

Dissipative Simulation of Quantum Dynamics with Tensor Networks

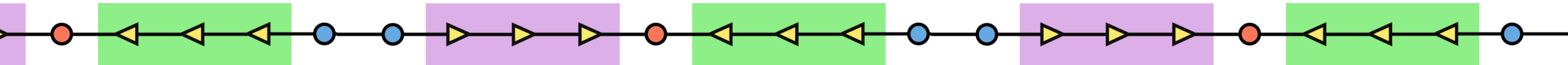
Carlos Ramos Marimón



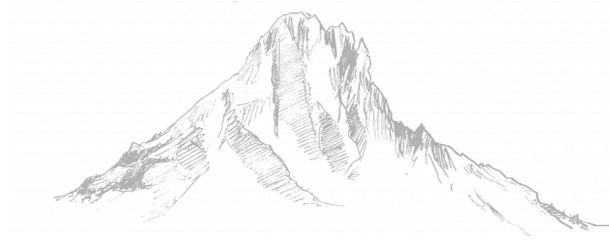
Institut de Ciències del Cosmos
UNIVERSITAT DE BARCELONA



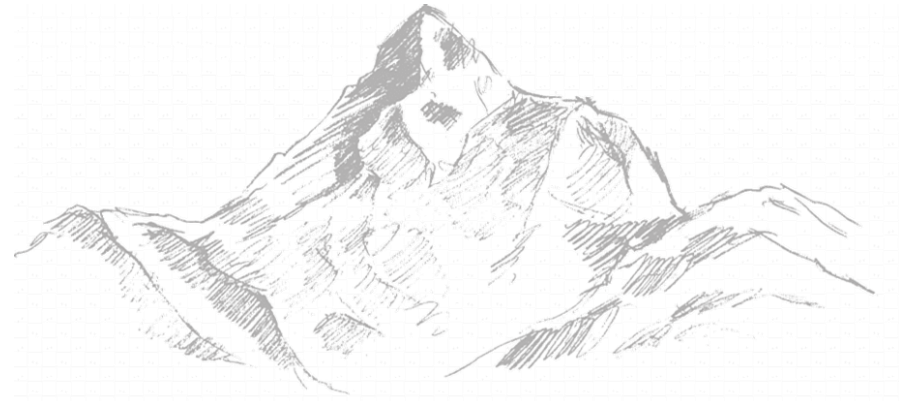
"la Caixa" Foundation



Roadmap



1. Quantum evolution
in tensor network language

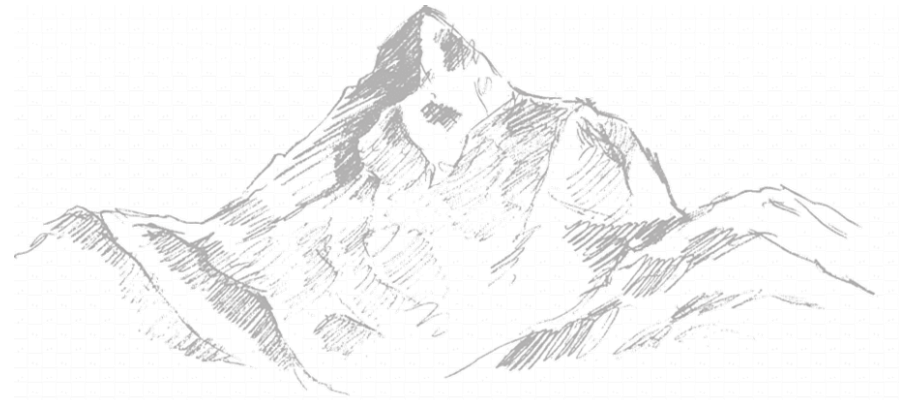


Roadmap



1. Quantum evolution
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2. Entanglement barrier

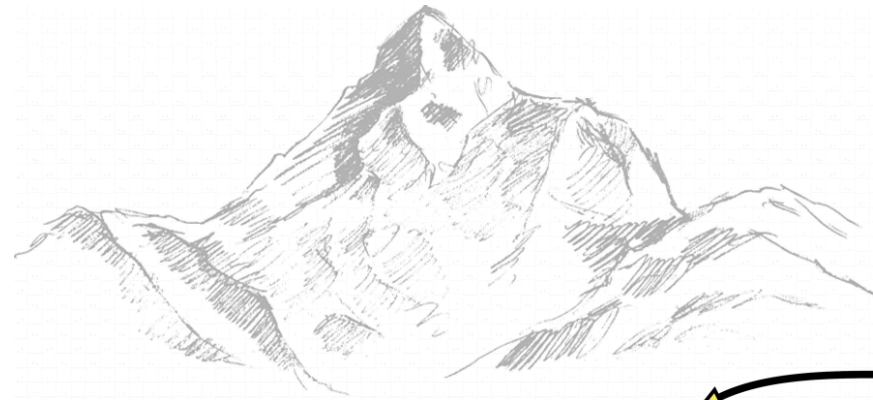


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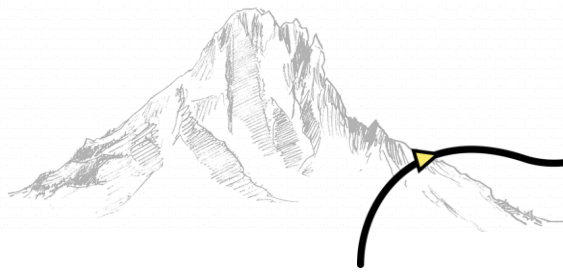


3. Dissipative evolution

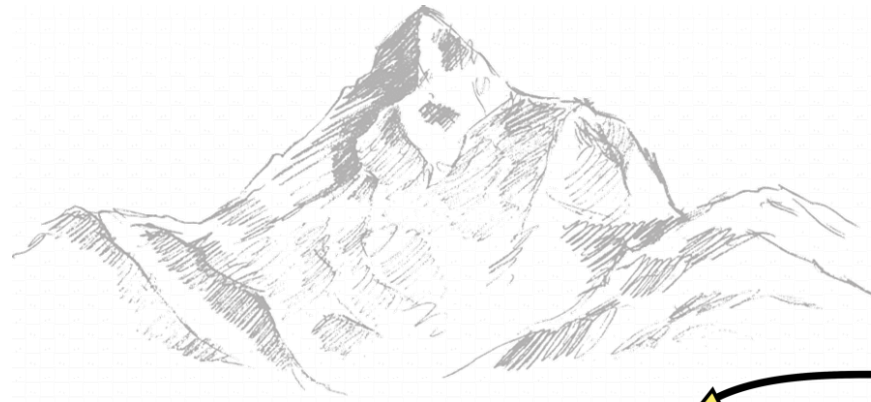


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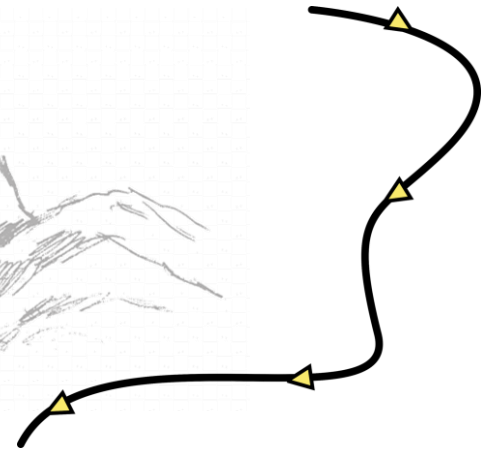
1. Quantum evolution
in tensor network language



2. Entanglement barrier



3. Dissipative evolution




4. Preliminary results




1. Quantum Evolution in Tensor Network Language

Given an initial state/density matrix



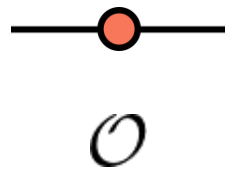
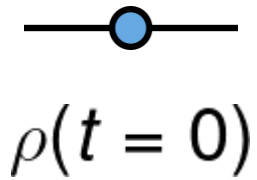
$\rho(t = 0)$




\mathcal{O}

1. Quantum Evolution in Tensor Network Language


Given an initial state/density matrix



the evolution is generated by


$$\delta\mathcal{U} = e^{-i \cdot \delta t \cdot \mathcal{H}}$$

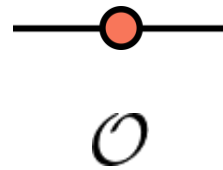
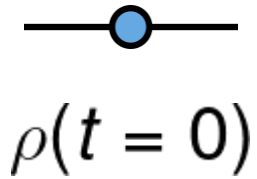
forward


$$\delta\mathcal{U}^\dagger = e^{i \cdot \delta t \cdot \mathcal{H}}$$

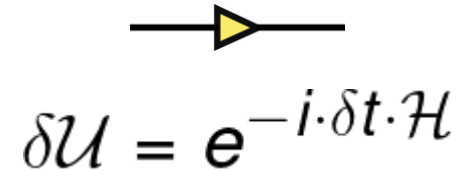
backward

1. Quantum Evolution in Tensor Network Language

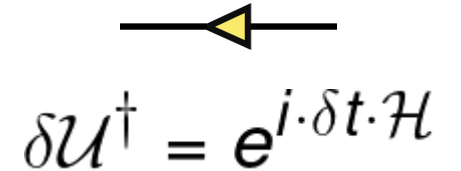
Given an initial state/density matrix



the evolution is generated by



forward



backward

The expectation value of the operator in the evolved state reads

$$\mathcal{O}(t = T) = \text{tr}\{\mathcal{U} \rho(t = 0) \mathcal{U}^\dagger \mathcal{O}\}$$

1. Quantum Evolution in Tensor Network Language

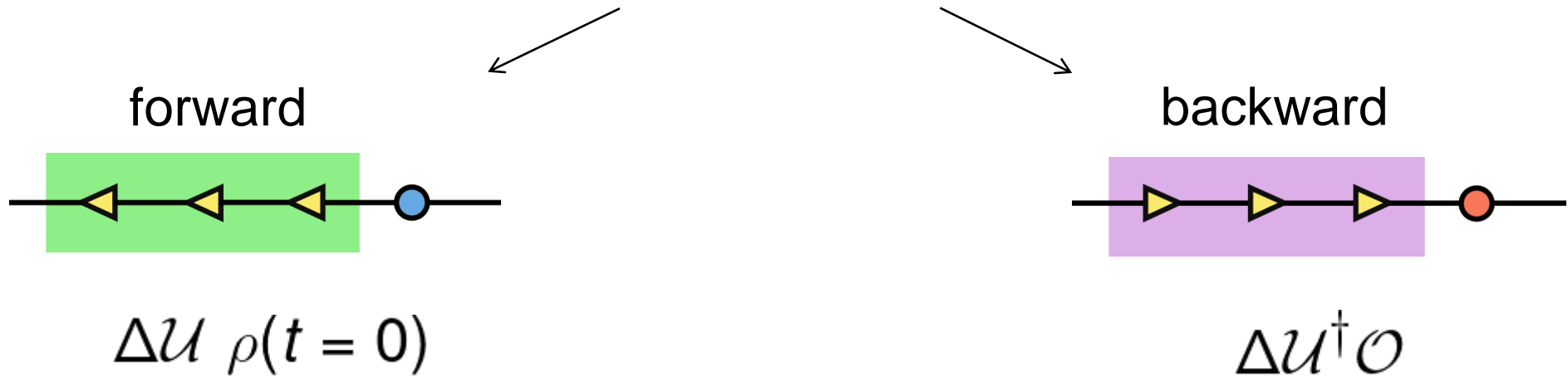
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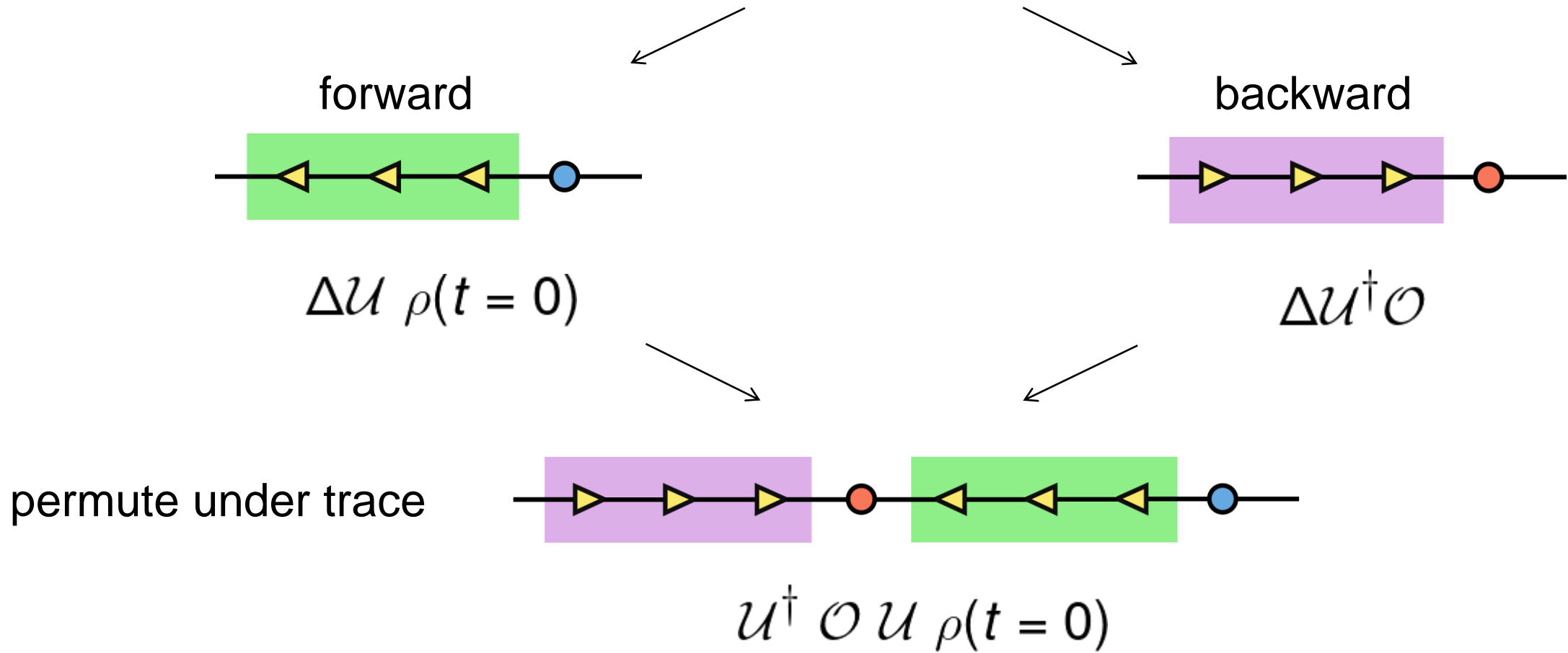
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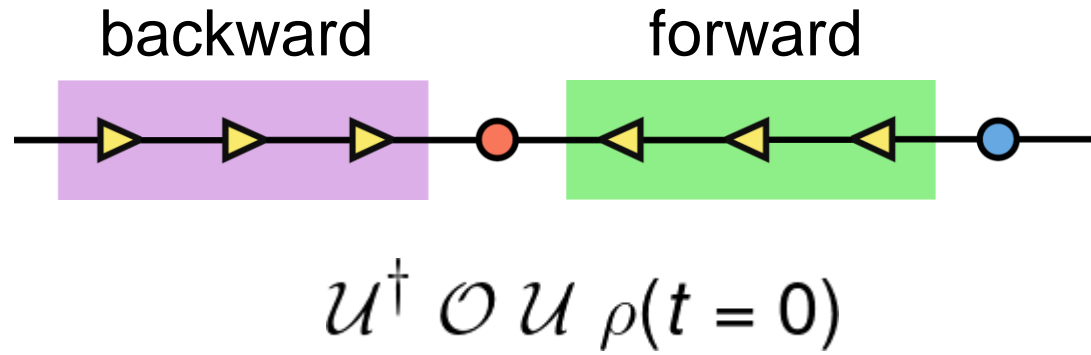


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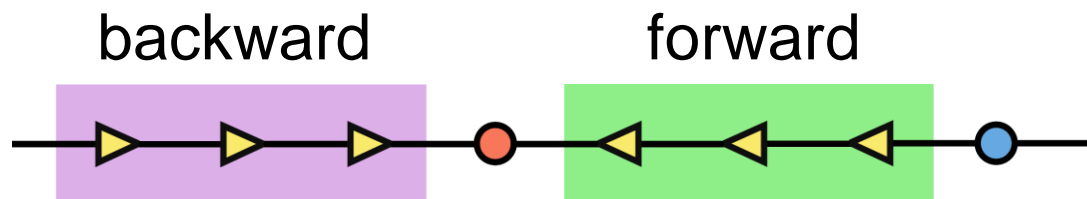
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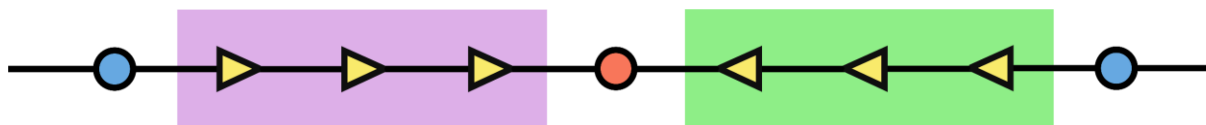
1. Quantum Evolution in Tensor Network Language



$$U^\dagger \circ U \rho(t=0)$$

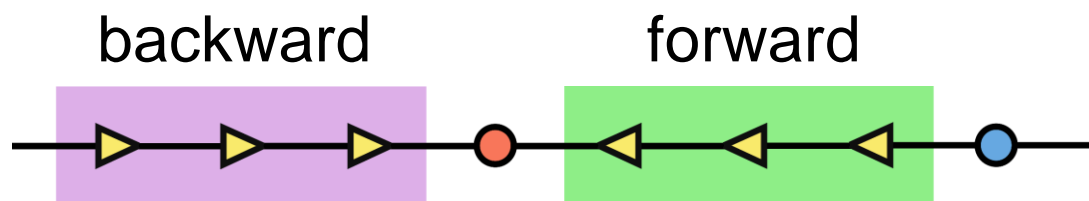


pure initial state



$$U^\dagger \circ U |\Psi(t=0)\rangle\langle\Psi(t=0)|$$

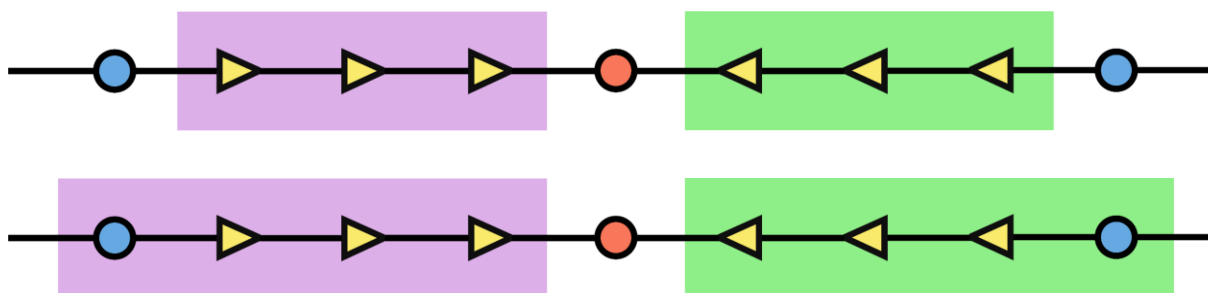
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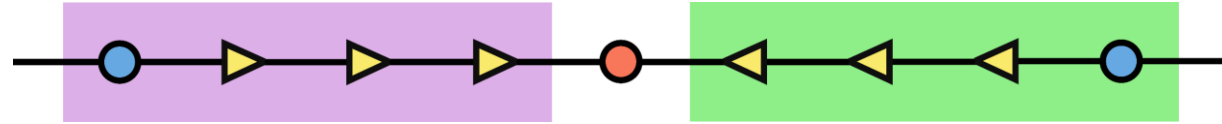


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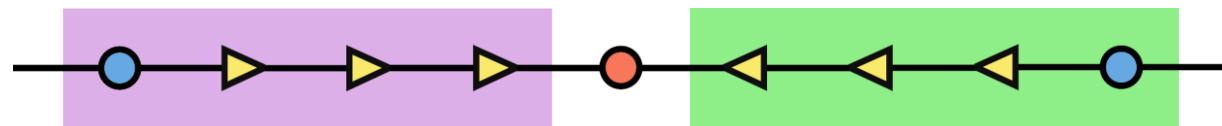
$$U^\dagger \circ U |\Psi(t=0)\rangle\langle\Psi(t=0)|$$

1. Quantum Evolution in Tensor Network Language



$$\mathcal{U}^\dagger \mathcal{O} \mathcal{U} |\psi(t=0)\rangle \langle\psi(t=0)|$$

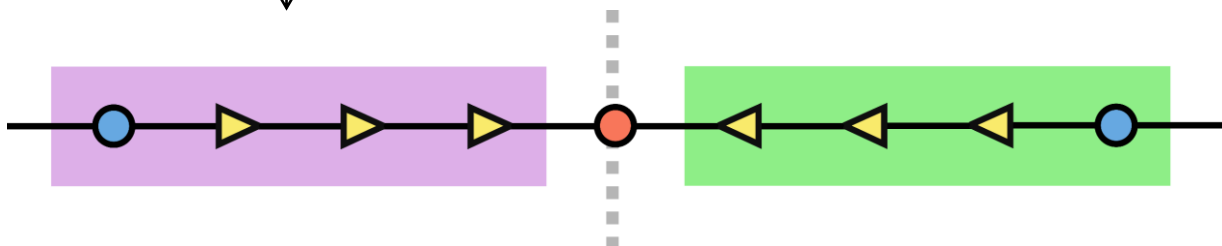
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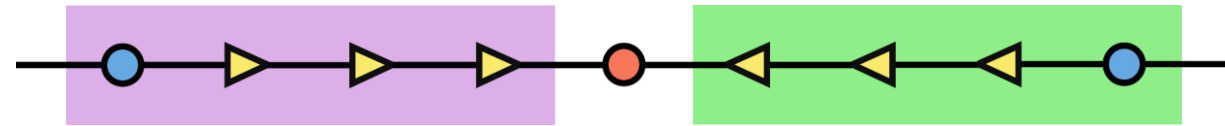
$$U^\dagger \circ U |\Psi(t=0)\rangle \langle \Psi(t=0)|$$



fold the network

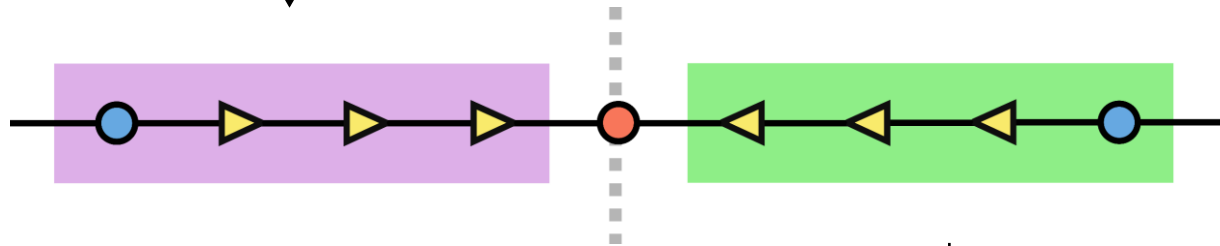


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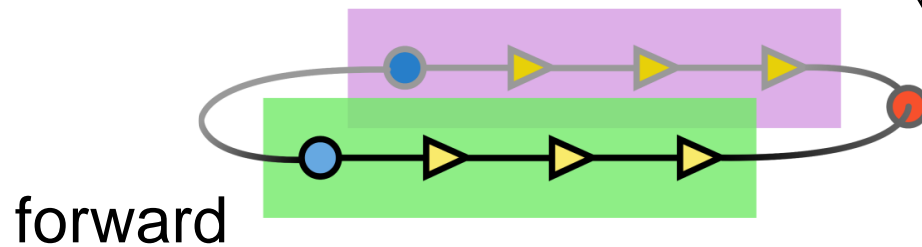


$$U^\dagger \circ U |\Psi(t=0)\rangle \langle \Psi(t=0)|$$

fold the network



backward ↓



trace (connect edges)

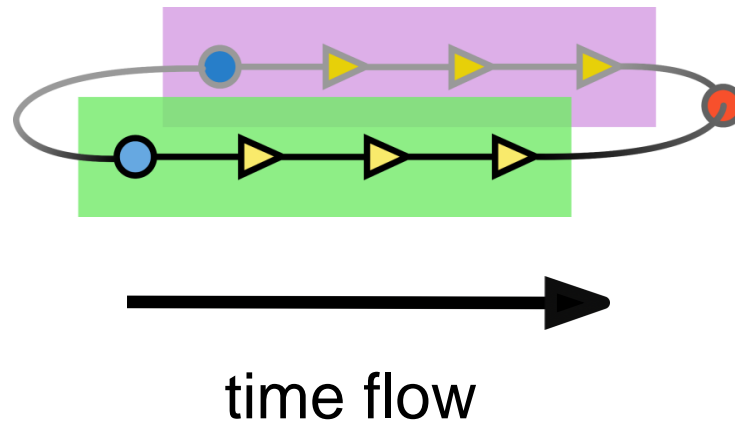
forward

$$\langle \Psi(t=0) | U^\dagger \circ U | \Psi(t=0) \rangle$$

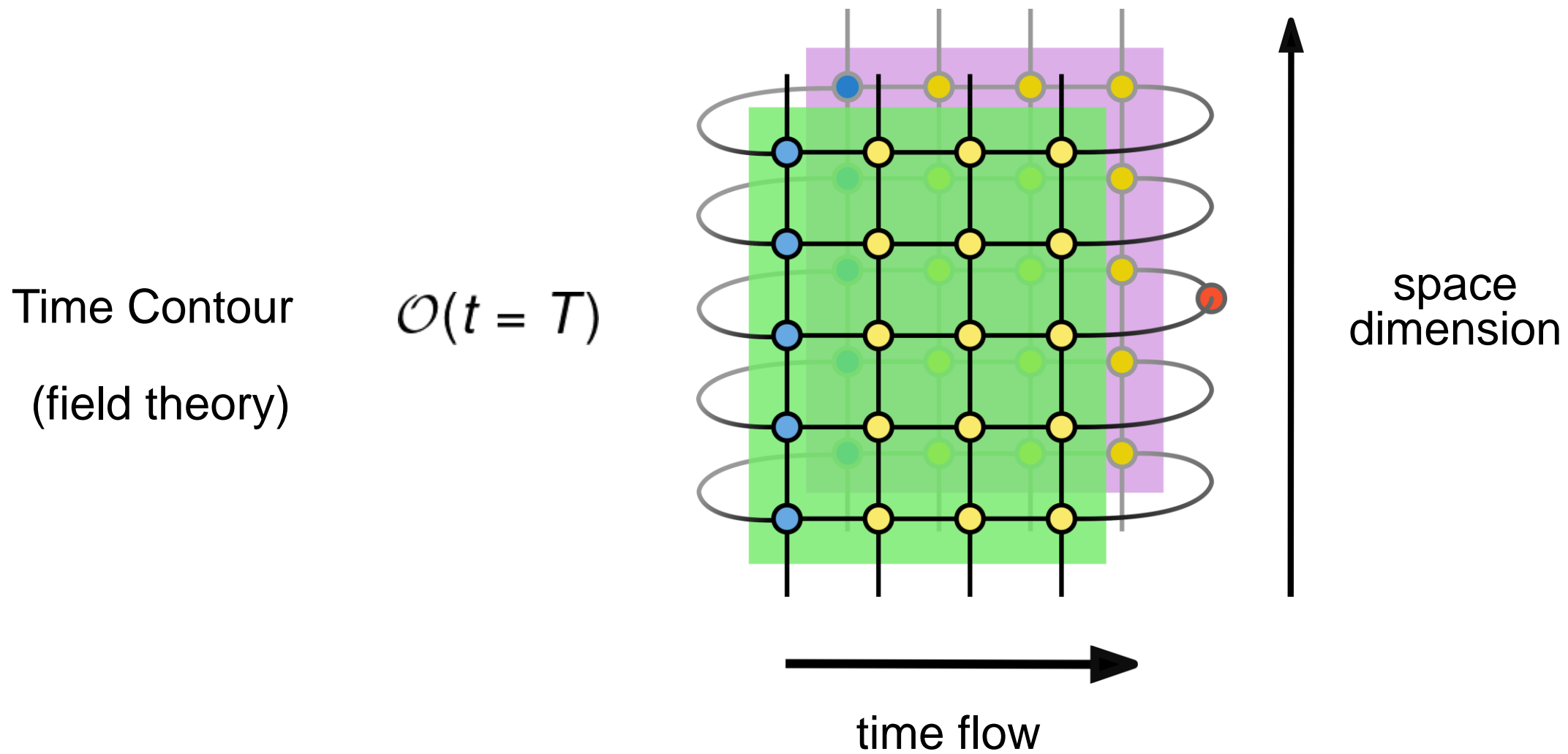
1. Quantum Evolution in Tensor Network Language

Time Contour
(field theory)

$$\mathcal{O}(t = T)$$



1. Quantum Evolution in Tensor Network Language

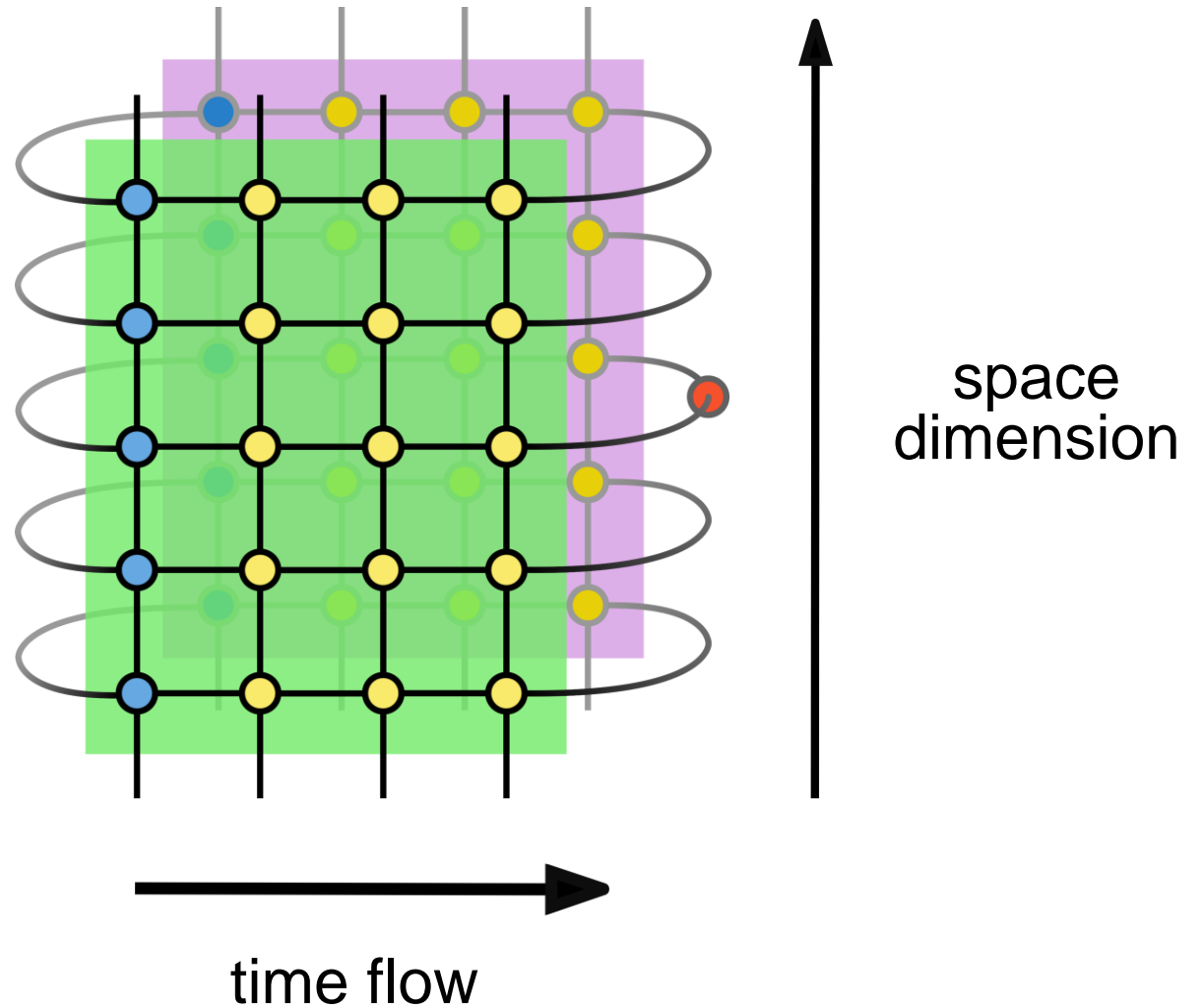


1. Quantum Evolution in Tensor Network Language

Time Contour
(field theory)

$$\mathcal{O}(t = T)$$

Can reach big sizes by **efficiently factoring** the operators into connected **local** tensors



2. Entanglement Barrier

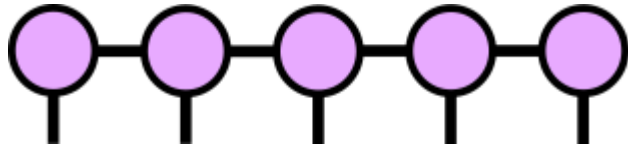
Swept under the carpet:

TNs work while **spatial entanglement is low**

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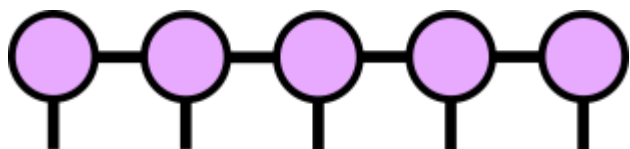


Matrix Product State

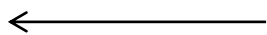
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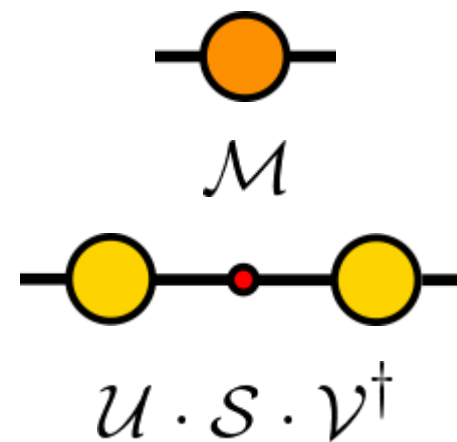
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Matrix Product State



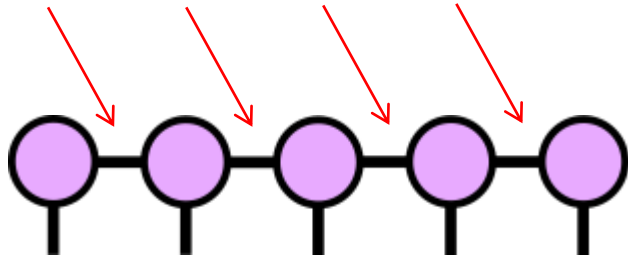
Singular
Value
Decomposition



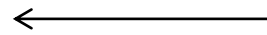
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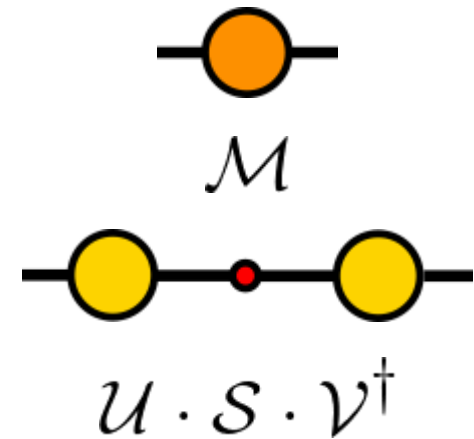
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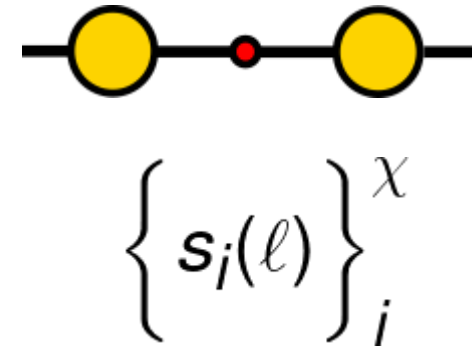


numerically: need **bounded number of singular values** per link

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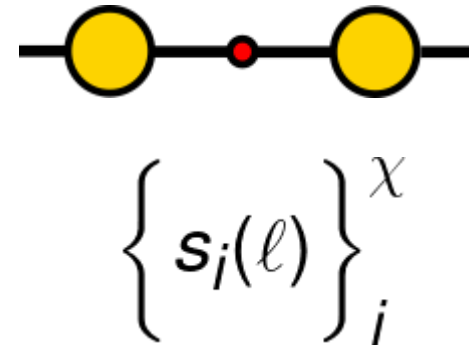
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Bipartite Entropy

$$S(\ell) = - \sum_i^{\chi} s_i^2(\ell) \cdot \log \left(s_i^2(\ell) \right)$$



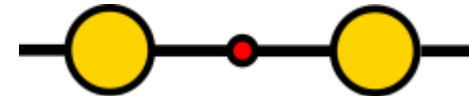
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$$\left\{ s_i(\ell) \right\}_i^{\chi}$$

numerically: need **bounded number of singular values** per link

physically: **subextensive entropy, no volume law**

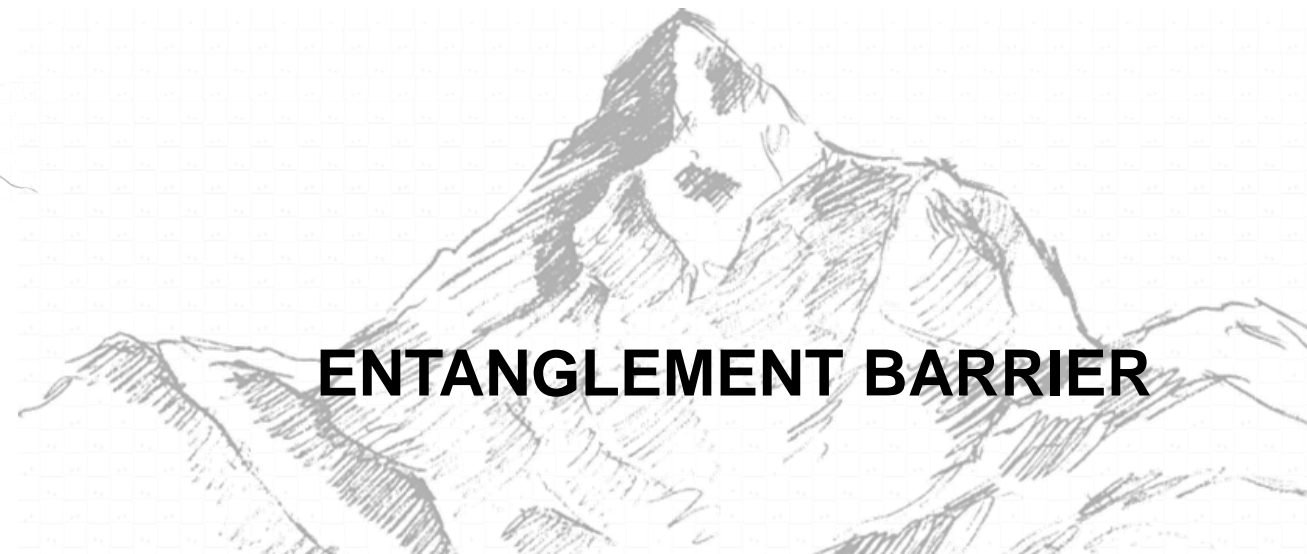
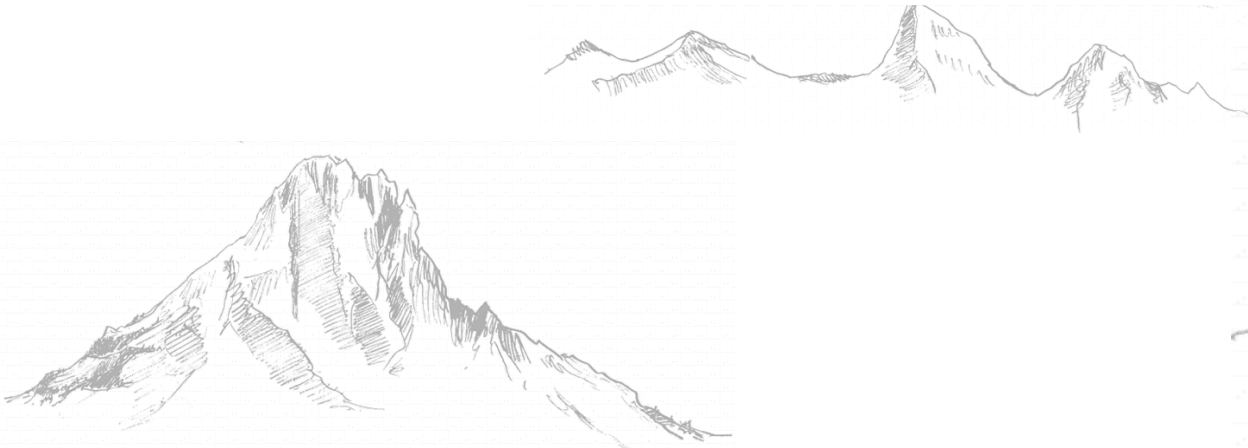
2. Entanglement Barrier

Nevertheless, evolution generically entangles

i.e. the sum of singular values **grows exponentially**

arXiv:0706.2480

arXiv:0903.2432



ENTANGLEMENT BARRIER

3. Dissipative Evolution

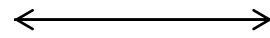
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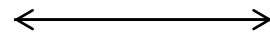
Our proposal is based on
decoherence

$$\begin{pmatrix} \rho_{00} & \rho_{01} \\ \rho_{10} & \rho_{11} \end{pmatrix}$$

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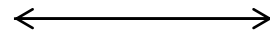


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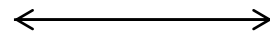
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We introduce a quantum channel targetting **unnecessary coherence**
for the description of **local observables**

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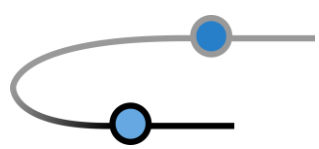
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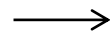
Connection to **thermalization**: reduced system is expected to decohere

3. Dissipative Evolution

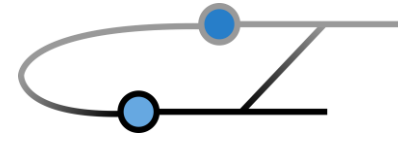
Engineer the quantum channel:



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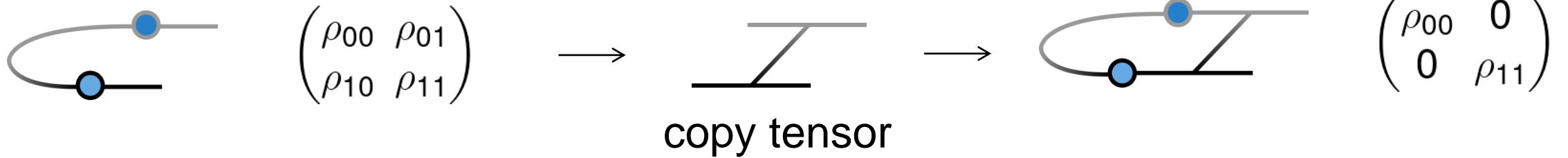
copy tensor



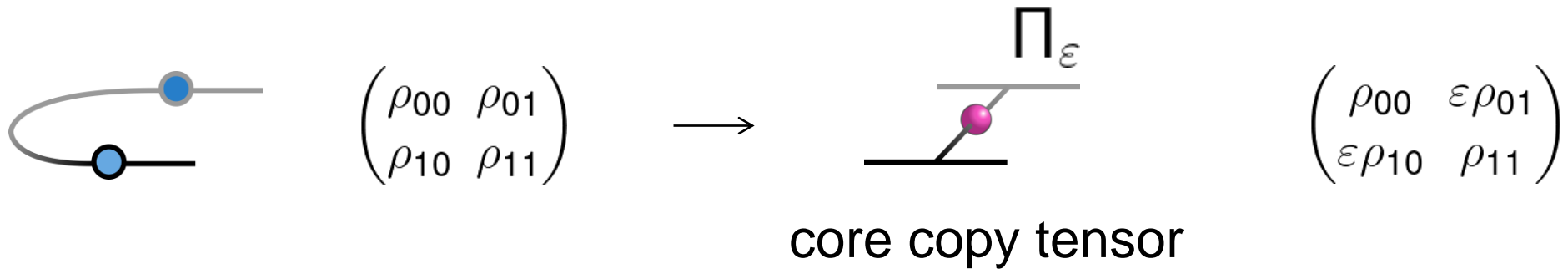
$$\begin{pmatrix} \rho_{00} & 0 \\ 0 & \rho_{11} \end{pmatrix}$$

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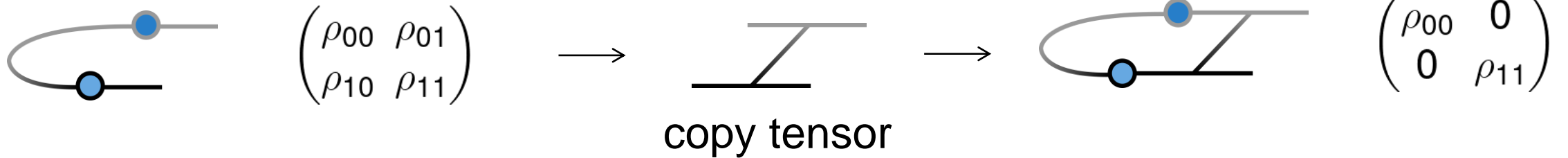


Can decide its intensity:

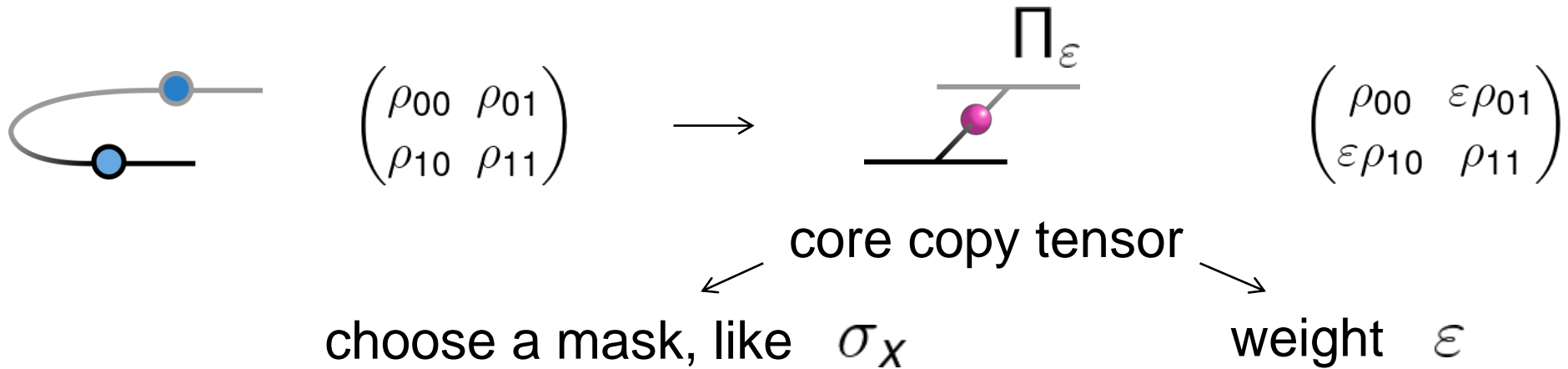


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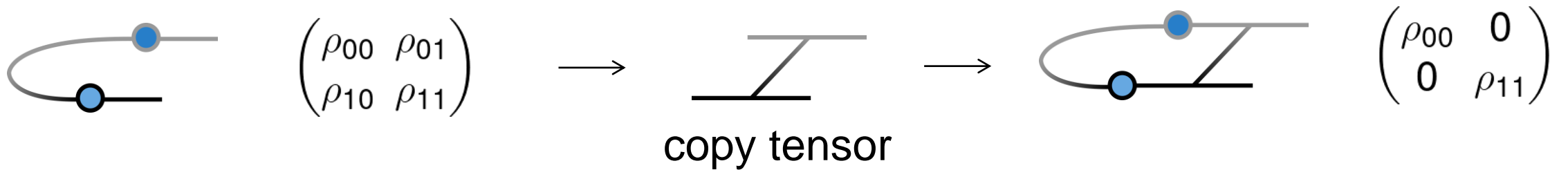


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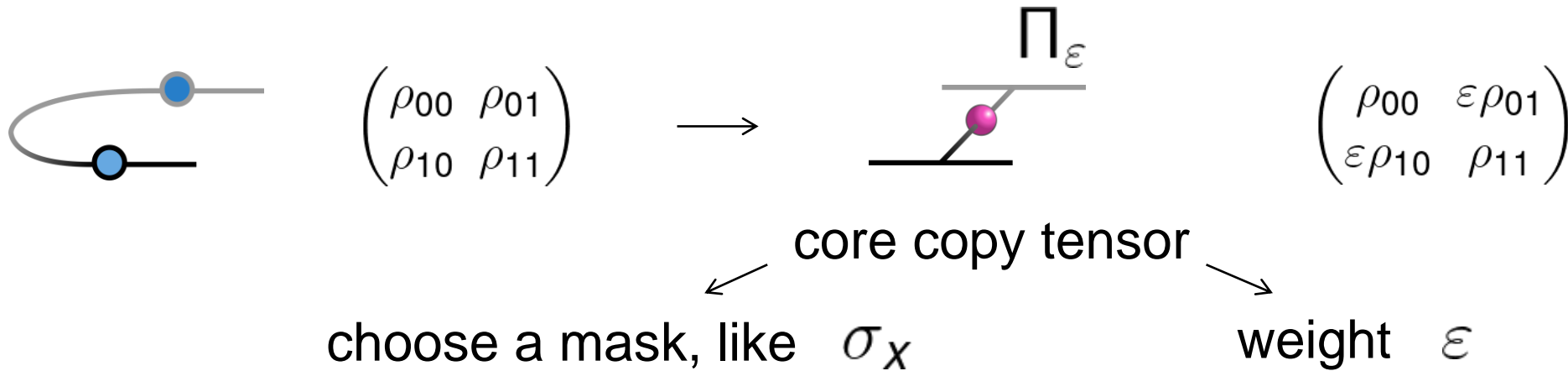


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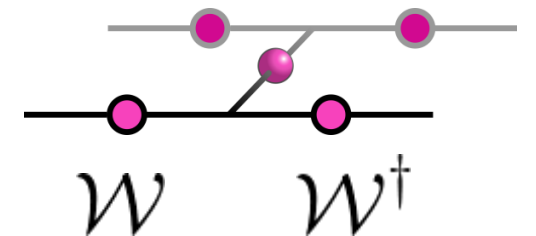
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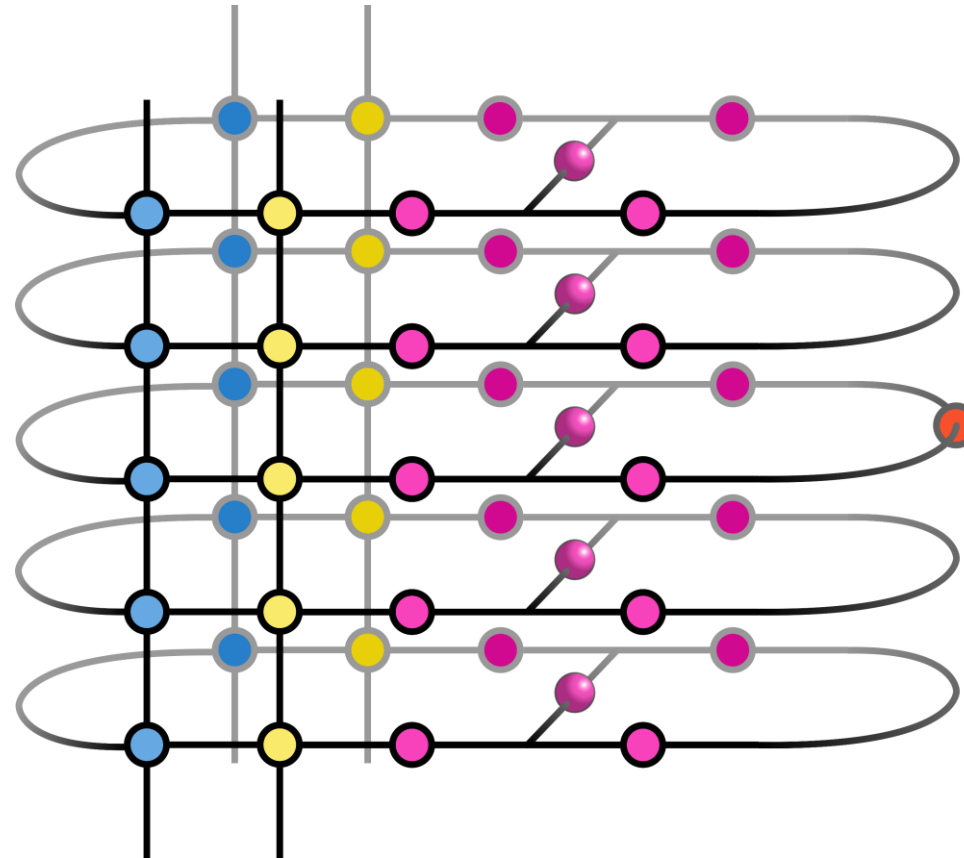
More freedom: the copy tensor is **not basis covariant**,
can introduce rotations



3. Dissipative Evolution

Design an auxiliary evolution expected to retrieve the right local observables:

one-step contour:

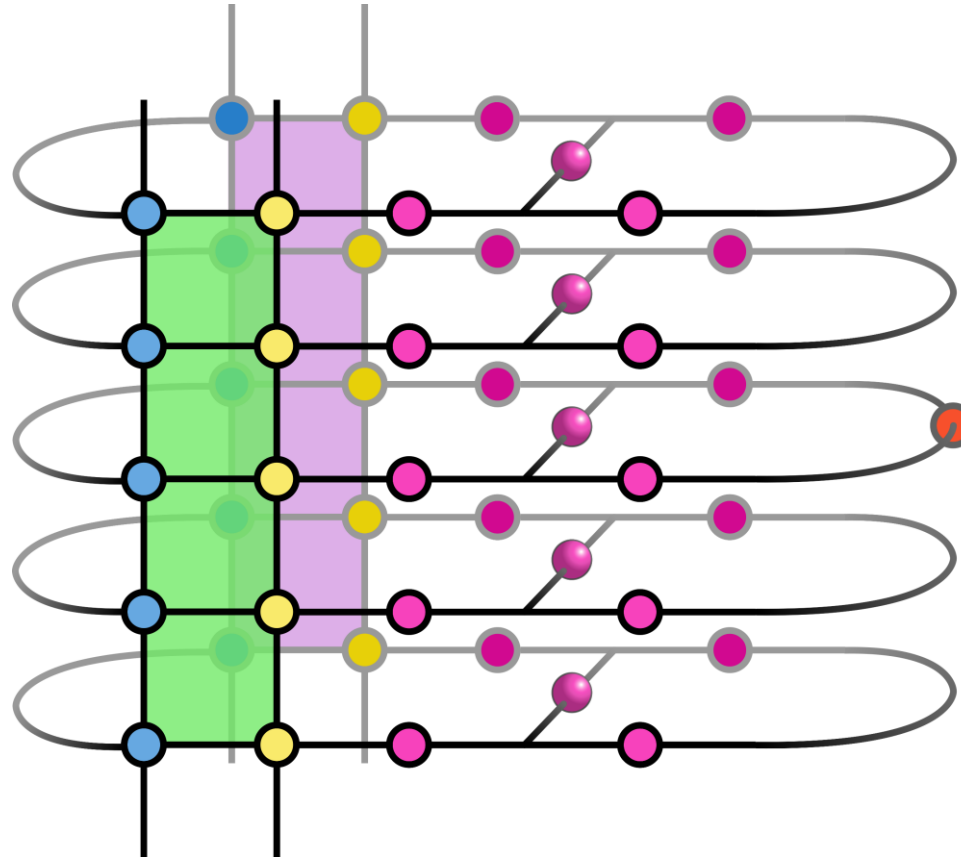


$$|\psi(t = 0)\rangle$$

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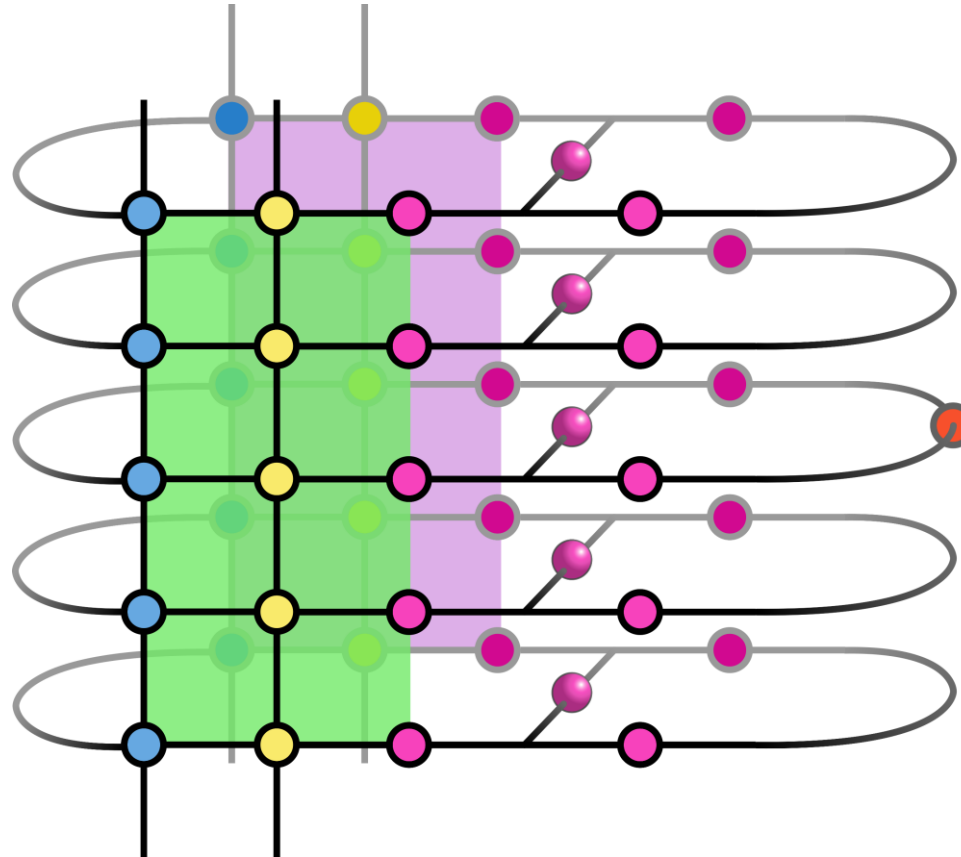


$$\delta\mathcal{U} |\psi(t=0)\rangle$$

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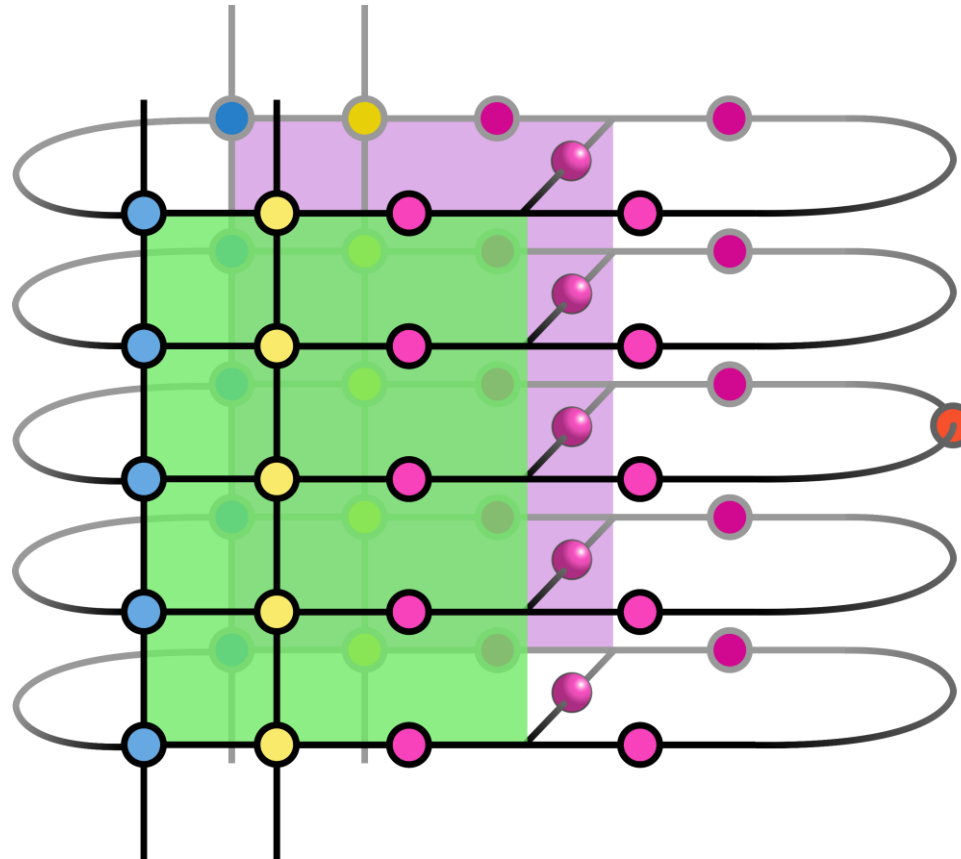


$$\mathcal{W} \delta \mathcal{U} |\psi(t=0)\rangle$$

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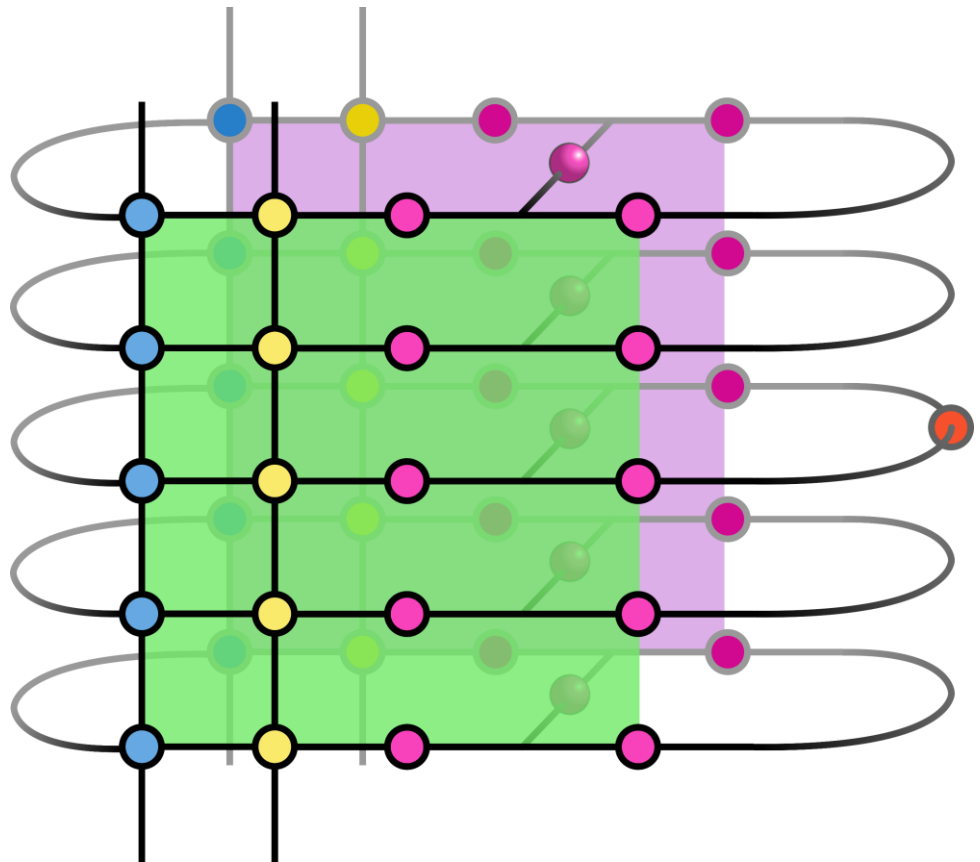


$$\Pi_{\varepsilon} \{ \mathcal{W} \delta \mathcal{U} | \Psi(t=0) \rangle \}$$

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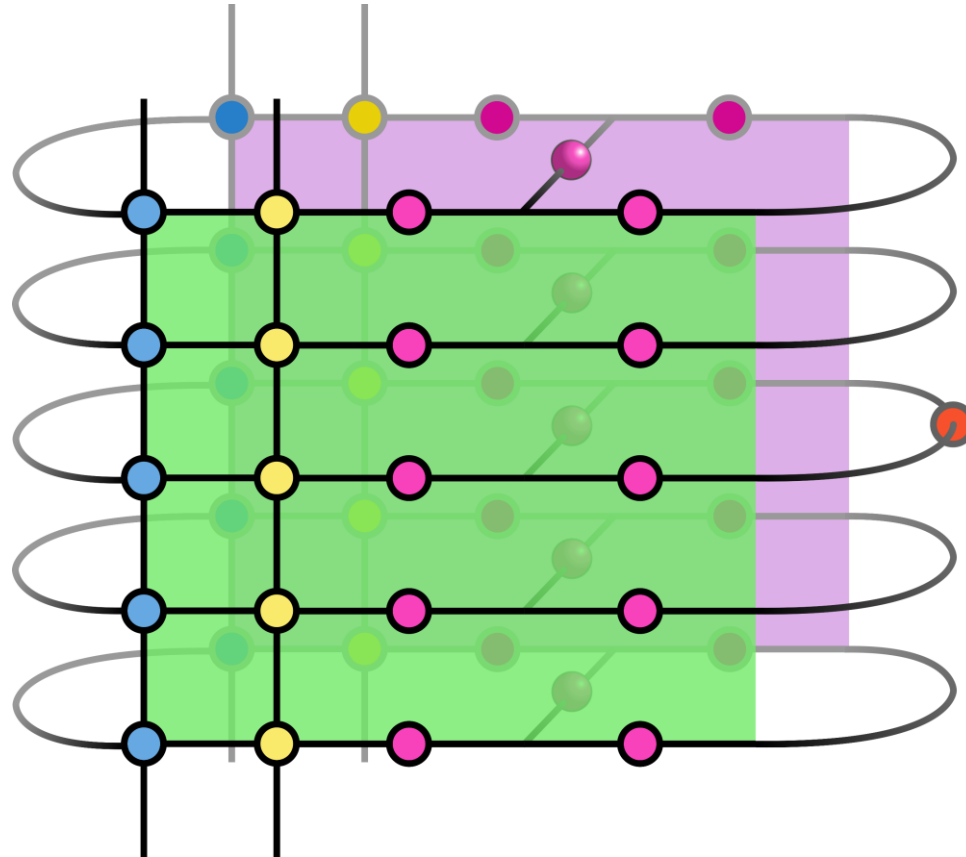


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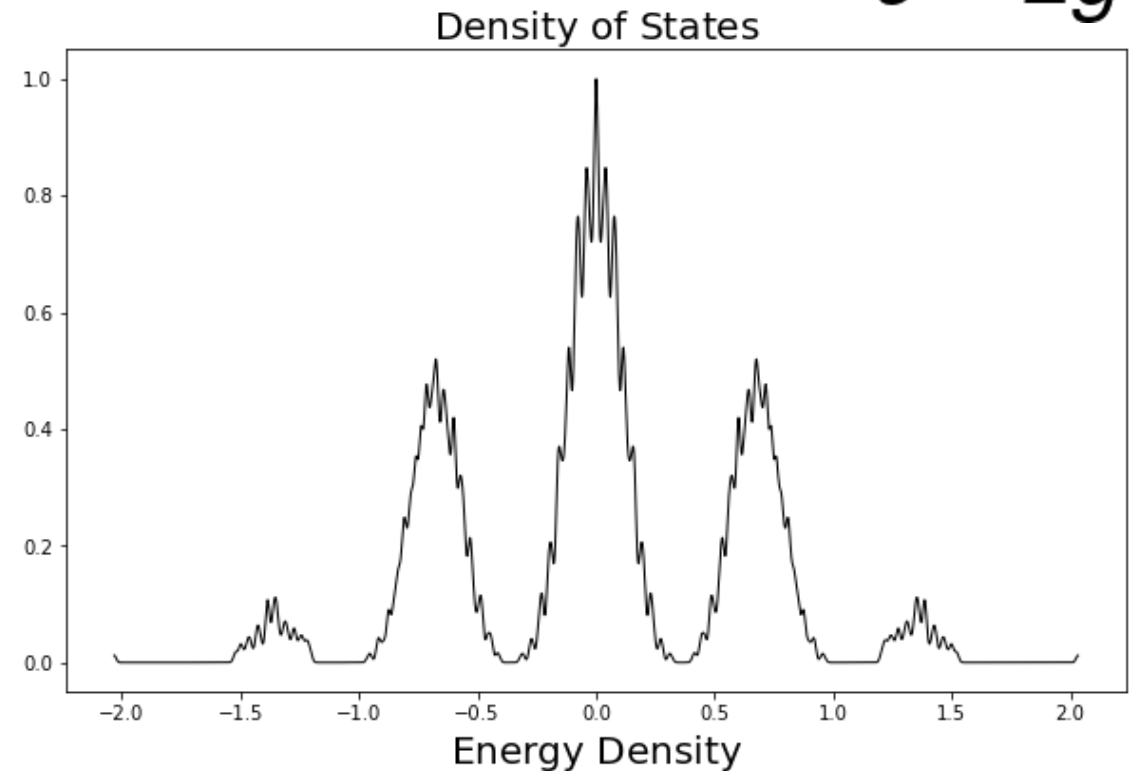
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4. Preliminary Results

Currently exploring **small systems** $L \leq 12$ for the **Ising model**

$$\mathcal{H} = -J \sum_i \sigma_i^z \sigma_{i+1}^z - g \sum_i \sigma_i^x$$

$$J = 2g$$



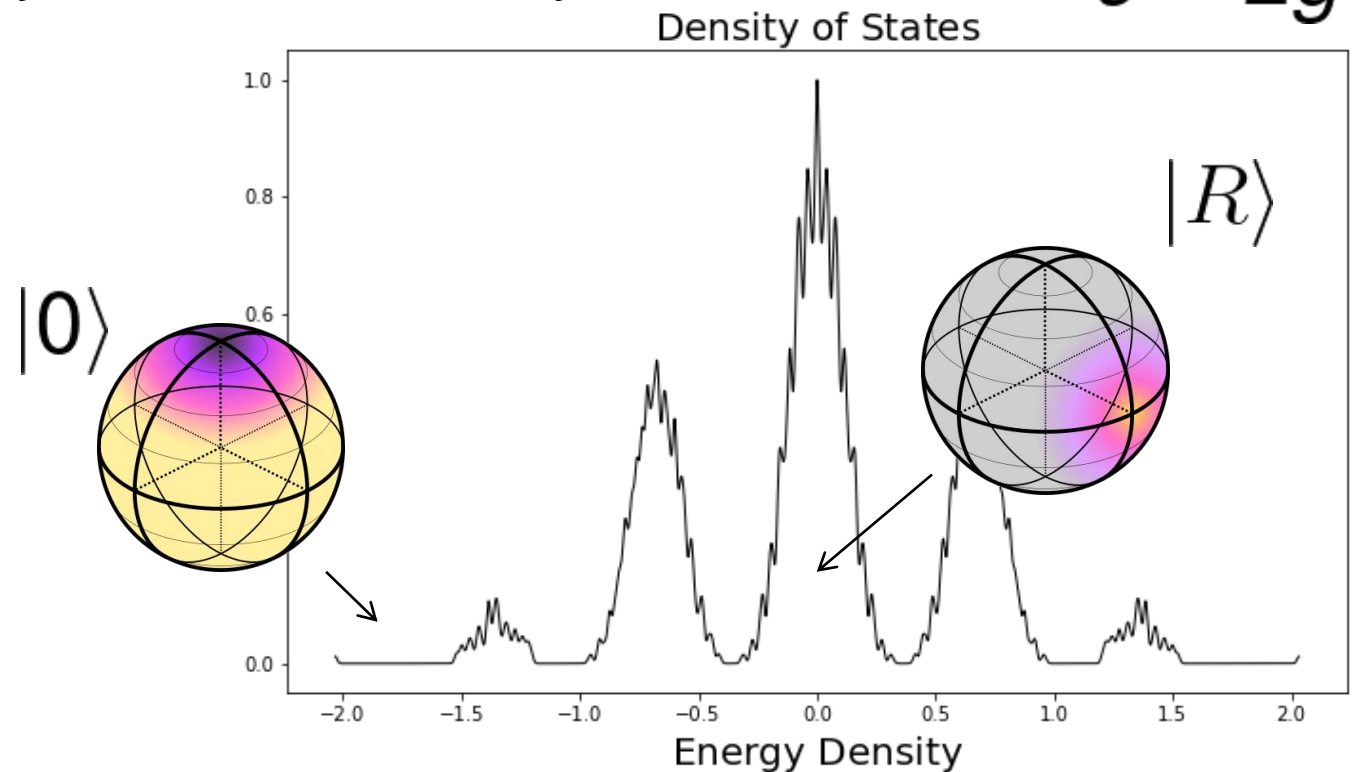
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and initial **Product States**
parametrized by a direction in
the Bloch sphere



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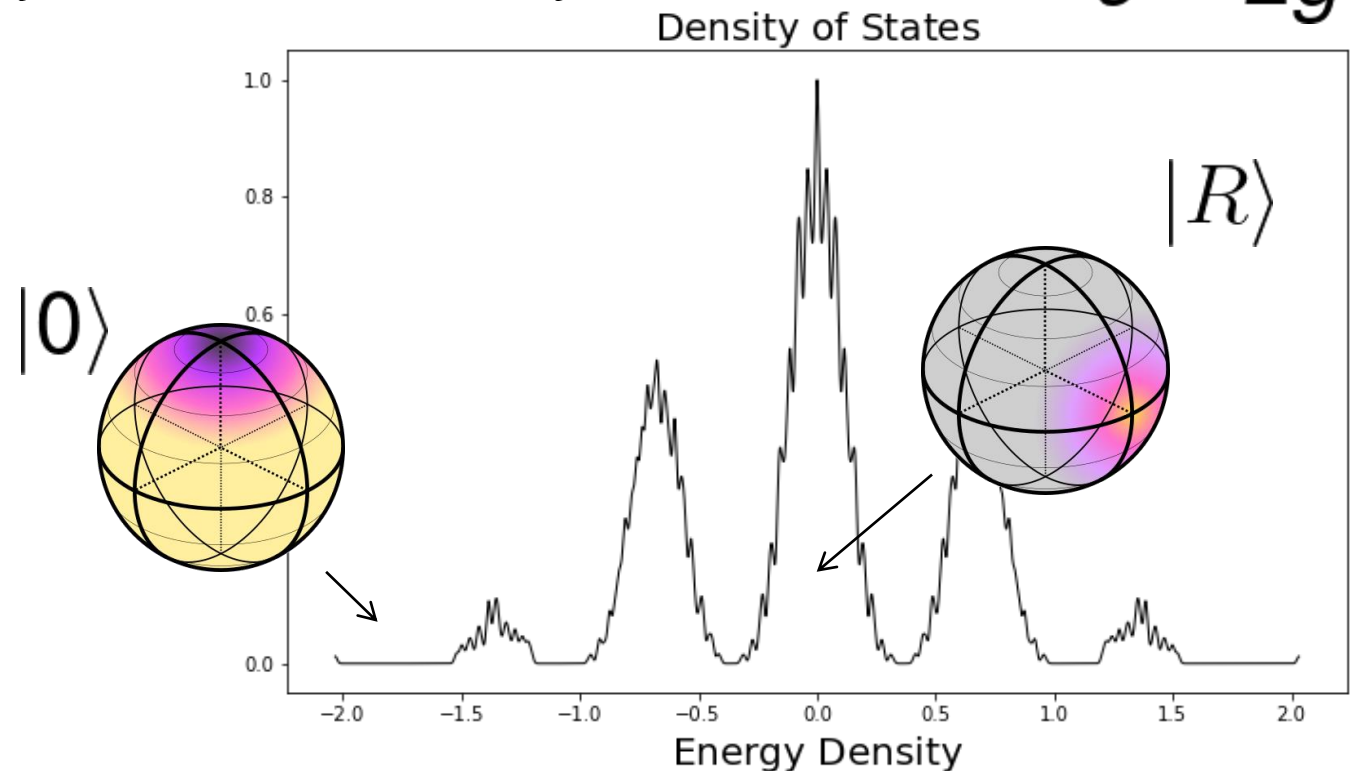
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parametrized by a direction in
the Bloch sphere

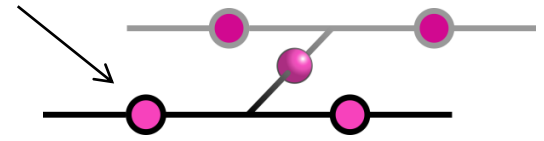
corresponding to many
energies (temperatures)



4. Preliminary Results

Basis election: **parallel magnetization**

$$M^{\parallel} = \sum_i \sigma_i^{\parallel}$$

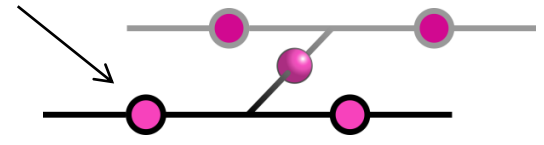


4. Preliminary Results

Basis election: **parallel magnetization**

$$M^{\parallel} = \sum_i \sigma_i^{\parallel}$$

(1) readily diagonal

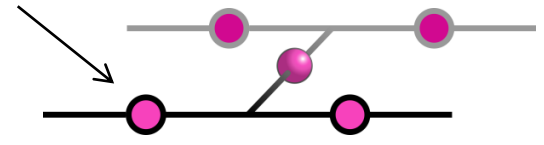


4. Preliminary Results

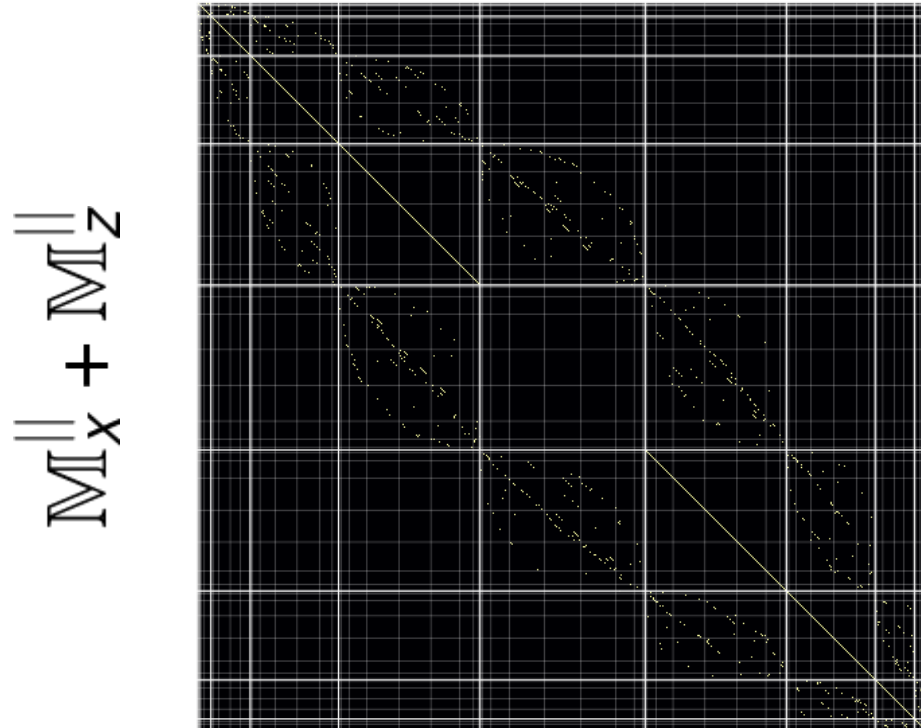
Basis election: **parallel magnetization**

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(1) readily diagonal



(2) magnetization sectors: **tidy and target**
particular coherences

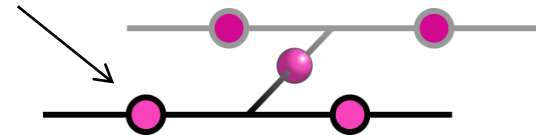


4. Preliminary Results

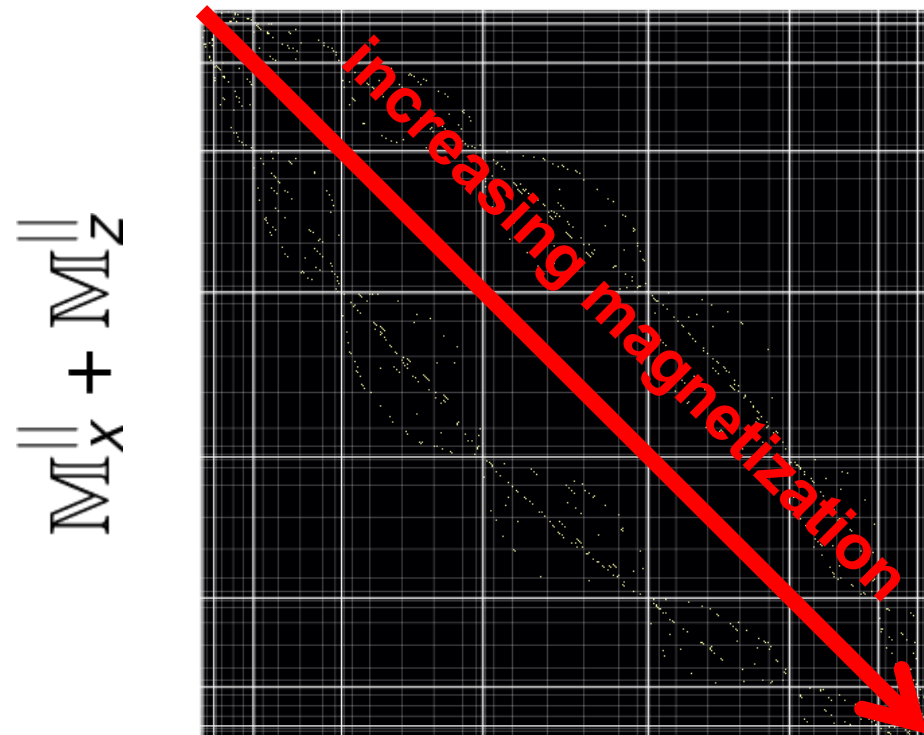
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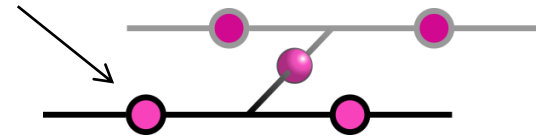
(2) magnetization sectors: **tidy and target**
particular coherences



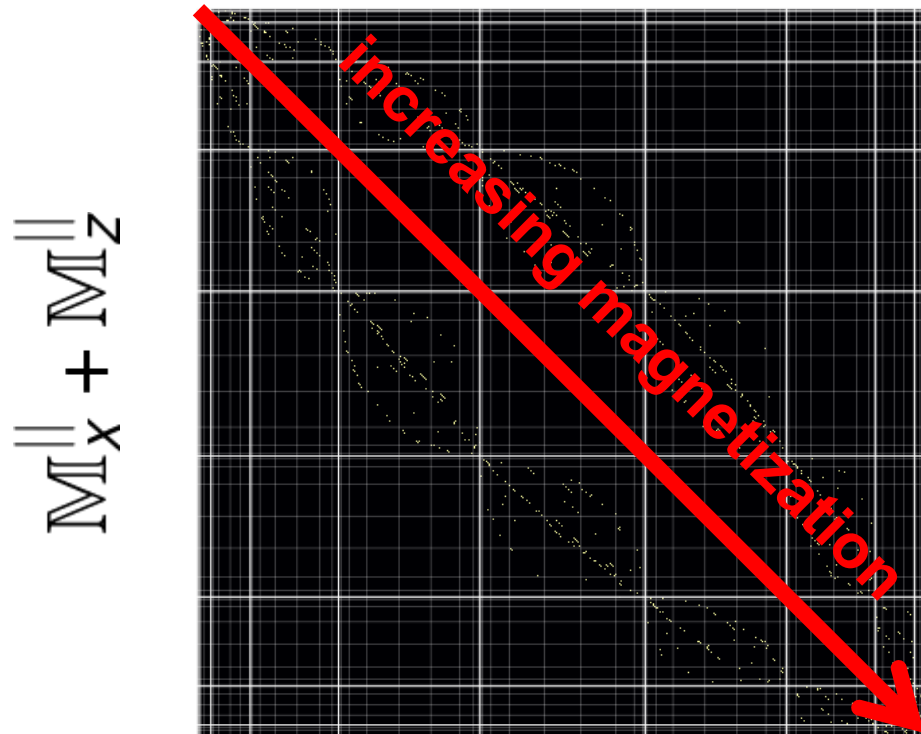
4. Preliminary Results

Basis election: **parallel magnetization**

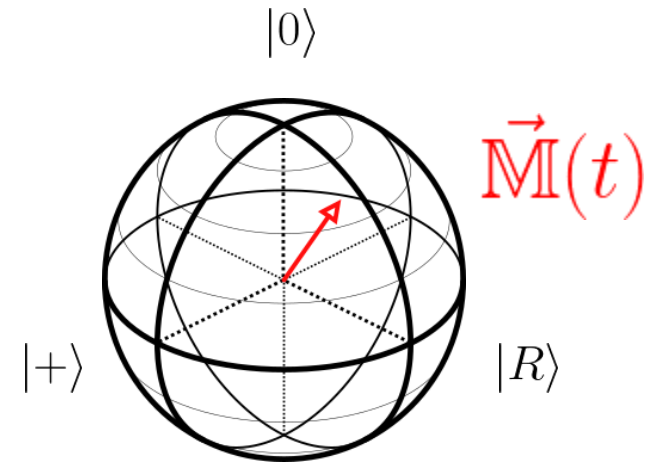
$$\mathbb{M}^{\parallel} = \sum_i \sigma_i^{\parallel} \quad (1) \text{ readily diagonal}$$



(2) magnetization sectors: **tidy and target**
particular coherences



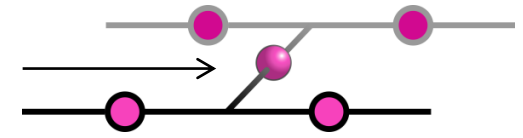
(3) expect strong overlap with
instantaneous state for short times



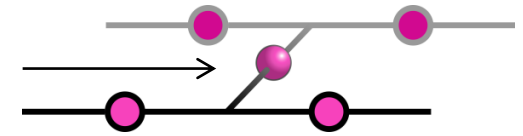
0-th order approximation,
system entangles

4. Preliminary results

Core election: action on magnetization eigenstate $|\Psi_i\rangle$



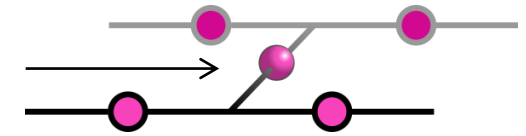
4. Preliminary results



Core election: action on magnetization eigenstate $|\Psi_i\rangle$

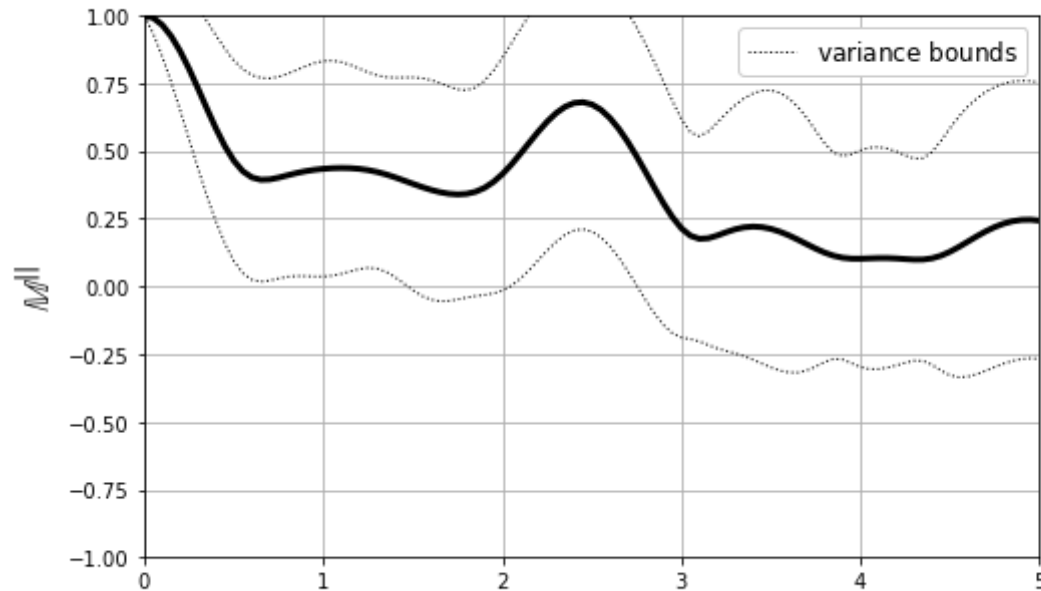
$$\begin{cases} 1 & \text{generically,} \\ e^{-\gamma(\mathbb{M} - \mathbb{I}\langle\mathbb{M}(t)\rangle)} & \text{if } |\langle\Psi_i|\mathbb{M}|\Psi_i\rangle - \langle\mathbb{M}(t)\rangle| > \text{Var}[\mathbb{M}(t)] \end{cases}$$

4. Preliminary results



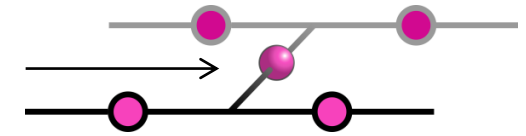
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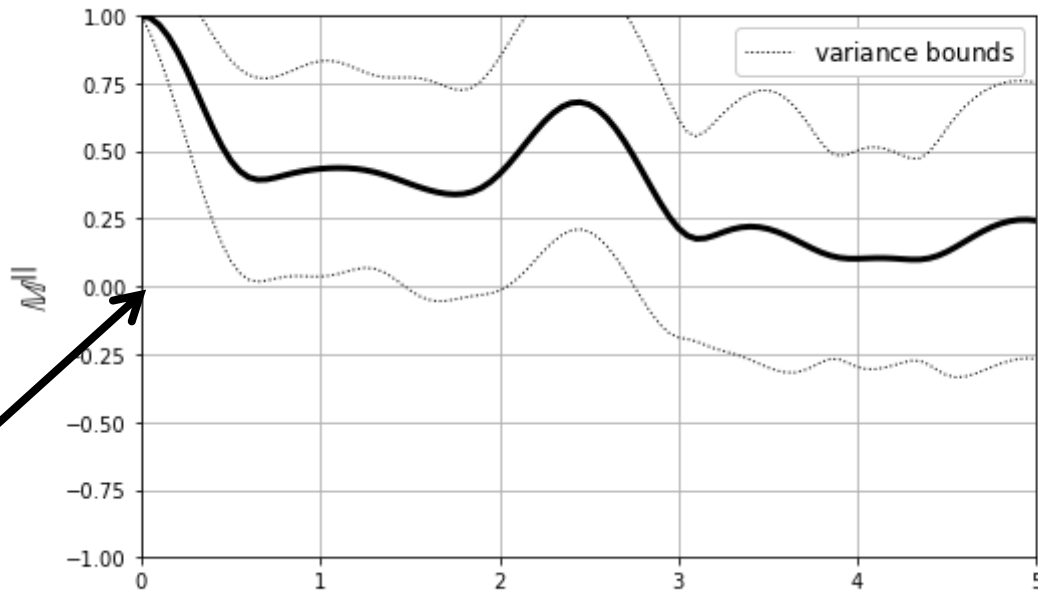
evolution of $|R\rangle$ for $J = 2g$

4. Preliminary results



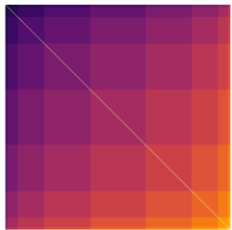
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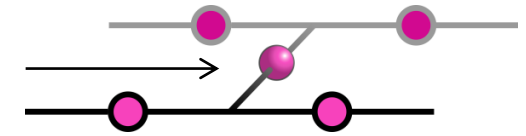


evolution of $|R\rangle$ for $J = 2g$

the quantum channel **masks** the coherences

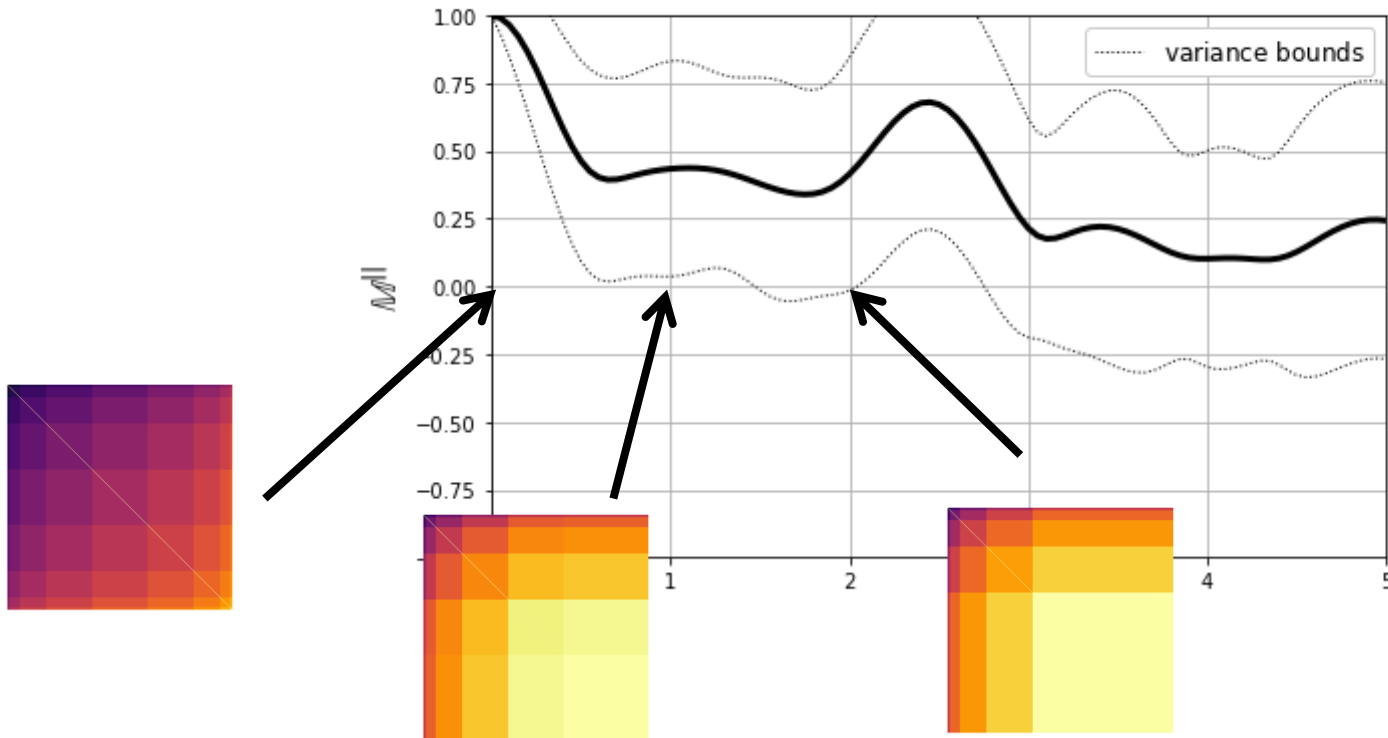


4. Preliminary results



Core election: action on magnetization eigenstate $|\Psi_i\rangle$

$$\begin{cases} 1 & \text{generically,} \\ e^{-\gamma(\mathbb{M} - \mathbb{I}\langle\mathbb{M}(t)\rangle)} & \text{if } |\langle\Psi_i|\mathbb{M}|\Psi_i\rangle - \langle\mathbb{M}(t)\rangle| > \text{Var}[\mathbb{M}(t)] \end{cases}$$



evolution of $|R\rangle$ for $J = 2g$

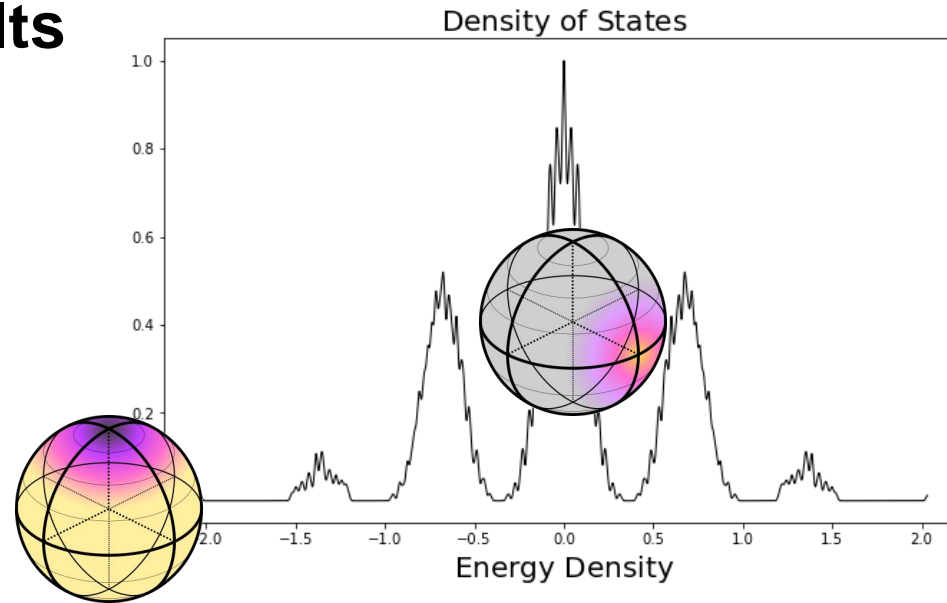
the quantum channel **masks** the coherences

time dependent!

follow the change in variance

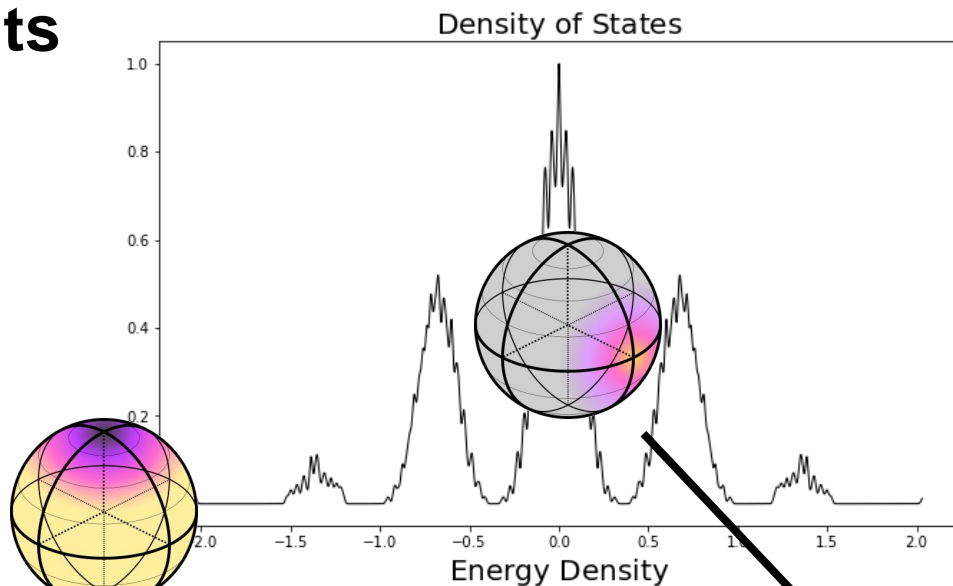
4. Preliminary Results

Two extreme cases:
close to **GS** and to
infinite temperature

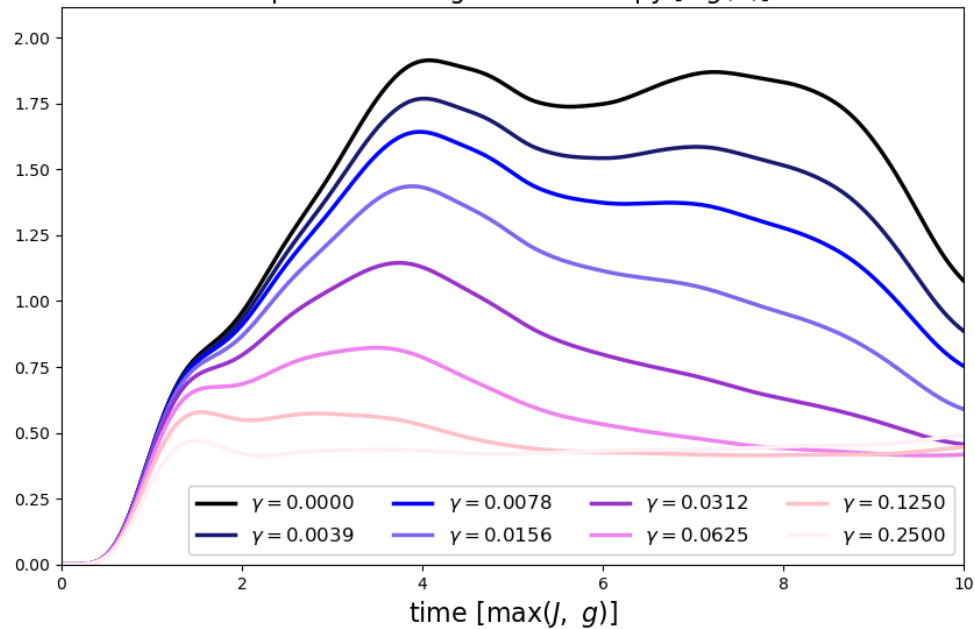


4. Preliminary Results

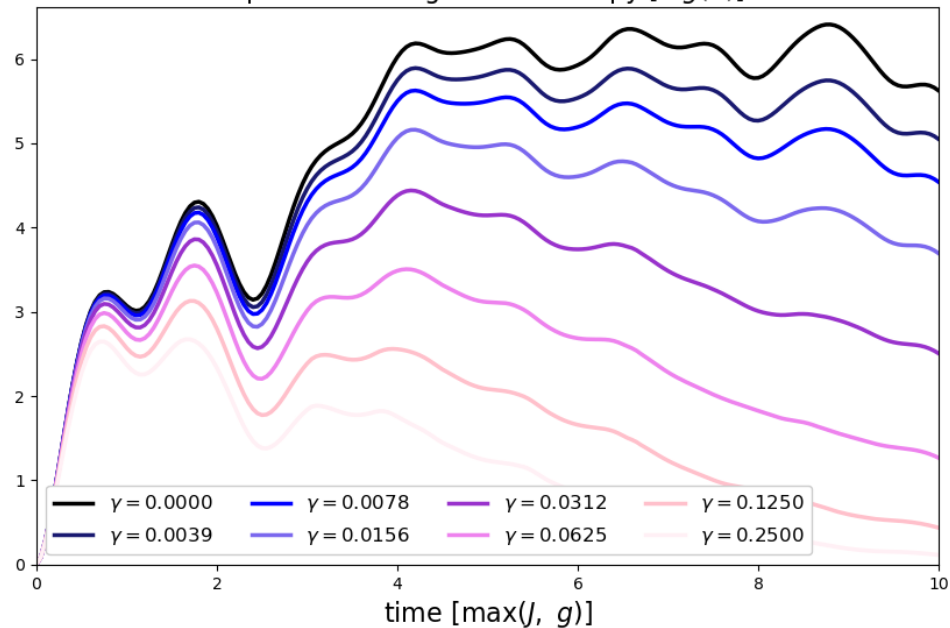
Two extreme cases:
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operator entanglement entropy [$\log(2)$]

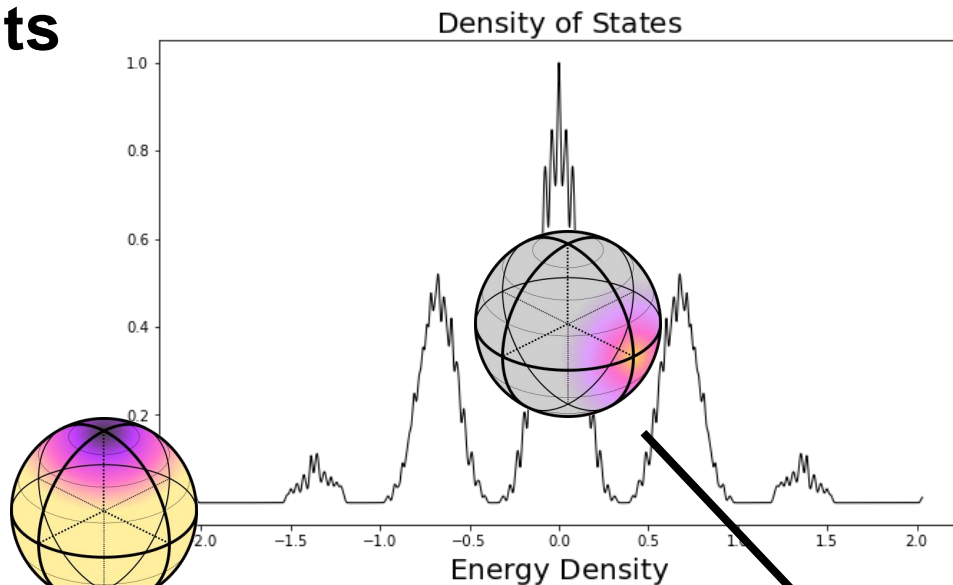


operator entanglement entropy [$\log(2)$]



4. Preliminary Results

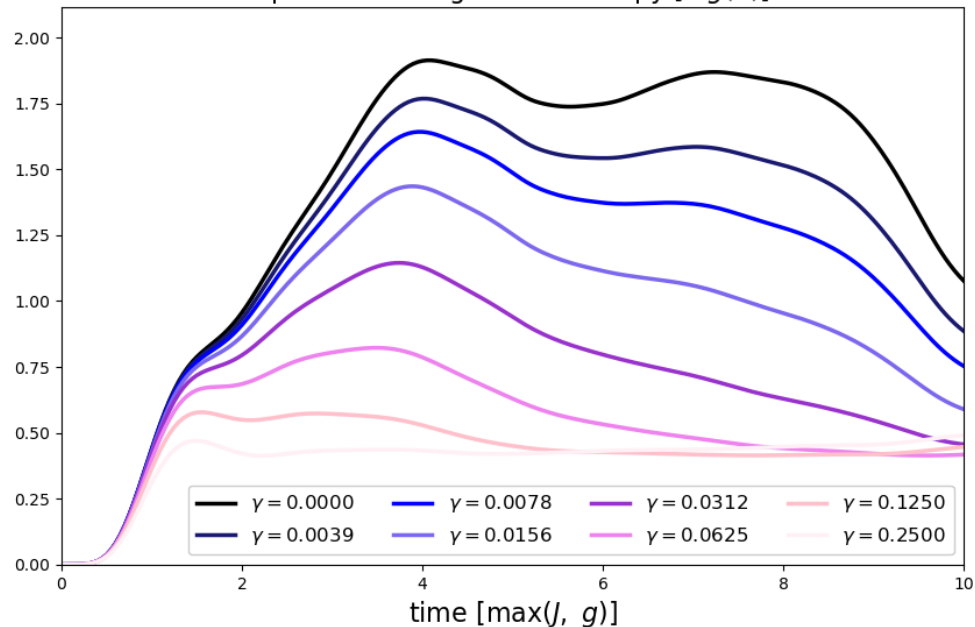
Two extreme cases:
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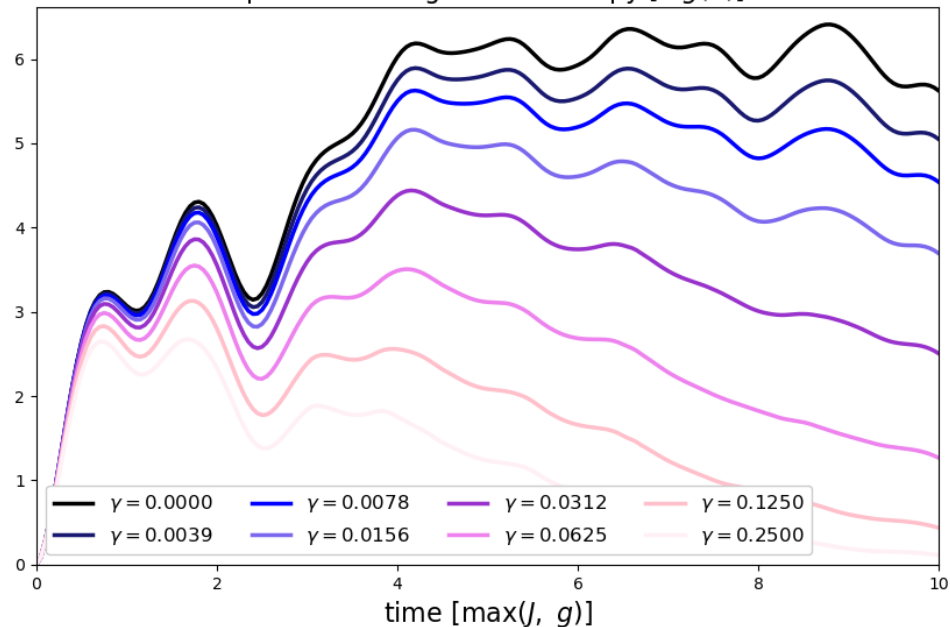
Observe a **reduction of entanglement** with increasing strength of decoherence

might improve simulability

operator entanglement entropy [$\log(2)$]

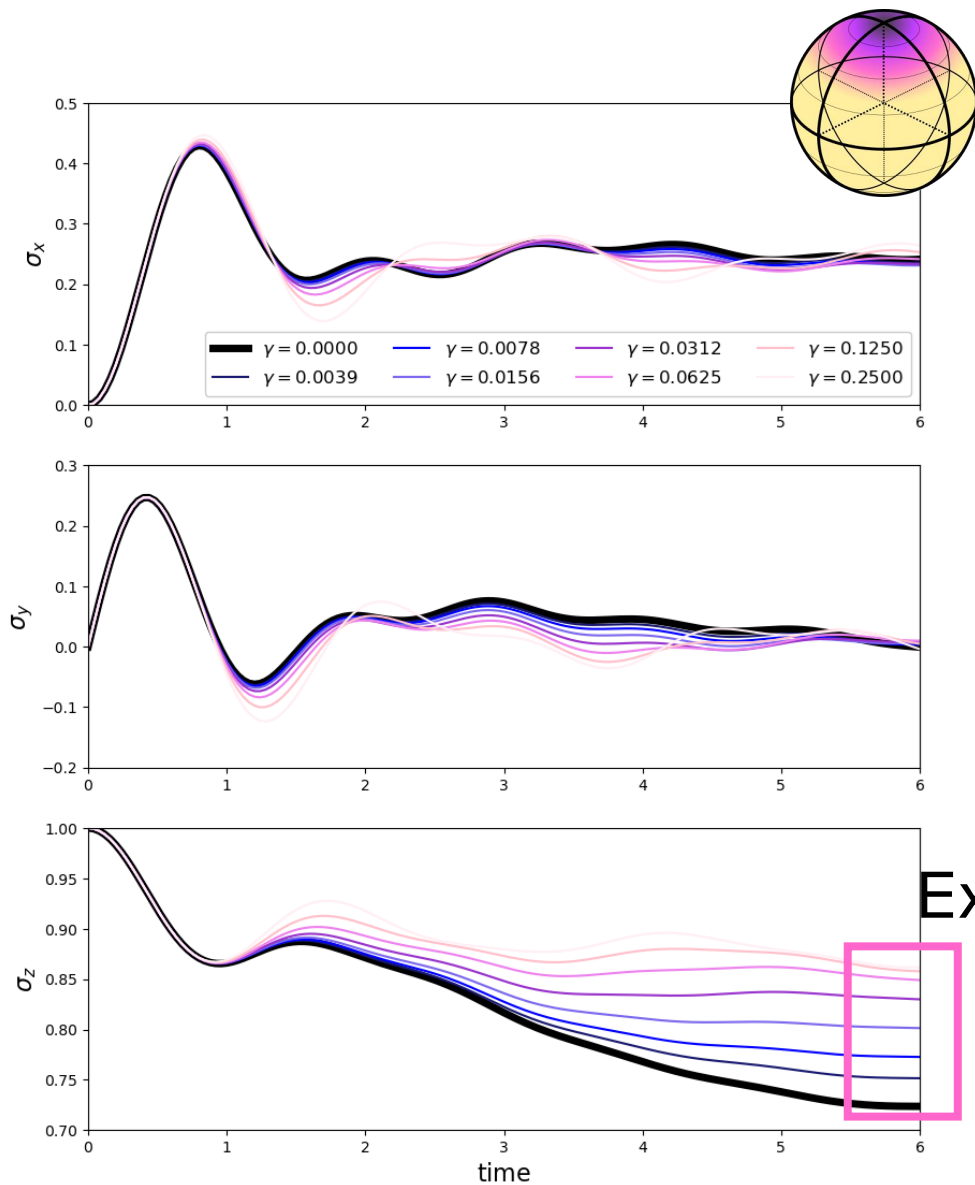


operator entanglement entropy [$\log(2)$]



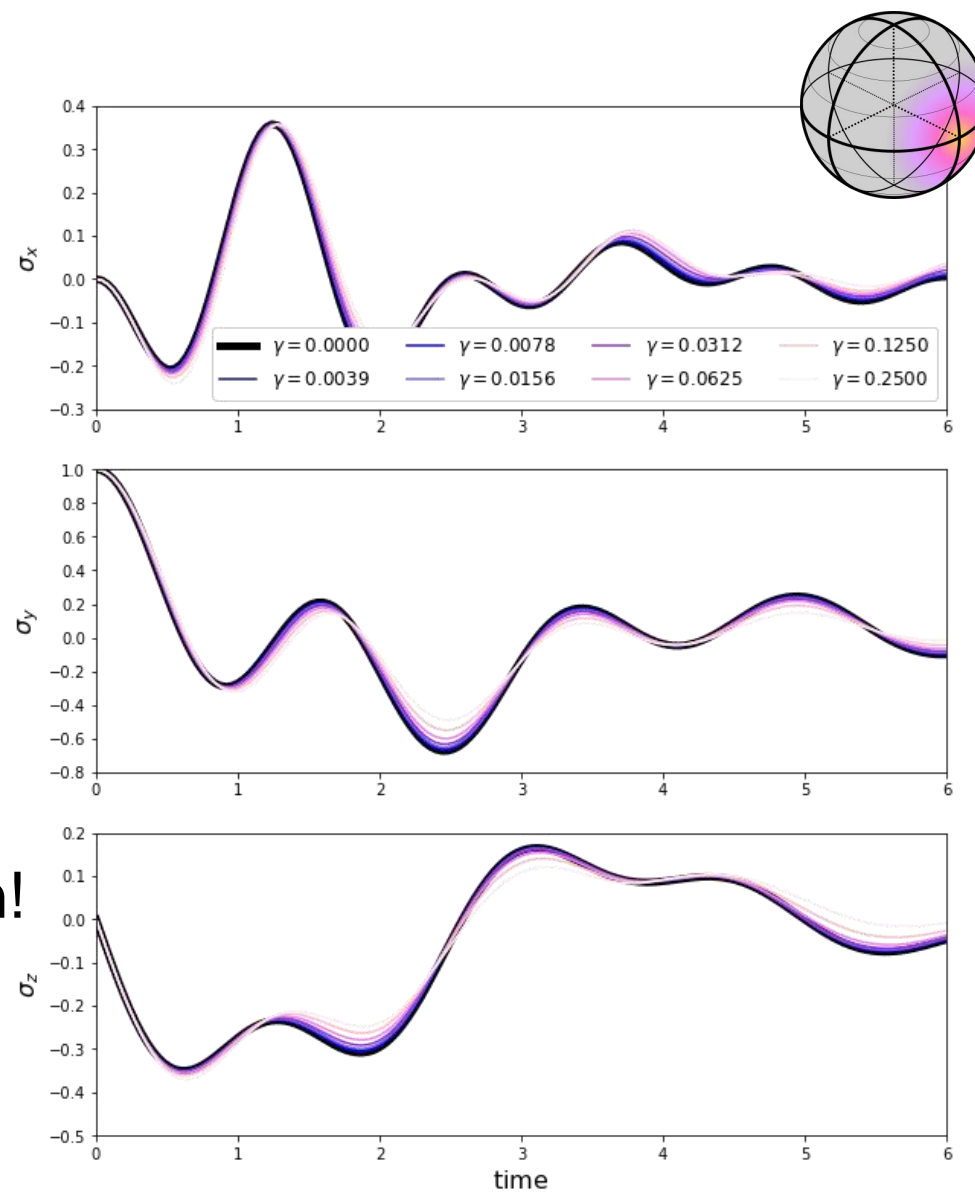
4. Preliminary Results

'Good performance'



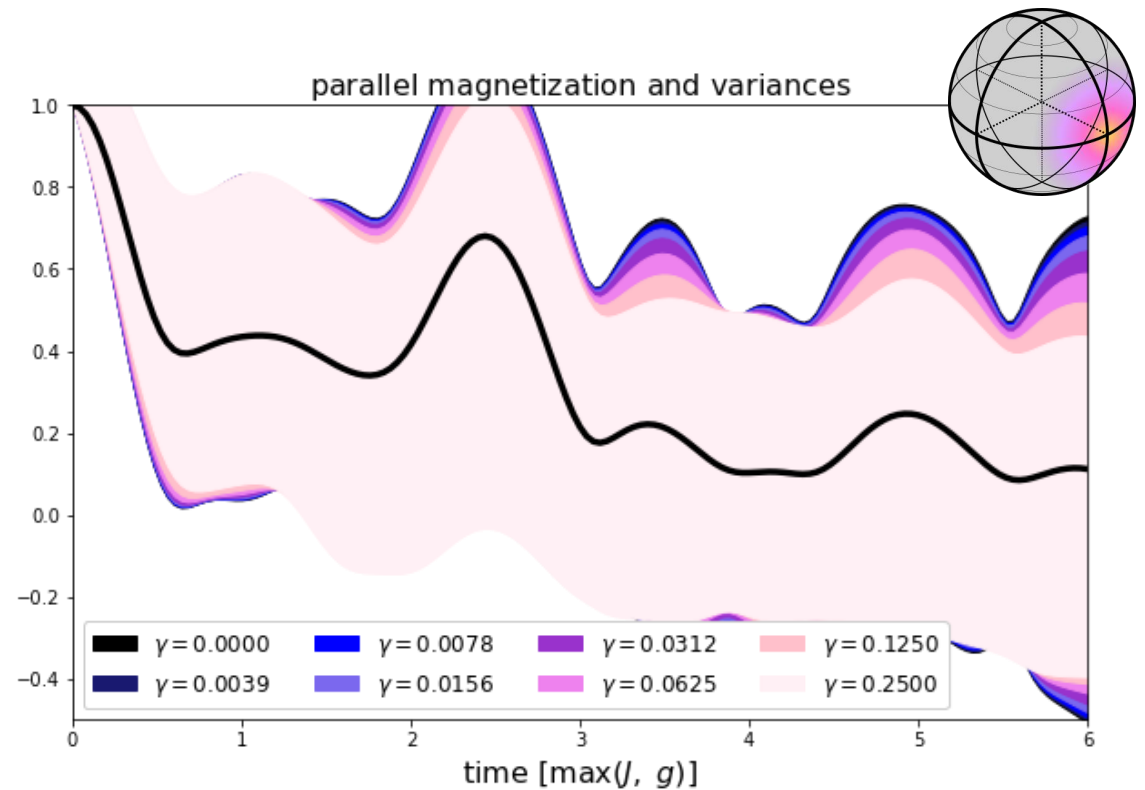
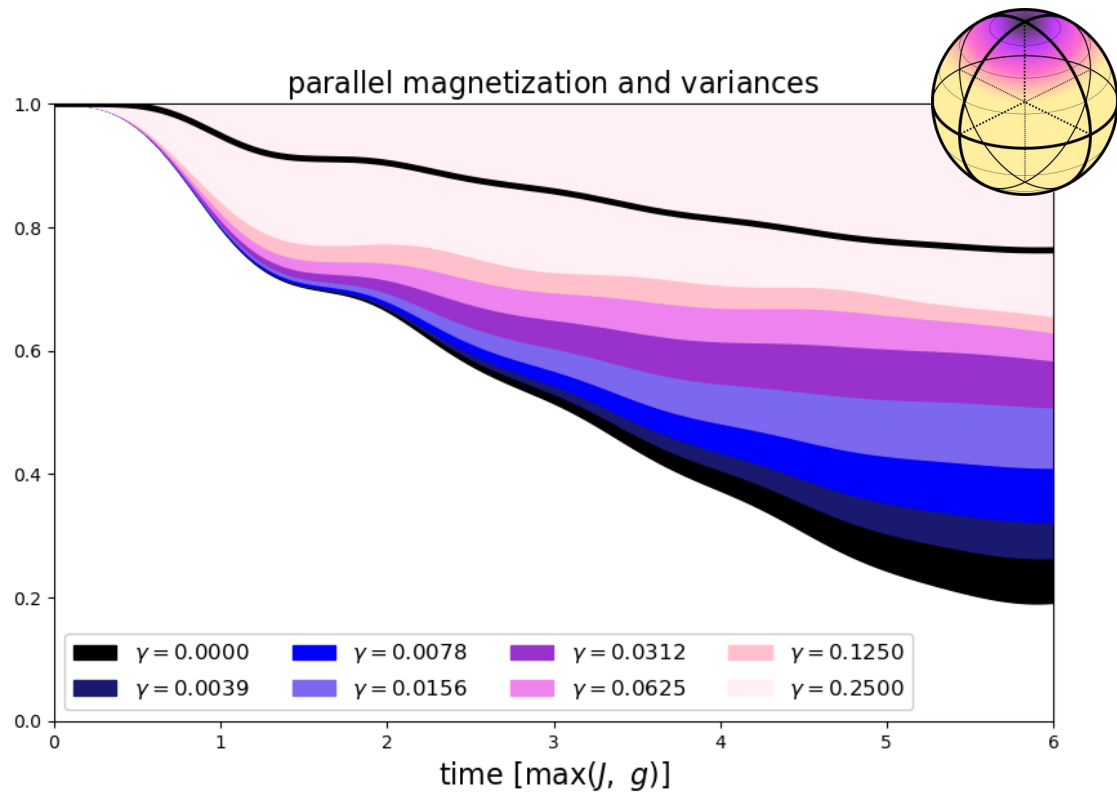
Extrapolation!

$$\gamma \rightarrow 0$$

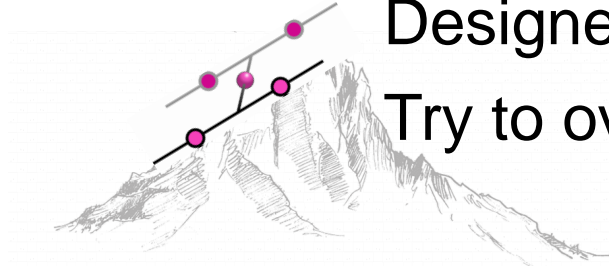


4. Preliminary Results

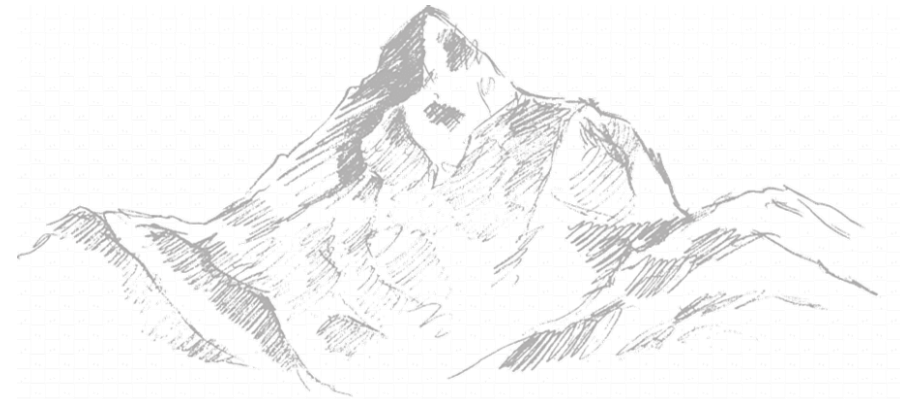
Physically: filtering with operator having a **small variance** on the current state



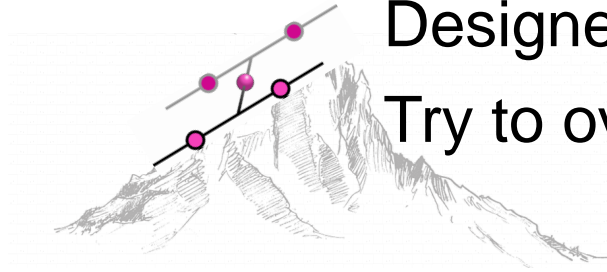
Concluding remarks



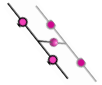
Designed a **flexible algorithm** targeting local dynamics.
Try to overcome the entanglement barrier.



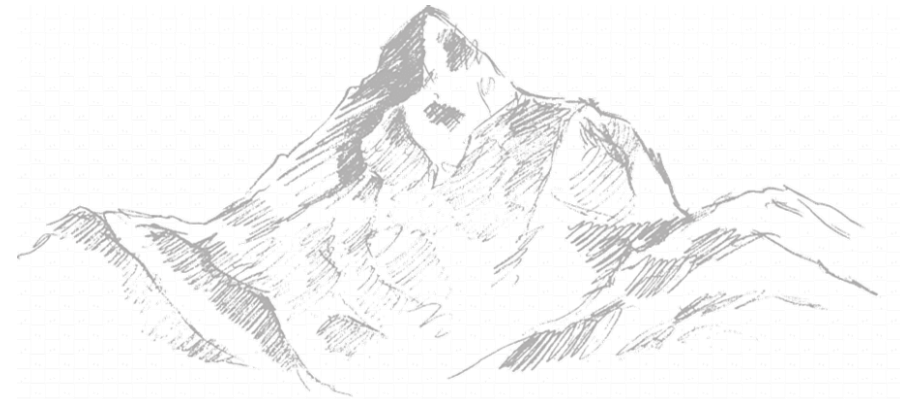
Concluding remarks



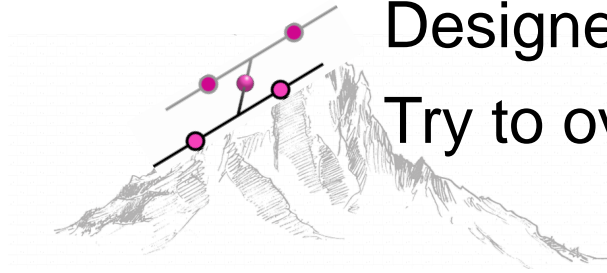
Designed a **flexible algorithm** targeting local dynamics.
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Not all coherence would be necessary to
predict local observables.



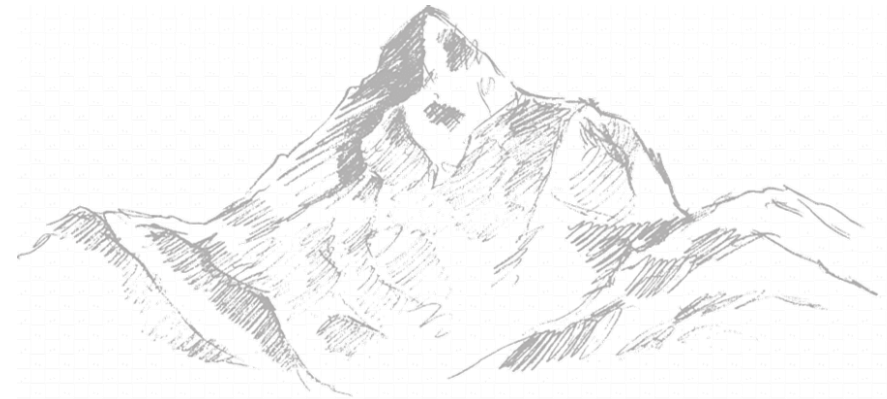
Concluding remarks



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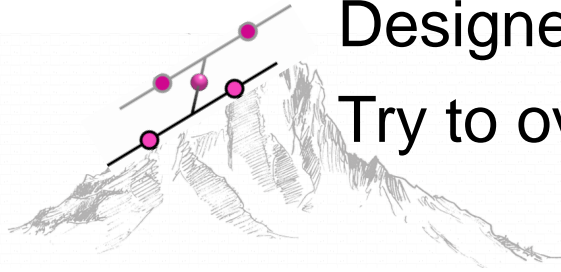
Not all coherence would be necessary to predict local observables.



The use of dissipative coupling can **decrease the operator entanglement** (range of applicability to be determined)



Concluding remarks



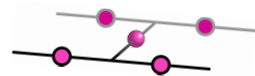
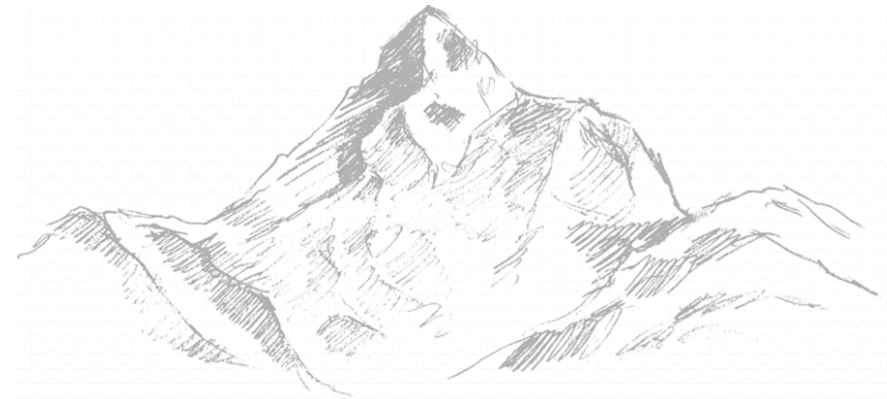
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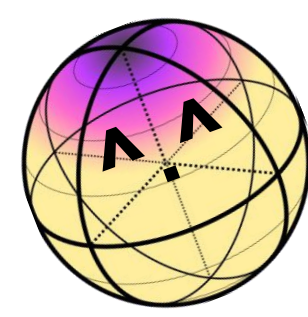
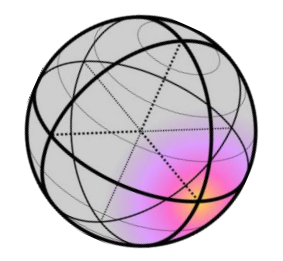
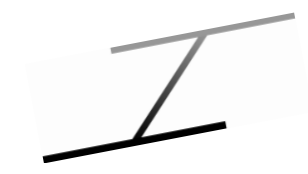
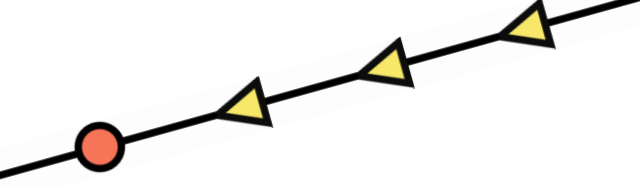


The use of dissipative coupling can **decrease the operator entanglement** (range of applicability to be determined)

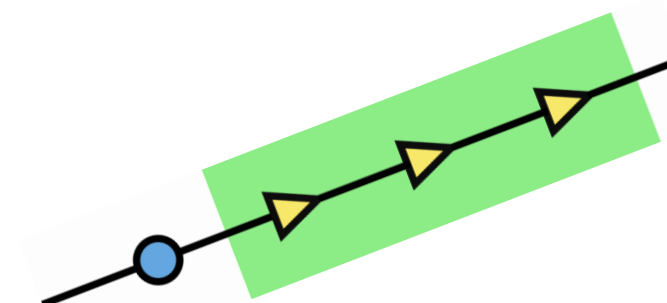
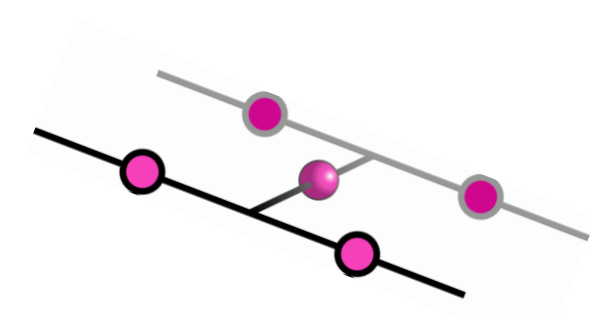
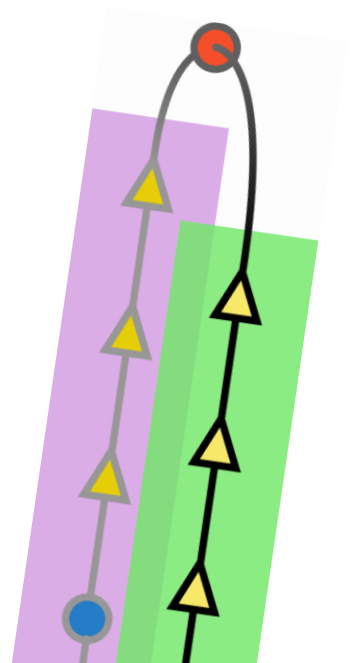


In case of proven efficiency: clear **extensions to non-local** observables (yet to come)



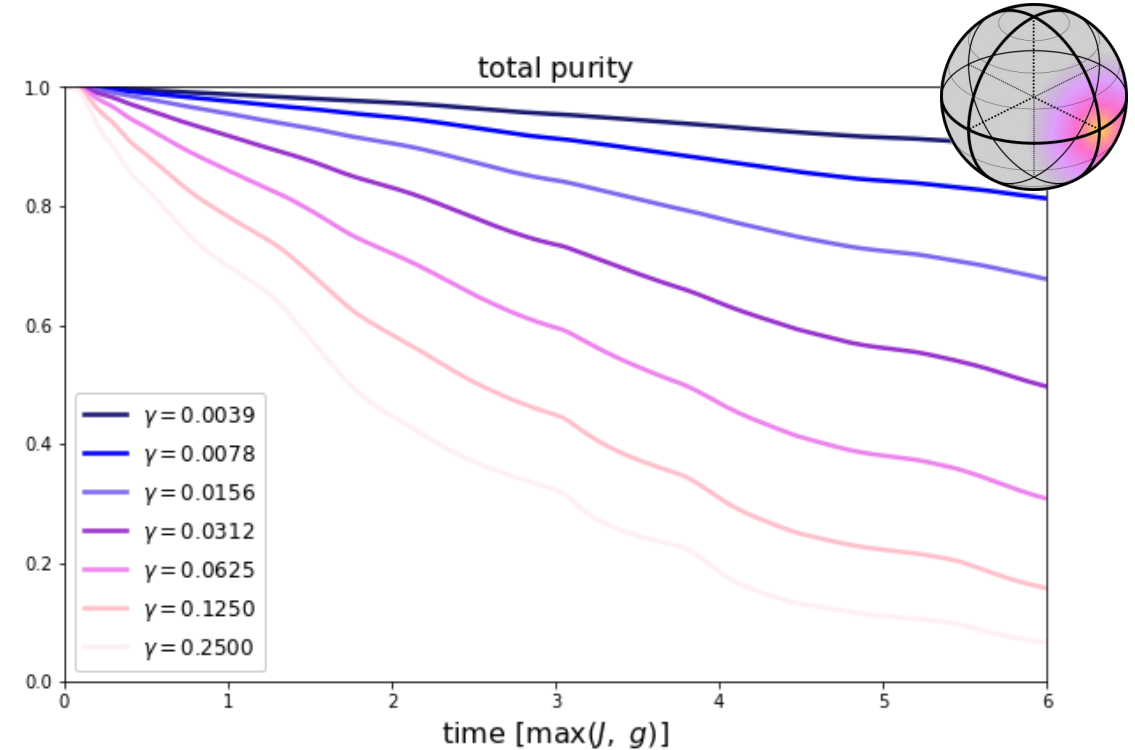
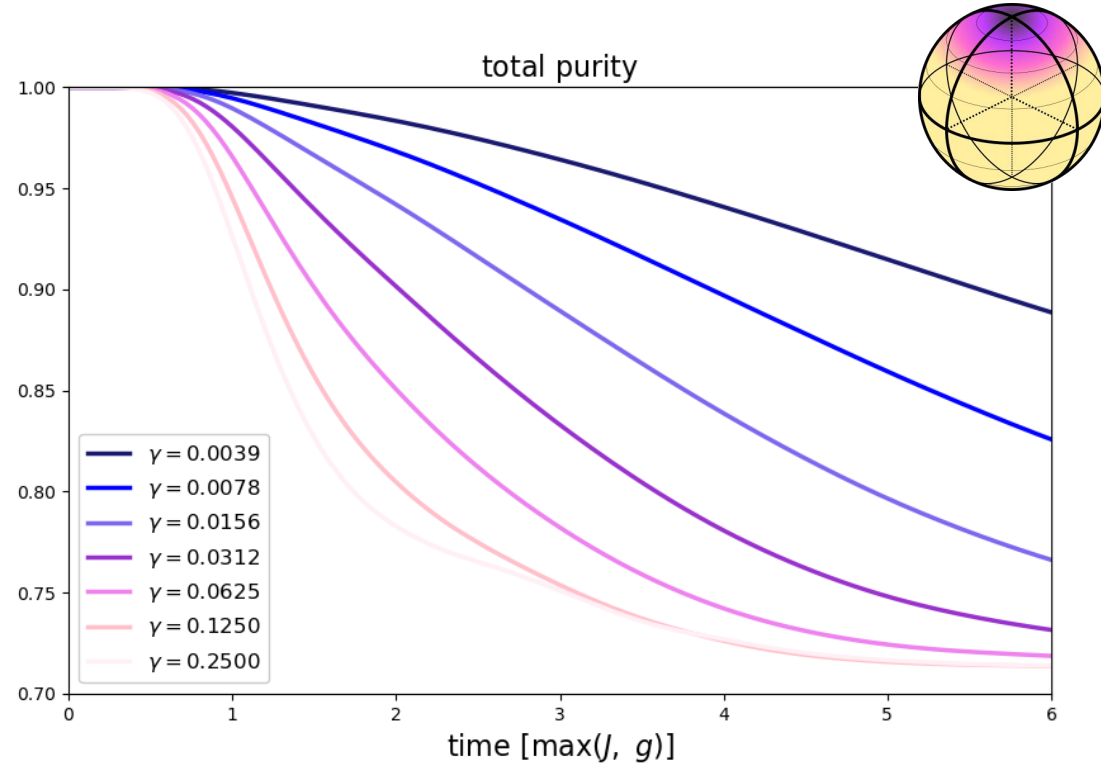


Thank you!



5. Backup: depurification

Dissipation depurifies total state...

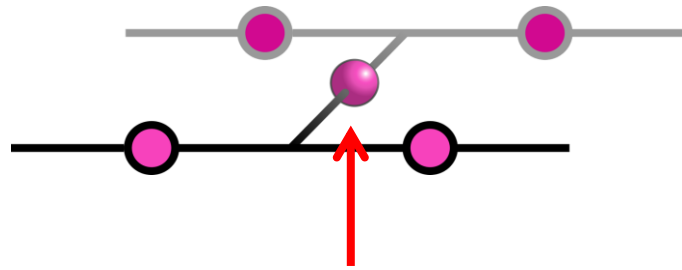


...this hinders the simulability in case of excessive loss of purity.

This bounds dissipation intensity from above (while memory does it from below).

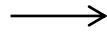
5. Backup: connection to measurement

Reshaping the core copy tensor:

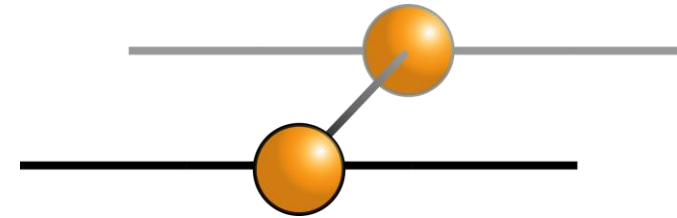


SVD

if weak: could think on
thermalization by
connection to a reservoir



$$\sum_q^{\chi_Q} M_q \{ \cdot \} M_q^\dagger$$



$$\mathbb{I} = \sum_q^{\chi_Q} M_q M_q^\dagger$$

generalized quantum
measurement



5. Backup: fine decoherence


$$\mathbb{M}_{(1)}^{\parallel} = \sum_i \sigma_i^{\parallel}$$

generate extensive set of commuting operators in which we can decohere

$$\mathbb{M}_{(2)}^{\parallel} = \sum_i \sigma_i^{\parallel} \sigma_{i+1}^{\parallel}$$

$$\mathbb{M}_{(3)}^{\parallel} = \sum_i \sigma_{i-1}^{\parallel} \sigma_i^{\parallel} \sigma_{i+1}^{\parallel}$$

new filter with extensive number of rates ~ Lagrange multipliers ~ temperatures

$$e^{-\sum_{i=1}^L \lambda_i \left(\mathbb{M}_{(i)}^{\parallel} - \langle \mathbb{M}_{(i)}^{\parallel}(t) \rangle \right)^2}$$


numerically: allow for **multidimensional extrapolation** $\lambda_i \rightarrow 0$

physically: **generalized thermalization?**

5. Backup: variational formulation

Generic formulation of the algorithm:

(1) find local operators whose variance is minimal at instantaneous state

$$\mathcal{C}(\{\theta_i\}) = \text{tr} \left\{ \rho^\dagger(t) Q^2(\{\theta_i\}) \right\} - \text{tr}^2 \left\{ \rho^\dagger(t) Q(\{\theta_i\}) \right\}$$

written as MPO

(2) filter in the operator eigenbasis

frustration free operator!

generate mixture of similar operators
with the proper local information

$$\mathbb{M}_{(1)}^{\parallel} = \sum_i \sigma_i^{\parallel}$$