

Test of lepton flavour universality in $B^+ \rightarrow K^+ l^+ l^-$ decays in high di-lepton invariant mass squared region

Zurich PhD Seminars 2022 (2023) 26 January 2023

Speaker: Vadym Denysenko

vadym.denysenko@physik.uzh.ch vadym.denysenko@cern.ch University of Zurich Zurich, Switzerland







- A theory that:
 - Contains the knowledge of all the elementary blocks
 - Explains the interaction between those elementary particles
 - Works everywhere and all the time

If only there was THE theory

Zurich PhD Seminars 2022 (2023) | Test of LFU in $B^+ \to K^+ l^+ l^-$ decays in high q^2 region | Vadym Denysenko



The Standard Model (SM) combines the vast majority of our knowledge of Particle Physics.

The Standard Model







- The Standard Model (SM) combines the vast majority of our knowledge of Particle Physics.
- Luckily, it is **excellent** at explaining an extensive range of physics processes.

The Standard Model





- The Standard Model (SM) combines the vast majority of our knowledge of Particle Physics.
- Luckily, it is excellent at explaining an extensive range of physics processes.
- Unfortunately, it does not explain all of them and does not answer **ALL** our questions.

The Standard Model







Looking Beyond the Standard Model with the SMEFT

"...the direct method may be used...but indirect methods will be needed in order to secure victory....' "The direct and the indirect lead on to each other in turn. It is like moving in a circle....' Who can exhaust the possibilities of their combination?"

Sun Tzu, The Art of War





ElliS ecture



The search for New Physics (NP) can be done in a couple of ways:

Direct searches

Focus on searches for particles that are directly created in high-energy beams

Б.

Beyond the SM

Indirect searches

Focus on searches for inconsistencies between experimental results and theoretical predictions

Zurich PhD Seminars 2022 (2023) | Test of LFU in $B^+ \to K^+ l^+ l^-$ decays in high q^2 region | Vadym Denysenko

Б.







Beyond the SM





26/01/2023

The LHCb experiment

Exploiting the indirect ways to search for New Physics, LHCb is primarily focused on beauty quark decays.

Zurich PhD Seminars 2022 (2023) | Test of LFU in $B^+ \rightarrow K^+ l^+ l^-$ decays in high q^2 region | Vadym Denysenko



Forward-arm spectrometer;



The LHCb experiment

Exploiting the indirect ways to search for New Physics, LHCb is primarily focused on beauty quark decays:

Zurich PhD Seminars 2022 (2023) | Test of LFU in $B^+ \to K^+ l^+ l^-$ decays in high q^2 region | Vadym Denysenko



- Forward-arm spectrometer;
- Excellent vertexing and PID



The LHCb experiment

Exploiting the indirect ways to search for New Physics, LHCb is primarily focused on beauty quark decays:

Zurich PhD Seminars 2022 (2023) | Test of LFU in $B^+ \to K^+ l^+ l^-$ decays in high q^2 region | Vadym Denysenko



Lepton flavour universality test

In the Standard Model (SM), the coupling of gauge bosons to leptons is independent of lepton flavour. • To test it, the ratio of $B^+ \to K^+ l^+ l^-$ branching fractions (a theoretically clean observable):

$$R_{K} = \frac{\int_{q_{min}^{2}}^{q_{max}^{2}} \frac{d\mathscr{B}(B^{+} \to K^{+}\mu^{+}\mu^{-})}{dq^{2}} dq^{2}}{\int_{q_{min}^{2}}^{q_{max}^{2}} \frac{d\mathscr{B}(B^{+} \to K^{+}e^{+}e^{-})}{dq^{2}} dq^{2}} \cong 1, \quad q^{2}$$

- Any significant deviation a hint of New Physics (NP):
 - The latest results show $< 1\sigma$ compatibility with SM [arXiv:2212.09152]

 \equiv dilepton invariant mass squared



• The ratio of $B^+ \to K^+ l^+ l^-$ branching fractions R_K as an experimental observable:

Problems due to the difference in leptons detection at LHCb.





Zurich PhD Seminars 2022 (2023) | Test of LFU in $B^+ \to K^+ l^+ l^-$ decays in high q^2 region | Vadym Denysenko



• The ratio of $B^+ \to K^+ l^+ l^-$ branching fractions (a theoretically clean observable):

$$\begin{split} R_{K} &= \left(\frac{N(B^{+} \to K^{+}\mu^{+}\mu^{-})}{\varepsilon(B^{+} \to K^{+}\mu^{+}\mu^{-})}\right) \left/ \left(\frac{N(B^{+} \to K^{+}e^{+}e^{-})}{\varepsilon(B^{+} \to K^{+}e^{+}e^{-})}\right) = \\ &= \frac{\left(\frac{N(B^{+} \to K^{+}\mu^{+}\mu^{-})}{\varepsilon(B^{+} \to K^{+}\mu^{+}\mu^{-})}\right) \left/ \left(\frac{N(B^{+} \to K^{+}e^{+}e^{-})}{\varepsilon(B^{+} \to K^{+}e^{+}e^{-})}\right)}{\left(\frac{N(B^{+} \to K^{+}J/\psi(\to \mu^{+}\mu^{-}))}{\varepsilon(B^{+} \to K^{+}J/\psi(\to e^{+}e^{-}))}\right)}\right) \left/ \left(\frac{N(B^{+} \to K^{+}J/\psi(\to e^{+}e^{-}))}{\varepsilon(B^{+} \to K^{+}J/\psi(\to e^{+}e^{-}))}\right) \right. \end{split}$$

$$= \left(\frac{N(B^+ \to K^+ \mu^+ \mu^-)}{\varepsilon(B^+ \to K^+ \mu^+ \mu^-)}\right) \left/ \left(\frac{N(B^+ \to K^+ e^+ e^-)}{\varepsilon(B^+ \to K^+ e^+ e^-)}\right) = \left(\frac{N(B^+ \to K^+ \mu^+ \mu^-)}{\varepsilon(B^+ \to K^+ \mu^+ \mu^-)}\right) \left/ \left(\frac{N(B^+ \to K^+ e^+ e^-)}{\varepsilon(B^+ \to K^+ e^+ e^-)}\right) - \left(\frac{N(B^+ \to K^+ J/\psi(\to \mu^+ \mu^-))}{\varepsilon(B^+ \to K^+ J/\psi(\to e^+ e^-))}\right) \right/ \left(\frac{N(B^+ \to K^+ J/\psi(\to e^+ e^-))}{\varepsilon(B^+ \to K^+ J/\psi(\to e^+ e^-))}\right)$$

modes allows us to have control over electron/muon differences.



□ Using a "double ratio" between non-resonant (rare) $B^+ \to K^+ l^+ l^-$ and resonant $B^+ \to K^+ J/\psi$ (→ $l^+ l^-$)



• The ratio of $B^+ \to K^+ l^+ l^-$ branching fractions (a theoretically clean observable):

$$\begin{split} R_{K} &= \left(\frac{N(B^{+} \rightarrow K^{+}\mu^{+}\mu^{-})}{\varepsilon(B^{+} \rightarrow K^{+}\mu^{+}\mu^{-})}\right) \left/ \left(\frac{N(B^{+} \rightarrow K^{+}e^{+}e^{-})}{\varepsilon(B^{+} \rightarrow K^{+}e^{+}e^{-})}\right) = \\ &= \frac{\left(\frac{N(B^{+} \rightarrow K^{+}\mu^{+}\mu^{-})}{\varepsilon(B^{+} \rightarrow K^{+}\mu^{+}\mu^{-})}\right) \left/ \left(\frac{N(B^{+} \rightarrow K^{+}e^{+}e^{-})}{\varepsilon(B^{+} \rightarrow K^{+}e^{+}e^{-})}\right)}{\left(\frac{N(B^{+} \rightarrow K^{+}J/\psi(\rightarrow \mu^{+}\mu^{-}))}{\varepsilon(B^{+} \rightarrow K^{+}J/\psi(\rightarrow \mu^{+}\mu^{-}))}\right) \right/ \left(\frac{N(B^{+} \rightarrow K^{+}J/\psi(\rightarrow e^{+}e^{-}))}{\varepsilon(B^{+} \rightarrow K^{+}J/\psi(\rightarrow e^{+}e^{-}))}\right) \\ &= r_{J/\psi} \quad \text{LFU within 0.4\%} \end{split}$$

$$= \left(\frac{N(B^+ \to K^+ \mu^+ \mu^-)}{\varepsilon(B^+ \to K^+ \mu^+ \mu^-)}\right) \left/ \left(\frac{N(B^+ \to K^+ e^+ e^-)}{\varepsilon(B^+ \to K^+ e^+ e^-)}\right) = \left(\frac{N(B^+ \to K^+ \mu^+ \mu^-)}{\varepsilon(B^+ \to K^+ \mu^+ \mu^-)}\right) \left/ \left(\frac{N(B^+ \to K^+ e^+ e^-)}{\varepsilon(B^+ \to K^+ e^+ e^-)}\right) \right.$$

$$= \left(\frac{N(B^+ \to K^+ J/\psi(\to \mu^+ \mu^-))}{\varepsilon(B^+ \to K^+ J/\psi(\to e^+ e^-))}\right) \left| \left(\frac{N(B^+ \to K^+ J/\psi(\to e^+ e^-))}{\varepsilon(B^+ \to K^+ J/\psi(\to e^+ e^-))}\right)\right| = r_{J/\psi} \quad \text{LFU within 0.4\%}$$

modes allows us to have control over electron/muon differences.



□ Using a "double ratio" between non-resonant (rare) $B^+ \to K^+ l^+ l^-$ and resonant $B^+ \to K^+ J/\psi (\to l^+ l^-)$



However, the result is statistically limited







Only have measurements from the *B*-factory experiments at high q^2 (such as Belle and BaBar) **Question:** Why hasn't LHCb made measurements of R_K at high q^2 ?







'High q^2 ' measurements:

Statistically independent measurements with a large yield;

Question: Why hasn't LHCb made measurements of R_K at high q^2 ?







'High q^2 ' measurements:

Statistically independent measurements with a large yield;

Question: Why hasn't LHCb made measurements of R_K at high q^2 ?

Bremsstrahlung effects are significant, thus **Answer**: the analysis becomes more difficult to perform.







Bremsstrahlung recovery

Even after Bremsstrahlung recovery, we see large differences between di-electron and di-muon final states:





Due to imperfect Bremsstrahlung recovery signal is 'washed out' for the electron channel.



Bremsstrahlung recovery





Narrow charmonium resonances have larger tails for the electron channel due to the imperfect Bremsstrahlung recovery and the energy resolution of the ECAL.



Bremsstrahlung recovery







Partially reconstructed backgrounds are not well separated from the signal for the electron case.



Bremsstrahlung recovery





The available phase space distorts combinatorial and partially reconstructed backgrounds.



Bremsstrahlung recovery





- **Di-electron** final state is complex due to the $\psi(2S)$ resonance leakage (Bremsstrahlung recovery smears out resonances)
- Di-muon final state is less problematic.



High q^2 problem





• 'Central q^2 ' measurements

LHCb [arXiv:2212.09153] R_K central- q^2 Counts Counts $\sqrt{32}$ MeV/ c^2) 200 100 100LHCb Data $9\,\mathrm{fb}^{-1}$ Total Signal Combinatorial Misidentification Partially reconstructed $B^+ \to K^+ J/\psi (\to e^+ e^-)$ $\left(\right)$ 5000 6000 5500 $m(K^+e^+e^-)$ [MeV/ c^2]

High q^2 problem

• 'High q^2 ' measurements: It becomes challenging to statistically separate signal from background.





Resonances are smeared due to the wrong Bremsstrahlung recovery.



High q^2 problem





Since resonances are smeared due to the wrong Bremsstrahlung recovery:



 $\square q_{no Brem.}^2 \equiv$ dilepton invariant mass (without adding Bremsstrashlung photons) squared

High q^2 solution

Since resonances are smeared due to the wrong Bremsstrahlung recovery:



 10^{-1}

 10^{-2}

 10^{-3}

 10^{-4}

10-5

1.2

1.4

1.6

1.8

2.0

*q*² [MeV² / c⁴]

 q^2

High q^2 solution



 $q_{no Brem.}^2$ cut based signal selection:

Loose ~ 50% of signal compared to q^2 cut.



High q^2 solution



Status: Starting to put different components together Expect around 800 signal events (~ $\frac{1}{2}$ of Central q^2)

- Signal: sum of 3 DCBs
 - Signal shape parameters fixed from simulation;
 - Brem. fractions are gaussian constrained to the fraction observed in MC;
 - Mean shift and scales are fixed from the simulation;
- Part. Reco.: KDE
- Comb.: Exponential (shape to be studied)

The total model was obtained by summing the PDFs with the relative fraction obtained from the simulation.

Electron Sensitivity studies



- PID cut corrected with weights (DLLe > 4)
- Combinatorial BDT cut and fit range

Zurich PhD Seminars 2022 (2023) | Test of LFU in $B^+ \rightarrow K^+ l^+ l^-$ decays in high q^2 region | Vadym Denysenko

The latest results in $B^+ \to K^+ l^+ l^-$ decays have been **shown** by LHCb to be consistent with Standard

• Measurement of the LFU ratio R_K in the 'high q^2 ' region:

- **first** LHCb measurement in this particular region;
- using the **entire** available dataset (9 fb⁻¹);
- statistically independent result;
- **high** yield;
- **complementary** phase space;

Work in progress. Stay tuned for further updates :)



Model, but the result is statistically limited. To further investigate further, more studies are required.





The latest results in $B^+ \rightarrow$ Model, but the result is sta

- Measurement of the LFU
 - first LHCb measurement
 - using the entire available
 - statistically independent
 - high yield;
 - complementary phase s

Work in progress. Sta



Summary

to be consistent with Standard er, more studies are required.

Zurich PhD Seminars 2022 (2023) | Test of LFU in $B^+ \to K^+ l^+ l^-$ decays in high q^2 region | Vadym Denysenko





Thanks for your attention

Combining the measured
$$R_K$$
 value with $\mathcal{B}(B^+ \rightarrow B)$

$$\int_{q^2=1.1 \text{ GeV}^2}^{q^2=6 \text{ GeV}^2} \frac{d\mathcal{B}(B^+ \to K^+ e^+ e^-)}{dq^2}$$

• Suggesting that electrons are more SMlike than muons.

D. Lancierini (Universität Zürich) talk





r	h
۰.	