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NextGenerationEU



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S. Carignano (BSC), C. Ramos (UB)

previous work with: M. Piani (evolutionQ) Niall Robertson
(IBM) Jacopo Surace (ICFO) Erik Tonni (SISSA)

Ongoing effort to simulate the out-of-equilibrium dynamics with TN

Entanglement and coherence out of equilibrium



Plan de Recuperación,
Transformación y Resiliencia

Phys. Rev. B 99, 235115 (2019)

M. Frias-Perez (LT) MC.- Bañuls
arXiv:2306.xxxx



PGC2018-095862-B-C22
PID2021-127968NB-I00
TED2021-130552B-C22

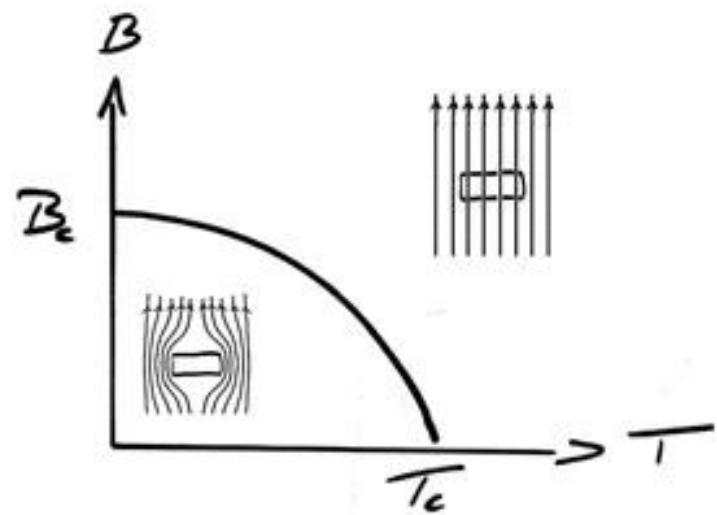
- Out of equilibrium
- Entanglement barrier
- Quasi-particles
- Trading entanglement for mixture
 - Unveiling local equilibration
 - Identifying fast degrees of freedom
 - Relation with decoherence

Classical:

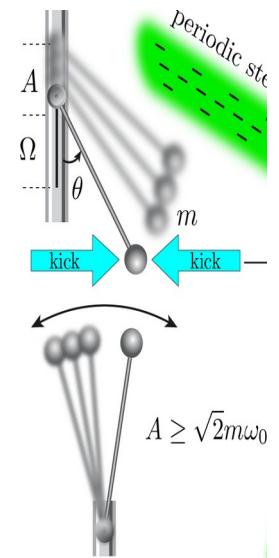
Kapitsa pendulum

Quantum:

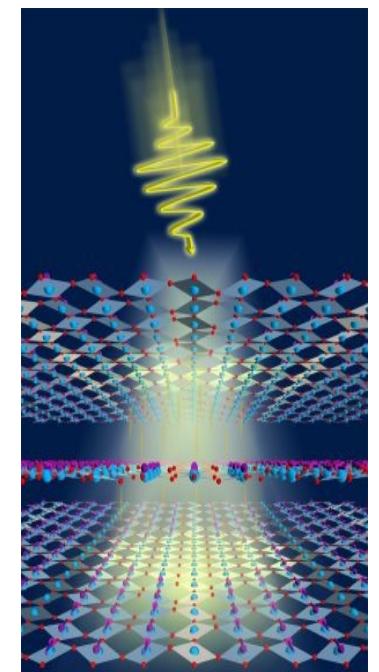
Light induced superconductor



Bukov



Cavallieri

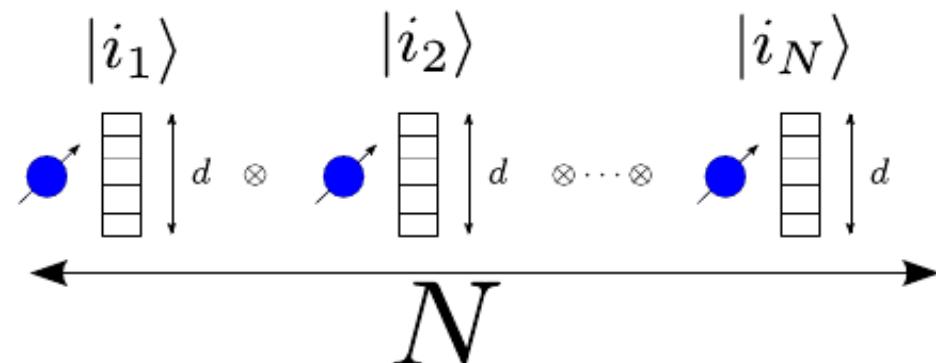


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A state of a many body systems

$$|\psi\rangle = \sum_{i_1 \cdots i_N} c^{i_1 \cdots i_N} |i_1 \cdots i_N\rangle$$

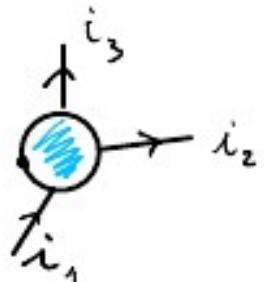
$$\mathcal{H} = \mathcal{H}_1 \otimes \mathcal{H}_2 \cdots \otimes \mathcal{H}_N$$



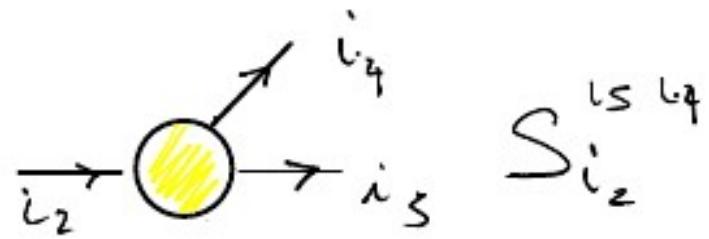
$c^{i_1 \cdots i_N}$ contains d^N parameters

- $2^2 = 4$
- $2^3 = 8$
- $2^4 = 16$
- $2^5 = 32$
- $2^6 = 64$
- ...
- $2^{40} = 1.0995116e+12$

Tensors:

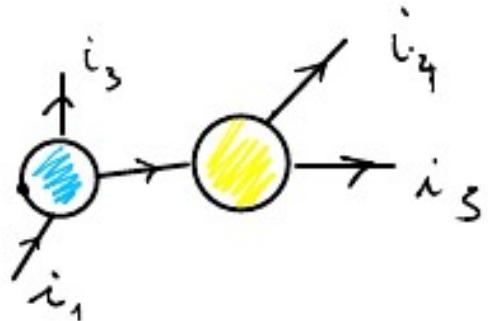


$$T_{i_1}^{i_2 i_3}$$



$$S_{i_2}^{i_5 i_4}$$

Operations among tensors:

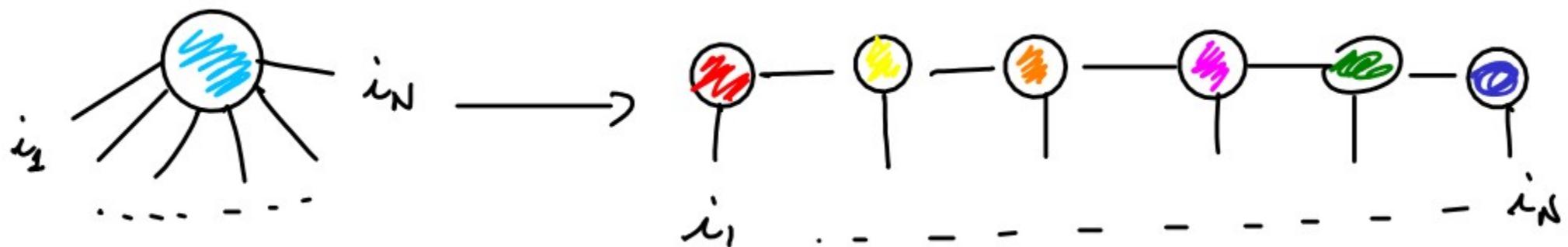


$$(TS)_{i_1}^{i_3 i_4 i_5} = \sum_{i_2} T_{i_1}^{i_3 i_2} S_{i_2}^{i_4 i_5}$$

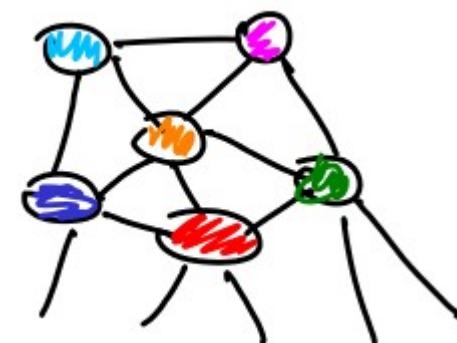
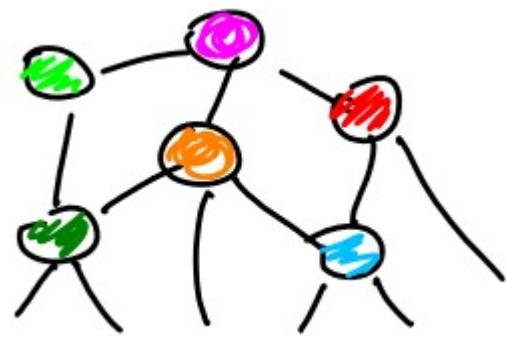
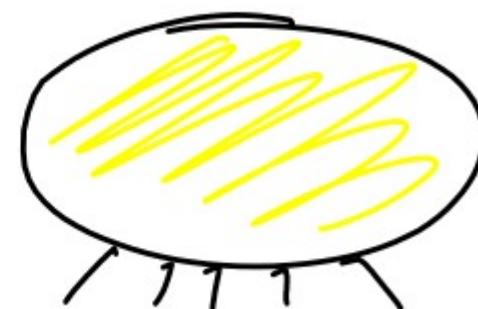
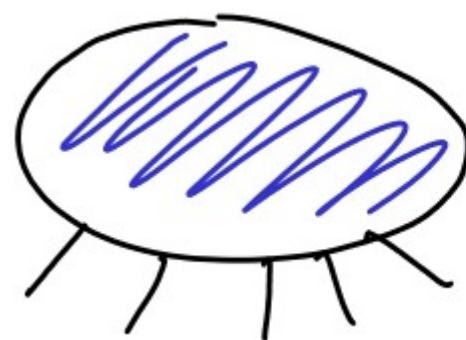
Express exponentially large tensors

$$c^{i_1 \cdots i_N}$$

as contraction of small elementary tensors, **TENSOR NETWORK**

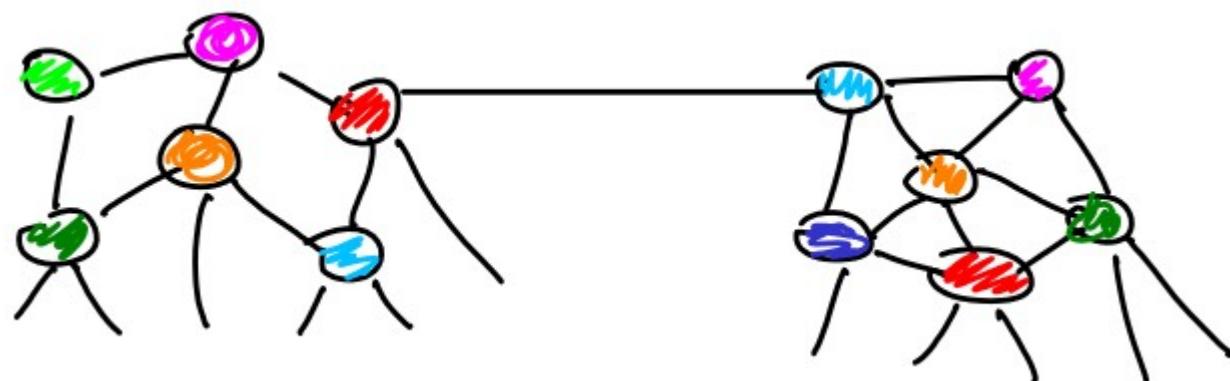
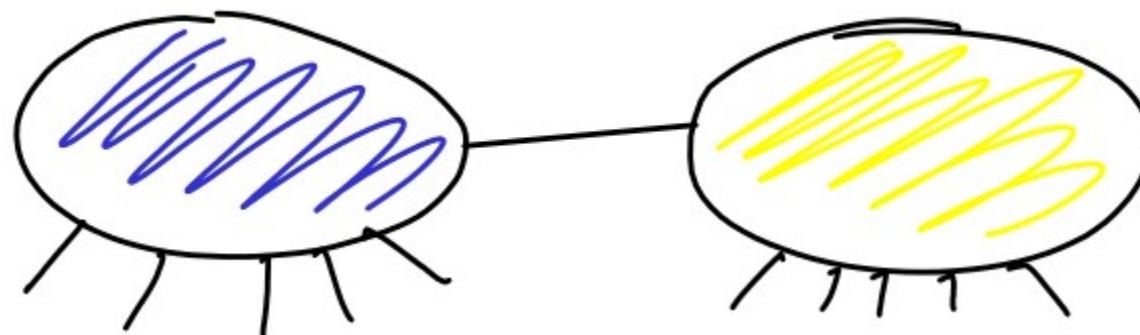


$$|\psi\rangle_A \otimes |\psi\rangle_B$$

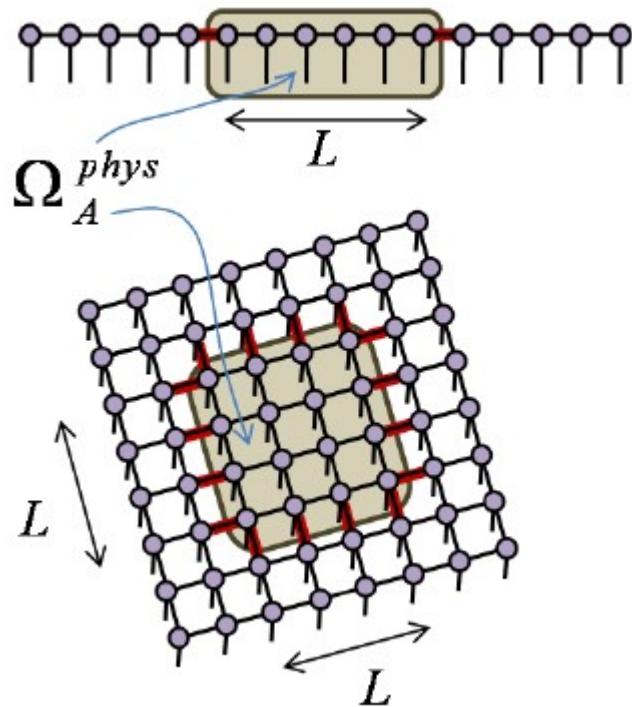


$$|\psi\rangle_{AB} \neq |\psi\rangle_A \otimes |\psi\rangle_B$$

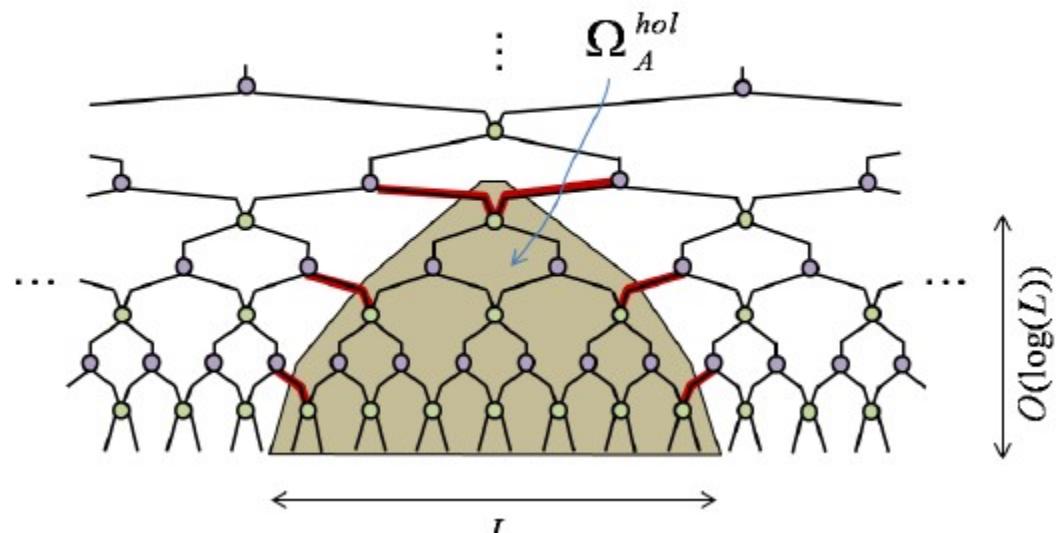
$$|\psi\rangle_{AB} = \sum_i c^i |\psi_i\rangle_A \otimes |\psi_i\rangle_B$$



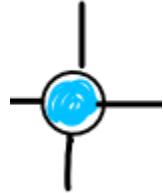
$$S \leq n_{AB} \log(\chi)$$



Entanglement area law states



Entanglement beyond
area law states
(Evenly Vidal 2011)



States and maybe space-time are Tensor Networks

Hilbert space is a “convenient illusion”

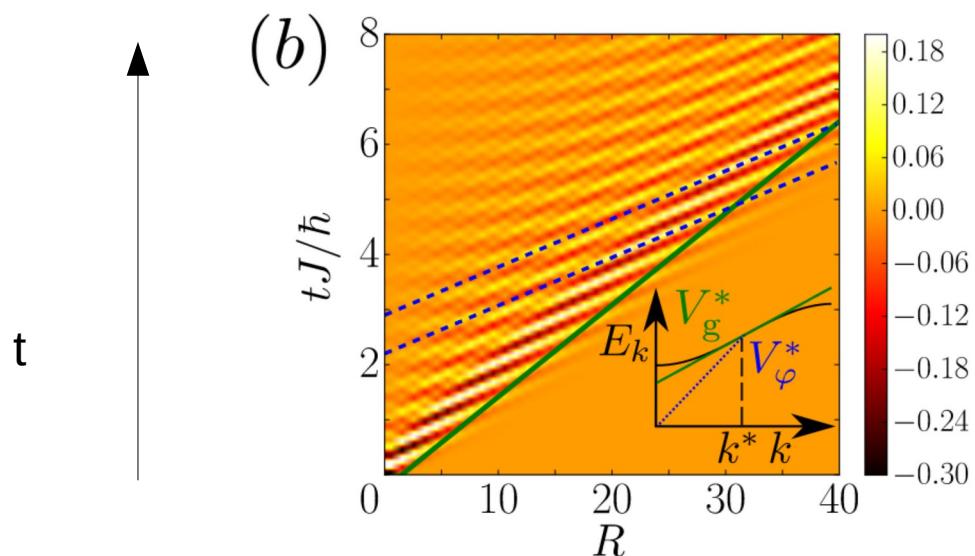
Large many-body systems become

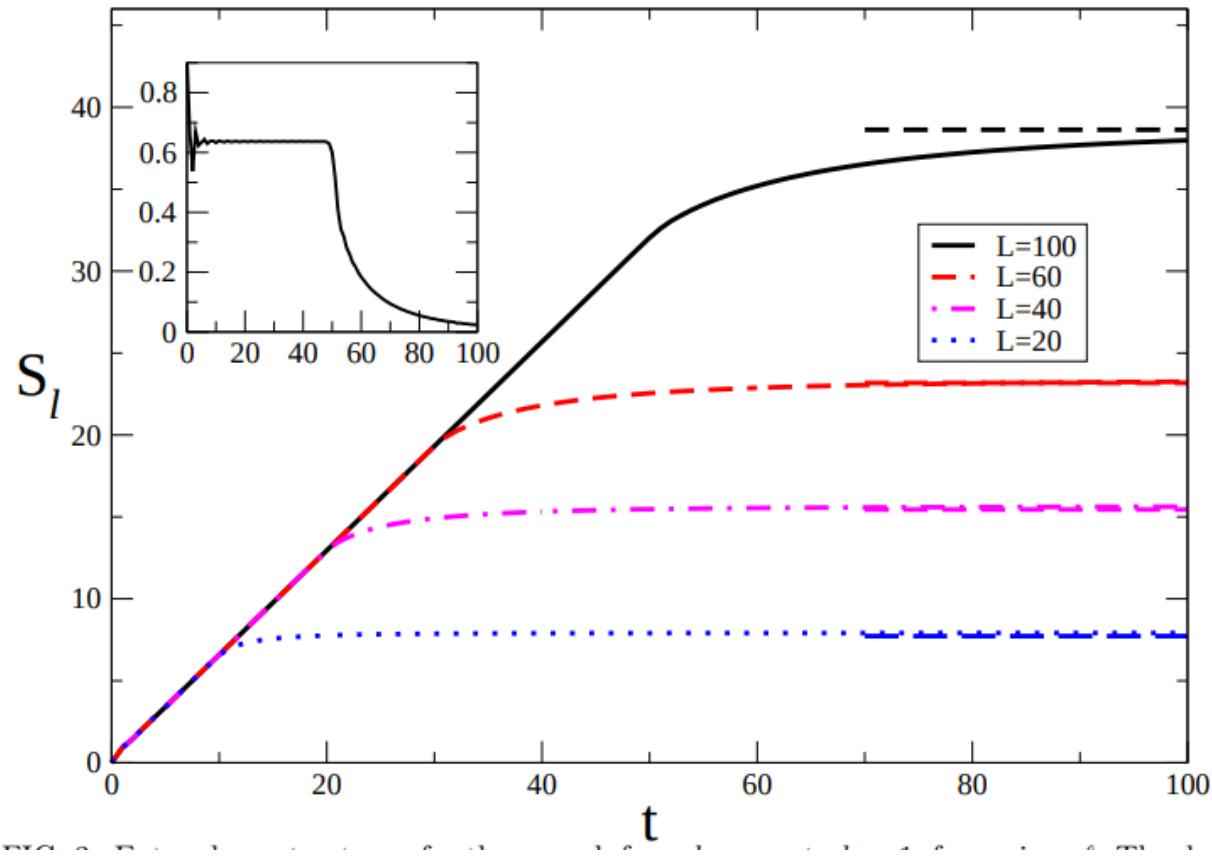
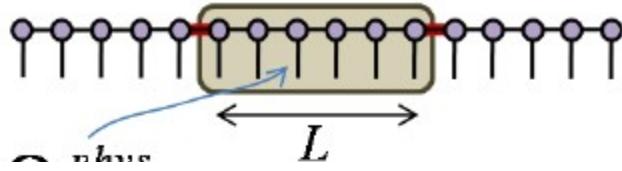
NUMERICALLY TRACTABLE

LT Optimal simulation of quantum dynamics
Nat Phys 2022

ground state $|\psi_0\rangle$ of H_0
out of equilibrium by **quenching**
the Hamiltonian to H_1

$$|\psi(t)\rangle = \exp[-iH_1 t] |\psi_0\rangle$$





from Calabrese Cardy 2005

$$S \leq n_{AB} \log(\chi)$$

$$\chi^{n_{AB}} \geq \exp S \propto \exp(t)$$

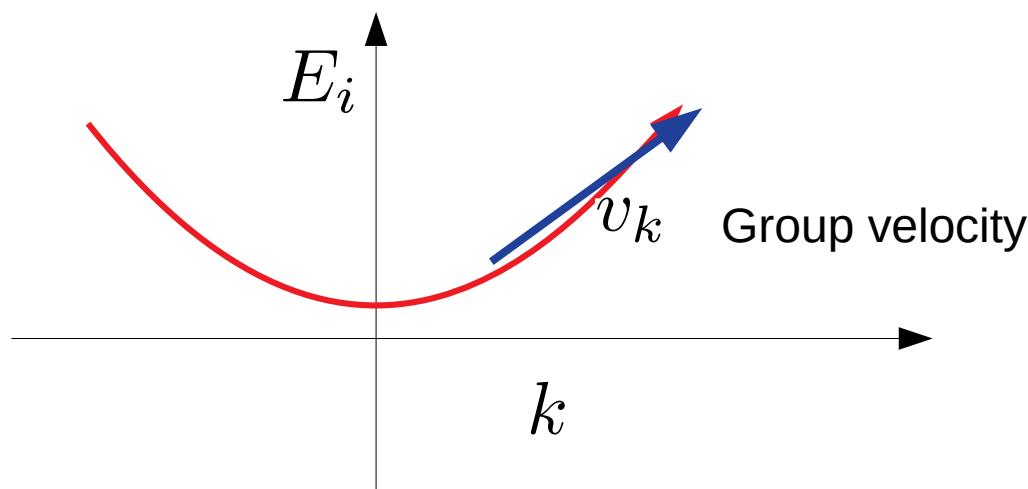
- Out of equilibrium
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- The Hamiltonian is translational invariant,

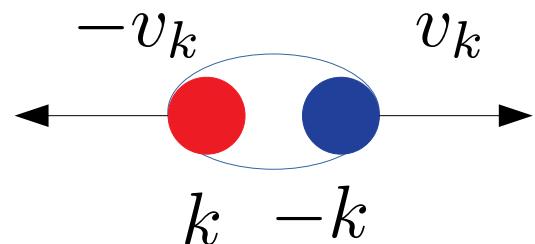
- Momentum is a good quantum number

$$H = \bigoplus_k H_k$$

- We have bands, **quasi-particles** $E_i(k)$

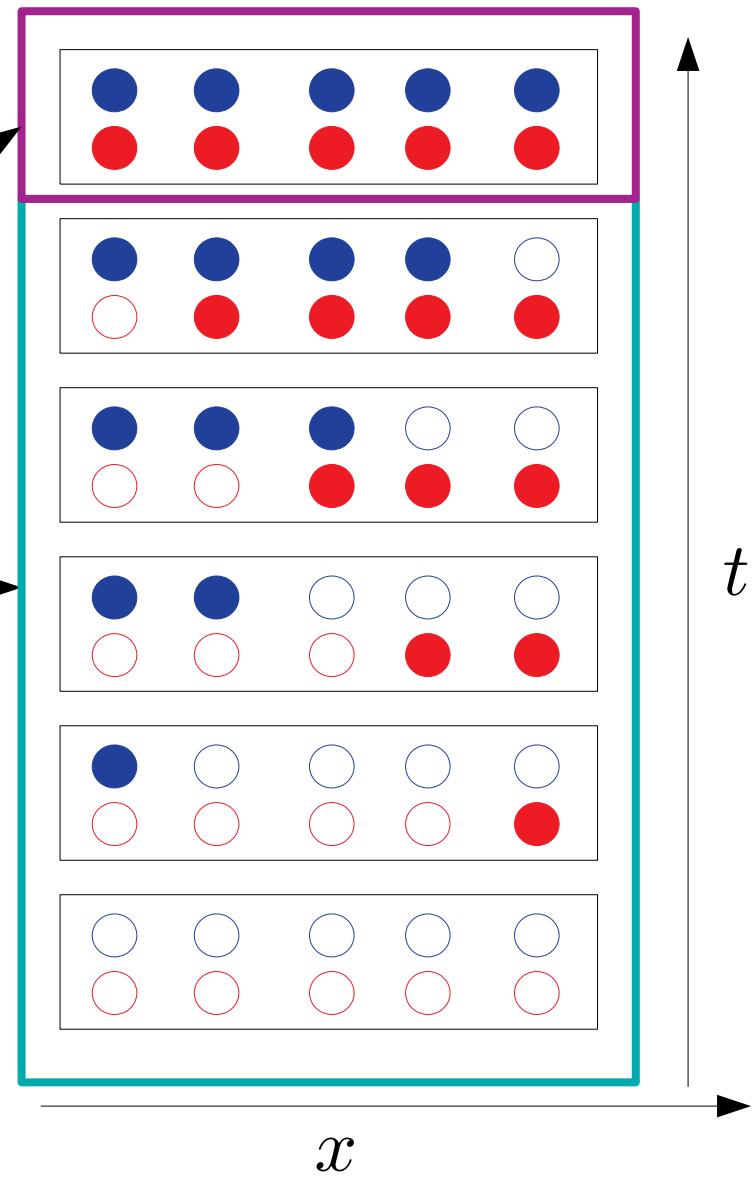
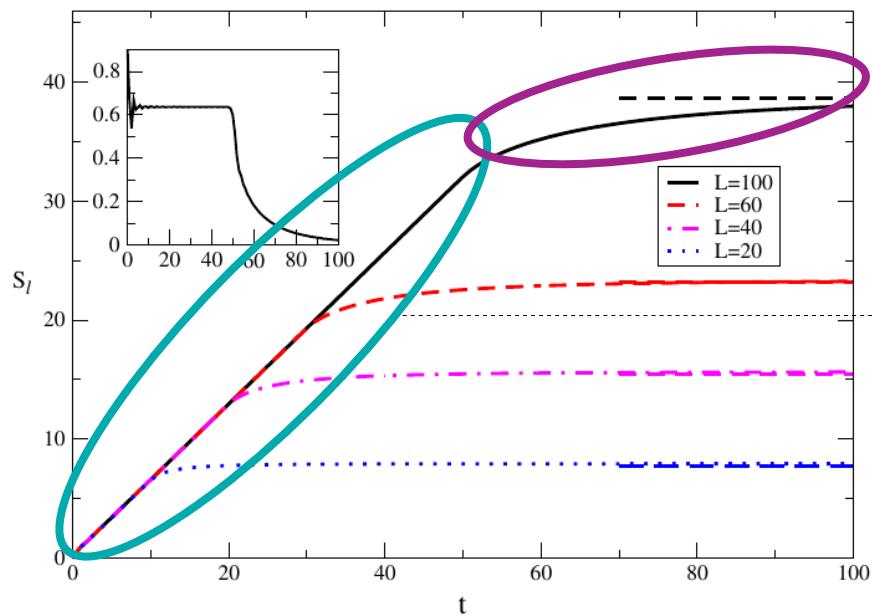


The **extra energy** in the quench is radiated by
entangled **quasi-particle** with opposite
momenta k and velocities v_k



Entanglement after a quench

Calabrese Cardy 05



- Improve TN ansatz
- Fight entanglement with universality and scaling
- Trade entanglement for mixture (decoherence)

- Out of equilibrium
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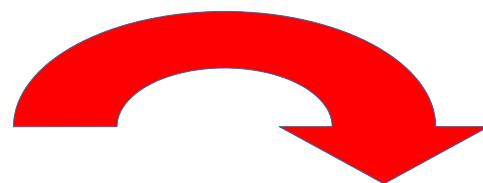
Trading entanglement for mixture, decoherence

- Surace Piani Tagliacozzo PRB2019



The quench protocol

$$H(\theta) = -\sin(\theta) \sum_{i=0}^{N-1} \sigma_i^x \sigma_{i+1}^x - \cos(\theta) \sum_{i=0}^{N-1} \sigma_i^z;$$



$\pi/4$



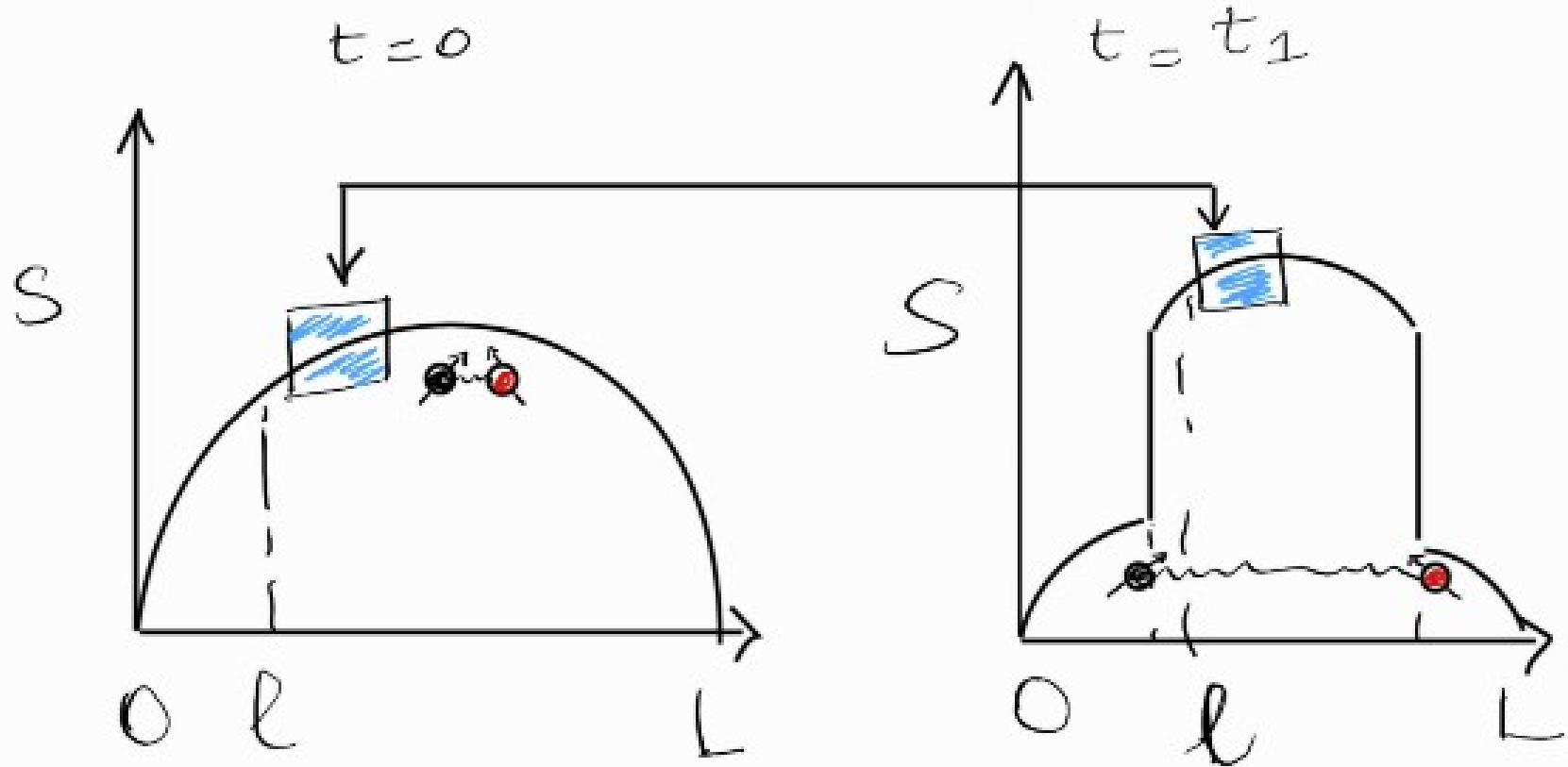
$$\rho = \sum_i p_i |\psi_i\rangle\langle\psi_i|$$

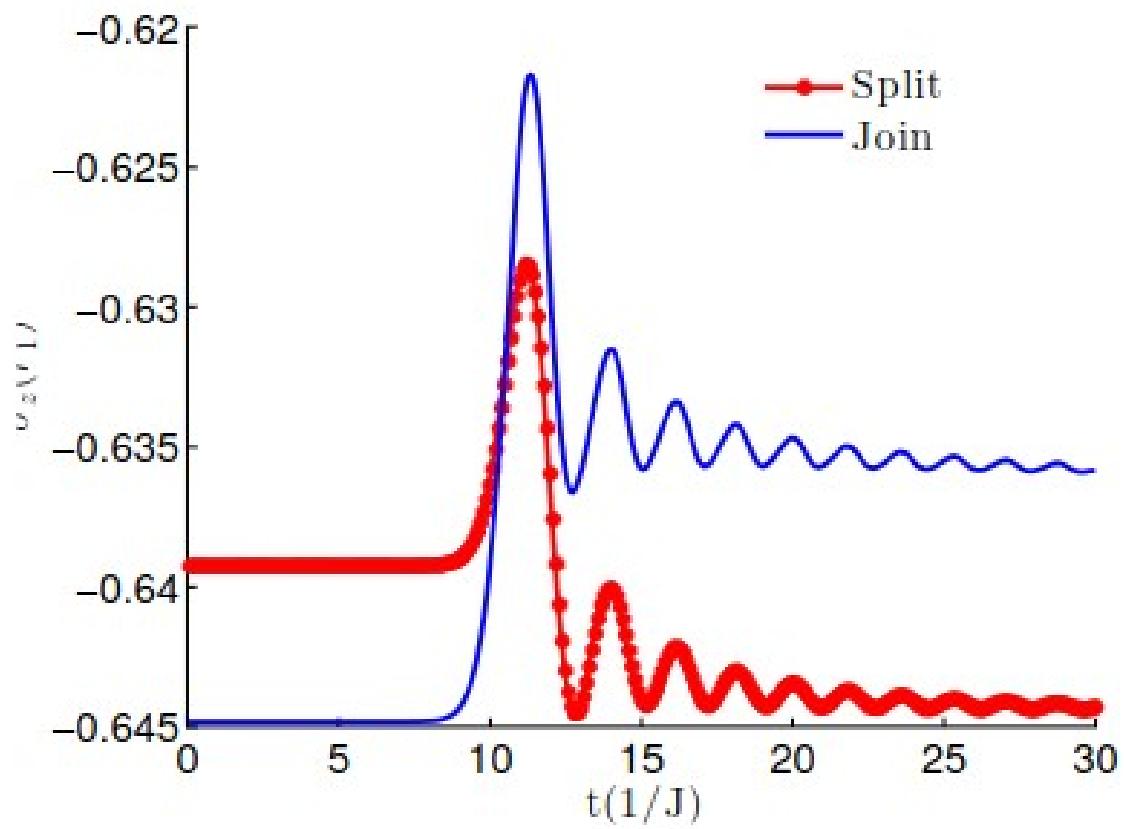
$[\mathcal{O}, \rho] \neq 0 \iff \mathcal{O}$ has quantum fluctuations
if the system is in the state ρ

$$\lambda_{dB}=\frac{h}{mv}$$

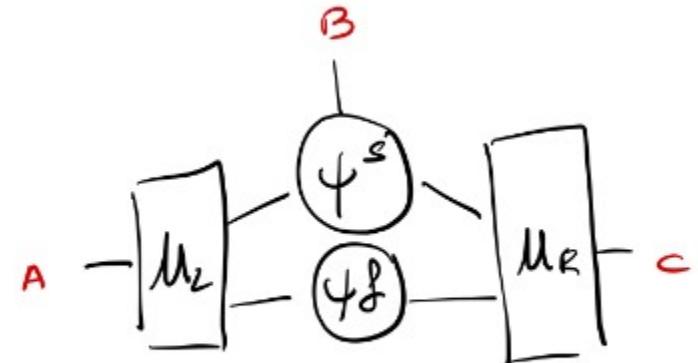
$$mv^2 \sim k_BT$$

$$\lambda_{dB}\sim\frac{h}{\sqrt{mk_BT}}~.$$





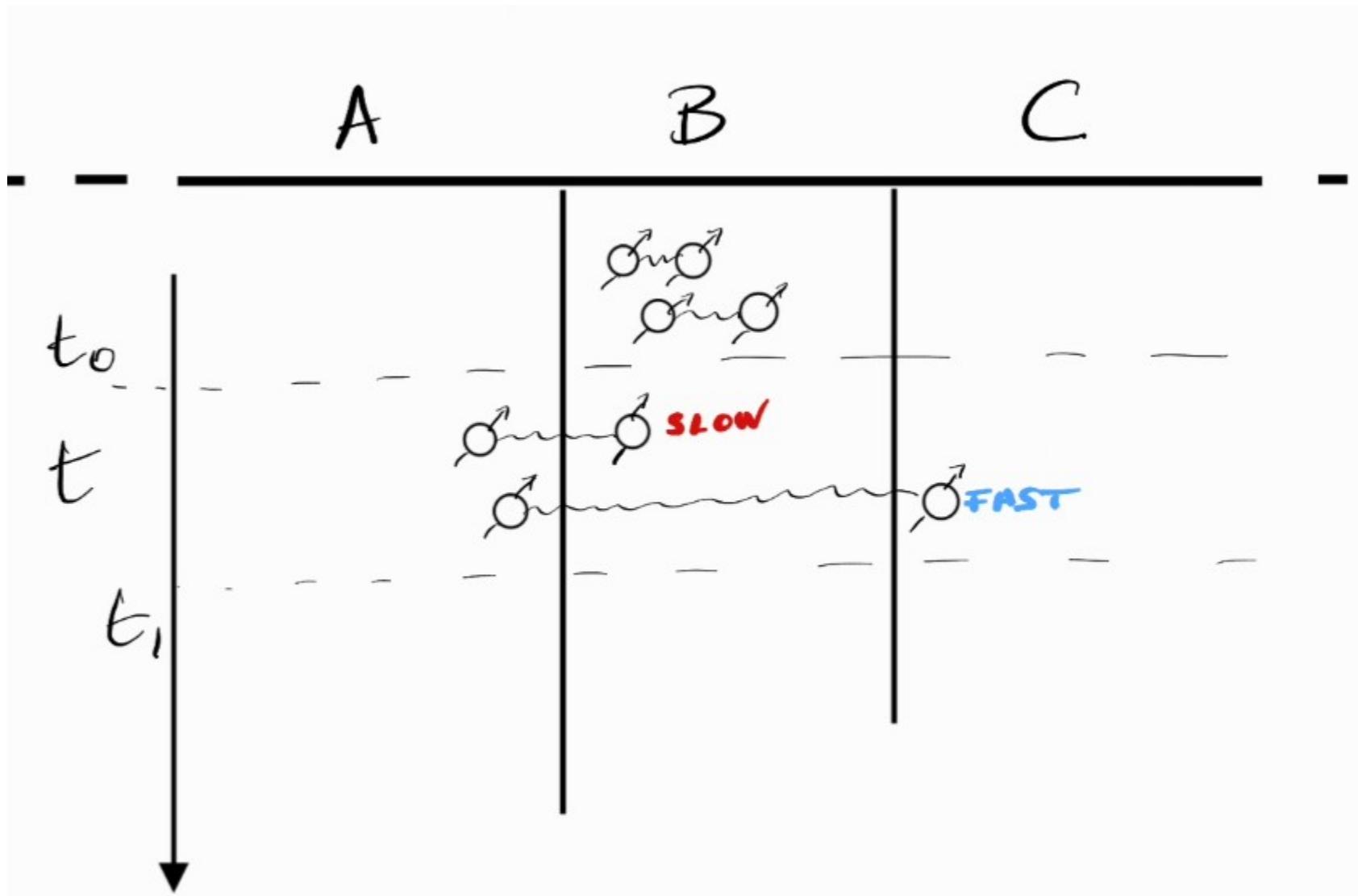
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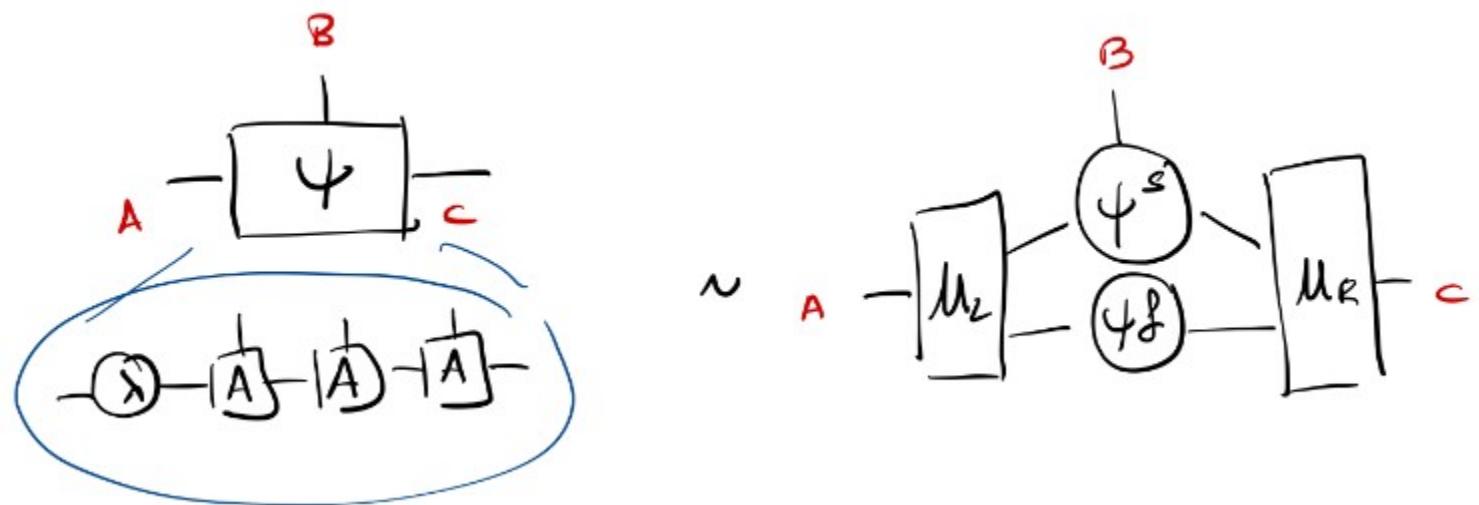


Unveiling local equilibration in TN

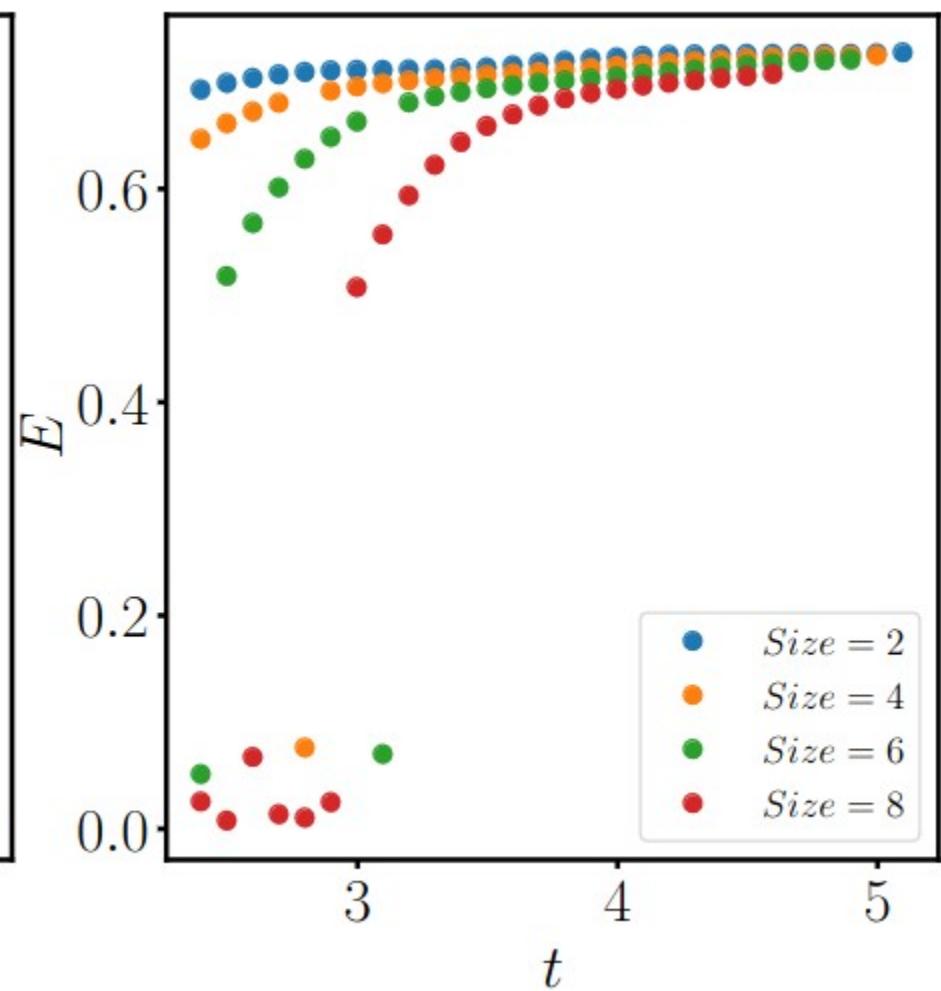
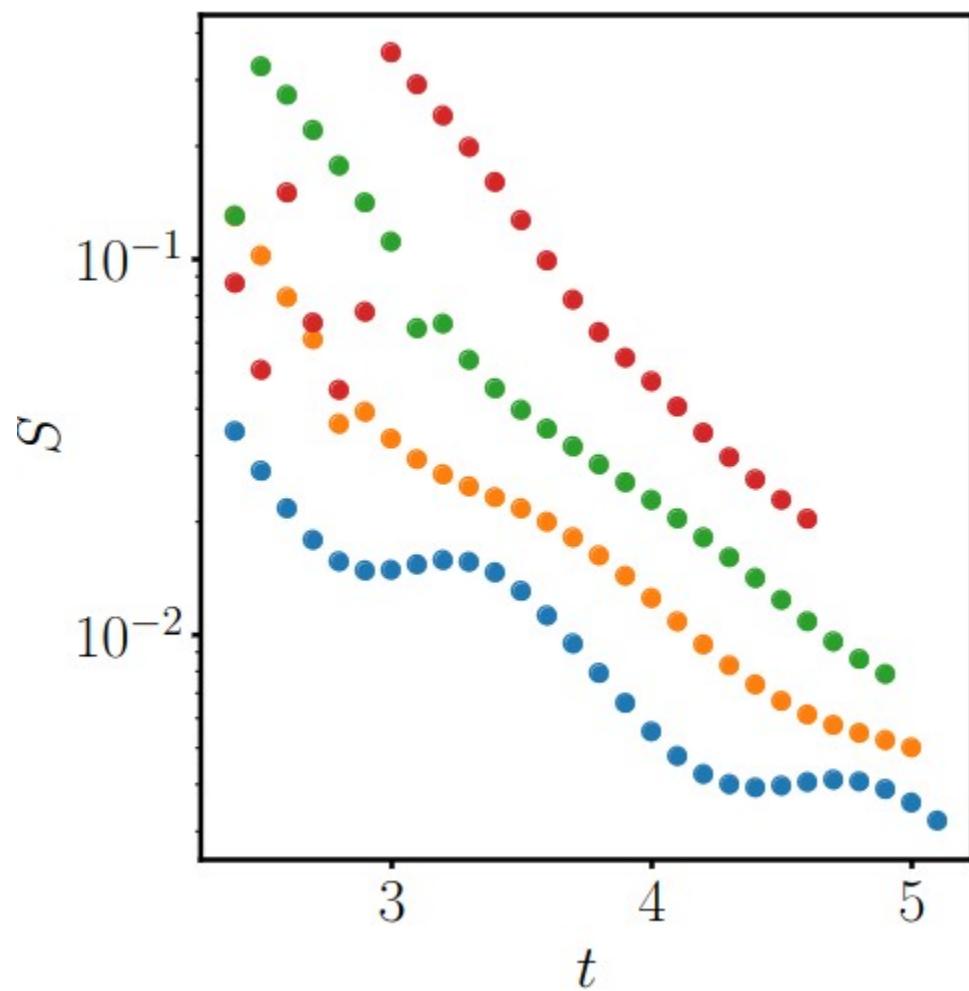
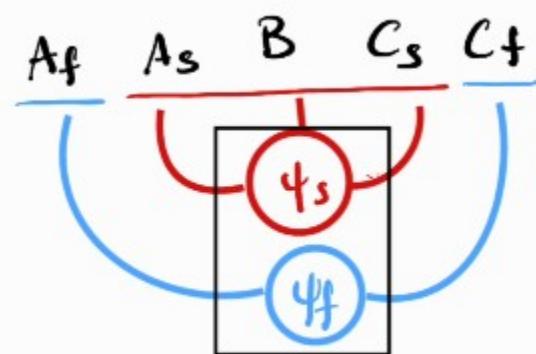
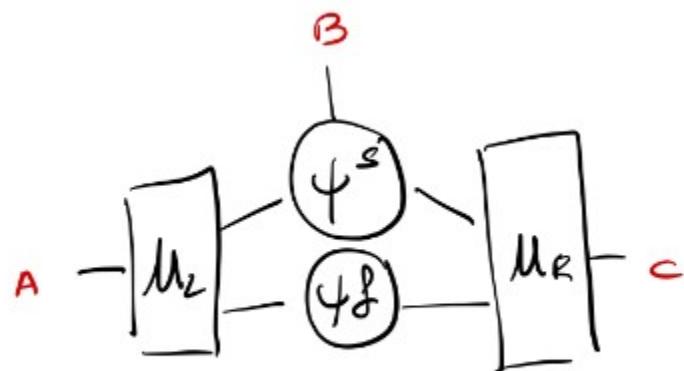
Work in progress with M Frias-Perez and MC Bañuls
(MPQ), arXiv2306xxxx

$$|\psi(t)\rangle \simeq |\psi_{A_s BC_s}\rangle \otimes |\psi_{A_f BC_f}\rangle$$



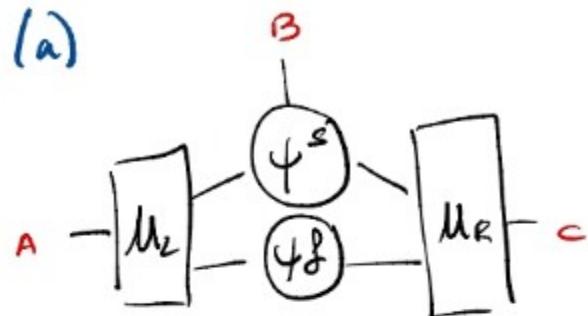


Frias-Perez (LT) Bañuls in prep



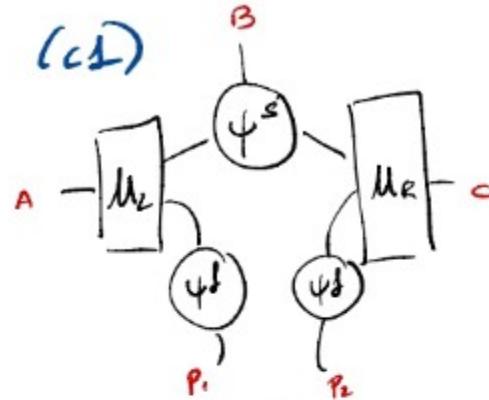
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(a)

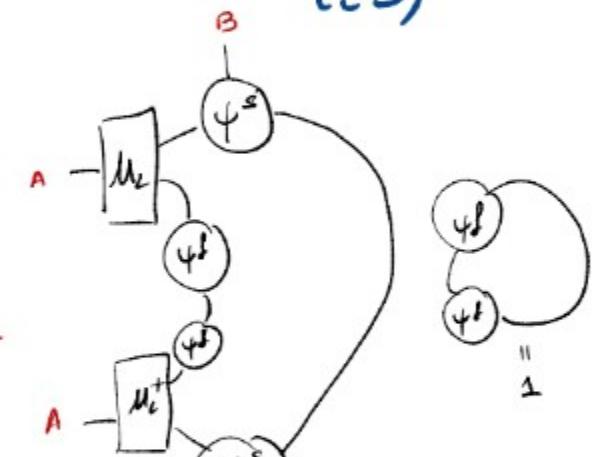


$|\Psi_{ABC}\rangle$

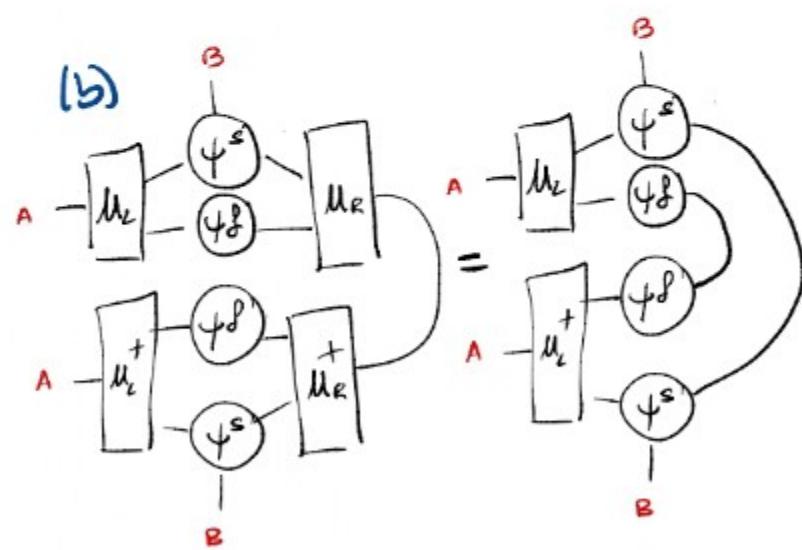
(c1)



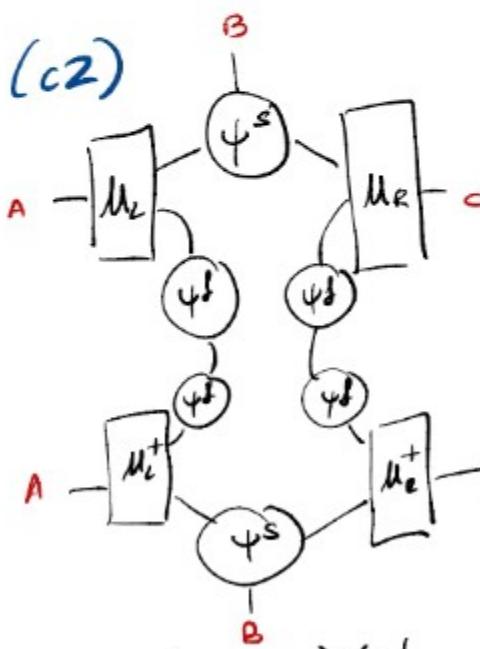
(c3)



(b)



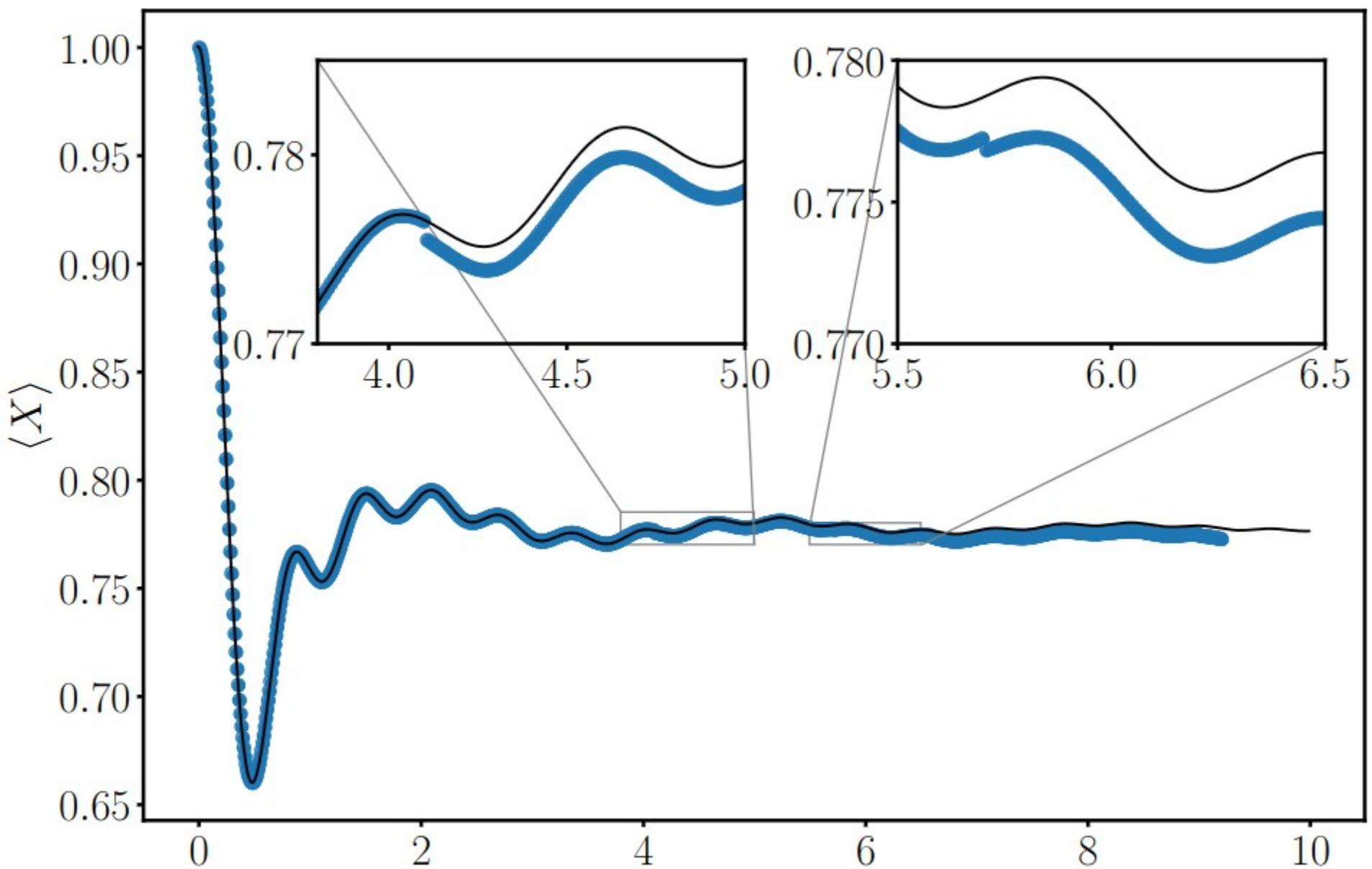
(c2)



$$\rho_{AB} = \text{Tr}_C |\Psi_{ABC}\rangle \langle \Psi_{ABC}|$$

$$\rho_{ABC} = \text{Tr}_{P_1 P_2} |\Psi_{ABC}, p_1 p_2\rangle \langle \Psi_{ABC}, p_1 p_2|$$

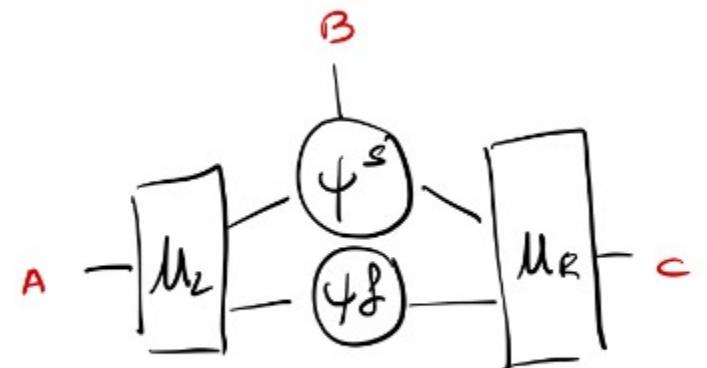
$$\rho_{AB} = \text{Tr}_{C, p_1 p_2} |\Psi_{ABC}, p_1 p_2\rangle \langle \Psi_{ABC}, p_1 p_2|$$



Frias-Perez (LT) Bañuls in prep



Different patterns of decoherences

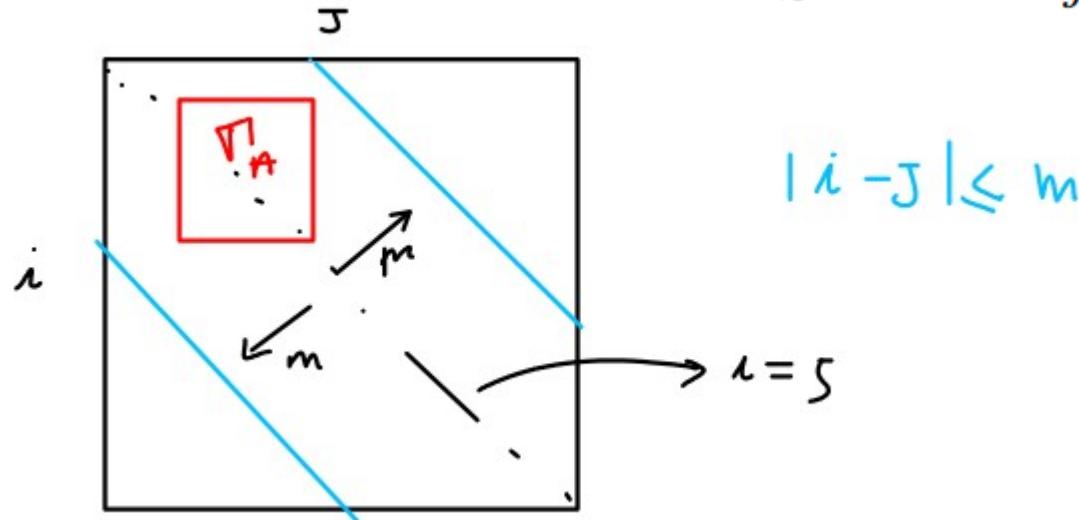


Frias-Perez (LT) Bañuls in prep

Gaussian states, Surace (LT) 2019

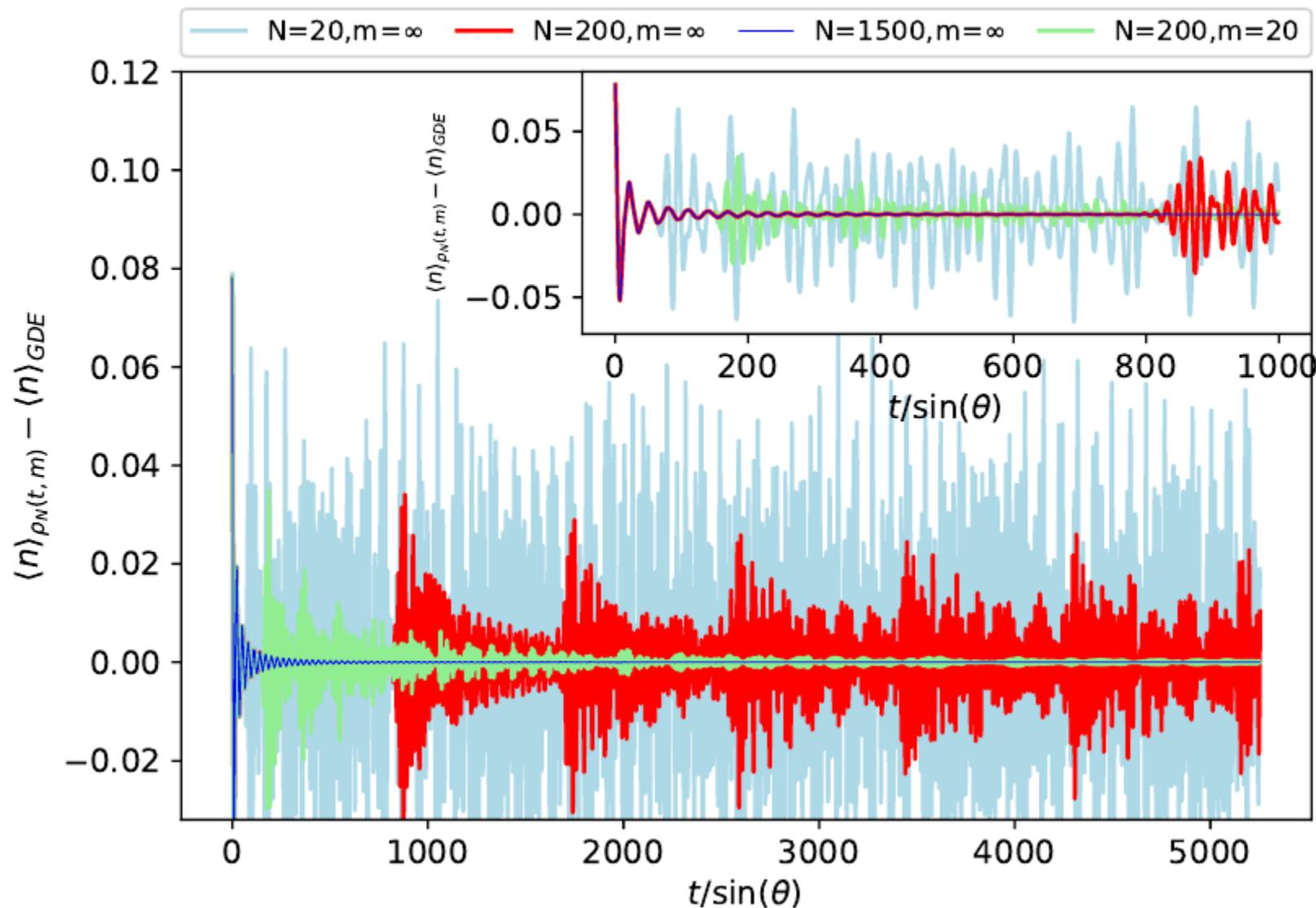
- The relevant correlations are inside a band of size m in the correlation matrix

$$\Gamma_{i,j} = \langle \vec{\alpha}_i \vec{\alpha}_j^\dagger \rangle,$$

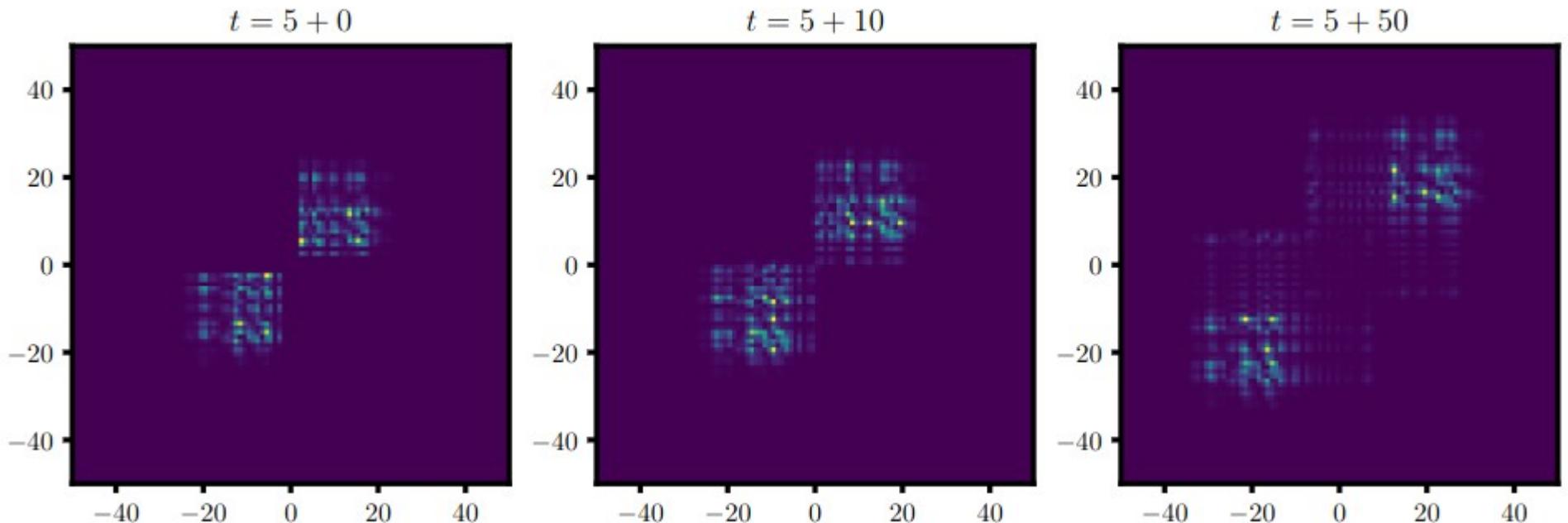


- At the truncation stage all correlation that are outside the band are zeroed

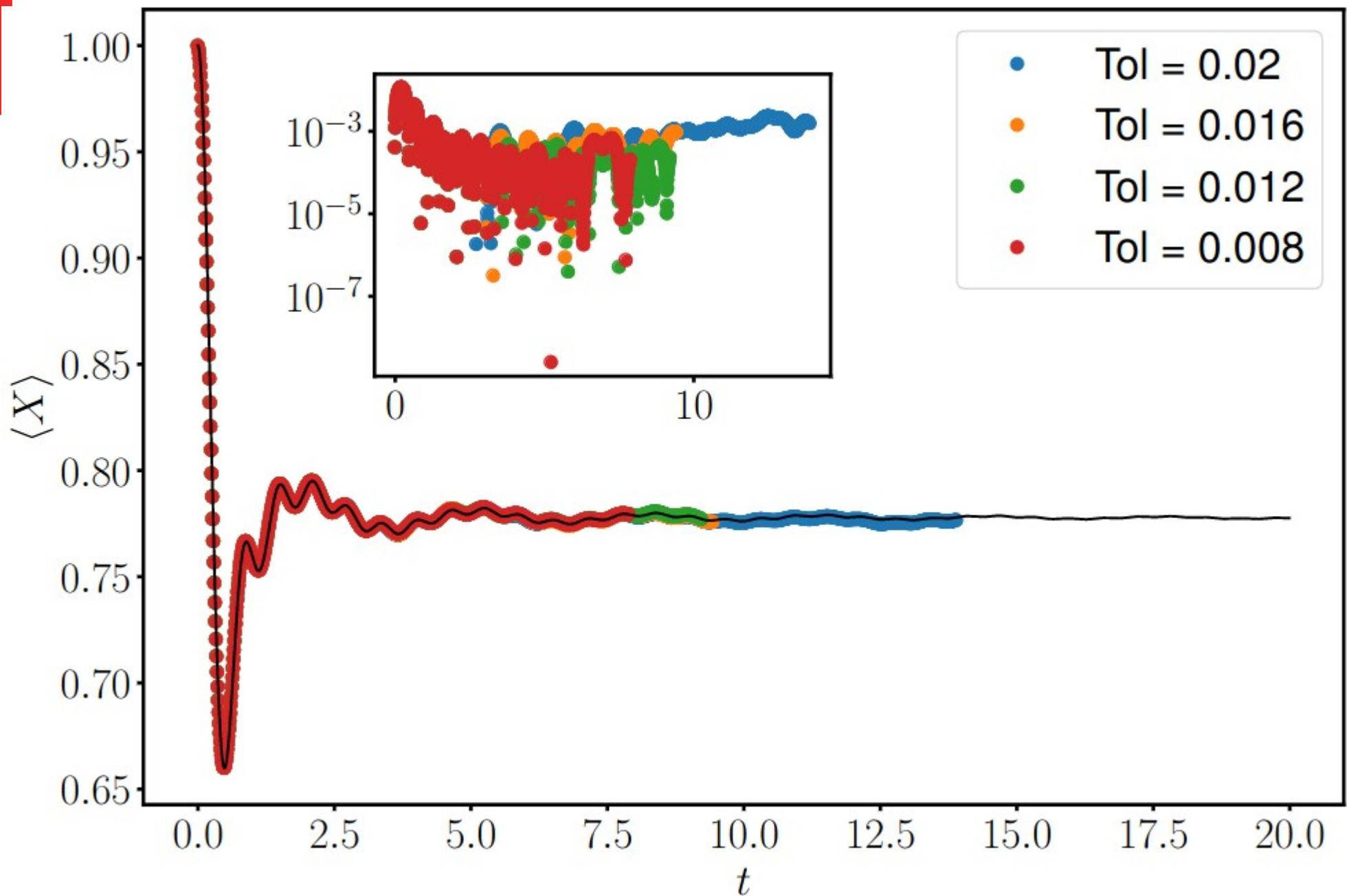
Quenches in Ising, using only local correlations



de-cohere only fast degrees of freedom



Frias-Perez (LT) Bañuls in prep



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- The entanglement barrier **can be circumvented**
- We need to **better understand the physics** of quenches to design better tensor network algorithms,
- **Decoherence** as the main ingredients allows to reach longer times

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