

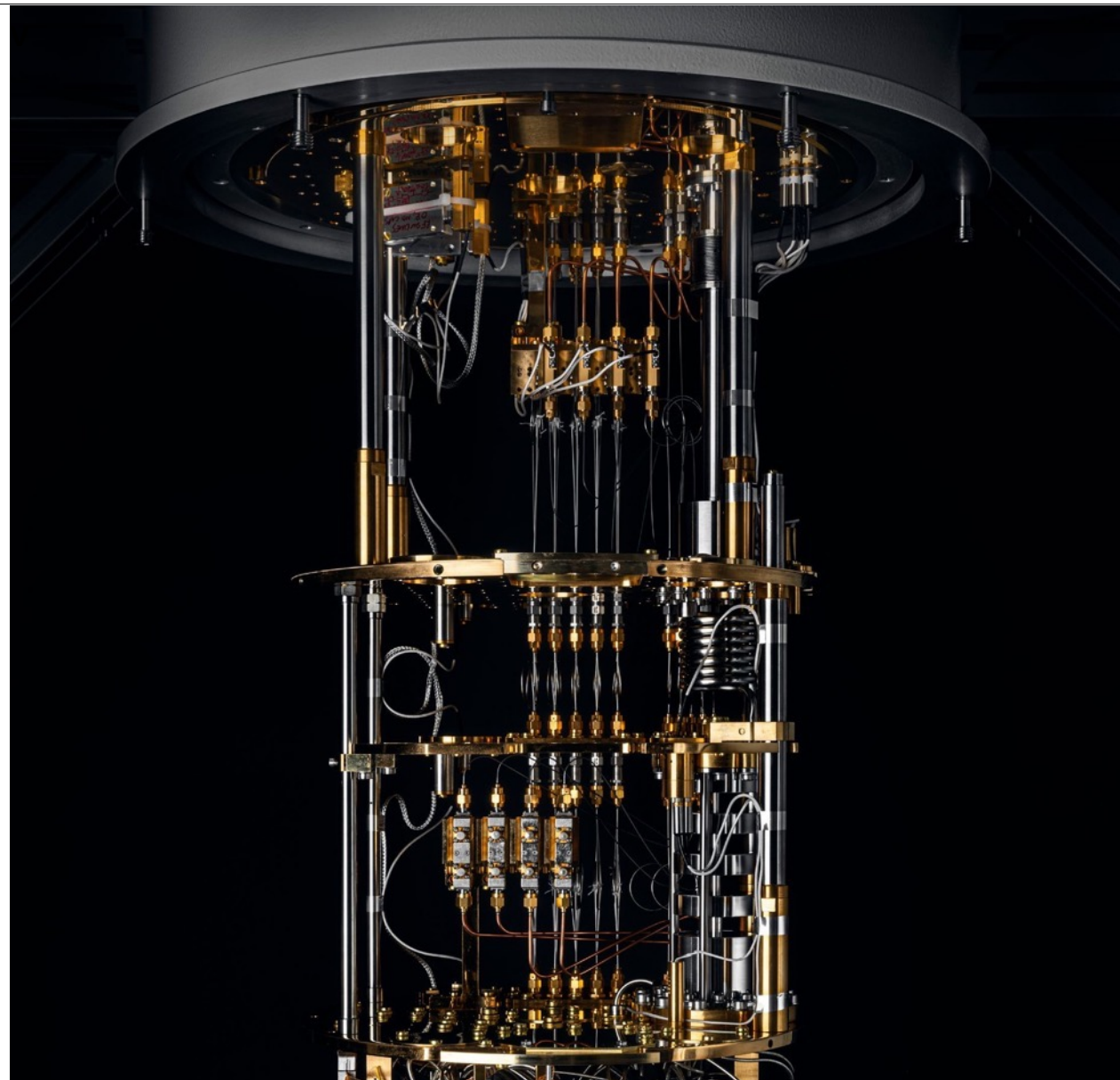
WE BUILD QUANTUM COMPUTERS

Low-depth simulations of fermionic systems on square-grid quantum hardware

Manuel Algaba,

P.V. Sriluckshmy, M. Leib, F. Šimkovic

arXiv:2302.01862 arXiv:2303.04498

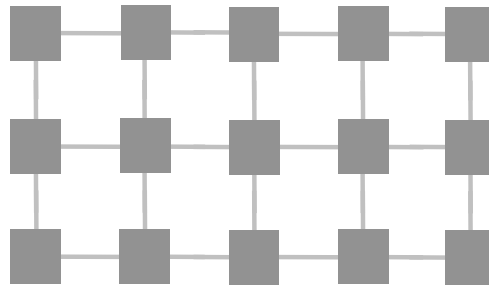


IQM

Objective:

Simulate a **single Trotter step of a fermionic system** on a **realistic quantum computer** in minimal circuit depth.

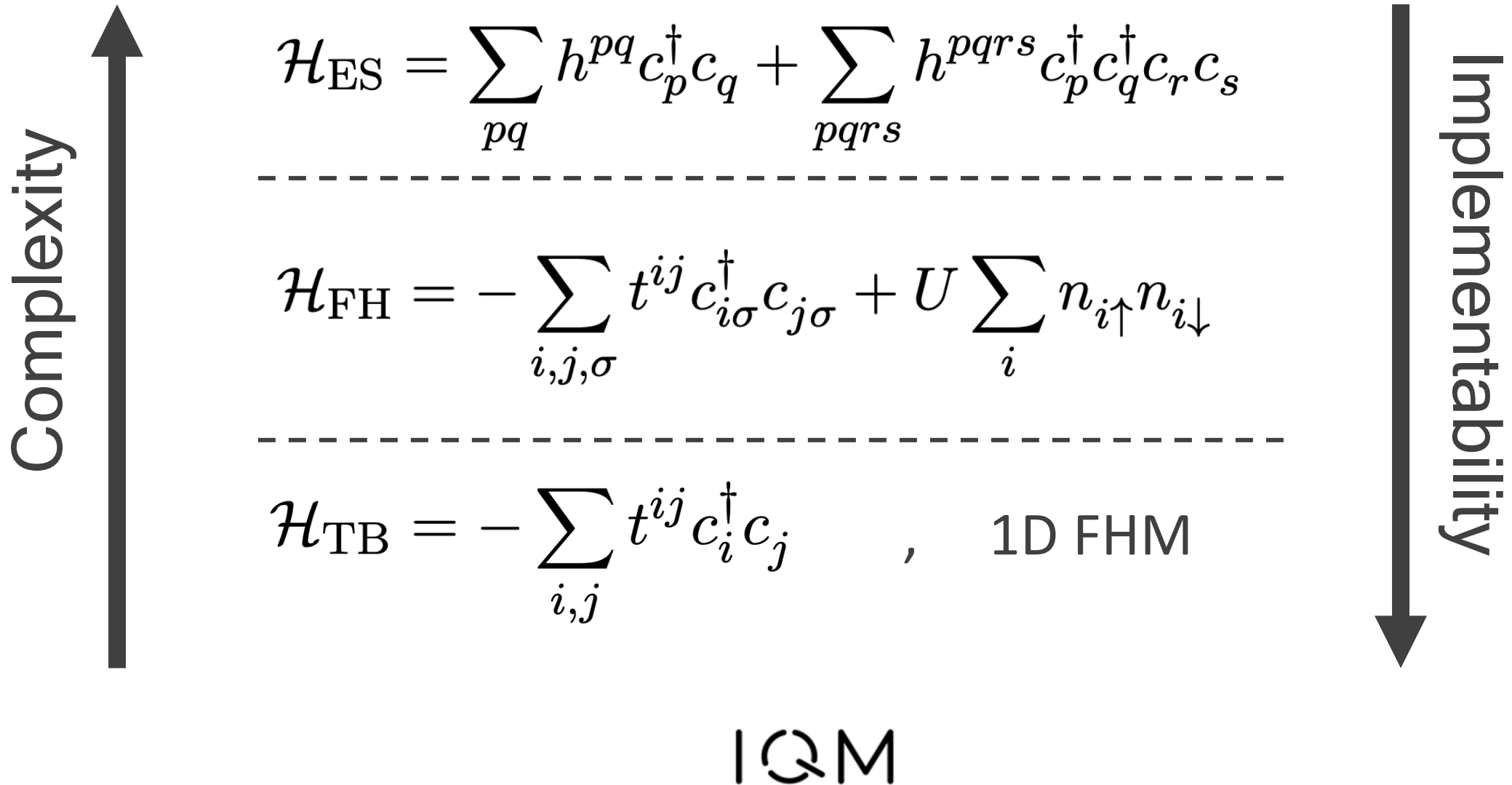
$$e^{-it\mathcal{H}} = \left(\prod_j e^{-i\frac{t}{N}\mathcal{H}_j} \right)^N$$



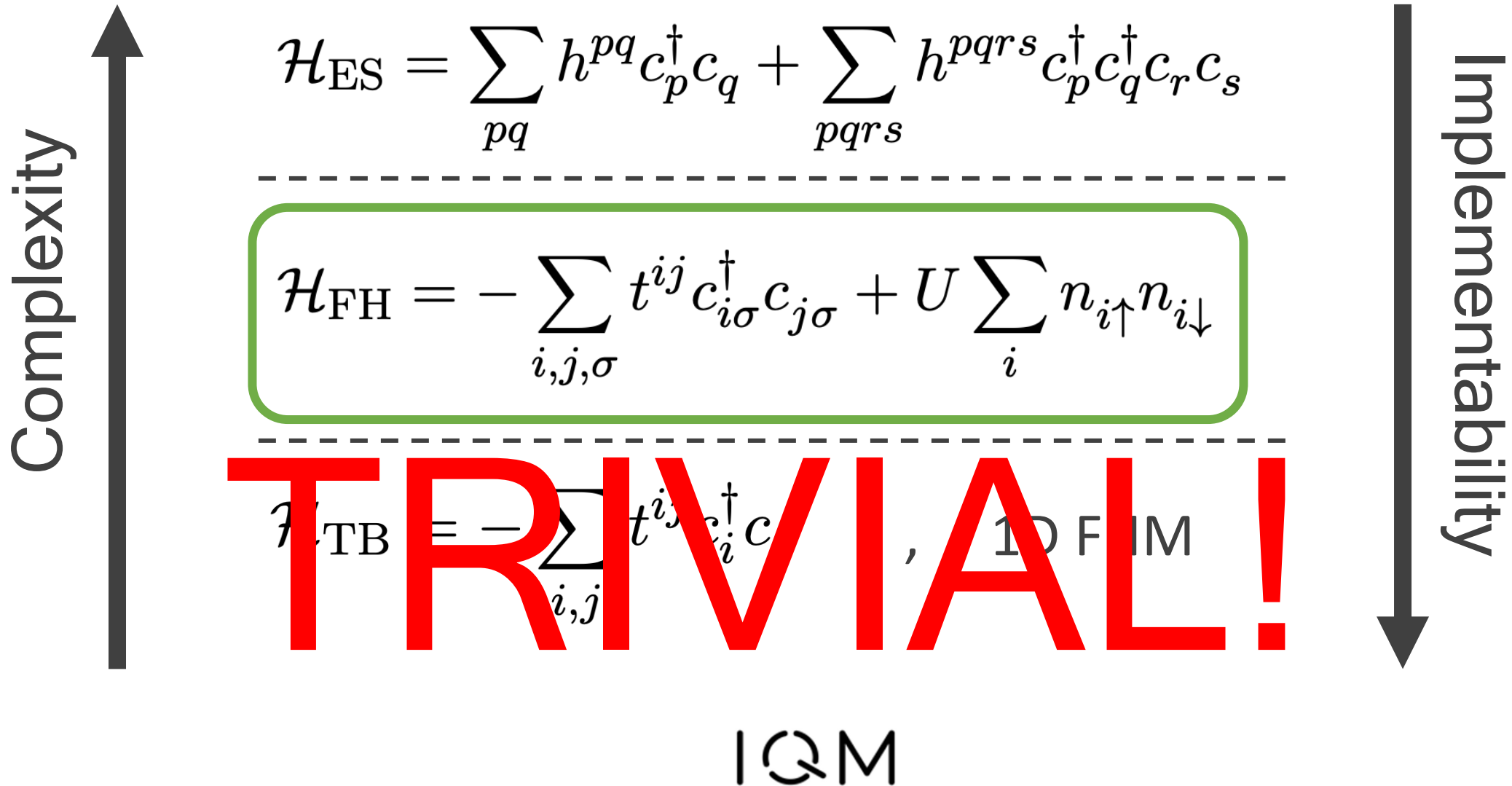
$$\text{fSIM}_{ij}(\theta, \phi) = e^{i\frac{\theta}{2}(X_i X_j + Y_i Y_j) + i\frac{\phi}{4}(Z_i + Z_j - Z_i Z_j)}$$

IQM

1. Fermionic Models



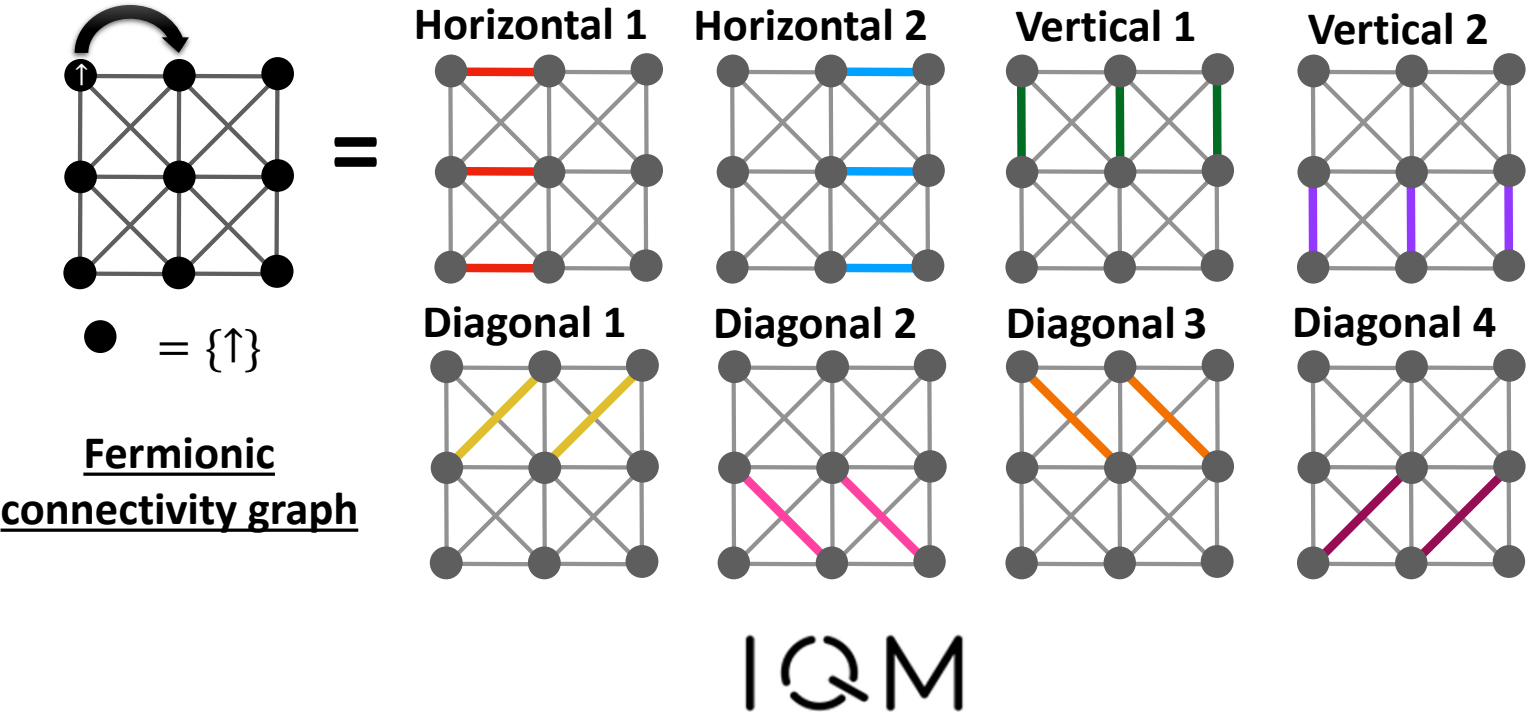
1. Fermionic Models



2. Fermion-to-qubit mappings

Tight Binding Hamiltonian:

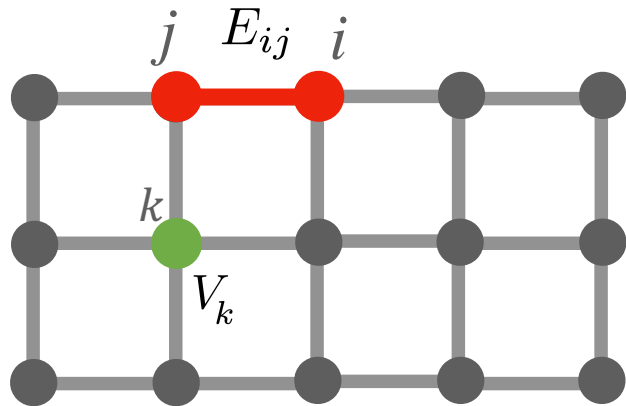
$$\mathcal{H}_{\text{TB}} = \sum_{i,j} t^{ij} c_i^\dagger c_j$$



2. Fermion-to-qubit mappings

$$\mathcal{H}_{\text{TB}} = \sum_{i,j} t^{ij} c_i^\dagger c_j$$

$$\begin{aligned} \{c_i, c_j^\dagger\} &= c_i c_j^\dagger + c_j^\dagger c_i = \delta_{ij} \\ \{c_i^\dagger, c_j^\dagger\} &= \{c_i, c_j\} = 0 \end{aligned}$$



Edge and vertex operators:

$$\begin{aligned} \{E_{ij}, V_i\} &= \{E_{ij}, E_{jk}\} = 0 \\ [E_{ij}, E_{kl}] &= [E_{ij}, V_k] = [V_i, V_j] = 0 \end{aligned}$$

Hopping operators:

$$\boxed{c_j^\dagger c_k + c_k^\dagger c_j} \rightarrow \boxed{\frac{i}{2}(V_k - V_j)E_{jk}}$$

Fermions Qubits

2. Fermion-to-qubit mappings

$$\mathcal{H}_{\text{TB}} = \sum_{i,j} t^{ij} c_i^\dagger c_j$$

$$\{c_i, c_j^\dagger\} = c_i c_j^\dagger + c_j^\dagger c_i = \delta_{ij}$$

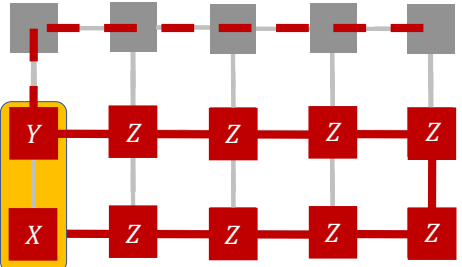
$$\{c_i^\dagger, c_j^\dagger\} = \{c_i, c_j\} = 0$$

$$\{E_{ij}, V_i\} = \{E_{ij}, E_{jk}\} = 0$$

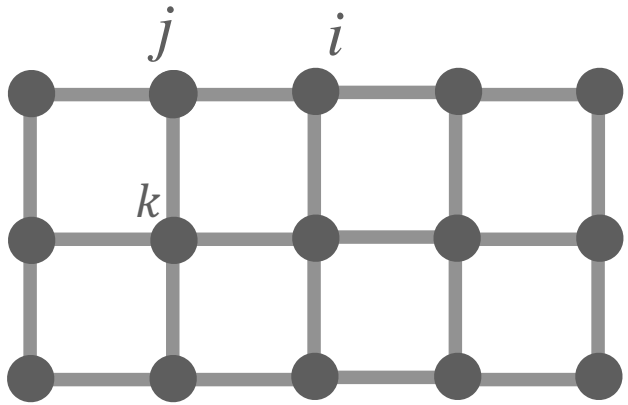
$$[E_{ij}, E_{kl}] = [E_{ij}, V_k] = [V_i, V_j] = 0$$

$$c_j^\dagger c_k + c_k^\dagger c_j \rightarrow \frac{i}{2}(V_k - V_j)E_{jk}$$

Jordan-Wigner mapping



$$c_i^\dagger c_j + c_j^\dagger c_i = X_i Z_{i-1} \dots Z_{j+1} Y_j + Y_i Z_{i-1} \dots Z_{j+1} X_j$$



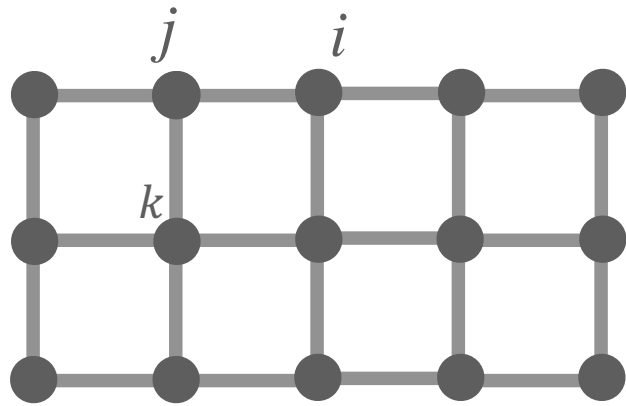
IQM

2. Fermion-to-qubit mappings

$$\mathcal{H}_{\text{TB}} = \sum_{i,j} t^{ij} c_i^\dagger c_j$$

$$\{c_i, c_j^\dagger\} = c_i c_j^\dagger + c_j^\dagger c_i = \delta_{ij}$$

$$\{c_i^\dagger, c_j^\dagger\} = \{c_i, c_j\} = 0$$



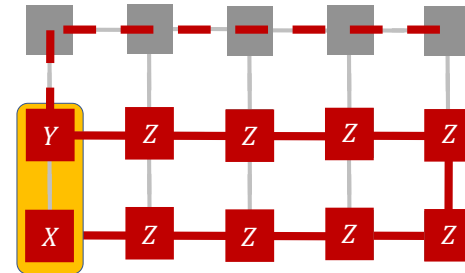
Black arrows

$$\{E_{ij}, V_i\} = \{E_{ij}, E_{jk}\} = 0$$

$$[E_{ij}, E_{kl}] = [E_{ij}, V_k] = [V_i, V_j] = 0$$

$$c_j^\dagger c_k + c_k^\dagger c_j \rightarrow \frac{i}{2}(V_k - V_j)E_{jk}$$

Jordan-Wigner mapping



$$c_i^\dagger c_j + c_j^\dagger c_i = X_i Z_{i-1} \dots Z_{j+1} Y_j + Y_i Z_{i-1} \dots Z_{j+1} X_j$$

Local mappings

- Bravyi-Kitaev mapping
- Verstraete-Cirac mapping
- Derby-Klassen mapping

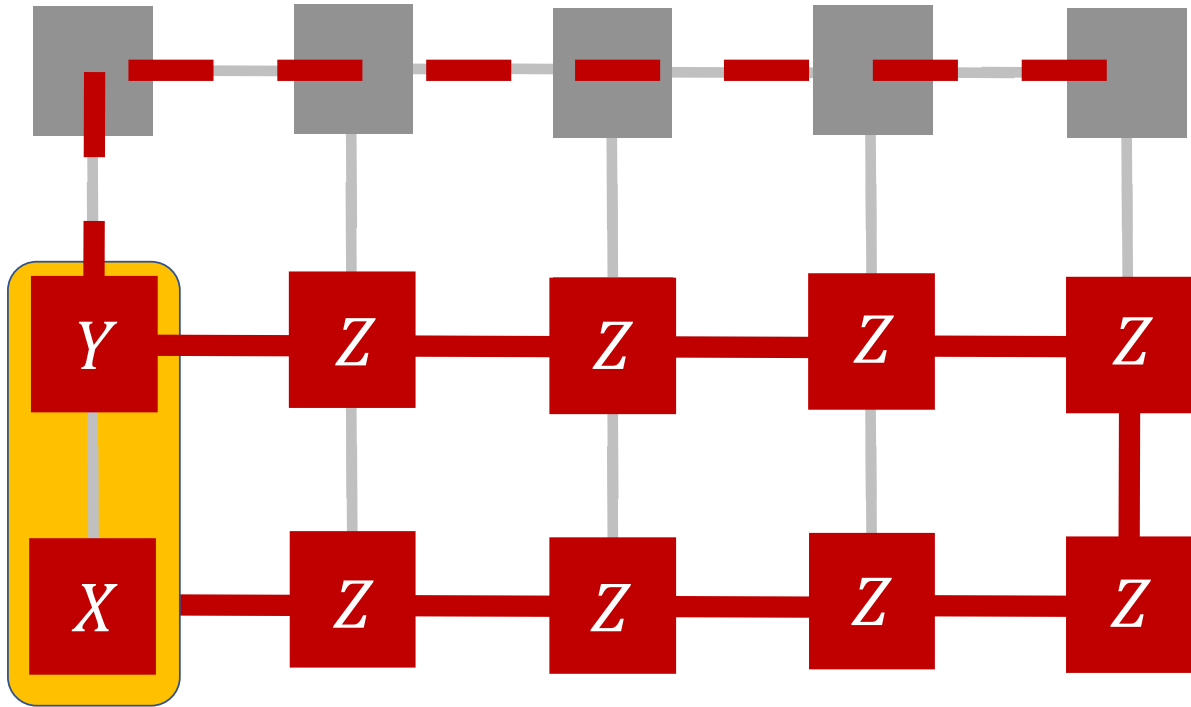
⋮

IQM

2. Fermion-to-qubit mappings

$$\begin{aligned} \{E_{ij}, V_i\} &= \{E_{ij}, E_{jk}\} = 0 \\ [E_{ij}, E_{kl}] &= [E_{ij}, V_k] = [V_i, V_j] = 0 \\ c_j^\dagger c_k + c_k^\dagger c_j &\rightarrow \frac{i}{2}(V_k - V_j)E_{jk} \end{aligned}$$

- Local mappings
- Bravyi-Kitaev mapping
- Verstraete-Cirac mapping
- Derby-Klassen mapping
- ⋮



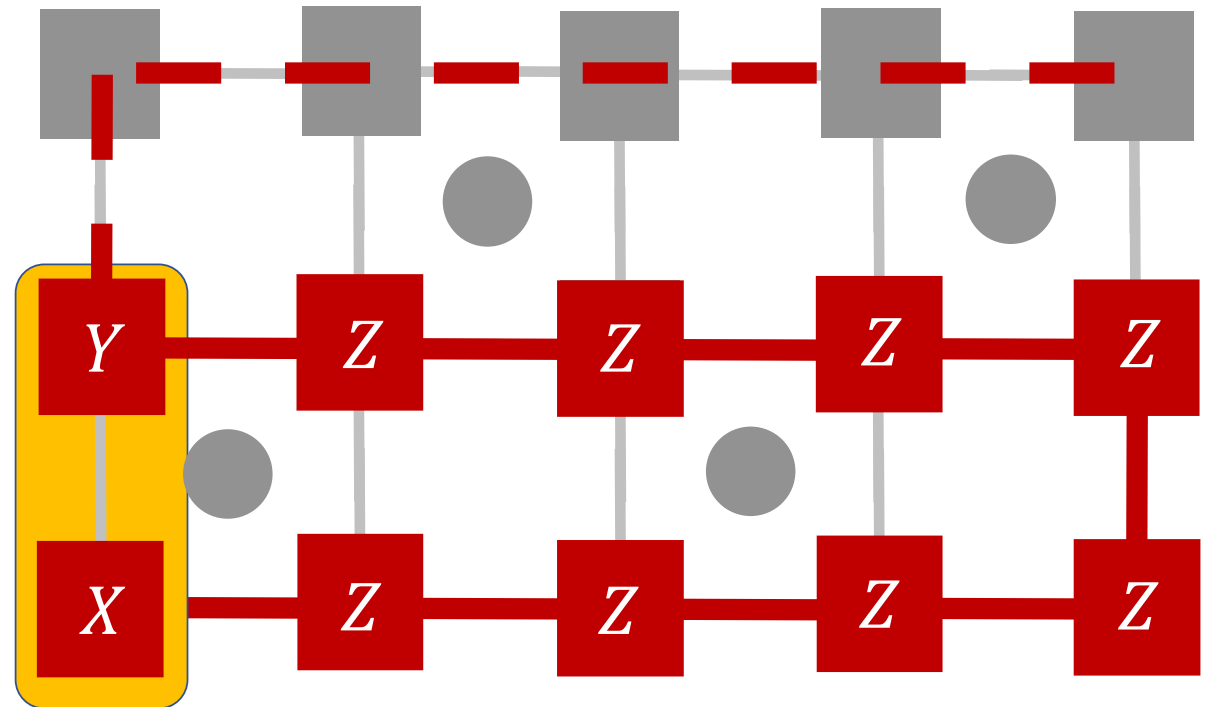
IQM

2. Fermion-to-qubit mappings

$$\begin{aligned} \{E_{ij}, V_i\} &= \{E_{ij}, E_{jk}\} = 0 \\ [E_{ij}, E_{kl}] &= [E_{ij}, V_k] = [V_i, V_j] = 0 \\ c_j^\dagger c_k + c_k^\dagger c_j &\rightarrow \frac{i}{2}(V_k - V_j)E_{jk} \end{aligned}$$

Local mappings

- Bravyi-Kitaev mapping
- Verstraete-Cirac mapping
- Derby-Klassen mapping
- ⋮



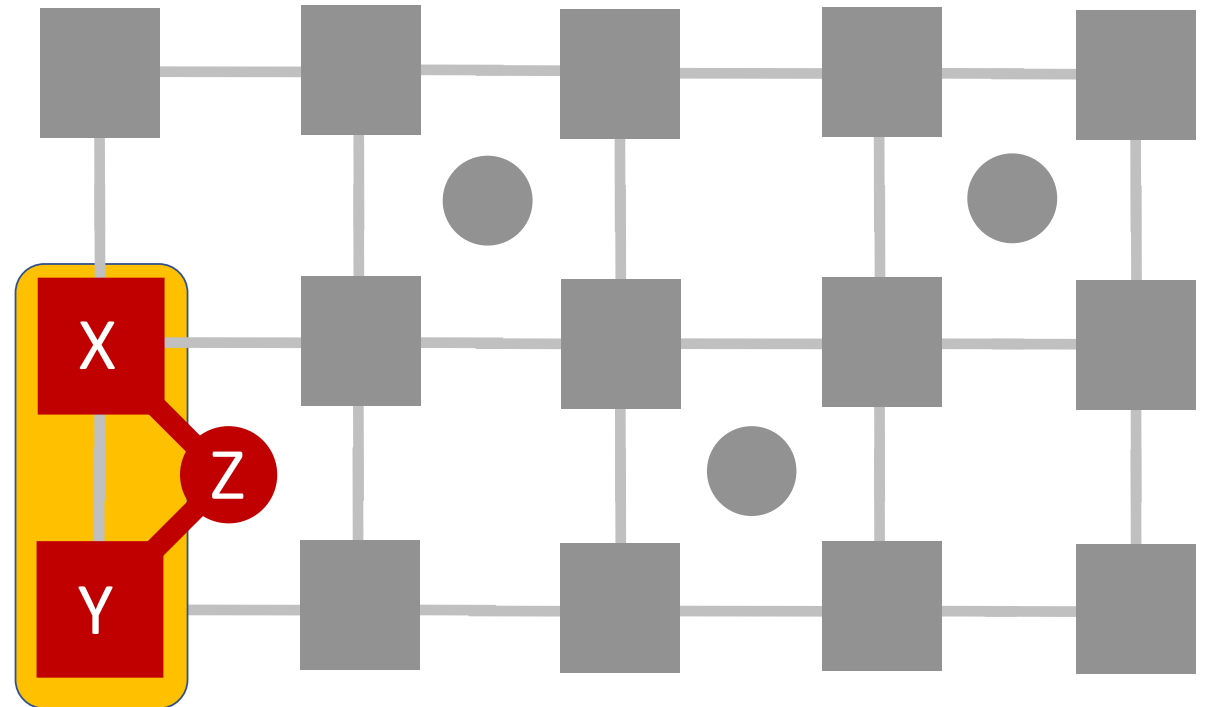
IQM

2. Fermion-to-qubit mappings

$$\begin{aligned} \{E_{ij}, V_i\} &= \{E_{ij}, E_{jk}\} = 0 \\ [E_{ij}, E_{kl}] &= [E_{ij}, V_k] = [V_i, V_j] = 0 \\ c_j^\dagger c_k + c_k^\dagger c_j &\rightarrow \frac{i}{2}(V_k - V_j)E_{jk} \end{aligned}$$

Local mappings

- Bravyi-Kitaev mapping
- Verstraete-Cirac mapping
- Derby-Klassen mapping
- ⋮



IQM

C. Derby *et al.* Phys. Rev. Res. 104 (2021)

2. Fermion-to-qubit mappings

Why local mappings are better when no ATA couplings?

- Local fermionic operators \longrightarrow Local two-qubit gates
- Less depth
- Less number of gates
- Error correction/mitigation properties

2. Fermion-to-qubit mappings

Why local mappings are better when no ATA couplings?

- Local fermionic operators \longrightarrow Local two-qubit gates
- Less depth
- Less number of gates
- Error correction/mitigation properties

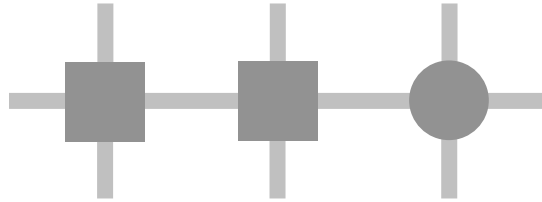
But more ancillas? I prefer Jordan-Wigner

- Number of qubits are not the limiting step nowadays

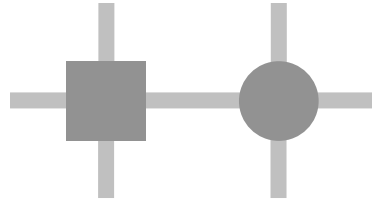
I. D. Kivlichan *et al.* PRL (2018)

2. Fermion-to-qubit mappings

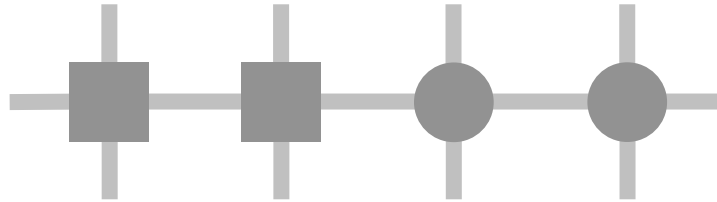
PPA



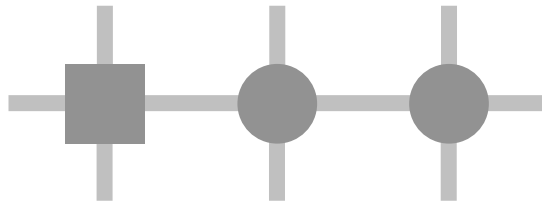
PA



PPAA



PAA



IQM

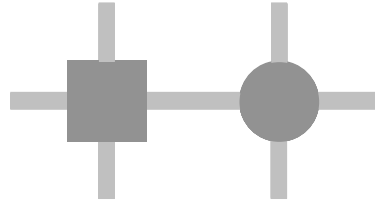
$$\begin{aligned} \{E_{ij}, V_i\} &= \{E_{ij}, E_{jk}\} = 0 \\ [E_{ij}, E_{kl}] &= [E_{ij}, V_k] = [V_i, V_j] = 0 \\ c_j^\dagger c_k + c_k^\dagger c_j &\rightarrow \frac{i}{2}(V_k - V_j)E_{jk} \end{aligned}$$

MA et al. (2023), arXiv:2302.01862

2. Fermion-to-qubit mappings

$$\begin{aligned} \{E_{ij}, V_i\} &= \{E_{ij}, E_{jk}\} = 0 \\ [E_{ij}, E_{kl}] &= [E_{ij}, V_k] = [V_i, V_j] = 0 \\ c_j^\dagger c_k + c_k^\dagger c_j &\rightarrow \frac{i}{2}(V_k - V_j)E_{jk} \end{aligned}$$

PA



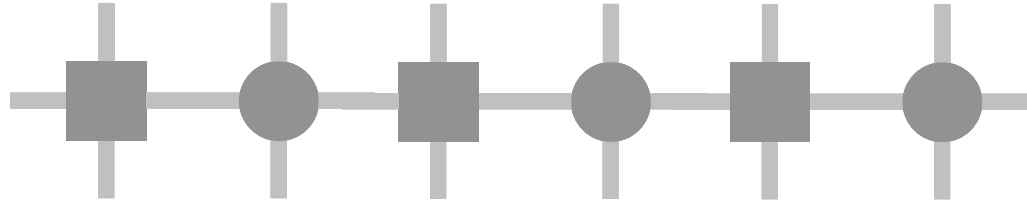
MA et al. (2023), arXiv:2302.01862

IQM

2. Fermion-to-qubit mappings

$$\begin{aligned} \{E_{ij}, V_i\} &= \{E_{ij}, E_{jk}\} = 0 \\ [E_{ij}, E_{kl}] &= [E_{ij}, V_k] = [V_i, V_j] = 0 \\ c_j^\dagger c_k + c_k^\dagger c_j &\rightarrow \frac{i}{2}(V_k - V_j)E_{jk} \end{aligned}$$

PA

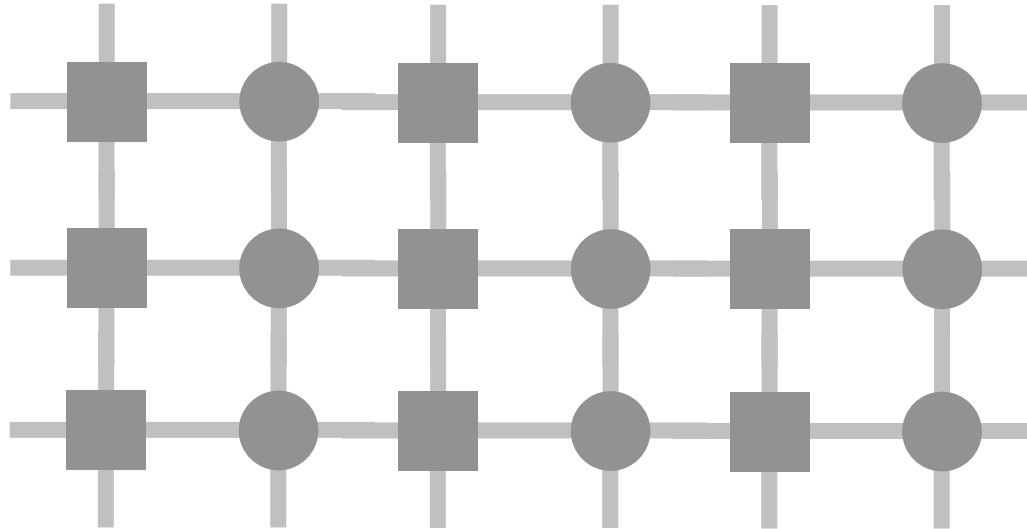


IQM

MA et al. (2023), arXiv:2302.01862

2. Fermion-to-qubit mappings

PA



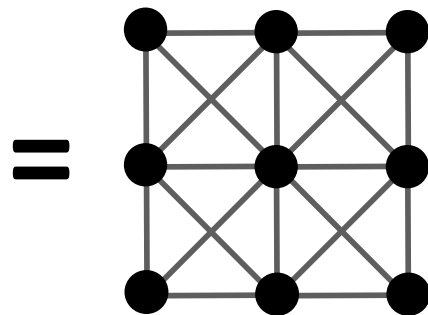
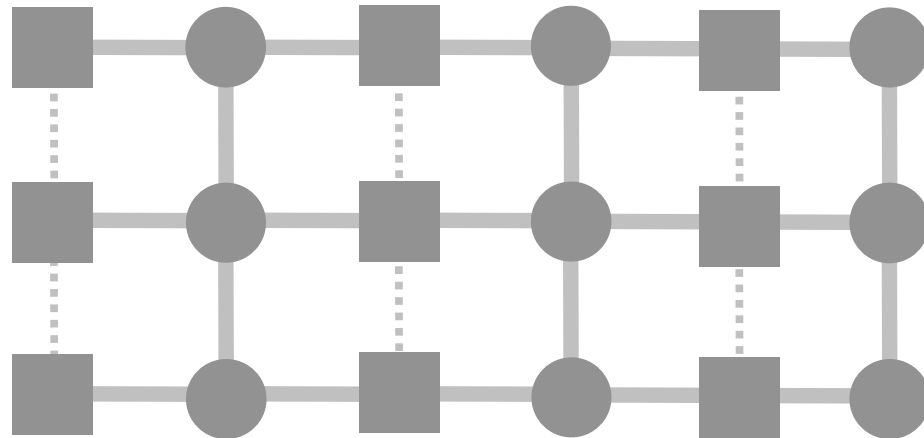
$$\begin{aligned} \{E_{ij}, V_i\} &= \{E_{ij}, E_{jk}\} = 0 \\ [E_{ij}, E_{kl}] &= [E_{ij}, V_k] = [V_i, V_j] = 0 \\ c_j^\dagger c_k + c_k^\dagger c_j &\rightarrow \frac{i}{2}(V_k - V_j)E_{jk} \end{aligned}$$

MA et al. (2023), arXiv:2302.01862

IQM

2. Fermion-to-qubit mappings

PA

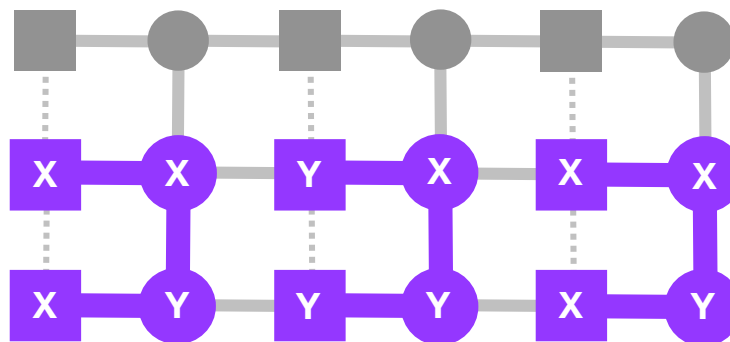
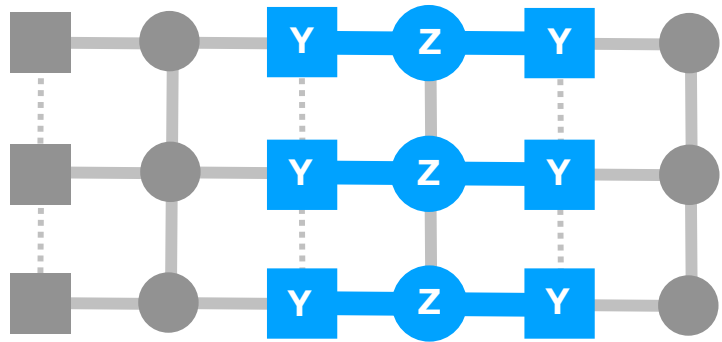
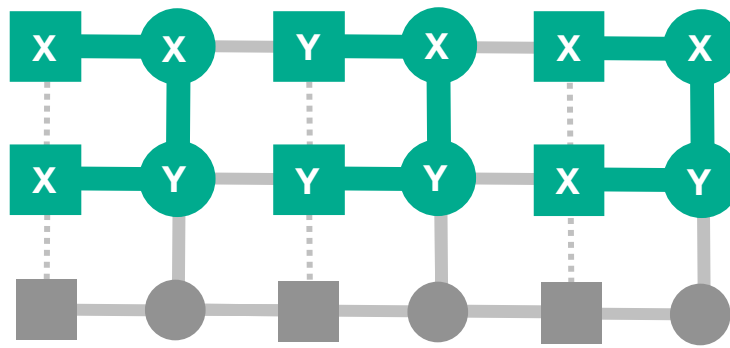
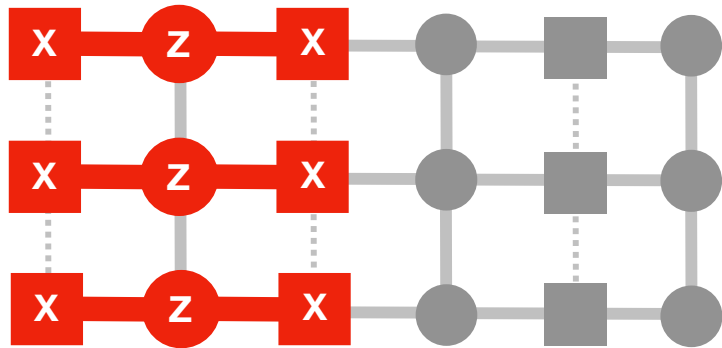


IQM

$$\begin{aligned} \{E_{ij}, V_i\} &= \{E_{ij}, E_{jk}\} = 0 \\ [E_{ij}, E_{kl}] &= [E_{ij}, V_k] = [V_i, V_j] = 0 \\ c_j^\dagger c_k + c_k^\dagger c_j &\rightarrow \frac{i}{2}(V_k - V_j)E_{jk} \end{aligned}$$

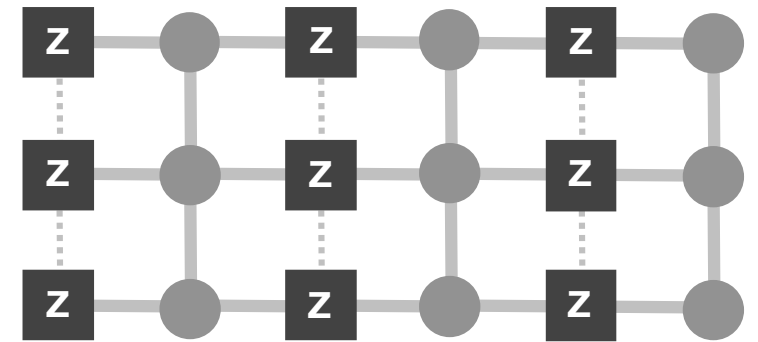
MA et al. (2023), arXiv:2302.01862

2. Fermion-to-qubit mappings



Edge operators

$$\begin{aligned} \{E_{ij}, V_i\} &= \{E_{ij}, E_{jk}\} = 0 \\ [E_{ij}, E_{kl}] &= [E_{ij}, V_k] = [V_i, V_j] = 0 \\ c_j^\dagger c_k + c_k^\dagger c_j &\rightarrow \frac{i}{2}(V_k - V_j)E_{jk} \end{aligned}$$

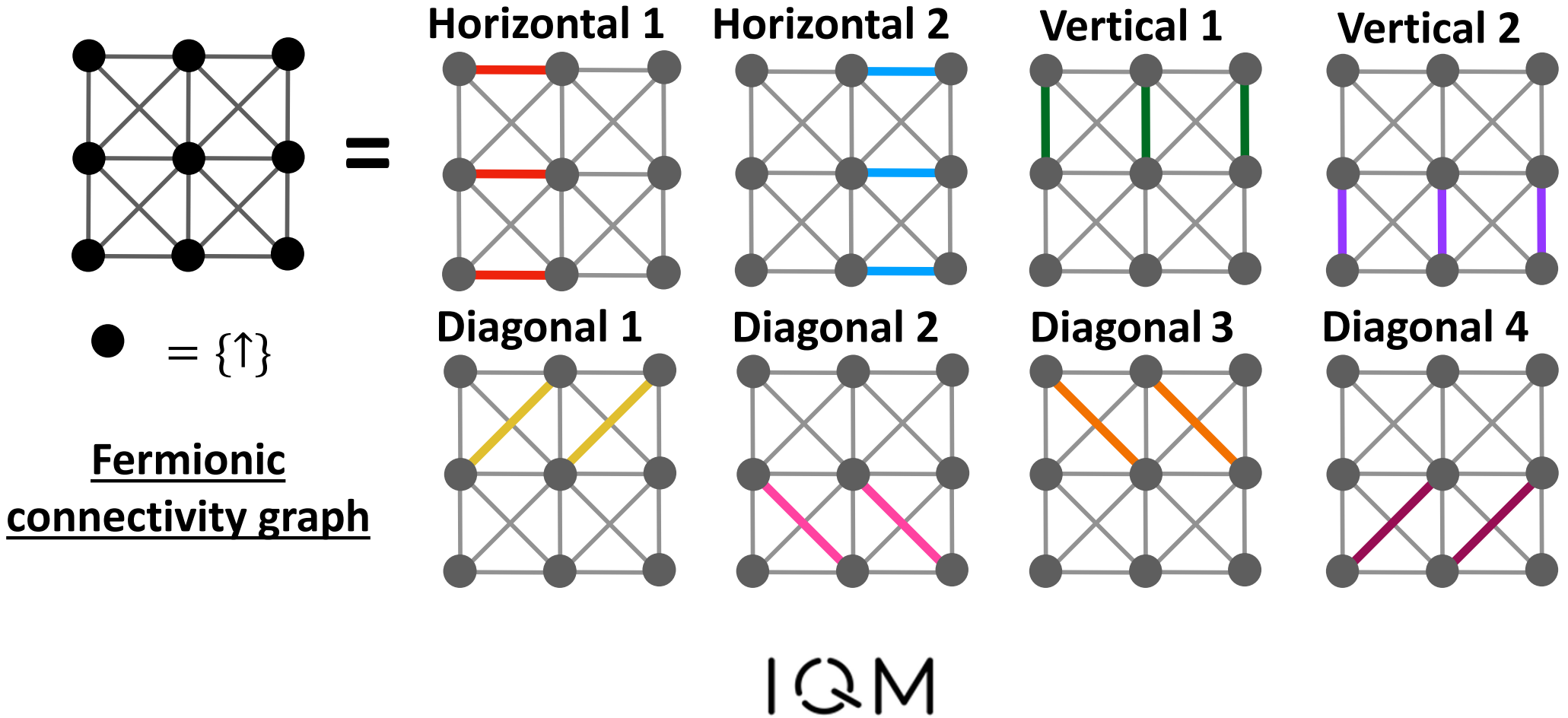


Vertex operators

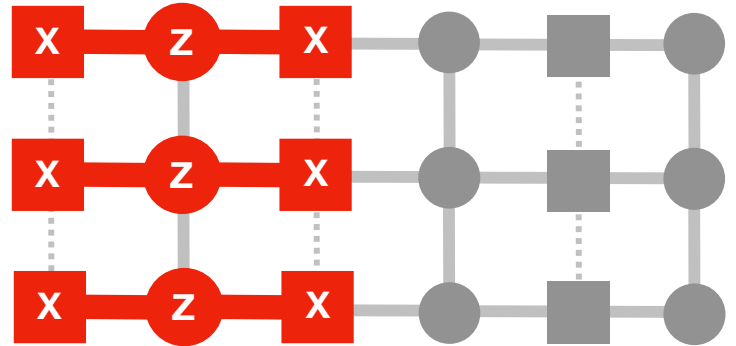
IQM

2. Fermion-to-qubit mappings

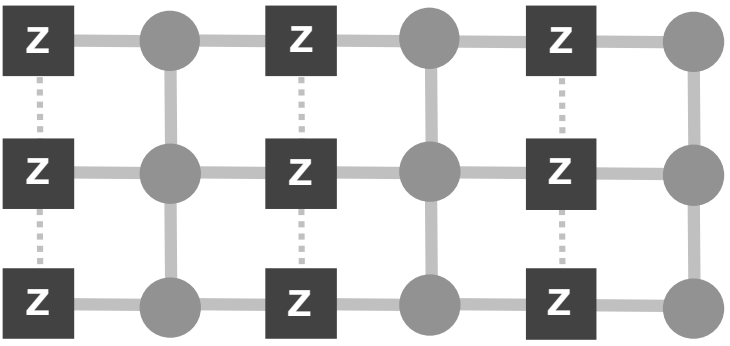
$$\begin{aligned} \{E_{ij}, V_i\} &= \{E_{ij}, E_{jk}\} = 0 \\ [E_{ij}, E_{kl}] &= [E_{ij}, V_k] = [V_i, V_j] = 0 \\ c_j^\dagger c_k + c_k^\dagger c_j &\rightarrow \frac{i}{2}(V_k - V_j)E_{jk} \end{aligned}$$



2. Fermion-to-qubit mappings

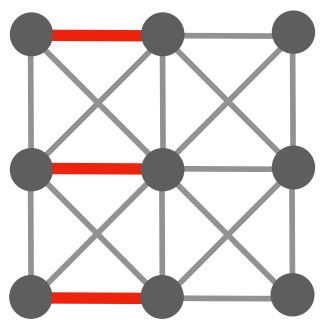


Edge operators

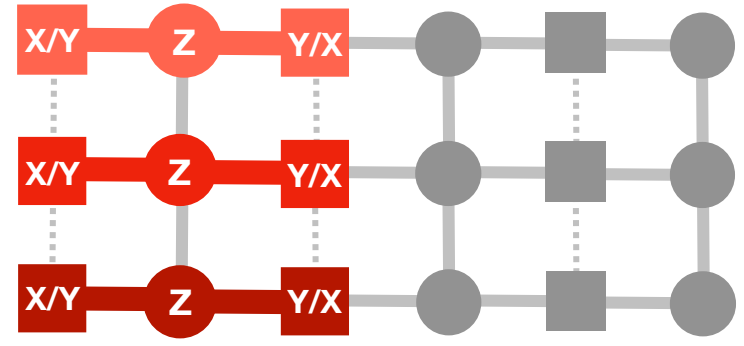


Vertex operators

$$\begin{aligned} \{E_{ij}, V_i\} &= \{E_{ij}, E_{jk}\} = 0 \\ [E_{ij}, E_{kl}] &= [E_{ij}, V_k] = [V_i, V_j] = 0 \\ c_j^\dagger c_k + c_k^\dagger c_j &\rightarrow \frac{i}{2}(V_k - V_j)E_{jk} \\ &\downarrow \\ &(Z_k - Z_j)X_k Z_a X_j \\ &\downarrow \\ &Y_k Z_a X_j + X_k Z_a Y_j \end{aligned}$$



=

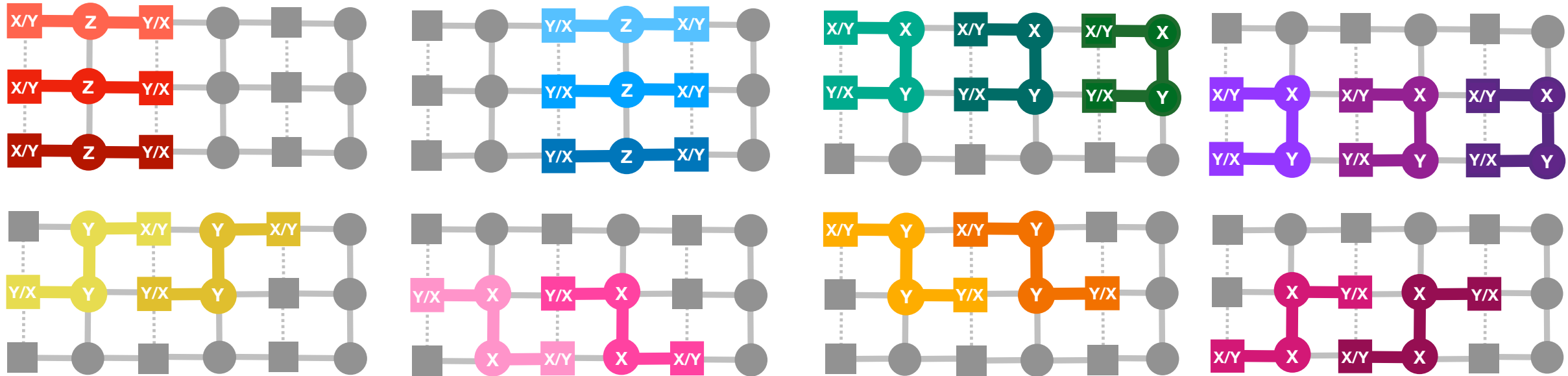


Hopping operators

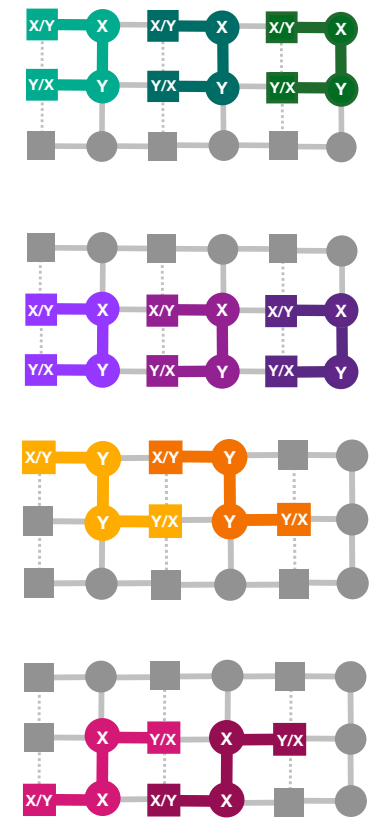
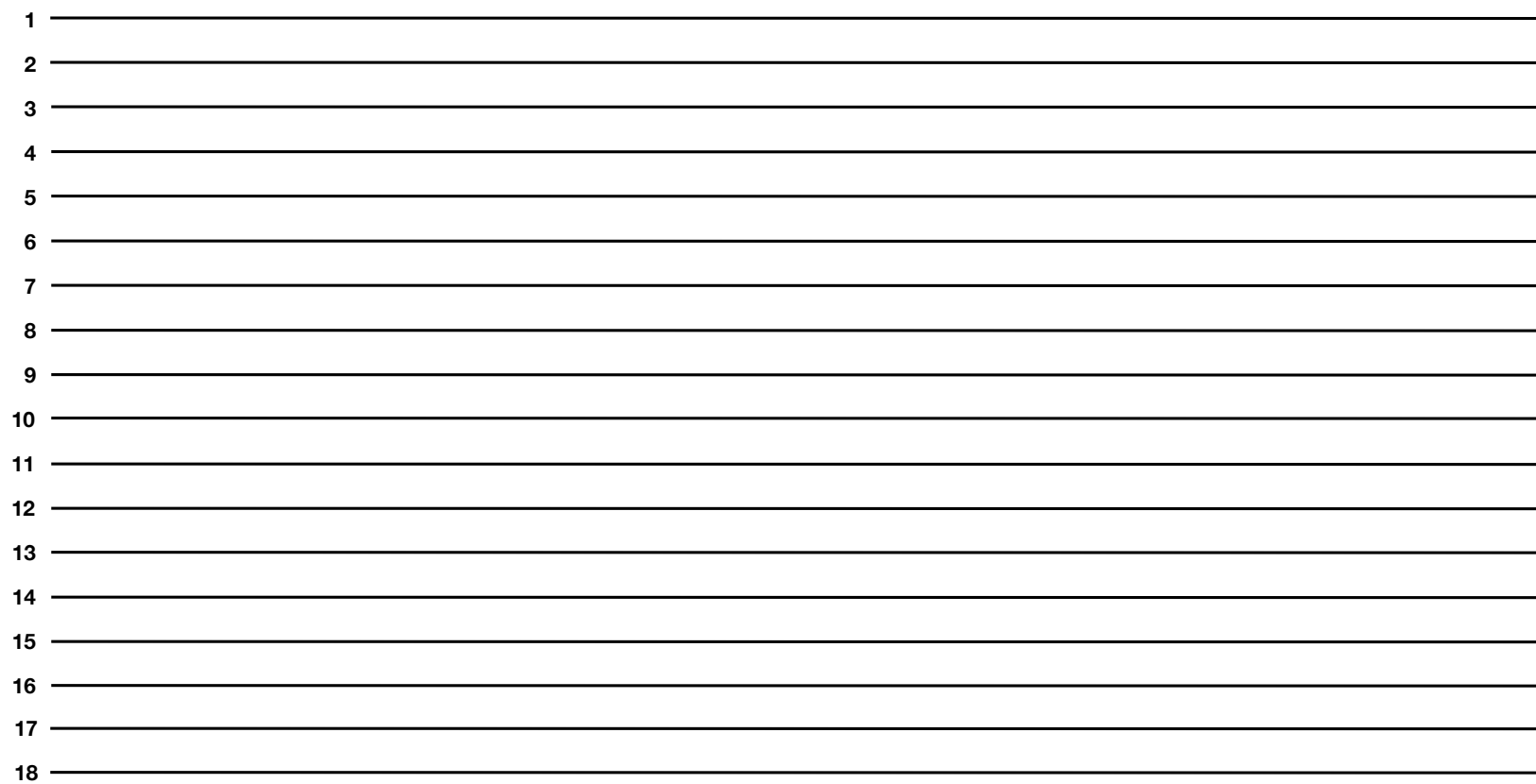
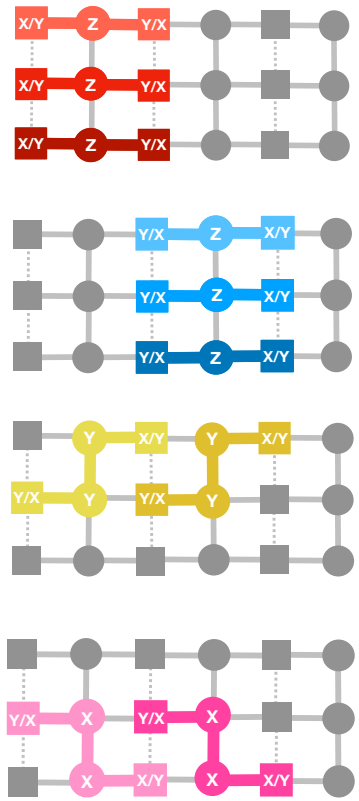
IQM

2. Fermion-to-qubit mappings

$$c_j^\dagger c_k + c_k^\dagger c_j \rightarrow \frac{i}{2}(V_k - V_j)E_{jk}$$

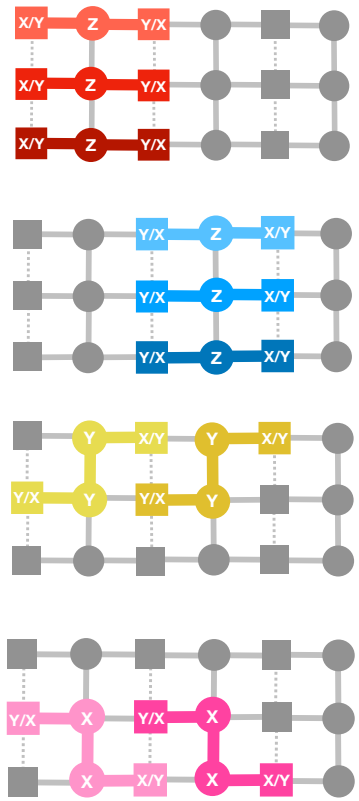


2. Fermion-to-qubit mappings

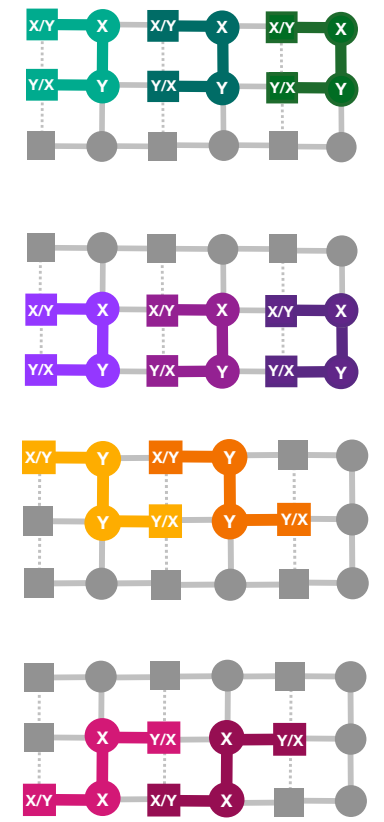


IQM

2. Fermion-to-qubit mappings

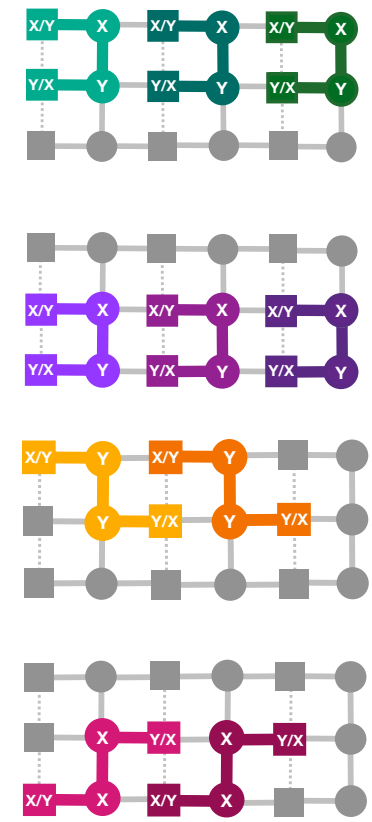
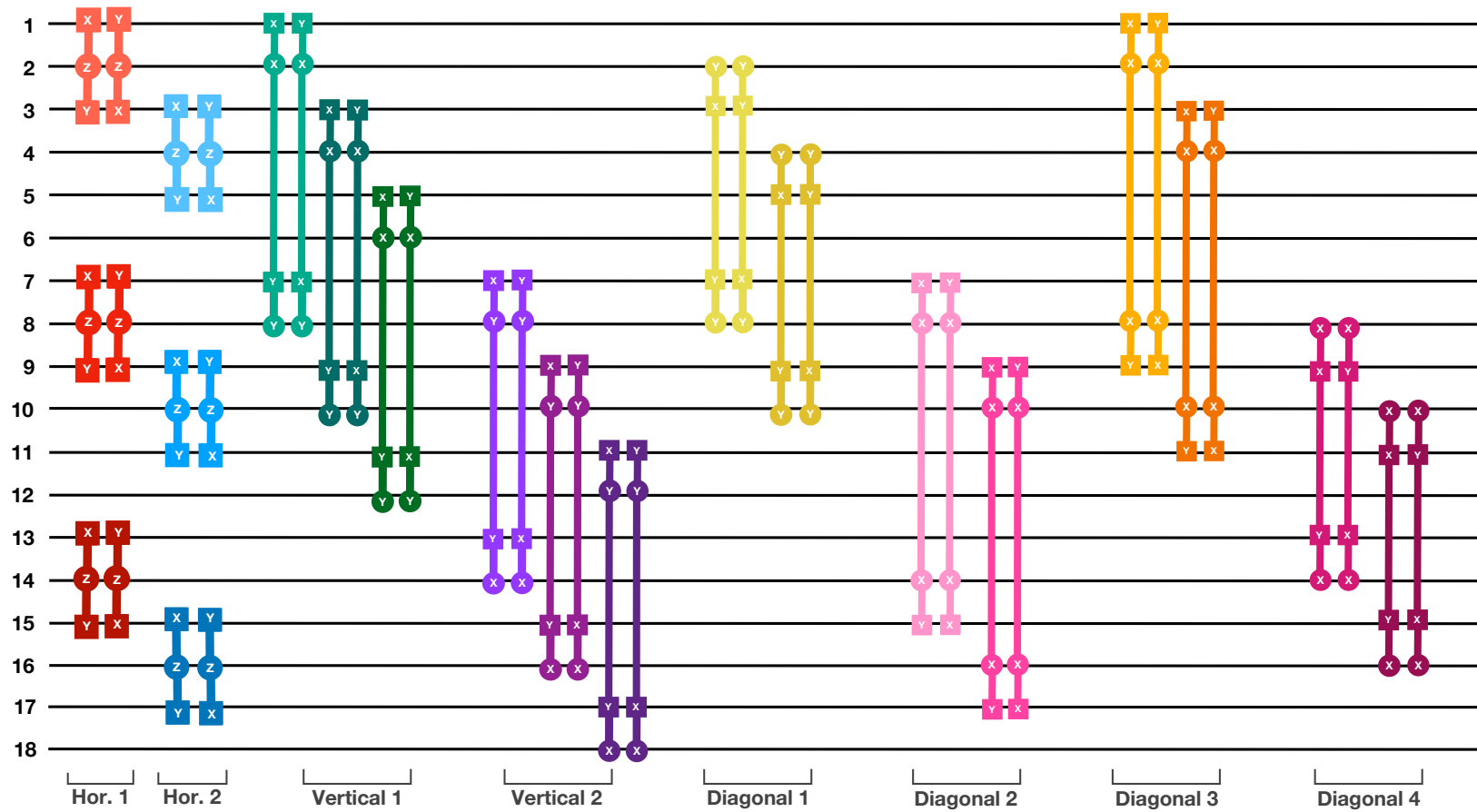
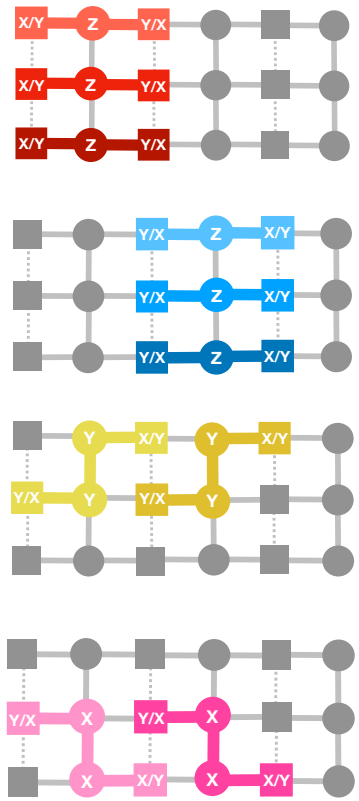


Hor. 1 Hor. 2 Vertical 1 Vertical 2 Diagonal 1 Diagonal 2 Diagonal 3 Diagonal 4



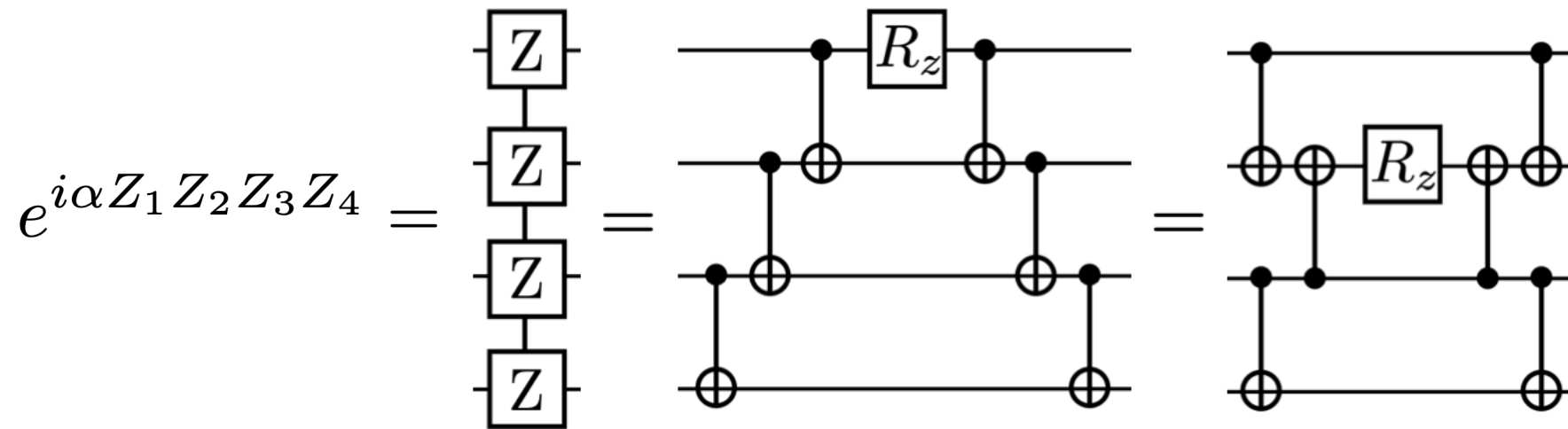
IQM

2. Fermion-to-qubit mappings



IQM

3. XYZ decomposition



Standard decomposition

IQM

3. XYZ decomposition

P. V. Sriluckshmy, MA et al. (2023) arXiv:2303.04498

XYZ decomposition

$$e^{i\alpha\mathcal{O}} = e^{i\frac{\pi}{4}\mathcal{O}_1} e^{i\alpha\mathcal{O}_2} e^{-i\frac{\pi}{4}\mathcal{O}_1}$$

$$\mathcal{O} = \frac{i}{2}[\mathcal{O}_1, \mathcal{O}_2]$$

$$\mathcal{O}^2 = \mathbb{1}$$

Graphical notation

$$\begin{array}{ccc} \text{---} \sigma \text{---} \equiv e^{i\frac{\pi}{4}\sigma} & , & \text{---} \sigma \text{---} \equiv e^{i\alpha\sigma} & , & \text{---} \sigma \text{---} \equiv e^{-i\frac{\pi}{4}\sigma} \\ \text{---} \sigma \text{---} & \equiv e^{i\frac{\pi}{4}\sigma_1\rho_2} & , & \text{---} \sigma \text{---} & \equiv e^{i\alpha\sigma_1\rho_2} & , & \text{---} \sigma \text{---} & \equiv e^{-i\frac{\pi}{4}\sigma_1\rho_2} \\ \text{---} \rho \text{---} & & \text{---} \rho \text{---} & & \text{---} \rho \text{---} & & \text{---} \rho \text{---} & \end{array}$$

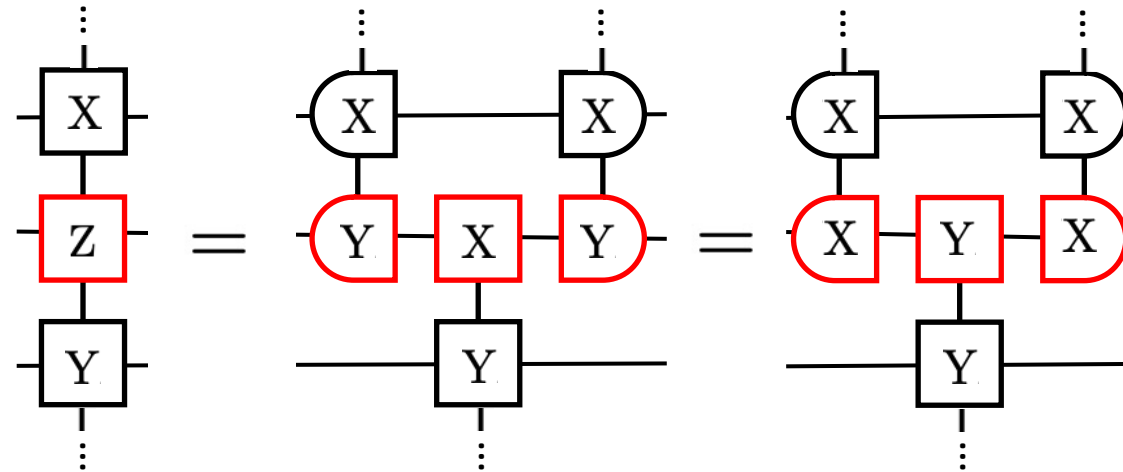
IQM

3. XYZ decomposition

$$e^{i\alpha O} = e^{i\frac{\pi}{4} O_1} e^{i\alpha O_2} e^{-i\frac{\pi}{4} O_1}$$

$$O = \frac{i}{2} [O_1, O_2]$$

$$O^2 = \mathbb{1}$$



IQM

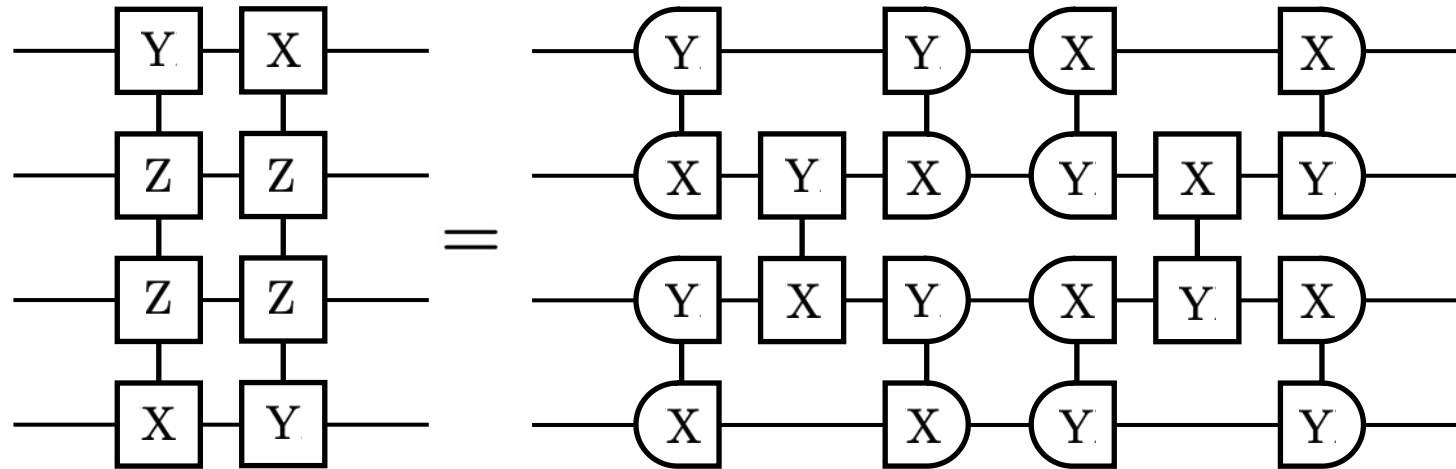
4. Fermionic Simulation

Hopping operators $c_j^\dagger c_k + c_k^\dagger c_j$

$$e^{i\alpha\mathcal{O}} = e^{i\frac{\pi}{4}\mathcal{O}_1} e^{i\alpha\mathcal{O}_2} e^{-i\frac{\pi}{4}\mathcal{O}_1}$$

$$\mathcal{O} = \frac{i}{2}[\mathcal{O}_1, \mathcal{O}_2]$$

$$\mathcal{O}^2 = \mathbb{1}$$



Remember we are using fSIM: $e^{i\frac{\theta}{2}(X_i X_j + Y_i Y_j) + i\frac{\phi}{4}(Z_i + Z_j - Z_i Z_j)}$

IQM

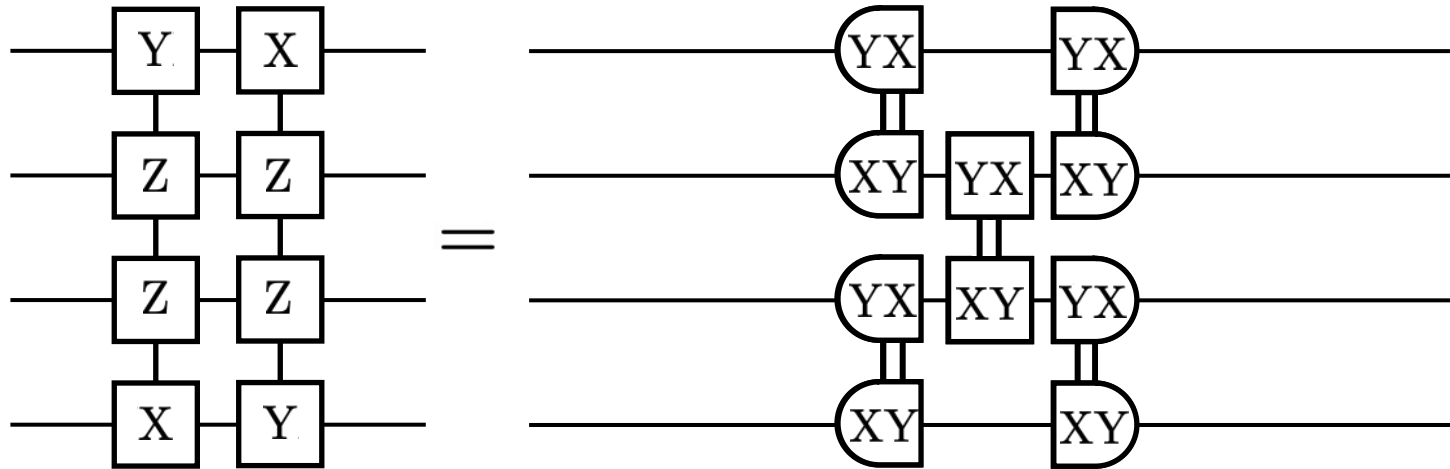
4. Fermionic Simulation

Hopping operators $c_j^\dagger c_k + c_k^\dagger c_j$

$$e^{i\alpha\mathcal{O}} = e^{i\frac{\pi}{4}\mathcal{O}_1} e^{i\alpha\mathcal{O}_2} e^{-i\frac{\pi}{4}\mathcal{O}_1}$$

$$\mathcal{O} = \frac{i}{2}[\mathcal{O}_1, \mathcal{O}_2]$$

$$\mathcal{O}^2 = \mathbb{1}$$



Remember we are using fSIM: $e^{i\frac{\theta}{2}(X_i X_j + Y_i Y_j) + i\frac{\phi}{4}(Z_i + Z_j - Z_i Z_j)}$

IQM

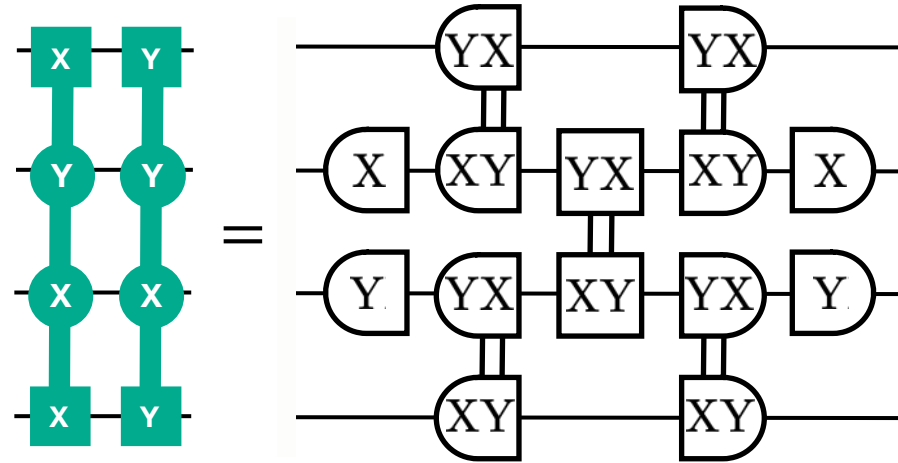
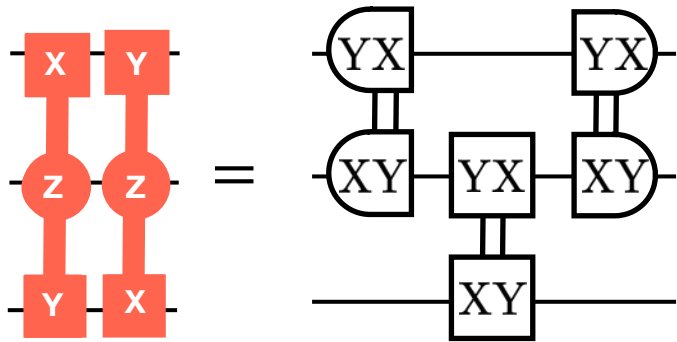
4. Fermionic Simulation

Hopping operators $c_j^\dagger c_k + c_k^\dagger c_j$

$$e^{i\alpha O} = e^{i\frac{\pi}{4} O_1} e^{i\alpha O_2} e^{-i\frac{\pi}{4} O_1}$$

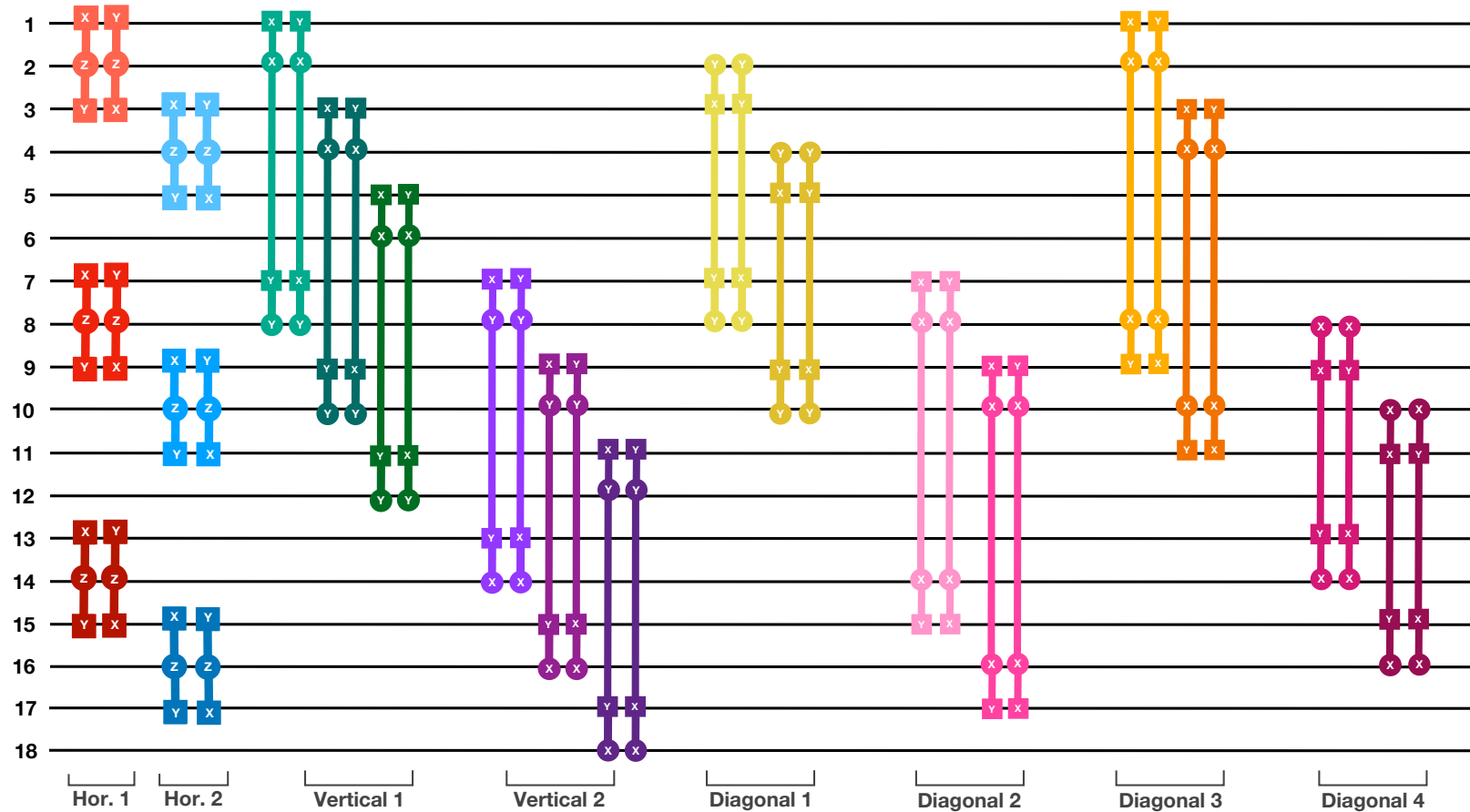
$$O = \frac{i}{2} [O_1, O_2]$$

$$O^2 = \mathbb{1}$$



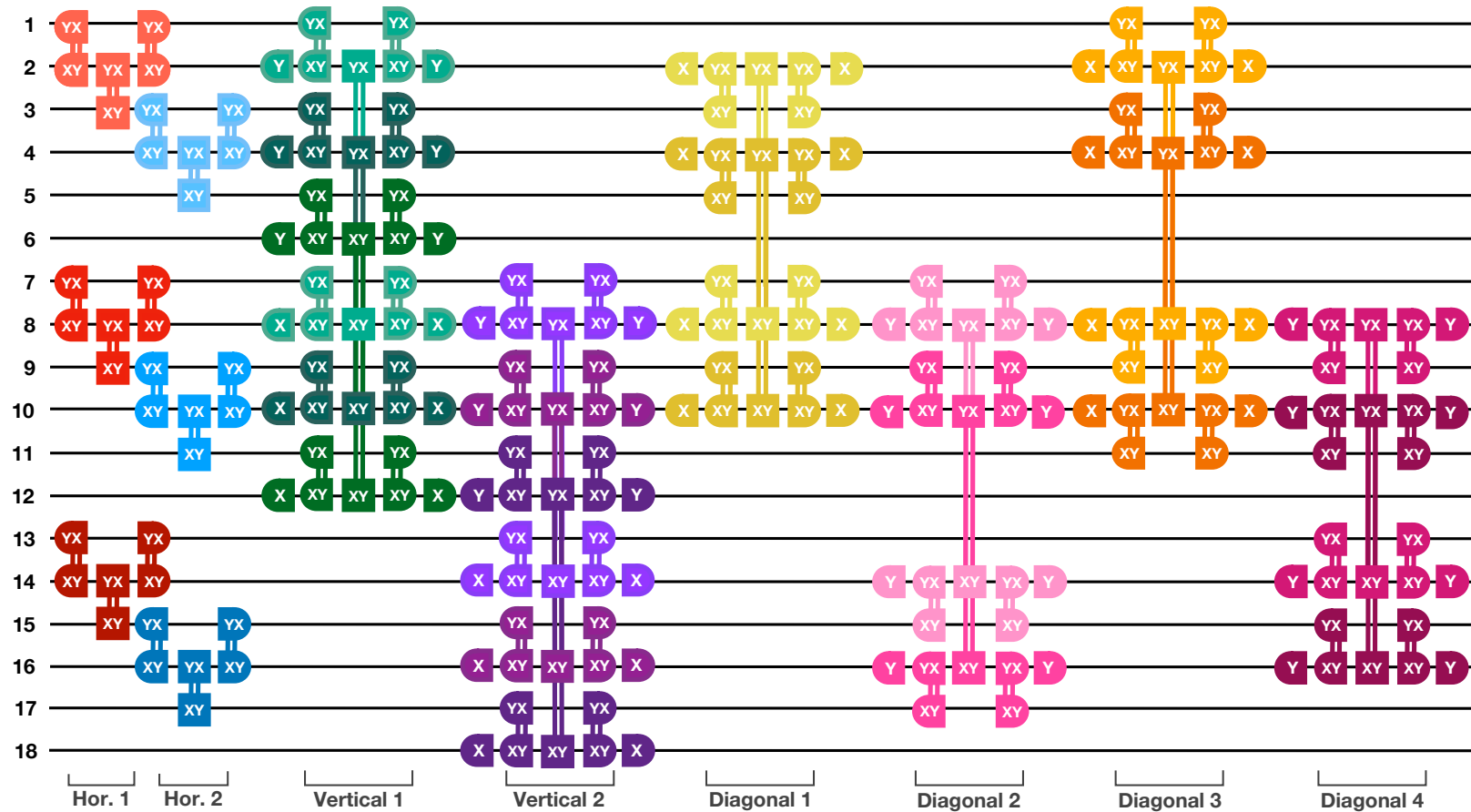
IQM

4. Fermionic Simulation



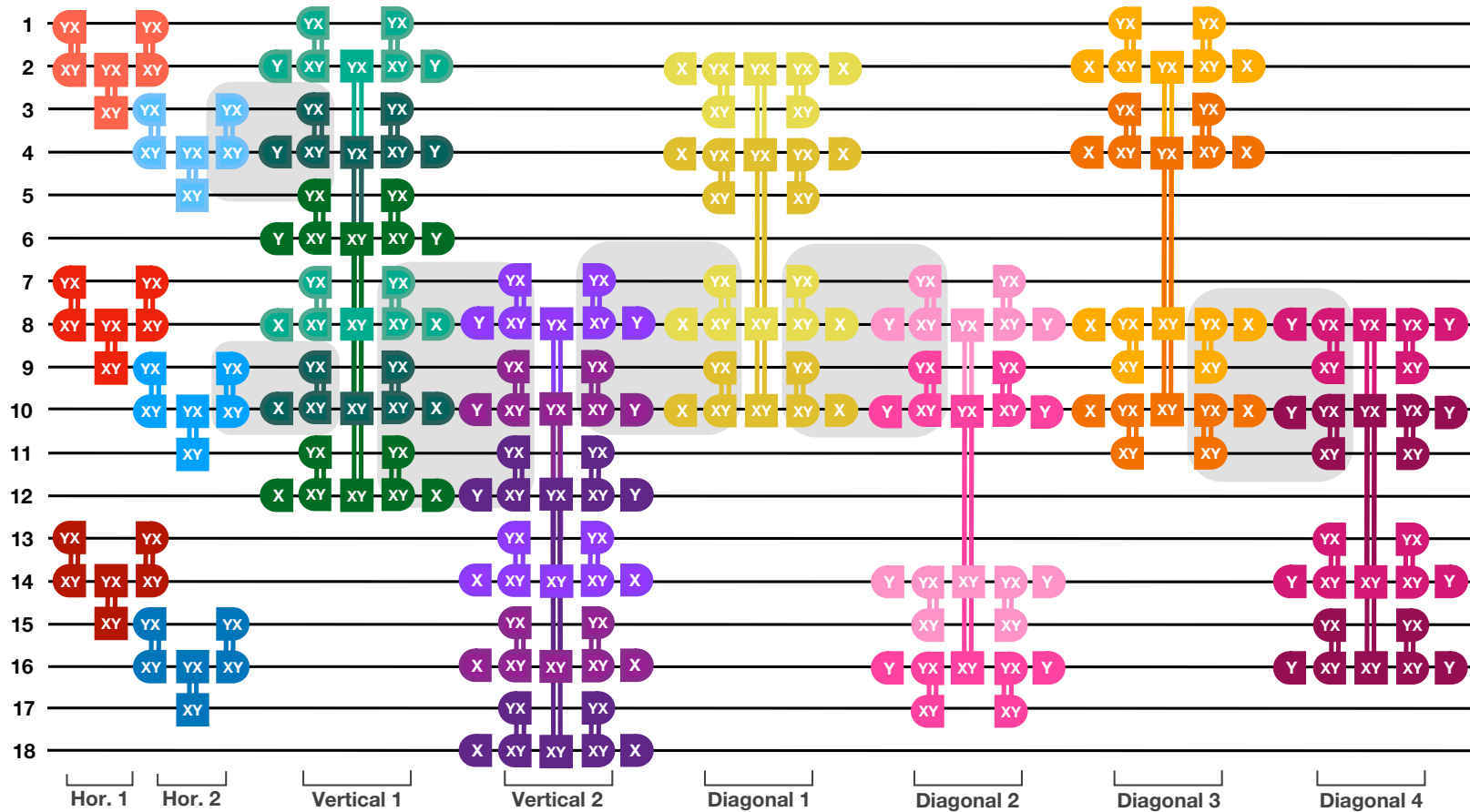
IQM

4. Fermionic Simulation



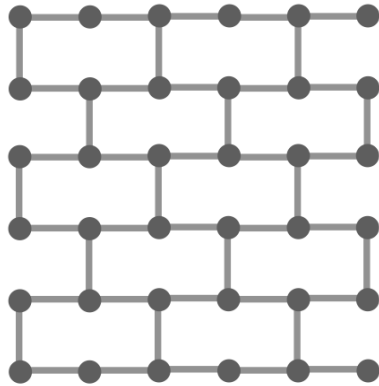
IQM

4. Fermionic Simulation

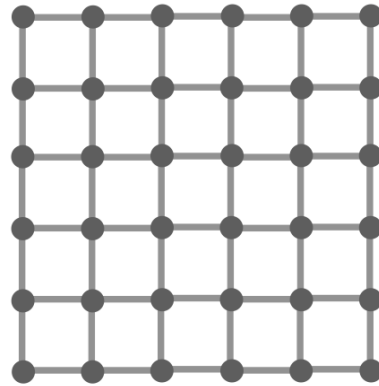


IQM

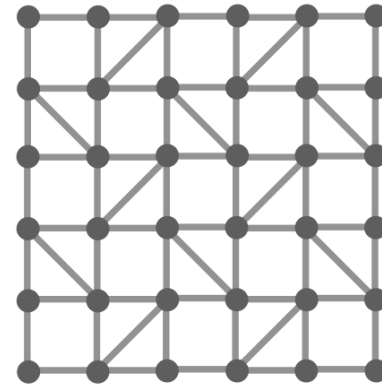
4. Fermionic Simulation



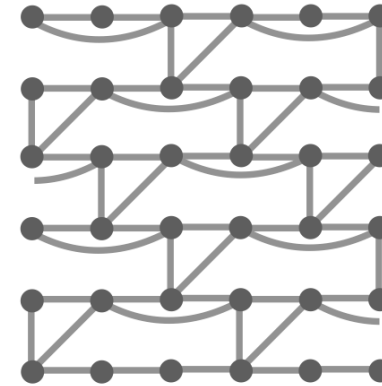
Honeycomb



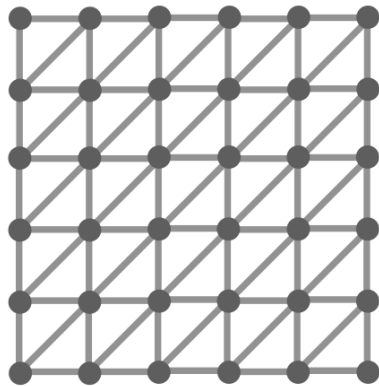
Square



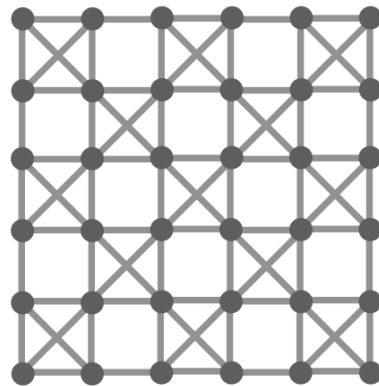
Shastry-Sutherland



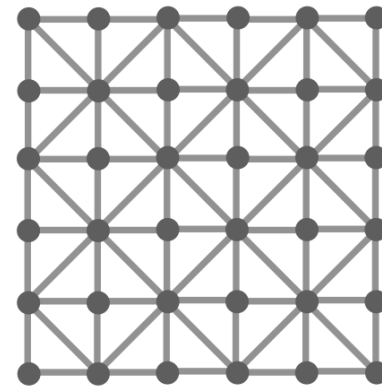
Kagome



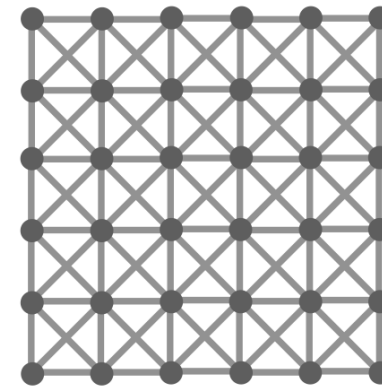
Triangular



Checkerboard



Tetrakis



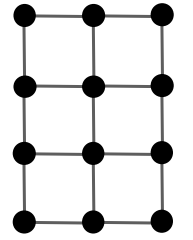
NNN Square

IQM

4. Fermionic Simulation

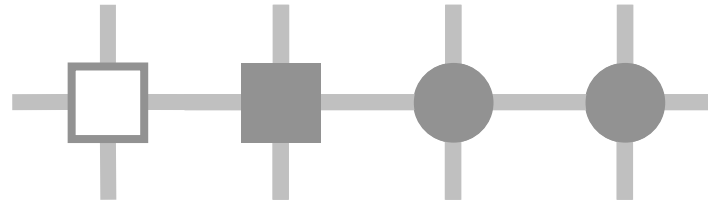
$$\mathcal{H}_{\text{FH}} = \sum_{i,j,\sigma} t^{ij} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Fermionic connectivity graph



● = {↑, ↓}

$\tilde{\text{PAA}}$

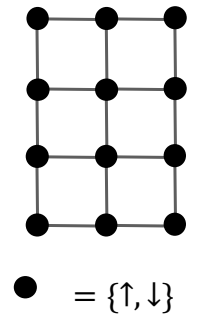


IQM

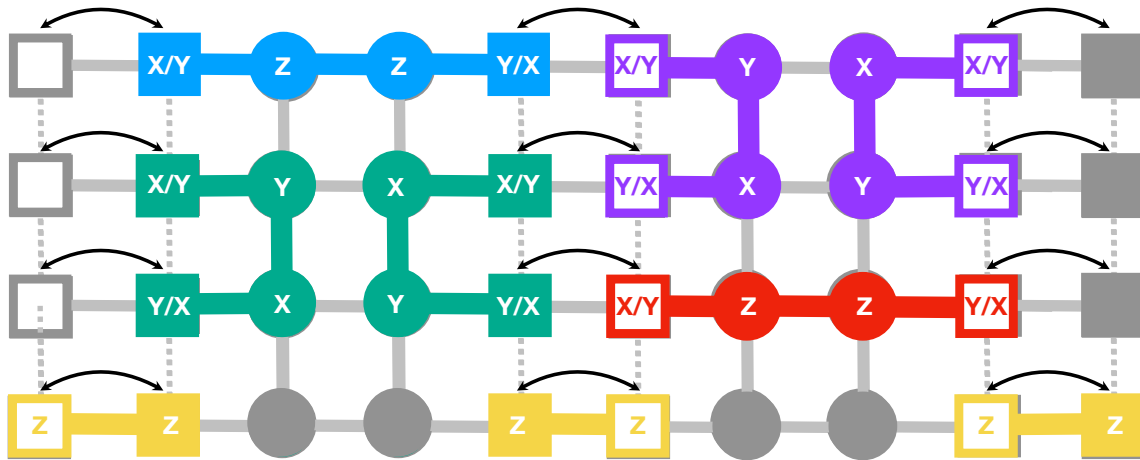
4. Fermionic Simulation

$$\mathcal{H}_{\text{FH}} = \sum_{i,j,\sigma} t^{ij} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Fermionic connectivity graph



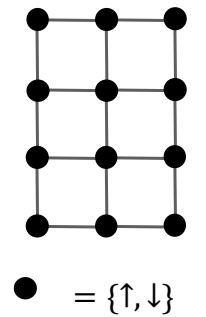
Qubit layout



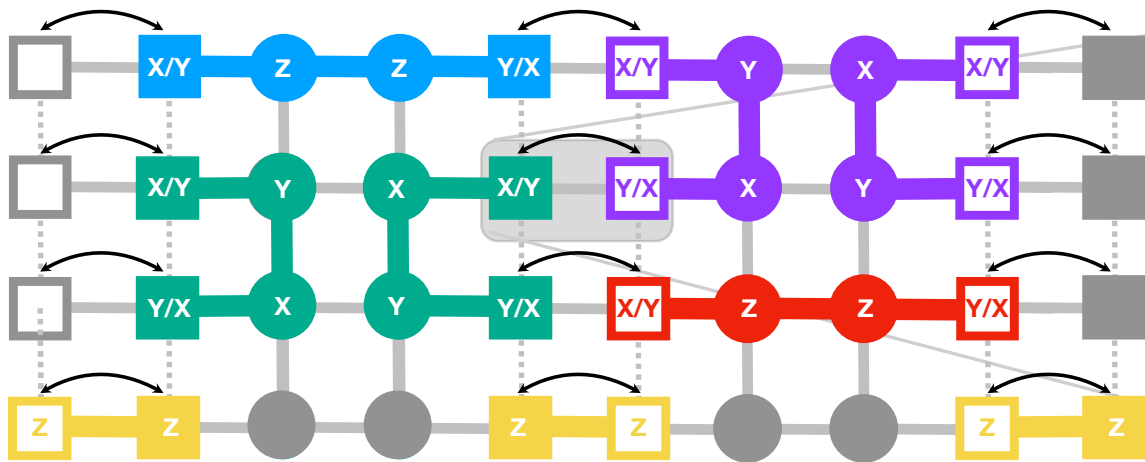
4. Fermionic Simulation

$$\mathcal{H}_{\text{FH}} = \sum_{i,j,\sigma} t^{ij} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

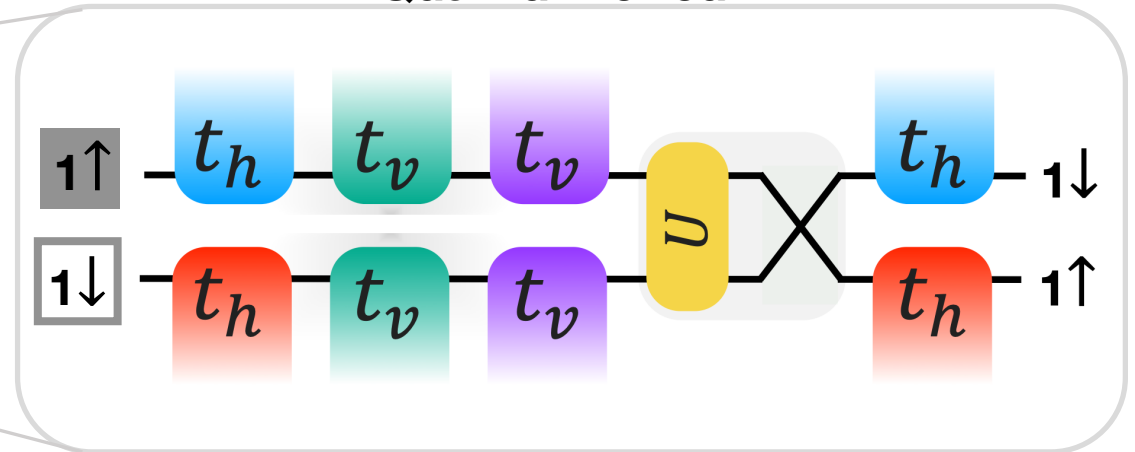
Fermionic connectivity graph



Qubit layout



Quantum circuit



$$\text{X} = \begin{array}{c} \text{---} \text{---} \\ \diagdown \diagup \\ \text{---} \text{---} \end{array} = \begin{array}{cc} \boxed{Z} & \boxed{YX} \\ \boxed{Z} & \boxed{XY} \end{array} = \text{fSWAP}_{ij}$$

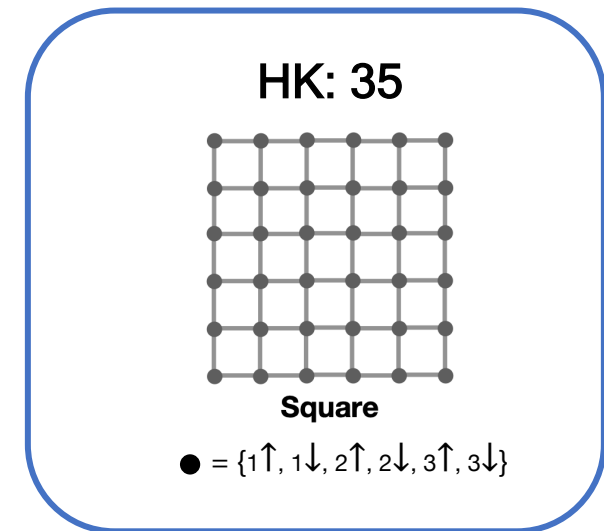
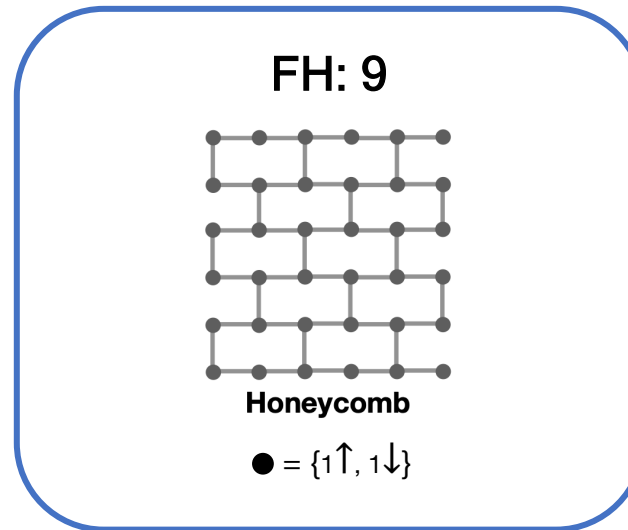
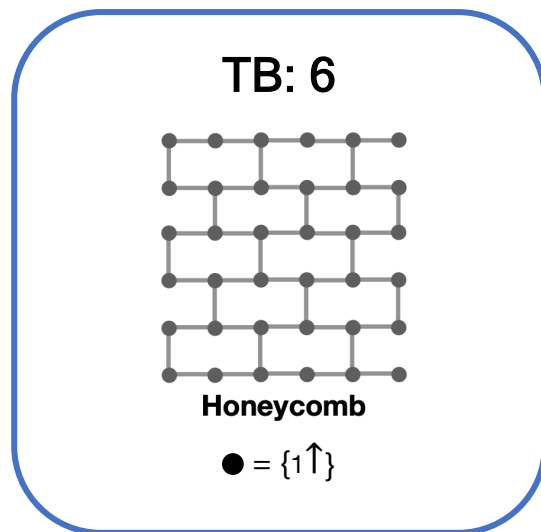
$$\boxed{U} = \begin{array}{cc} \boxed{Z} & \boxed{Z} \\ \boxed{Z} & \boxed{Z} \end{array}$$

IQM

4. Fermionic Simulation

So, what's the improvement?

- Least number of TQGs with DK + XYZ decomposition
- Up to 72% depth reduction (3.2x).
- Shallowest single-Trotter-step circuits for these condensed matter Hamiltonians in literature:



IQM

Thank you for your
attention!

 manuel.algaba@meetiqm.com

 [@ManuQPhys](https://twitter.com/ManuQPhys)

IQM

Backup I

But, where is the advantage coming from?

Decomp.	Native TQGs	TB NNN	FH NNN	FH NNN	HK NN
		PA	PAA	DK	PAA
XYZ	fSIM	18	30	74	55
Standard	fSIM	31	68	97	84
XYZ	CNOT	31	53	125	100
Standard	CNOT	47	93	130	132

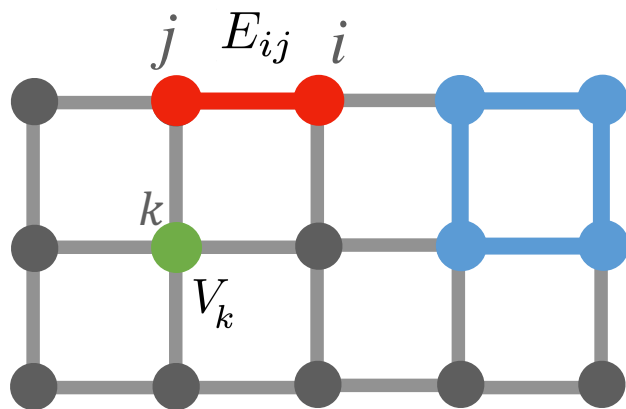
IQM

Backup II

$$\mathcal{H}_{\text{TB}} = \sum_{i,j} t^{ij} c_i^\dagger c_j$$

$$\{c_i, c_j^\dagger\} = c_i c_j^\dagger + c_j^\dagger c_i = \delta_{ij}$$

$$\{c_i^\dagger, c_j^\dagger\} = \{c_i, c_j\} = 0$$



Edge and vertex operators:

$$\{E_{ij}, V_i\} = \{E_{ij}, E_{jk}\} = 0$$

$$[E_{ij}, E_{kl}] = [E_{ij}, V_k] = [V_i, V_j] = 0$$

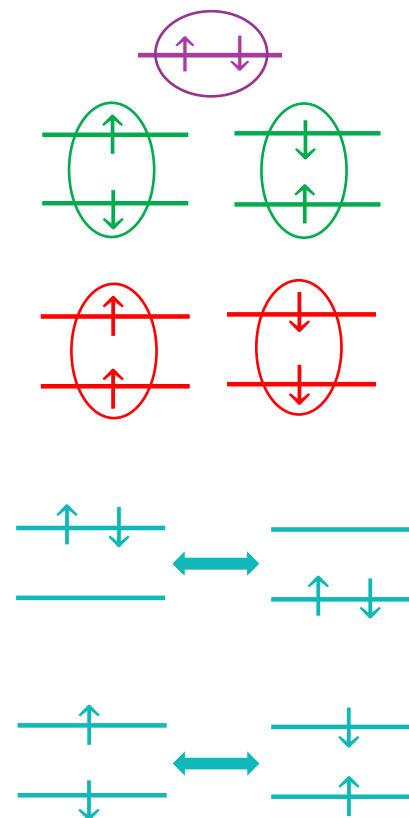
Eigenspace condition:

$$i^{(|p|-1)} \prod_j^{|p|-1} E_{p_j, p_{j+1}} = \mathbb{1}$$

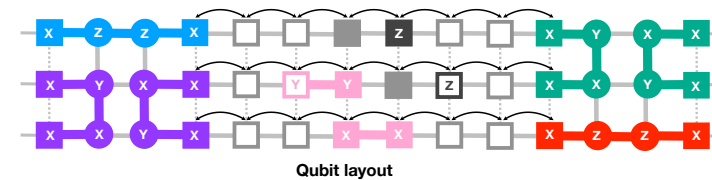
Backup III

Hubbard-Kanamori Hamiltonian:

$$\begin{aligned}
 \mathcal{H}_{\text{HK}} = & \sum_{i,j,m,\sigma} t^{ijm\sigma} c_{im\sigma}^\dagger c_{jm\sigma} + \sum_{i,m} U^{im} n_{im\uparrow} n_{im\downarrow} \\
 & + \sum_{i,m < \bar{m}} U_1^{im\bar{m}} \left(n_{im\uparrow} n_{i\bar{m}\downarrow} + n_{im\downarrow} n_{i\bar{m}\uparrow} \right) \\
 & + \sum_{i,m < \bar{m}} U_2^{im\bar{m}} \left(n_{im\uparrow} n_{i\bar{m}\uparrow} + n_{im\downarrow} n_{i\bar{m}\downarrow} \right) \\
 & + \sum_{i,m < \bar{m}} J^{im\bar{m}} \left(c_{im\uparrow}^\dagger c_{im\downarrow}^\dagger c_{i\bar{m}\downarrow} c_{i\bar{m}\uparrow} + c_{i\bar{m}\uparrow}^\dagger c_{i\bar{m}\downarrow}^\dagger c_{im\downarrow} c_{im\uparrow} \right. \\
 & \quad \left. + c_{im\uparrow}^\dagger c_{i\bar{m}\downarrow}^\dagger c_{im\downarrow} c_{i\bar{m}\uparrow} + c_{i\bar{m}\uparrow}^\dagger c_{im\downarrow}^\dagger c_{i\bar{m}\downarrow} c_{im\uparrow} \right)
 \end{aligned}$$



Backup IV



Quantum circuit

