

Efficient quantum simulation of jet evolution in a medium

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Quantum Information in Spain ICE-8, Santiago de Compostela, May 29 - Jun 1, 2023

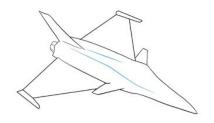
Based on 2104.04661, 2208.06750, 23XX.XXXX

In collaboration with João Barata, Meijian Li, Xiaojian Du, Carlos Salgado



What is jet quenching?

Slide from M Li's talk at Qiskit Fall Fest at USC (2022)



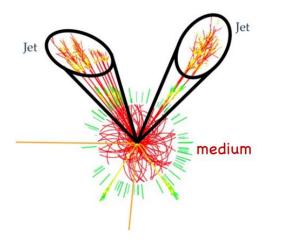


"Jet" a rapid stream "Quenching" a rapid cooling process



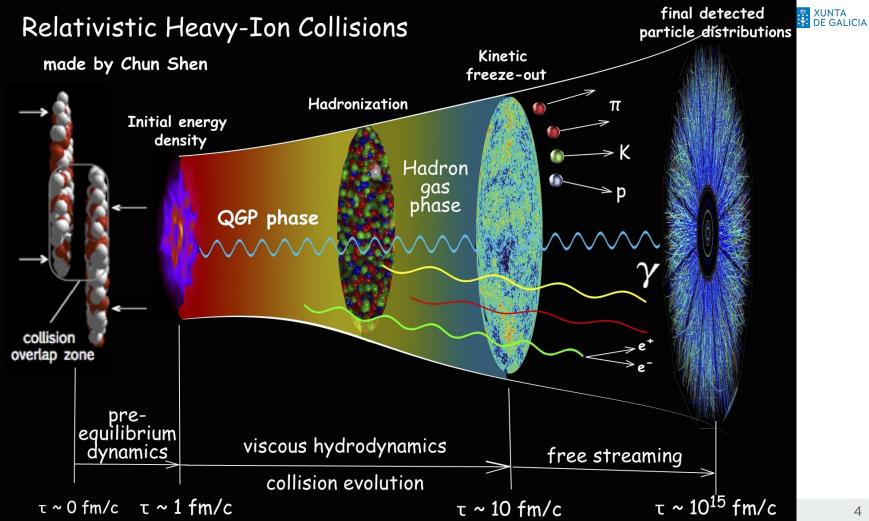
What is jet quenching?

Slide from M Li's talk at Qiskit Fall Fest at USC (2022)



In heavy ion collisions, a jet is a cone-shaped beam of energetic particles.

When propagating through the hot medium, it loses energy due to **jet-medium** interaction, a phenomenon known as "jet quenching".





Outline

- 1. Light-front Hamiltonian and physical setups
- 2. Quantum simulation algorithm
- 3. Results



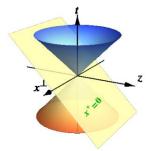
From LF Hamiltonian formalism to quantum simulation

Classical light-front Hamiltonian formalism

Scattering in Time-Dependent Basis Light-Front Quantization, PRD 88 (2013) 065014

Ultrarelativistic quark-nucleus scattering, PRD 101 (2020) 7, 076016

Scattering and gluon emission in a color field, PRD 104 (2021) 5, 056014



 $\frac{\text{front form}}{x^+ \triangleq x^0 + x^3}$

Quantum simulation

Single-particle digitization strategy for quantum computation of a ϕ 4 scalar field theory, PRA 103 (2021) 4, 042410

A quantum strategy to compute the jet quenching parameter, Eur.Phys.J.C 81 (2021) 10, 862

Quantum simulation of nuclear inelastic scattering, PRA 104 (2021) 1, 012611

Medium induced jet broadening in a quantum computer, PRD 106 (2022) 7, 074013

Quantum simulation of jet evolution in a medium (soon)



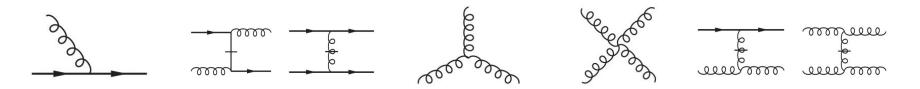
QCD Lagrangian

We start with the QCD lagrangian, with an external field

Li, Lappi, Zhao, PRD104.056014 (2021)

$$\mathcal{L} = -\frac{1}{4} F^{\mu\nu}{}_{a} F^{a}_{\mu\nu} + \overline{\Psi} (i\gamma^{\mu} D_{\mu} - m_{q}) \Psi$$
$$D^{\mu} \equiv \partial_{\mu} + ig(A^{\mu} + \mathcal{A}^{\mu})$$

The light-front Hamiltonian is obtained by the canonical light-front quantization via the standard Legendre transformation, Brodsky, Pauli, Pinsky, Phys.Rept.301: 299-486 (1998)





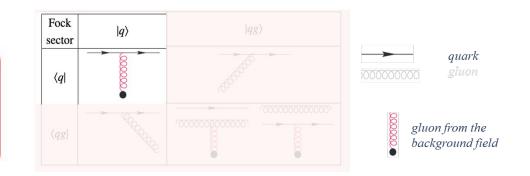
Physical set up, I

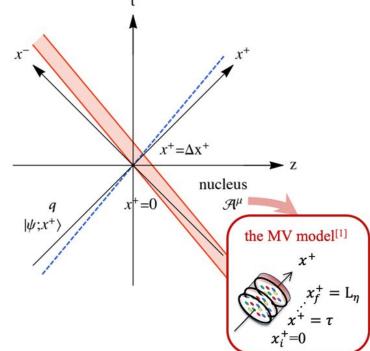
Li, Zhao, Maris, Chen, Li, Tuchin, Vary PRD101.076016 (2020) Barata, Du, Li, WQ, Salgado, PRD106,074013 (2022)

High-energy quark moving close to the light cone scattering on a dense nucleus medium

The light-front Hamiltonian in the $|q\rangle$ Fock sector:

$$P^{-}(x^{+}) = P^{-}_{\rm KE} + V_{\mathcal{A}}(x^{+})$$







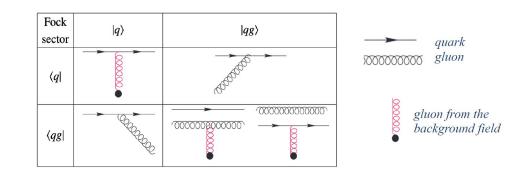
Physical set up, II

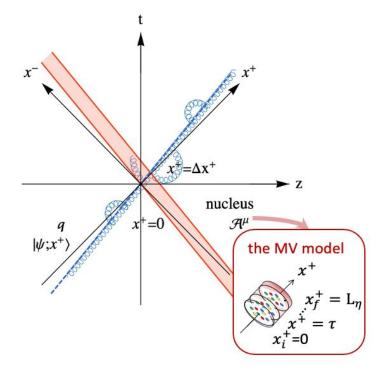
Li, Lappi, Zhao, PRD104.056014 (2021) Barata, Du, Li, WQ, Salgado (to appear)

High-energy quark moving close to the light cone scattering on a dense nucleus medium

The light-front Hamiltonian in the |q
angle+|qg
angle Fock sector:

$$P^{-}(x^{+}) = P_{\rm KE}^{-} + V(x^{+}) = P_{\rm KE}^{-} + \left\{ V_{qg} + V_{\mathcal{A}}(x^{+}) \right\}$$



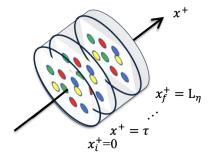




The medium and evolution

The stochastic background field uses the McLerran-Venugopalan (MV) model

McLerran, Venugopalan, PRD49, 2233; PRD49, 3352; PRD50, 2225 (1994)



$$\langle\!\langle \rho_a(\vec{x}_{\perp}, x^+) \rho_b(\vec{y}_{\perp}, y^+) \rangle\!\rangle = g^2 \tilde{\mu}^2 \delta_{ab} \delta^2(\vec{x}_{\perp} - \vec{y}_{\perp}) \delta(x^+ - y^+)$$

$$(m_g^2 - \nabla_{\perp}^2) \mathcal{A}_a^-(\vec{x}_{\perp}, x^+) = \rho_a(\vec{x}_{\perp}, x^+) \qquad Q_s^2 = C_F \frac{(g^2 \tilde{\mu})^2 L_\eta}{2\pi}$$
saturation scales

Light-front time evolution of the probe, decomposed as sequence of unitary operators

$$|\psi_{L_{\eta}}\rangle = U(L_{\eta}; 0) |\psi_{0}\rangle \equiv \mathcal{T}_{+} e^{-i \int_{0}^{L_{\eta}} dx^{+} P^{-}(x^{+})} |\psi_{0}\rangle \qquad \qquad U(L_{\eta}; 0) = \prod_{k=1}^{N_{t}} U(x_{k}^{+}; x_{k-1}^{+})$$

QI in Spain ICE-8, 2023



Quantum simulation algorithm

Wiesner, 9603028 (1996); Zalka, 9603026 (1996)

- 1. Define problem Hamiltonian
- 2. Encode Hamiltonian onto basis
- 3. Prepare initial states
- 4. Evolution
- 5. Measurement protocol

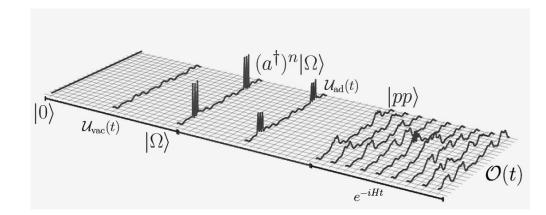


Image from Lamm's talk at Fermilab (2021)



Basis encoding

The general QCD quantum state in single-particle basis:

Barata, Mueller, Tarasov, Venugopalan, PRA103, 4, 042410 (2021)

$$|\psi
angle = |q
angle \cdots |q
angle \otimes |g
angle \cdots |g
angle \qquad \sim \prod_{i=1}^m \left(|e_{g_i}
angle \otimes |g_i
angle
ight) \otimes \prod_{i=1}^n \left(|e_{q_i}
angle \otimes |q_i
angle
ight)$$

In our cases of |q
angle+|qg
angle and |q
angle Fock spaces:

$$\begin{split} |\psi\rangle &= |\zeta\rangle \otimes \underbrace{\left(\left| g_x \right\rangle \left| g_y \right\rangle \left| c_g \right\rangle \right)}_{|g\rangle} \otimes \underbrace{\left(\left| q_x \right\rangle \left| q_y \right\rangle \left| c_q \right\rangle \right)}_{|q\rangle} \\ |\psi\rangle &= \underbrace{\left(\left| q_x \right\rangle \left| q_y \right\rangle \left| c_q \right\rangle \right)}_{|q\rangle} \\ \zeta\rangle &= |0\rangle \Leftrightarrow \operatorname{Fock} |q\rangle, \, k_g^+ = 0, k_q^+ = K \\ \zeta\rangle &= |1\rangle \Leftrightarrow \operatorname{Fock} |qg\rangle, k_g^+ = 1, k_q^+ = K - 1 \\ \cdots \\ \end{split}$$

$$N_{\text{tot}} \sim \lceil K \rceil N_{\perp}^4 \to n_Q \sim 4 \log N_{\perp} + \log \lceil K \rceil$$



Transverse & Longitudinal Lattice

• Transverse lattice is periodical, position & momentum related via quantum fourier transform (FT)

$$(n_x, n_y) \iff (n_x + i2N_\perp, n_y + j2N_\perp)$$

$$\vec{r}_{\perp} = (n_x, n_y)a_{\perp} \qquad \vec{p}_{\perp} = (k_x, k_y)b_{\perp}$$

• Longitudinal direction is periodical (anti-periodical) for bosons (fermions)

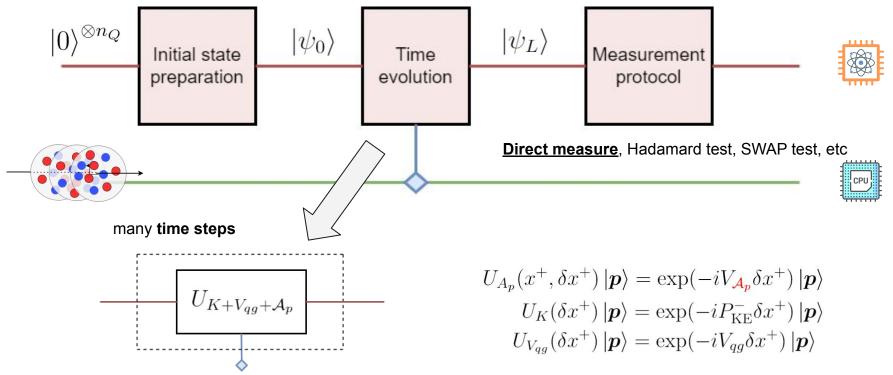
$$p_l^+ = \frac{2\pi}{L}k_l^+, \quad k_q^+ = \frac{1}{2}, \frac{3}{2}, \cdots, \quad k_g^+ = 1, 2, 3, \cdots$$

(-3, 3)	(-2, 3)	(-1, 3)	(0, 3)	(1, 3)	(2, 3)	(3, 3)	(4, 3)
(-3, 2)	(-2, 2)	(-1, 2)	(0, 2)	(1, 2)	(2, 2)	(3, 2)	(4, 2)
(-3, 1)	(-2, 1)	(-1, 1)	(0, 1)	(1, 1)	(2, 1)	(3, 1)	(4, 1)
(-3, 0)	(-2, 0)	(-1, 0)	(0, 0)	(1, 0)	(2, 0)	(3, 0)	(4, 0)
(-3, -1)	(-2, -1)	(-1, -1)	(0, -1)	(1, -1)	(2, -1)	(3, -1)	(4, -1)
(-3, 3)	(-2, -2)	(-1, -2)	(0, -2)	(1, -2)	(2, 2)	(3, 2)	(4, 2)
(-3, 3)	(-2, -3)	(-1, -3)	(0, -3)	(1, -3)	(2, 3)	(3, 3)	(4, 3)



Momentum-space simulation

Amplitude level

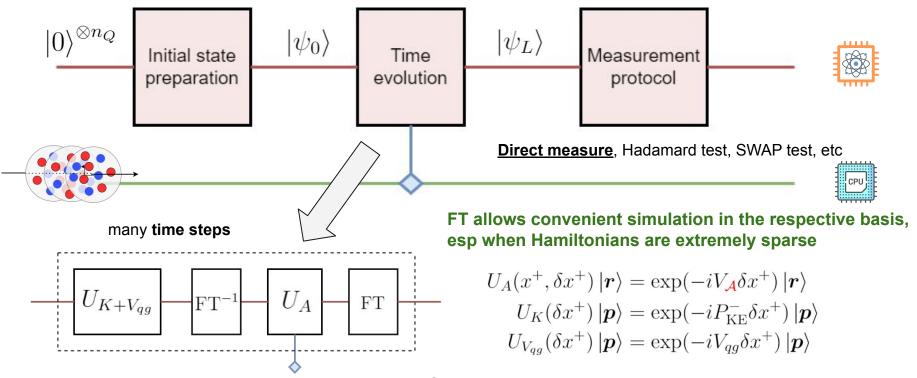


Ham => pauli strings => quantum gates



Mixed-space simulation

Amplitude level



Sparse Ham => pauli strings => quantum gates



Simulation strategies

Momentum-space simulation

$$U(x_k^+ + \delta x; x_k^+) = U_{K+V_{qg}+V_{\mathcal{A}_p}}(\delta x^+, x_k^+)$$

$$\equiv \exp\left\{-i\delta x^+ \left[\tilde{K} + \tilde{V}_{qg} + \tilde{V}_{\mathcal{A}_p}(x_k^+)\right]\right\},\$$

Pros Easy to trotter (no need of basis transform)

Mixed-space simulation

$$U(x_k^+ + \delta x; x_k^+) \approx U_{K+V_{qg}}(\delta x^+) U_{V_{\mathcal{A}}}(\delta x^+, x_k^+)$$

$$\equiv \exp\left\{-i\delta x^+ \left[\tilde{K} + \tilde{V}_{qg}\right]\right\} \exp\left\{-i\delta x^+ \left[\tilde{V}_{\mathcal{A}}(x_k^+)\right]\right\}$$

Need quantum Fourier Transforms (qFT)

Cons Harder to scale to larger problem size; Expensive to evaluate Pauli terms

Sparse Hamiltonian; better scaling at larger problem size



Simulation parameters

- Our initial state for the jet is a quark state with zero transverse momenta: $(p_x, p_y) = (0, 0)$
- Duration of static medium: $L_{\eta} = 50 \,\mathrm{GeV}^{-1} \approx 10 \,\mathrm{fm}$
- 5 stochastic fields are used for configuration average $\langle\!\langle \langle \hat{O} \rangle \rangle\!\rangle = \lim_{n \to \infty} \frac{1}{n} \sum_{i=1}^{n} \langle \hat{O} \rangle_i$
- Computational lattices:

 $|q\rangle + |qg\rangle$ Transverse: $N_{\perp} = 1$ Longitudinal: K = 3.5

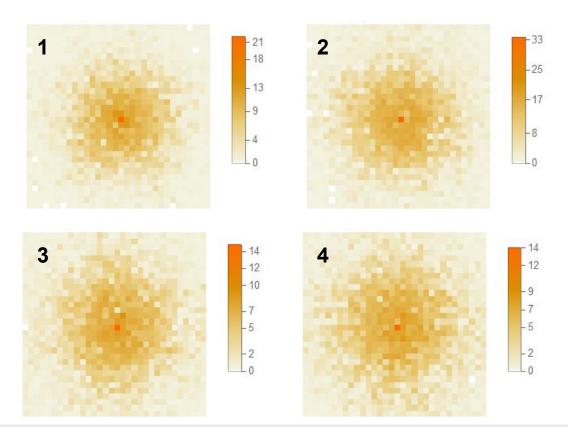
|q> Transverse: $N_{\perp} = 16$

- Non-abelian SU(2) color
- Non-flip case for spin $(\lambda_Q,\lambda_q,\lambda_g)=(\uparrow,\uparrow,\uparrow)$
- Sufficient time steps in trotterization and medium layers in background medium

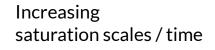
We present our preliminary results of jet evolution in vacuum and in medium, using IBM Qiskit quantum simulators



Momentum broadening Fock |q> only



Barata, Du, Li, WQ, Salgado, PRD106,074013 (2022)



transverse lattice: 32 X 32

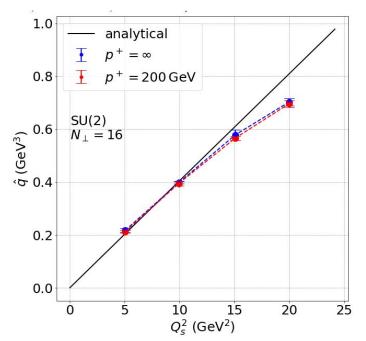
819200 shots, 11 qubits

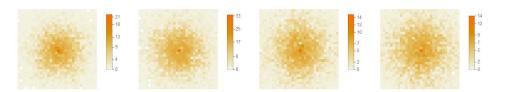


Momentum broadening Fock |q> only

Barata, Du, Li, WQ, Salgado, PRD106,074013 (2022)

quenching parameter





rtical
$$\hat{q} = \frac{g^4}{4\pi} C_F \tilde{\mu}^2 \left\{ \log \left(1 + \frac{\pi^2}{a_\perp^2}}{m_g^2} \right) - \frac{1}{1 + \frac{a_\perp^2 m_g^2}{\pi^2}} \right\} \sim Q_s^2 / L_\eta$$

Simulation

$$\hat{q} = rac{\Delta \left\langle p_{\perp}^2(x^+)
ight
angle}{\Delta x^+} \;\; = \left\langle oldsymbol{p}_{\perp}^2
ight
angle / L_\eta$$

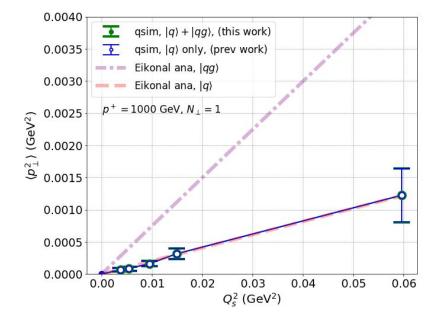
819200 shots, 11 qubits



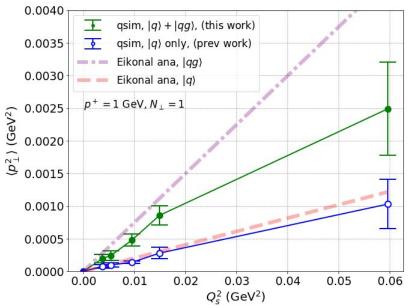
Momentum broadening Fock |q> + |qg>

Barata, Du, Li, WQ, Salgado (to appear)

$$\langle p_{\perp}^2 \rangle = \mathcal{P}_{|q\rangle} \langle p_{\perp}^2 \rangle_{|q\rangle} + \mathcal{P}_{|qg\rangle} \langle p_{\perp}^2 \rangle_{|qg\rangle}$$



$$\hat{q}_{Eik}(x = a_{\perp} m_g/\pi, N_{\perp} = 1) \big|_{\text{on lattice}}$$
$$= C_F g^4 \tilde{\mu}^2 \frac{1}{(2\pi)^2} \left[\frac{2}{(x^2 + 1)^2} + \frac{2}{(x^2 + 2)^2} \right] \sim Q_s^2 / L_\eta$$

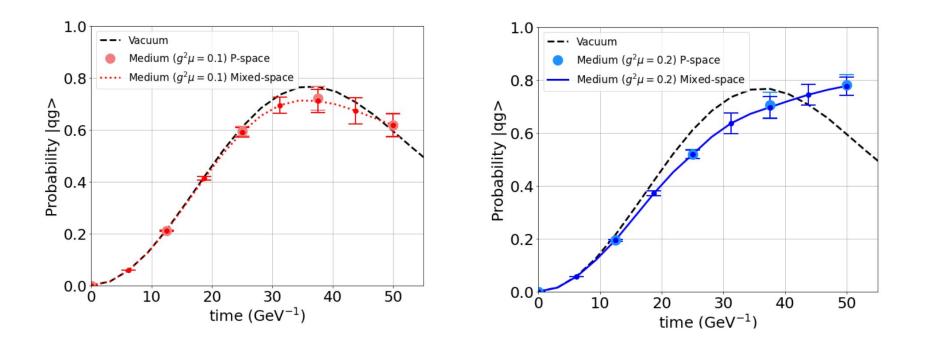


819200 shots, 9 qubits, Ndx=16, Neta=16



Quantum jet evolution

Barata, Du, Li, WQ, Salgado (to appear)



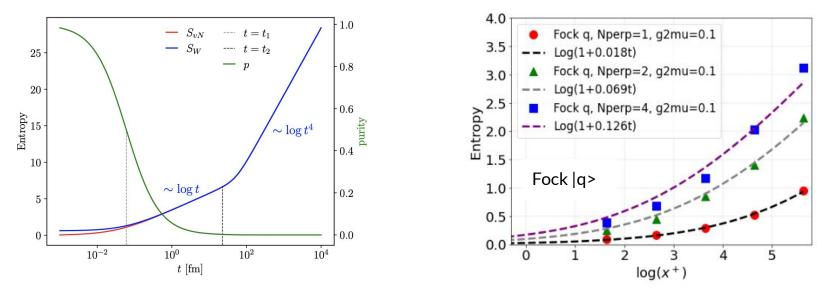
819200 shots, 9 qubits, Ndx=16, Neta=16



Entropy evolution of the quark jet

Barata, Du, Li, WQ, Salgado (to appear)

$$\rho_{\mathsf{N}_{\mathsf{m}}} = (1/N_m) \sum_{i=1}^{N_m} \rho_{\mathsf{i}} \qquad \langle S \rangle = -\mathrm{Tr}(\rho_{\mathsf{N}_{\mathsf{m}}} \log \rho_{\mathsf{N}_{\mathsf{m}}})$$



Barata, Blazizot, Mehtar-Tani, 2305.10476 (2023)

819200 shots, 9 qubits, Ndx=16, Neta=16



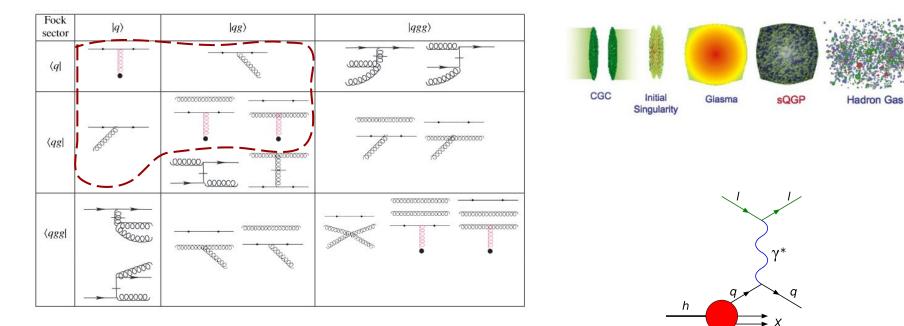
Summary and outlook

- We study multi-particle jet evolution in a medium within two Fock sectors |q> + |qg> using light-front QCD Hamiltonian approach on quantum simulator.
- Despite of having a small model space, we can study jet evolution using quantum simulators, in agreement with analytical result and classical diagonalization.
- Quantum simulation can be effective in reducing problem complexity faced in classical simulation. Our formalism can be extended to higher Fock sectors, various QCD mediums, different physical processes.





Future work



McLerran, 0812.4989 (2008)