

# Bounding the minimum time of a quantum measurement

*Nathan Shettel, Federico Centrone, Luis Pedro Garcia-Pintos*

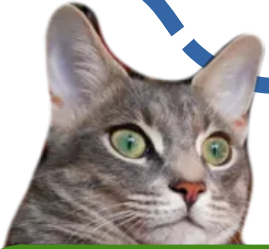


**ICE-8**

**Quantum Information in Spain**

Santiago de Compostela, May 29<sup>th</sup> to June 1<sup>st</sup> 2023

# Quantum mechanics is hard to digest

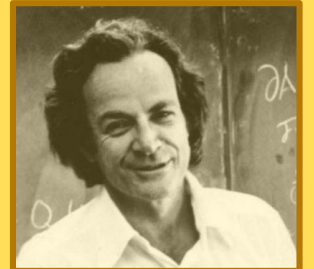


*"If all of this is true, then it means the end of physics"*  
– Albert Einstein

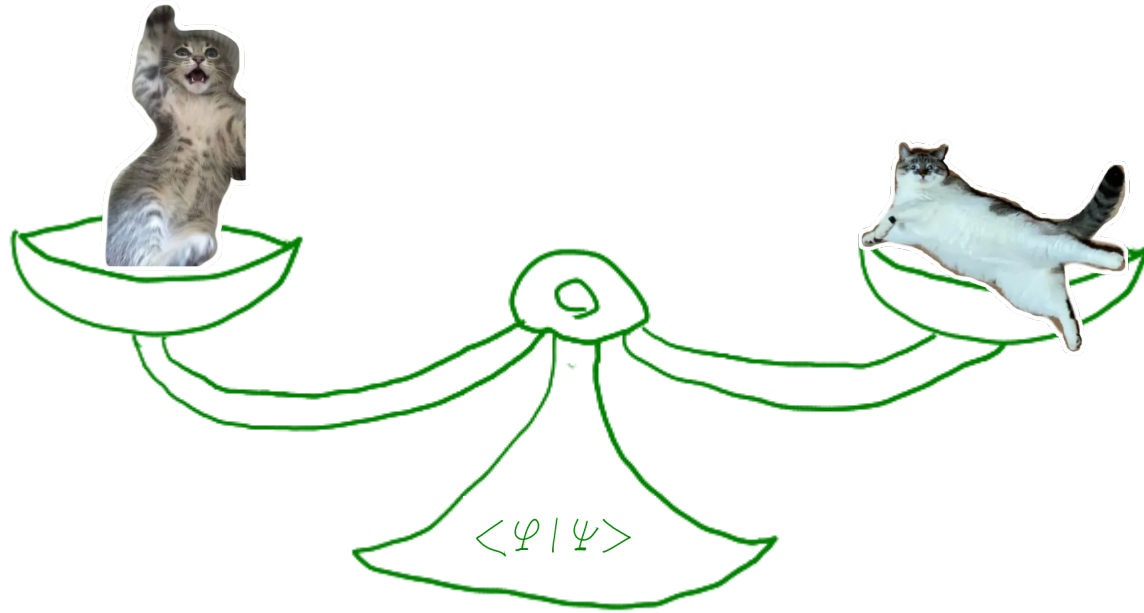


*"I don't like it, and I'm sorry I ever had anything to do with it."*  
- Erwin Schrödinger

*"I think I can safely say that nobody understands quantum mechanics."*  
- Richard Feynman

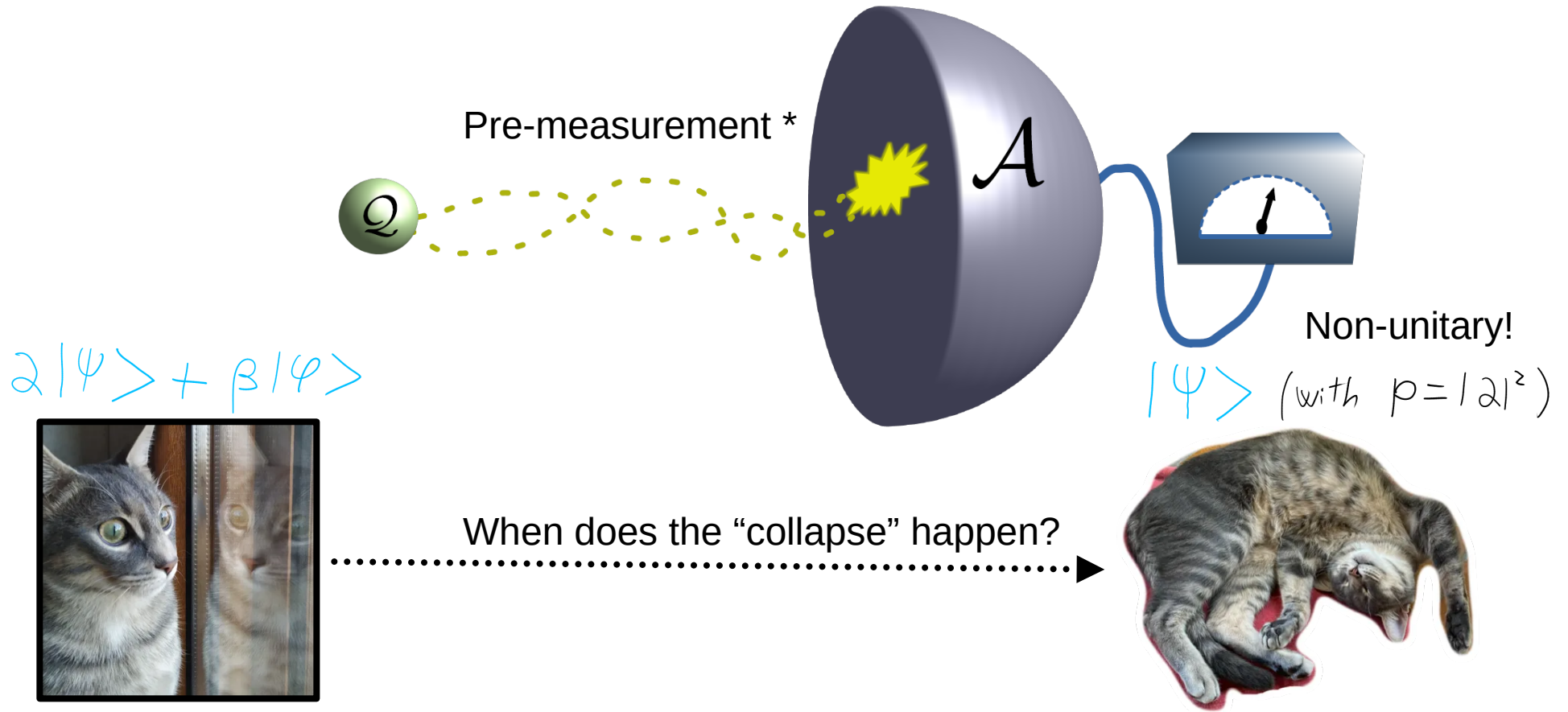


# The measurement problem



*Can measurement itself be described by a physical process?*

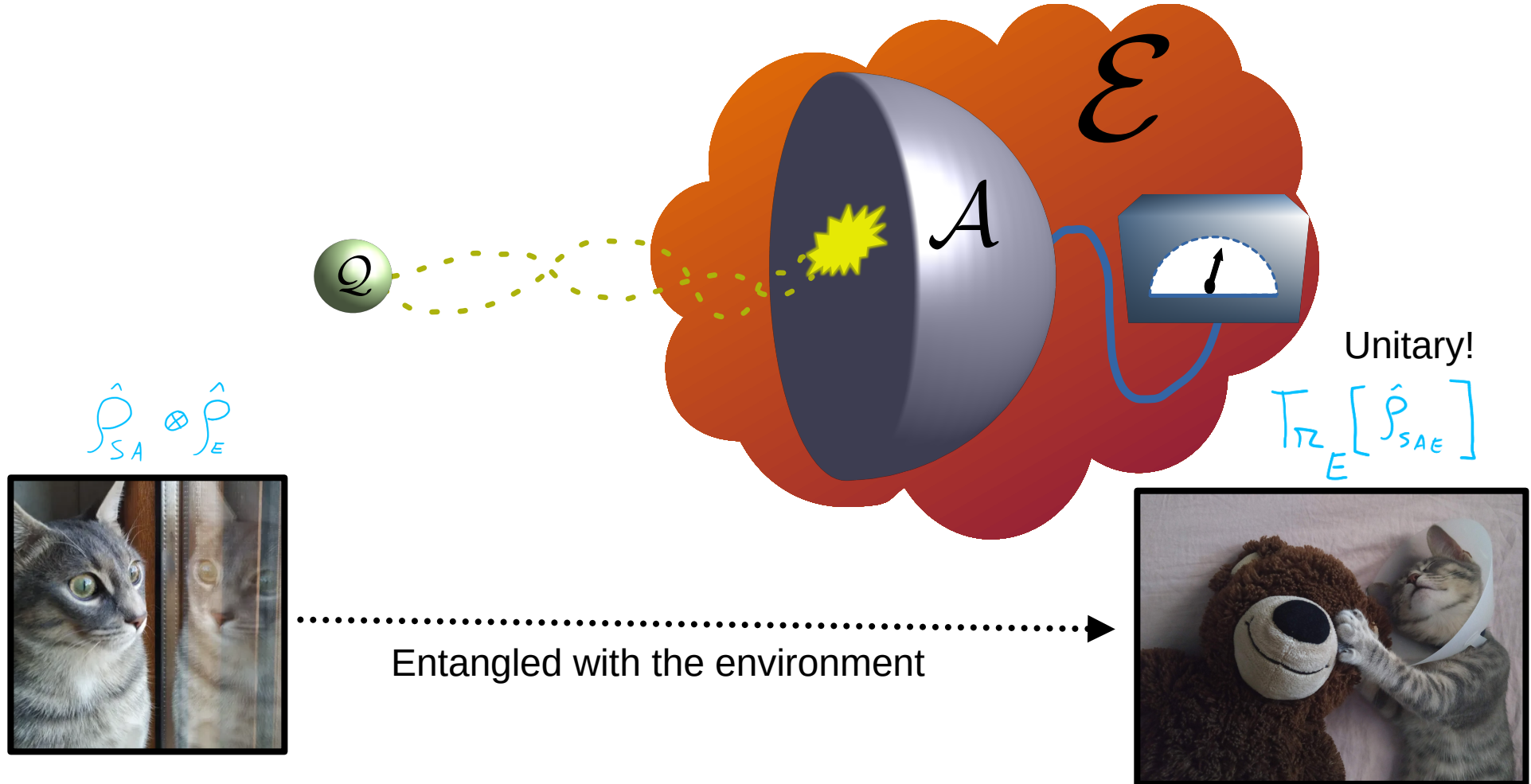
# Collapse of the wavefunction



\* Strasberg, Philipp, Kavan Modi, and Michalis Skotiniotis. "How long does it take to implement a projective measurement?." European Journal of Physics 43.3 (2022): 035404.



# Decoherence



**Heisenberg's uncertainty relations :**

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$

Conjugate observables  
(well understood!)

$$\Delta E \Delta t \geq \frac{\hbar}{2}$$

Energy-time  
(wtf?)



Say my name...

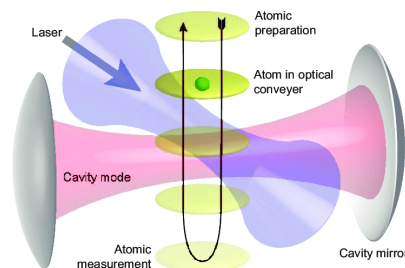
# Time scale of the evolution: **Quantum Speed Limits**

## Mandelstam and Tamm

$$\Delta_{\rho} H \Delta_{\rho} M \geq \frac{|\langle [\hat{H}, \hat{M}] \rangle_{\rho}|}{2} = \frac{\hbar}{2} \left| \frac{\partial}{\partial t} \langle \hat{M} \rangle_{\rho} \right|$$
$$\Rightarrow \Delta H \Delta \tau \geq \frac{\hbar}{2}$$



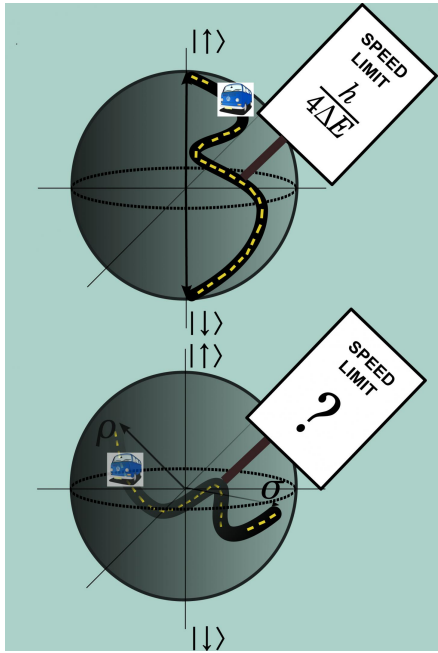
QSL have been used to explore the limits of computation and complexity. [1]



Non-Markovian dynamics can speed up quantum processes :  
Verified in a cavity QED experiment. [2]

[1] Lloyd, Seth. "Ultimate physical limits to computation." Nature 406.6799 (2000): 1047-1054.

[2] Deffner, Sebastian, and Eric Lutz. "Quantum speed limit for non-Markovian dynamics." Physical review letters 111.1 (2013): 010402.



### Geometric approach:

The shortest path between distinguishable quantum states.

$$v(t) := \lim_{\delta t \rightarrow 0} \frac{D_B(\rho(t + \delta t), \rho(t))}{\delta t}$$

$$= \frac{1}{2\hbar} \sqrt{I_Q(\rho(t))}$$

### Information-Time Uncertainty relation:

Fast observable dynamics requires large fluctuations.

$$\Delta\tau \sqrt{I_Q} \geq 1$$



# Timescale bound

How much time  $\tau$  does it take for a car to run a distance  $d$ ?



Need to make some assumptions and solve complex differential equation to know the dynamics

# Timescale bound

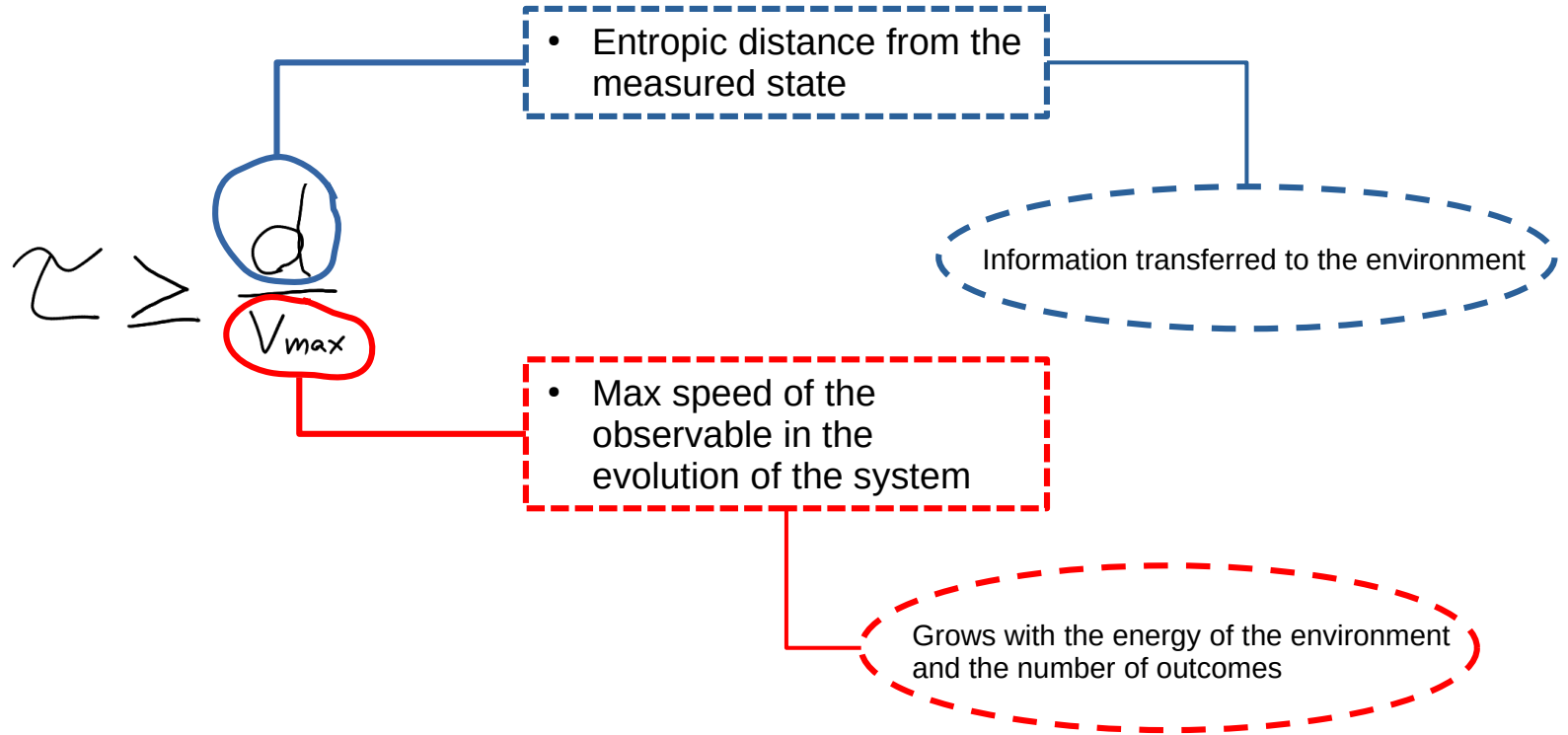
$\tau$  is at least  $d/v_{\max}$





# Timescale bound

This bound can be generalized from cars to quantum states and observables

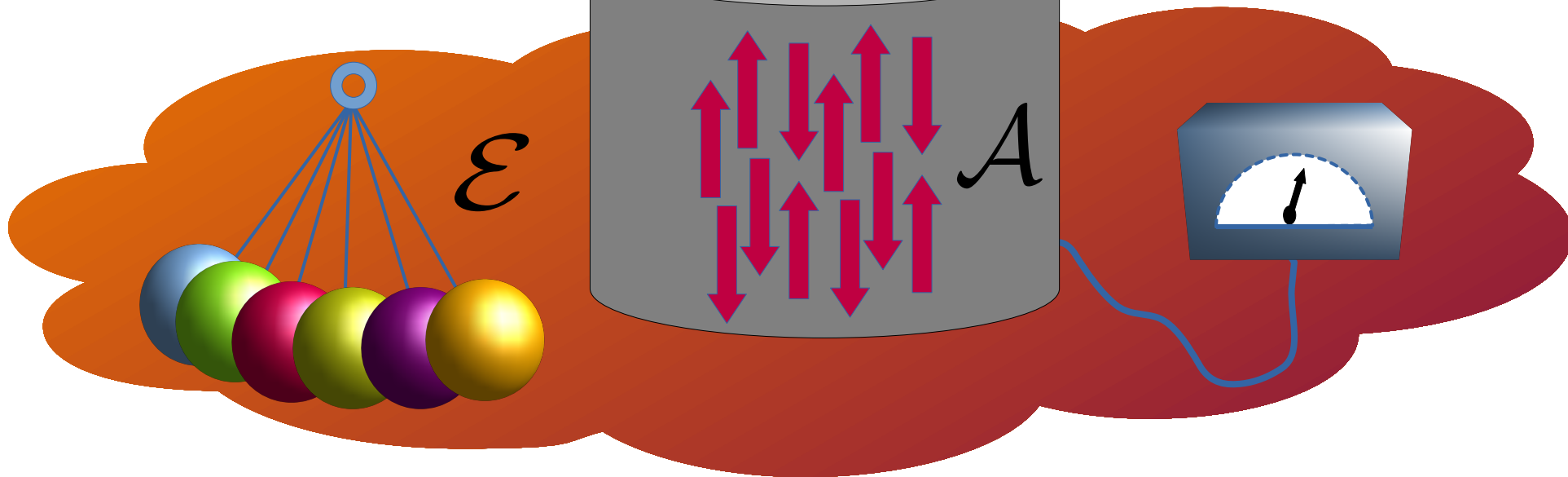
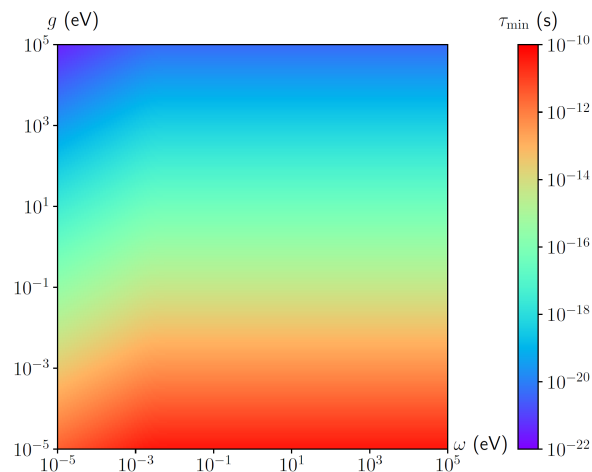
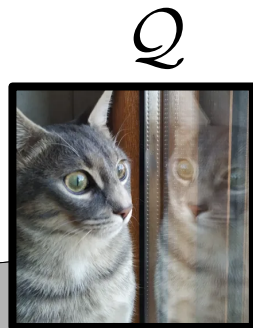


$$\tau = O\left(\frac{1}{N}\right)$$

State of the art quantum clocks  
~attosecond ( $10^{-18}$ s) range

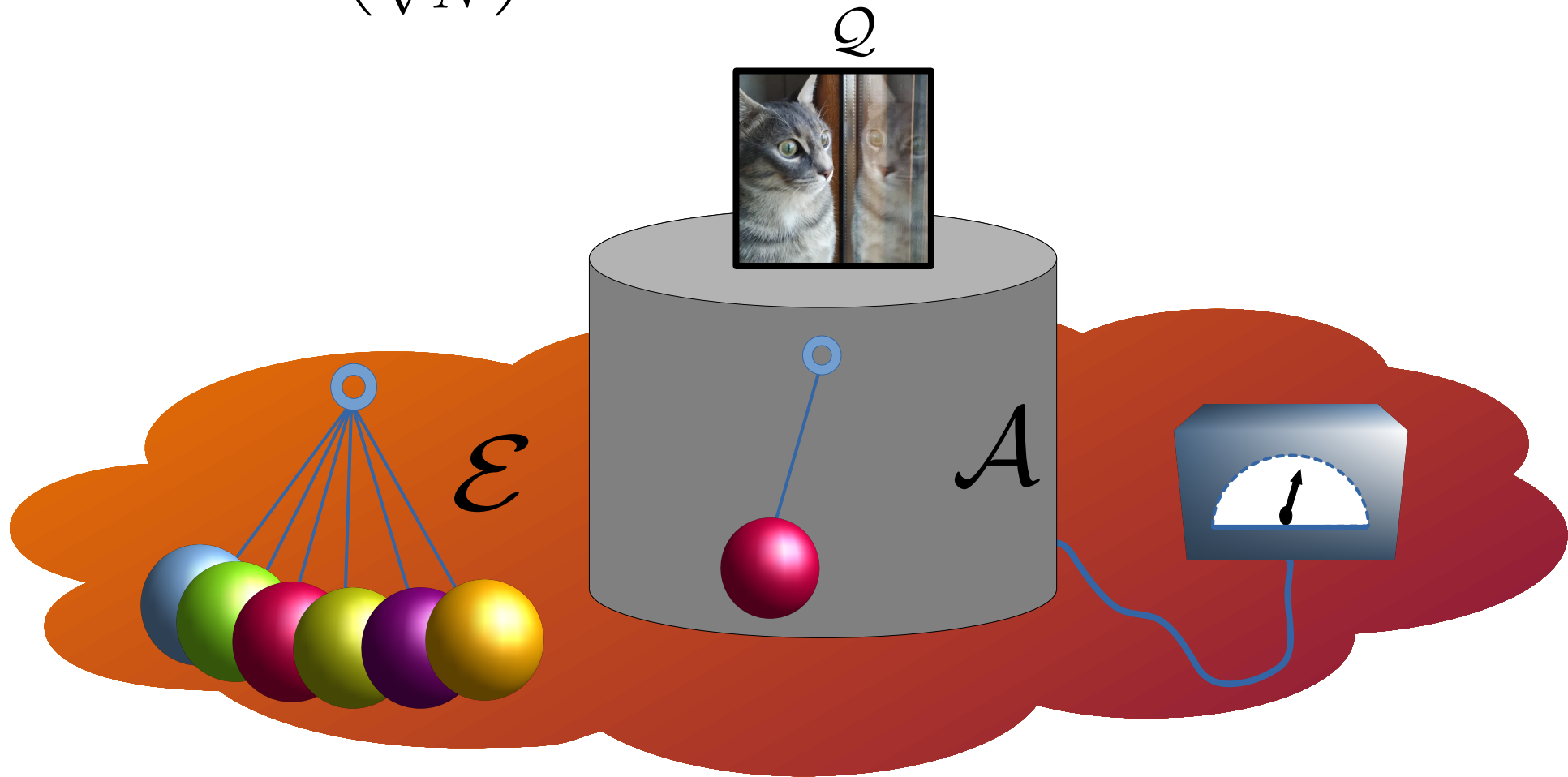
T. Gaumnitz, et al. *Streaking of 43-attosecond soft-X-ray pulses generated by a passively CEP-stable mid-infrared driver* (2017). *Optics express*, 25(22), 27506-27518.

## Example: spin-boson



**Example: boson-boson**

$$\tau = O\left(\frac{1}{\sqrt{N}}\right)$$



## To sum up:

- The decoherence model allows to compute a **non-vanishing** measurement time.
- Quantum speed limits can be used to find a bound in very **general scenarios**.
- Our bound does not require to solve **complex differential equations**.
- Applicable to **more than just measurement**.
- The bound is loose but **experimentally testable**.

## Outlook:

- Tightening the bound.
- Testbed for non-linear Schrodinger equation.
- Influence from exotic environments.
- Physics of quantum information (Bremmerman-Bekenstein bound).

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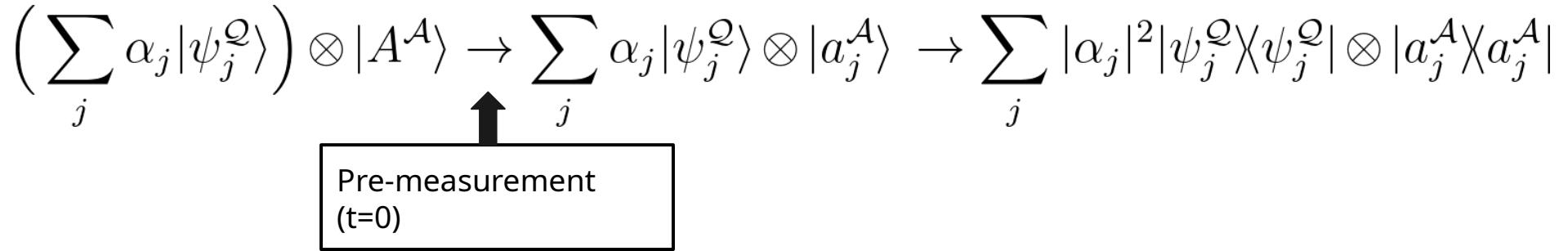


$$\langle \psi | \psi \rangle$$



## Decoherence measurement

$$\left( \sum_j \alpha_j |\psi_j^Q\rangle \right) \otimes |A^A\rangle \rightarrow \sum_j \alpha_j |\psi_j^Q\rangle \otimes |a_j^A\rangle \rightarrow \sum_j |\alpha_j|^2 |\psi_j^Q\rangle\langle\psi_j^Q| \otimes |a_j^A\rangle\langle a_j^A|$$



Pre-measurement  
(t=0)

# Time scale bound

$$\hbar \left| \frac{dS}{dt} \right| \leq 2\Delta S \Delta H_{\text{int}}^* \longrightarrow \tau \geq \frac{S(\rho^{\mathcal{Q}\mathcal{A}} || \rho_{\mathcal{M}}^{\mathcal{Q}\mathcal{A}})|_{t=0}}{\max_t |\partial_t S(\rho^{\mathcal{Q}\mathcal{A}} || \rho_{\mathcal{M}}^{\mathcal{Q}\mathcal{A}})|} \geq \frac{\hbar}{2f_A} \frac{S(\rho^{\mathcal{Q}\mathcal{A}} || \rho_{\mathcal{M}}^{\mathcal{Q}\mathcal{A}})|_{t=0}}{\max_t \Delta H_{\text{int}}}$$

Kullback–Leibler divergence

$$\Delta S \leq \sqrt{\frac{1}{4} \ln^2(A-1) + 1} = f_A.$$

Approximately time independent in the strong coupling regime

# Example: Spin-Boson

$$|\psi^{\mathcal{QA}}\rangle = x|0\rangle|\downarrow\rangle^{\otimes N} + y|1\rangle|\uparrow\rangle^{\otimes N}$$

$$\rho^{\mathcal{E}} = \bigotimes_k \frac{1}{Z_k} e^{-\beta\omega_k a_k^\dagger a_k}$$

- Q: 2-level system
- A: N spins
- E: multi-mode thermalized bosonic states (harmonic oscillators)

Continuum of modes:

$$\sum_k |g_k|^2 \rightarrow \int_0^\infty J(\omega) d\omega$$

Interaction couples spins to E-modes

$$H_{\text{int}}^{\mathcal{AE}} = \sum_{i=1}^N \sum_k \sigma_Z^{(i)} \otimes g_k (a_k + a_k^\dagger)$$

Time scales inversely with N and coupling strength

$$\tau \geq \frac{\hbar \ln 2}{2N} \left( \int_0^\infty J(\omega) \coth(\beta\omega/2) d\omega \right)^{-1/2}$$

## Spin-Boson: Analytic expression

Born-Markov approximation:

$$\rho^{\mathcal{QA}} = \begin{pmatrix} |x|^2 & xy^* e^{-\Gamma} \\ x^* y e^{-\Gamma} & |y|^2 \end{pmatrix}$$

$$\Gamma = 4N \int_0^\infty \frac{J(\omega)}{\omega^2} (1 - \cos(\omega t/\hbar)) \coth(\beta\omega/2) d\omega$$

**Solution for relative entropy to equal**

**Our bound**

$$t = \frac{\hbar \sqrt{\ln(1/\varepsilon)}}{2\sqrt{N}} \left( \int_0^\infty J(\omega) \coth(\beta\omega/2) d\omega \right)^{-1/2} \quad \tau \geq \frac{\hbar \ln 2}{2N} \left( \int_0^\infty J(\omega) \coth(\beta\omega/2) d\omega \right)^{-1/2}$$

## Example: Boson-Boson

$$|\psi^{QA}\rangle = x|0\rangle|\alpha\rangle + y|1\rangle|-\alpha\rangle$$

- Q: 2-level system
- A: Cat state
- E: multi-mode thermalized bosonic states (harmonic oscillators)

$$\rho^E = \bigotimes_k \frac{1}{Z_k} e^{-\beta\omega_k a_k^\dagger a_k}$$

Interaction couples A-modes to E-modes

$$H_{\text{int}}^{AE} = \sum_k g_k (a_k b^\dagger + a_k^\dagger b)$$

**Our bound:**

Scales inversely with  $|\alpha|$

**Experimental paper [1]:**

Scales inversely with  $|\alpha|^2$

Due to different Hamiltonian

## Tightening the bound

$$\tau \geq \frac{S(\rho^{\mathcal{Q}\mathcal{A}} || \rho_{\mathcal{M}}^{\mathcal{Q}\mathcal{A}})|_{t=0}}{\max_t |\partial_t S(\rho^{\mathcal{Q}\mathcal{A}} || \rho_{\mathcal{M}}^{\mathcal{Q}\mathcal{A}})|} \geq \frac{\hbar}{2f_A} \frac{S(\rho^{\mathcal{Q}\mathcal{A}} || \rho_{\mathcal{M}}^{\mathcal{Q}\mathcal{A}})|_{t=0}}{\max_t \Delta H_{\text{int}}}$$

Middle inequality 'fixes' N vs  $\sqrt{N}$  discrepancy in spin-boson example

$$|\partial_t S(\rho^{\mathcal{Q}\mathcal{A}} || \rho_{\mathcal{M}}^{\mathcal{Q}\mathcal{A}})| \leq \Delta S \sqrt{I} \leq 2\Delta S \Delta H_{\text{int}}^{\mathcal{A}\mathcal{E}} \quad [1]$$

$$\max(\Delta S) = f_A$$

Cannot be saturated