

Goldstino condensation

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Plan:

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 $\rightarrow~$ Goldstino condensation at large N

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Motivation

Dark Energy in String Theory/Supergravity means

 $\rho_{DE} \sim V_{4D} = f^2 - 3m_{3/2}^2 > 0 \implies \text{Supersymmetry Breaking.}$

We will focus on non-linear SUSY because:

1. NL-SUSY underlines many EFTs with broken SUSY. See e.g. Dudas, Dall'Agata, FF '16, Dall'Agata, FF, Cribiori '17

 For "anti-brane uplifts" the supersymmetry breaking is described by sectors with non-linear supersymmetry. See e.g. Kallosh, Wrase '14, Bergshoeff, Dasgupta, Kallosh, Van Proeyen, Wrase '15, Dasgupta, Emelin and McDonough '16 Goldstino: An indispensable part of uplifts

Salient features of NL-SUSY in Volkov–Akulov model ('73)

The (N=1) goldstino Lagrangian

$$\mathcal{L}_{VA} = -\mathbf{f}^2 - i\overline{G}\partial G + \frac{1}{4f^2}\overline{G}^2\partial^2 G^2 - \frac{1}{16f^6}G^2\overline{G}^2\partial^2 G^2\partial^2\overline{G}^2,$$

and generates the uplift when coupled to SG.

$$\delta G_{\alpha} = -f\xi_{\alpha} - (i/2f)\partial_{\alpha\dot{\alpha}} G^2 \overline{\xi}^{\dot{\alpha}} + \dots$$

We want to understand if the non-linear terms are spectators or they have any physical significance/impact. Goldstino condensation: A threat to uplifts

 If there is a condensate of the form

$$\langle i\overline{G}\partial G\rangle \neq 0$$
,

then this may potentially ruin uplifts.



(uplift Kachru, Kallosh, Linde, Trivedi '03)

- Due to the non-linear terms, the goldstini may produce such condensates.
- We will discuss the simplest model where goldstino condensation can be studied convincingly.

Goldstino condensation at large N

FF, Matteo Morittu '22



We can work with N copies of 4D Dirac goldstini

$$\mathcal{L} = -\mathsf{N}f^2 \det[\mathcal{A}_m{}^a] = -\mathsf{N}f^2 + i\overline{G}{}^A \partial \!\!\!/ G^A + \dots$$

with A = 1, ..., N, with cut-off $\Lambda \ll \sqrt{f}$, where

$$\mathcal{A}_m{}^a = \delta_m^a + \frac{i}{2Nf^2} \left(\overline{G}^A \gamma^a \partial_m G^A - \partial_m \overline{G}^A \gamma^a G^A \right) \,,$$

and
$$\delta G^A = -f\xi^A + \dots$$

$$\mathcal{L} = -\mathsf{N}f^2 \det[\mathbf{e}_m{}^a] + \mathsf{N}f^2 C_a{}^m(\mathbf{e}_m{}^a - \mathcal{A}_m{}^a),$$

making the path integral Gaussian in the fermions.

 We can perform the Gaussian integral over the fermions, to find formally

$$Z_{F} \sim \left(\det[\mathit{iC}_{a}{}^{m}\gamma^{a}\partial_{m}]
ight)^{\mathsf{N}} \, ,$$

and then the action for the bosons reads

$$S = \mathbf{N} \times \left\{ -f^2 \int d^4 x \left[\det[\mathbf{e}_n^{\ b}] - C_a^{\ m} (\mathbf{e}_m^{\ a} - \delta_m^{\ a}) \right] -i \operatorname{tr} \log \left[i C_a^{\ m} \gamma^a \partial_m \right] \right\}.$$

▶ Restore momentarily \hbar and notice an "effective" \hbar as

$$\hbar_{eff} = \hbar/N$$
,

infinitesimal at large N and theory behaves "classically".

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We want to find new stationary points therefore we can focus on

$$C_a^{\ m} = (1+h) \, \delta_a^{\ m}, \quad e_m^{\ a} = (1-\phi) \, \delta_m^{\ a}.$$

The potential reads

$$V(h,\phi) = \mathsf{N}\left\{f^{2}\left[(1-\phi)^{4} + 4(1+h)\phi\right] - \frac{\Lambda^{4}}{16\pi^{2}}\log\left[(1+h)^{2}\right]\right\}$$

We can readily eliminate h to get an effective V(\u03c6); an "effective potential" for the condensate to search for critical points.



 We interpret this as a goldstino condensate because classically

$$\phi = -\frac{i}{8Nf^2} \left(\overline{G}^A \gamma^m \partial_m G^A - \partial_m \overline{G}^A \gamma^m G^A \right) \,.$$

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Outlook

- → Large N: GC quite robust under higher derivative goldstino self-couplings / typical matter coupling of NL-SUSY.
- → Using "exact RG flow", we are studying the 4D goldstino condensate also at N=1, which is physically more relevant. Dall'Agata, Emelin, FF, Morittu '22, '23
- $\rightarrow\,$ Our findings apply to anti-branes and can ruin uplifts, and inflationary models.
- → Our results resonate with gravitino condensation. *E.g.* Jasinschi, Smith '83, Alexandre, Houston, Mavromatos '13-'15

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Thank you

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