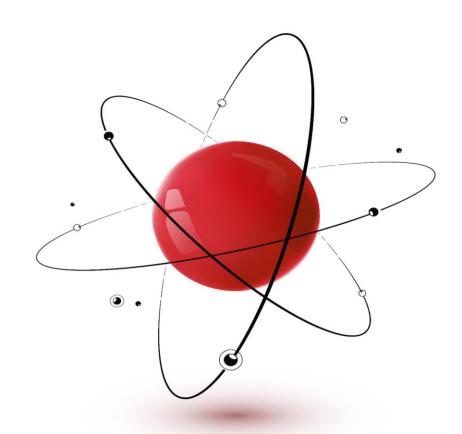
LOW-ENERGY NUCLEON STRUCTURE AT THE

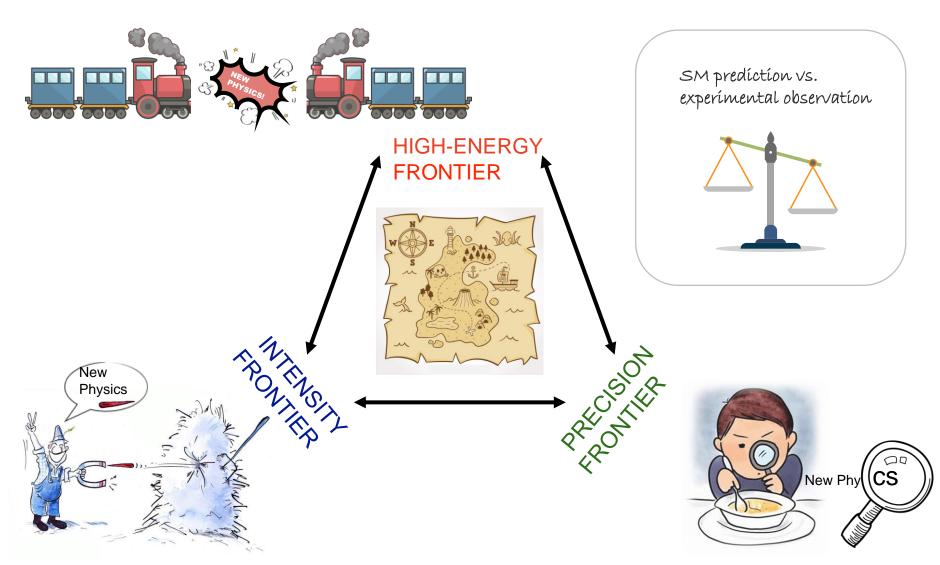
PRECISION FRONTIER



Vladyslava Sharkovska (PSI&UZH)

NEW PHYSICS SEARCHES

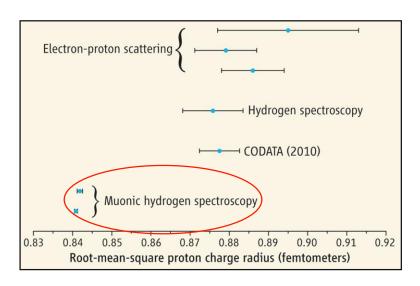
Lab searches for New Physics proceed along 3 frontiers:



PRECISION FRONTIER

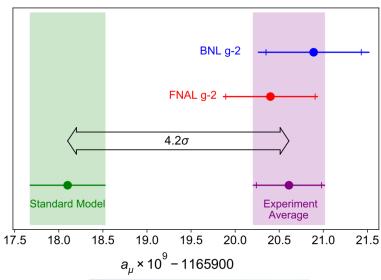
Low-energy observables measured to high precision provide stringent tests of the Standard Model (SM) of particle physics

Proton Radius





Muon g-2

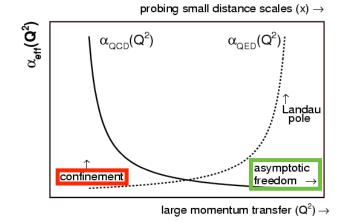




HADRONIC CORRECTIONS

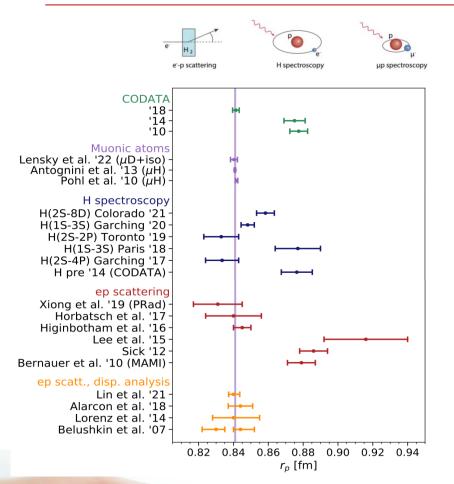
Hadrons (e.g., the proton) are often involved in "puzzles" — disagreements between SM predictions and experiments.

- Often limit precision of SM predictions
- QCD is non-perturbative at low energies



- Hadronic contributions are challenging to calculate:
 - Low-energy effective field theories, e.g., chiral perturbation theory (ChPT)
 - Lattice QCD
 - Dispersion relations basis for data-driven evaluations

STATUS OF THE PROTON RADIUS PUZZLE



- μ H Lamb shift measurements allowed for r_p extraction with unprecedented precision but gave a smaller value
- e-p scattering improved (PRad gives now the smaller value)
- H Lamb shift was remeasured and now agrees with the smaller value
- Still open issues: H(2S-8D), H(1S-3S) transitions (systematic effects?)



Considered resolved, but future efforts will still be necessary to deliver a clear verdict

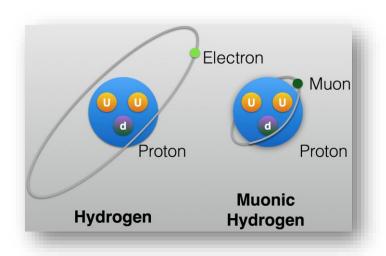
FROM PUZZLE TO PRECISION

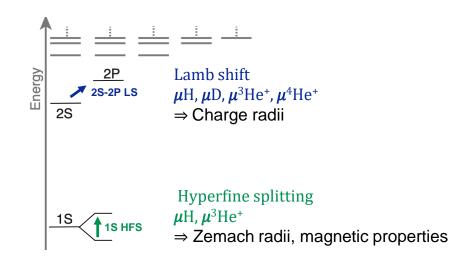
- Several experimental activities ongoing and proposed:
 - 1S hyperfine splitting in μ H (ppm accuracy) and μ He
 - Improved measurement of Lamb shift in μ H, μ D and μ He⁺ possible (× 5)
 - Medium- and high-Z muonic atoms
- ► Theory improvements needed!

$$r_p = 0.84087(12)_{\rm sys}(23)_{\rm stat}(29)_{\rm theory}$$
 fm
 $r_d = 2.12562(5)_{\rm sys}(12)_{\rm stat}(77)_{\rm theory}$ fm
 $r_{\alpha} = 1.67824(2)_{\rm sys}(13)_{\rm stat}(82)_{\rm theory}$ fm

- Charge radius extractions from Lamb shift in muonic atoms:
 - $-\mu$ H: present accuracy comparable with experimental precision
 - $-\mu D$, $\mu^3 He^+$ and $\mu^4 He^+$: present accuracy factor 5-10 worse than experimental precision

WHY MUONIC ATOMS?





Lamb shift:

$$\Delta E_{nl}(\text{LO+NLO}) = \delta_{l0} \frac{2\pi Z\alpha}{3} \frac{1}{\pi (an)^3} \left[R_E^2 - \frac{Z\alpha m_r}{2} R_{E(2)}^3 \right]$$

Fermi energy:

$$E_F(nS) = \frac{8}{3} \frac{Z\alpha}{a^3} \frac{1+\kappa}{mM} \frac{1}{n^3}$$

NLO becomes appreciable in µH

• Hyperfine splitting (HFS):



$$\Delta E_{nS}(\text{LO} + \text{NLO}) = E_F(nS) \left[1 - 2 \, Z \alpha m_r \, R_Z \right]$$

Bohr radius:

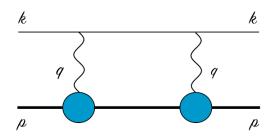
$$a = 1/(Z\alpha m_r)$$

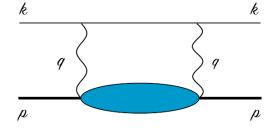
TRUCTURE EFFECTS THROUGH TPE

Proton-structure effects at subleading orders arise through multi-photon

processes

forward two-photon exchange (TPE)





Elastic finite-size recoil. 3rd Zemach moment (Lamb shift), Zemach radius (HFS)

Non-Born polarizability contribution

• "Blob" corresponds to doubly-virtual Compton scattering:

$$T^{\mu\nu}(q,p) = \left(-g^{\mu\nu} + \frac{q^{\mu}q^{\nu}}{q^2}\right) T_1(\nu,Q^2) + \frac{1}{M^2} \left(p^{\mu} - \frac{p \cdot q}{q^2} q^{\mu}\right) \left(p^{\nu} - \frac{p \cdot q}{q^2} q^{\nu}\right) T_2(\nu,Q^2) - \frac{1}{M^2} (\gamma^{\mu\nu}q^2 + q^{\mu}\gamma^{\nu\alpha}q_{\alpha} - q^{\nu}\gamma^{\mu\alpha}q_{\alpha}) S_2(\nu,Q^2)$$

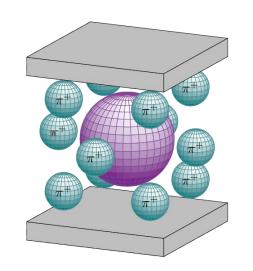
• Proton structure functions:
$$f_1(x,Q^2), f_2(x,Q^2), g_1(x,Q^2), g_2(x,Q^2)$$

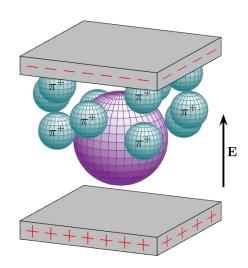
Lamb shift HFS

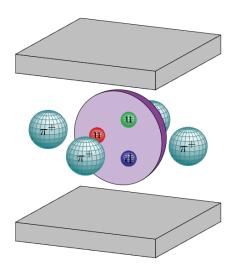
POLARIZABILITIES

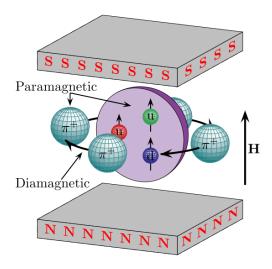
Electric dipole polarizability:

$$ec{P}=lpha_{E1}\,ec{E}$$
 induced electric dipole polarization







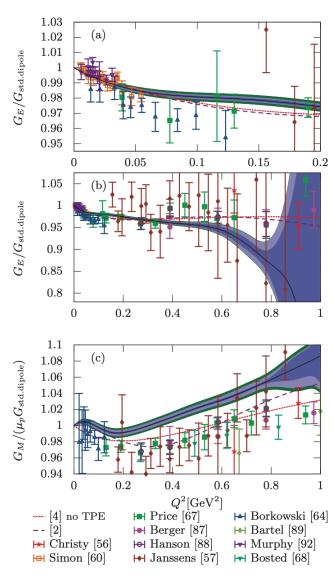


Magnetic dipole polarizability:

$$ec{P}=eta_{M1} ec{H}$$
 polarization induced by magnetic field

diamagnetic: $\beta_{M1} < 0$ paramagnetic: $\beta_{M1} > 0$

PROTON FORM FACTORS



J. C. Bernauer *et al.*, Phys. Rev. C **90**, 015206 (2014)

• Form factors (FF): Fourier transforms of charge and magnetization distributions:

$$\rho(r) = \int \frac{\mathrm{d}\boldsymbol{q}}{(2\pi)^3} G(\boldsymbol{q}^2) e^{-i\boldsymbol{q}\boldsymbol{r}}$$

Root-mean-square (rms) charge radius:

$$R_E = \sqrt{\langle r^2 \rangle_E}$$

$$\langle r^2 \rangle_E \equiv \int d\mathbf{r} \, r^2 \, \rho_E(\mathbf{r}) = -6 \frac{\mathrm{d}}{\mathrm{d}Q^2} G_E(Q^2) \Big|_{Q^2 = 0}$$

 Extraction of the proton charge radius from e-p scattering requires extrapolation of FF data to zero momentum transfer

HYPERFINE SPLITTING IN μ H

$$\Delta E_{\rm HFS}(nS) = \left[1 + \Delta_{\rm QED} + \Delta_{\rm weak} + \Delta_{\rm structure}\right] E_F(nS)$$
 with
$$\Delta_{\rm structure} = \Delta_Z + \Delta_{\rm recoil} + \Delta_{\rm pol}$$

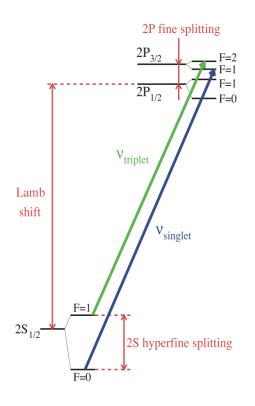
Zemach radius:

$$\Delta_Z = \frac{8Z\alpha m_r}{\pi} \int_0^\infty \frac{\mathrm{d}Q}{Q^2} \left[\frac{G_E(Q^2)G_M(Q^2)}{1+\kappa} - 1 \right] \equiv -2Z\alpha m_r R_\mathrm{Z}$$

experimental value:

$$R_Z = 1.082(37) \,\mathrm{fm}$$

A. Antognini, et al., Science 339 (2013) 417-420



- Measurements of the μH ground-state HFS planned by the CREMA, FAMU and J-PARC / Riken-RAL collaborations
- Zemach radius can help to pin down the magnetic properties of the proton

Theory: QED, ChPT, data-driven dispersion relations, ab-initio few-nucleon theories

Experiment: HFS in μ H, μ He⁺, ...

Guiding the exp.

find narrow 1S HFS transitions in μ H with the help of full theory predictions: QED, weak, finite size, polarizability

Interpreting the exp.

extract E^{TPE} , $E^{\text{pol.}}$ or R_Z

Input for datadriven evaluations form factors, structure functions, polarizabilities

Electron and Compton Scattering

Testing the theory

- discriminate between theory predictions for polarizability effect
 - disentangle R_Z & polarizability effect by combining HFS in H & μ H
- ► test HFS theory
 - combining HFS in H & μ H with theory prediction for polarizability effect
- test nuclear theories

Spectroscopy of ordinary atoms (H, He⁺)

Determine fundamental quantity

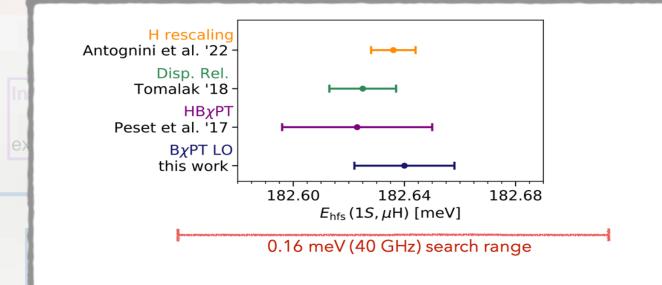
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A. Antognini, F. Hagelstein, V. Pascalutsa, Ann. Rev. Nucl. Part. 72 (2022)

Electron and Compton Scattering

The hyperfine splitting of μH (theory update):

A. Antognini, FH, V. Pascalutsa, Ann. Rev. Nucl. Part. 72 (2022)

$$E_{1S-\rm hfs} = \left[\underbrace{182.443}_{E_{\rm F}} \underbrace{+1.350(7)}_{\rm QED+weak} \underbrace{+0.004}_{\rm hVP} \underbrace{-1.30653(17) \left(\frac{r_{\rm Z}p}{\rm fm}\right) + E_{\rm F} \left(1.01656(4) \, \Delta_{\rm recoil} + 1.00402 \, \Delta_{\rm pol}\right)}_{2\gamma \; \rm incl. \; radiative \; corr.}\right] \text{meV}$$

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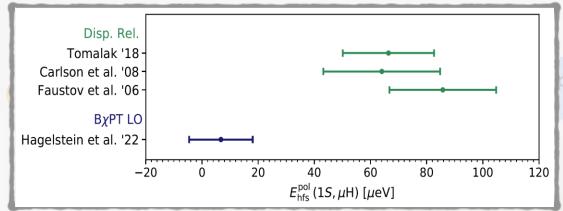
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Spectroscopy of ordinary atoms (H, He⁺(

Determine fundamental constants

Zemach radius R_2



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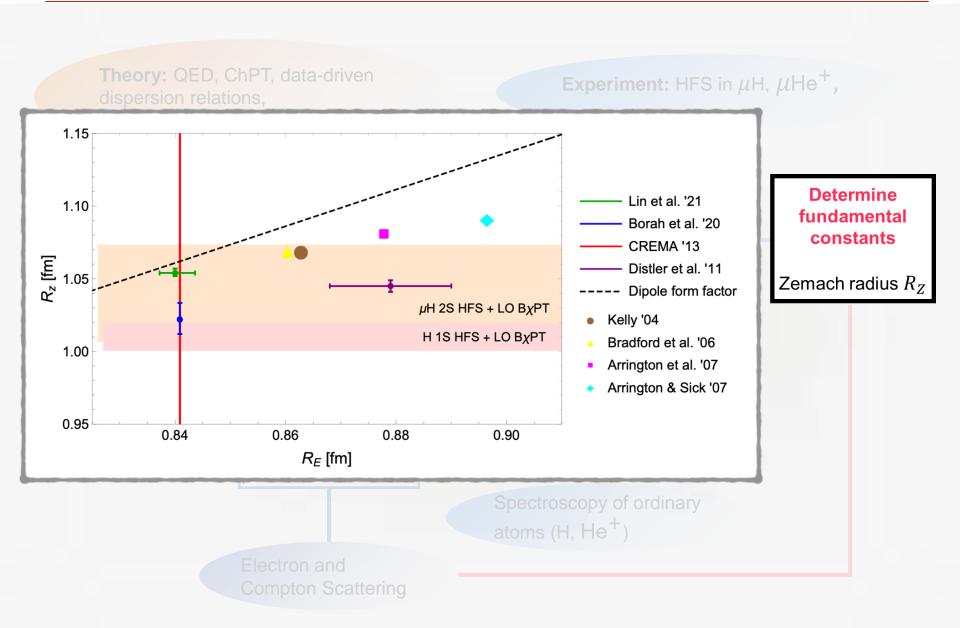
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Spectroscopy of ordinary atoms (H, He⁺)

Determine fundamental constants

Zemach radius R₂



CONCLUSIONS!

During my PhD the main interest is low-energy nucleon structure at the precision frontier

- Experimental precision presents a challenge for theory
- Further progress in studying the proton structure is important for matching the precision, but also as input for heavier nuclei
- Precise theoretical prediction for μ H 1S HFS is needed to guide the upcoming experiment
- Future plans: improve the (data-driven) dispersive prediction for Δ_{pol} contribution to HFS in (μ) H based on the global analysis of nuclear structure functions

Thank you for having me here!

