial \mathbf{r}_i} = -\frac{\partial \mathbf{V}}{\partial \mathbf{r}_i} \quad n_{\{i,j\}} \equiv n_B(r_{\{i,j\}}) \\ &= -\left(\frac{\partial V_i}{\partial n_i} \cdot \frac{\partial n_i}{\partial \mathbf{r}_i} \right) - \left(\sum_{j \neq i} \frac{\partial V_j}{\partial n_j} \cdot \frac{\partial n_j}{\partial \mathbf{r}_i} \right) \end{aligned}$$

The local interaction density n_B at r_k is calc by assuming each particle as gaussian wave packet

$$\begin{aligned} n_B(r_k) &= n_k = \sum_{j, j \neq k} n_{j,k} \\ &= \left(\frac{\alpha}{\pi} \right)^{3/2} \sum_{j, j \neq k} B_j \exp(-\alpha(\mathbf{r}_k - \mathbf{r}_j)^2) \\ \alpha &= 1/2L, \quad L = 2 \text{ fm}^2 \end{aligned}$$

Force on i^{th} baryon depends on change in potential energy at point r_i due to local gradient of $n_B(r_i)$ and change in potential at positions r_j of all baryons j due to change in r_i

-solved in timestep 0.2fm/c

$$P(n_B) = P_{id}(n_B) + \int_0^{n_B} n' \frac{\partial U(n')}{\partial n'} dn', \quad U(n_B) = \frac{\partial(n_B \cdot V(n_B))}{\partial n_B}$$

$$\mu'_B(n_B) = \mu_B^{id}(n_B) + U(n_B)$$

$$\epsilon(n_B) = -P(n_B) + \mu'_B n_B + sT$$