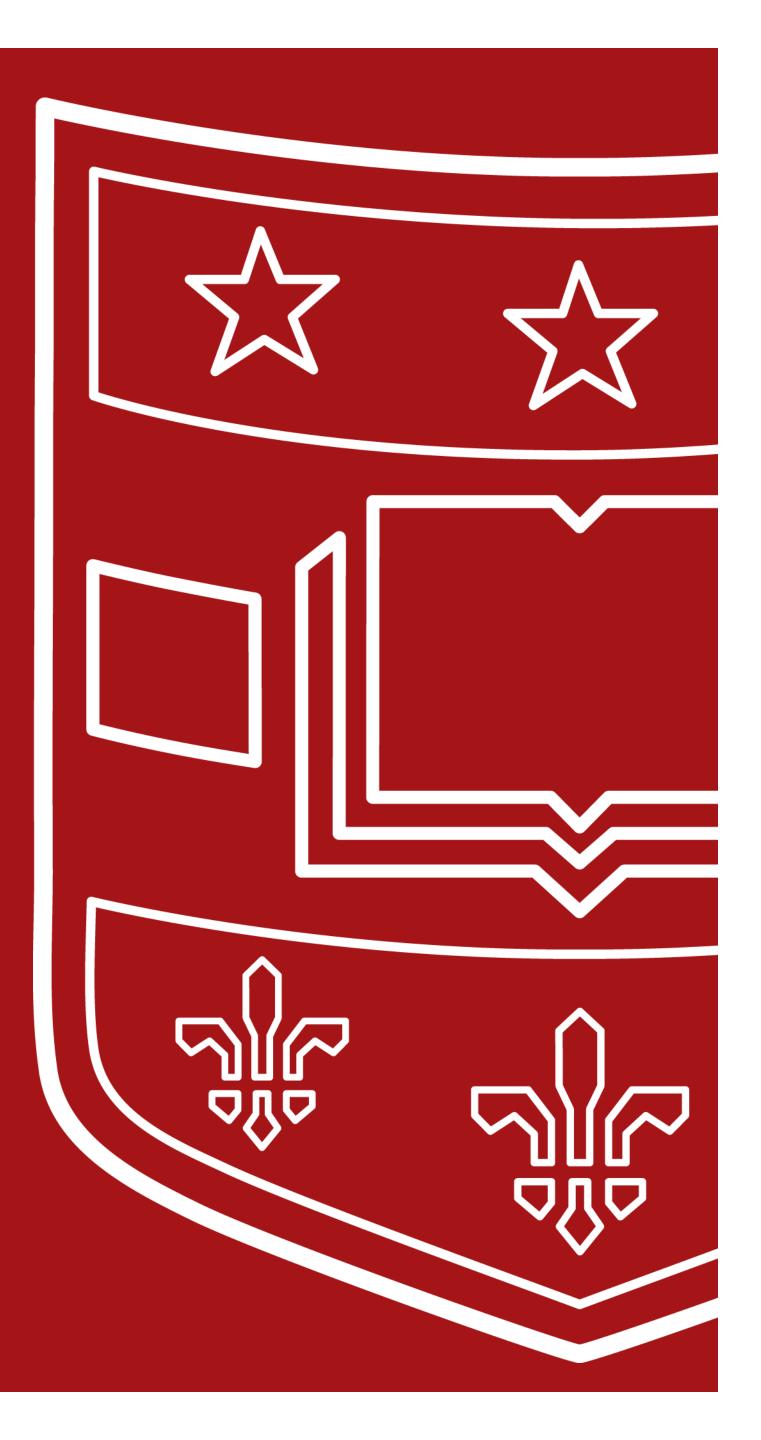
How Uncertain Am I? Theoretical errors in Bayesian model calibration for EFTs

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Effective Field Theory

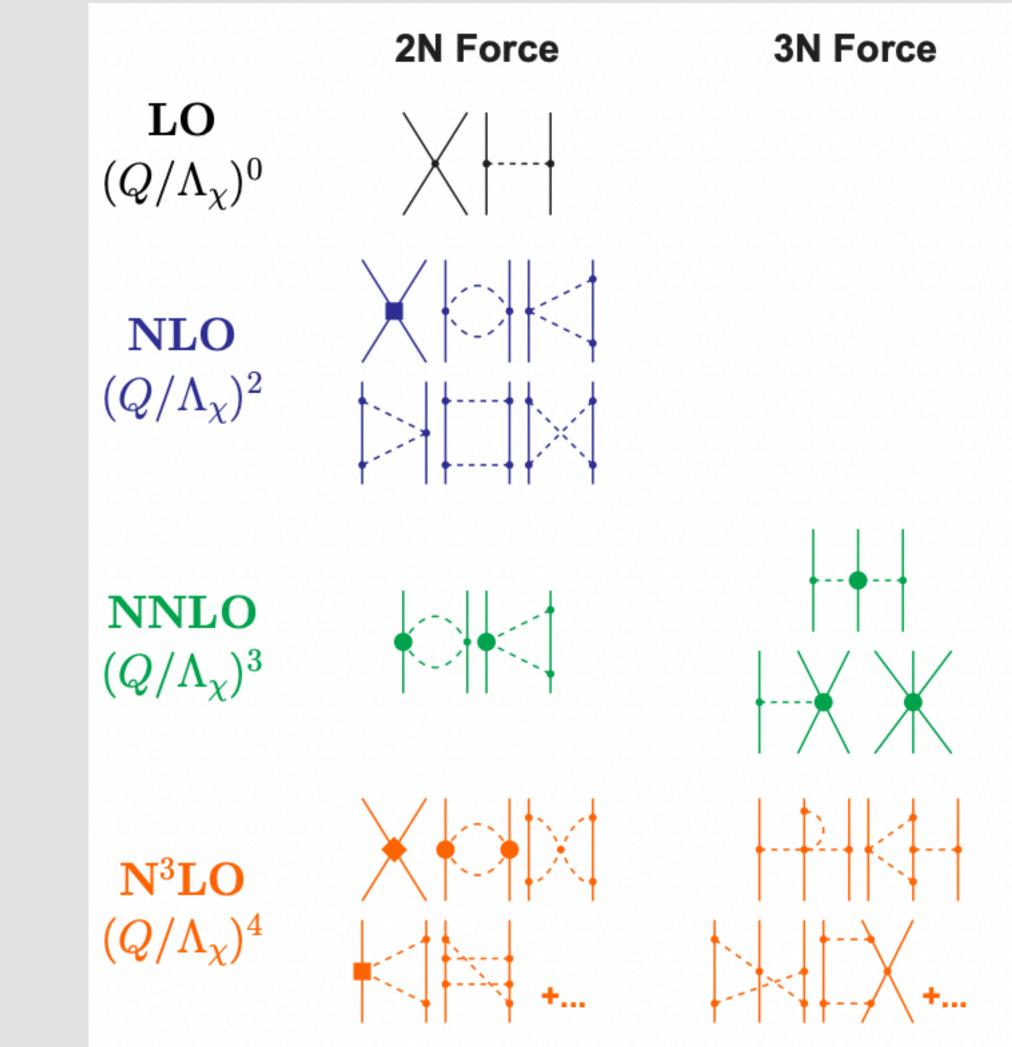
We take an effective expansion of QCD preserving chiral symmetry.

The interaction can be ordered in terms of powers of Q/Λ_{γ}

- Q is a momentum or pion mass
- Λ_{γ} is the symmetry breaking scale

Gives a systematic ordering to improve the interaction.





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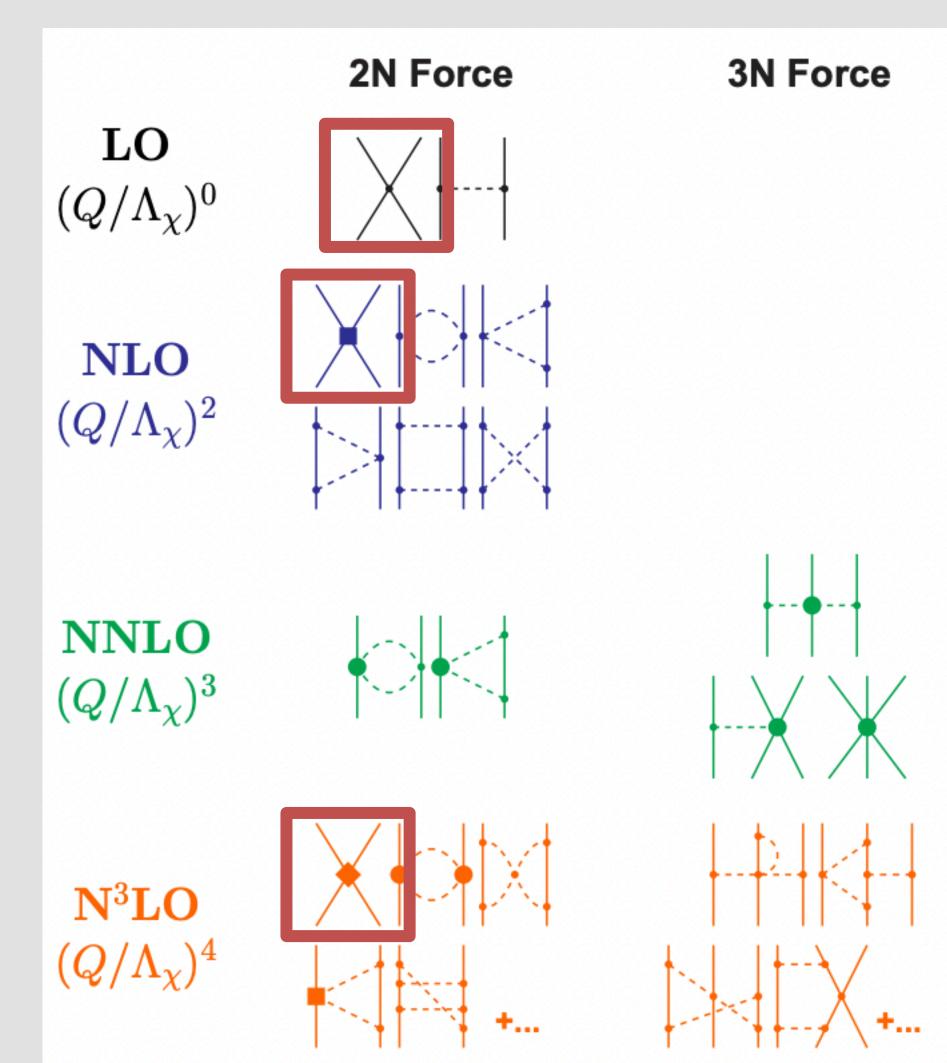
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Pionless EFT

We are working in an EFT framework without pions

Our interaction takes the form:

 $v^{(NLO)}(k) = C_1 k^2 + C_2 k^2 \sigma_1 \cdot \sigma_2 + C_3 S_{12}(k) + C_4 k^2 \tau_1 \cdot \tau_2$ $+iC_5S \cdot (K \times k) + C_6k^2\tau_1 \cdot \tau_2\sigma_1 \cdot \sigma_2 + C_7S_{12}(k)\tau_2 \cdot \tau_2$



 $v^{(LO)}(k) = C_{\rm S} + C_T \sigma_1 \cdot \sigma_2$



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Undetermined Low Energy Constants (LECs)

$$C_S + C_T \sigma_1 \cdot \sigma_2$$

+ $C_5 \delta \cdot (K \times k) + C_6 k^2 \tau_1 \cdot \tau_2 \sigma_1 \cdot \sigma_2 + C_7 \delta_{12}(k) \tau_2 \cdot \tau_2$



Regularization

must be local in coordinate space (for QMC).

We employ a Gaussian cutoff in coordinate space, which smears δ -functions upon Fourier transformation.

f(r) = -

space cutoff of 200 MeV.



- To use these interactions, they must be regularized in some fashion and

$$\frac{1}{3/2R_0^3}e^{-\left(\frac{r}{R_0}\right)^2}$$

For this work, we choose $R_0 = 2.0$ fm corresponding to a momentum



Model Calibration

To calibrate our EFT model, we use a Bayesian framework $\rho(\mathbf{a} | \mathbf{y}, I) \propto \rho(\mathbf{y} | \mathbf{a}, I) \rho(\mathbf{a} | I)$

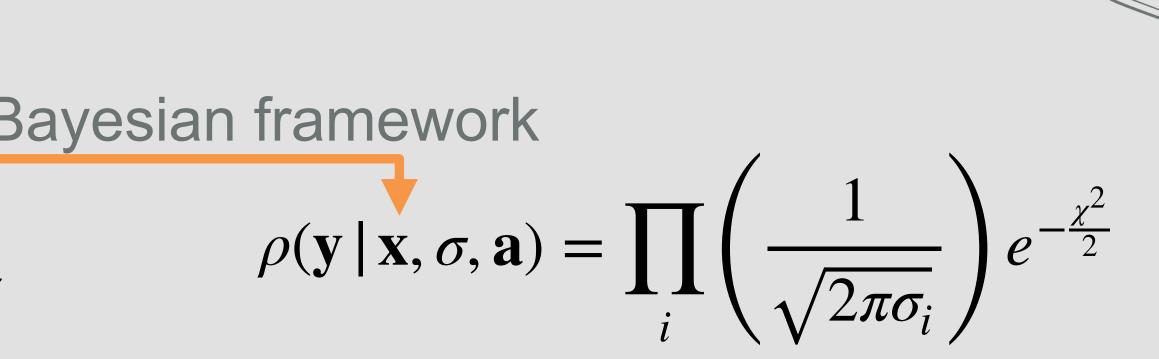
posterior

likelihood prior

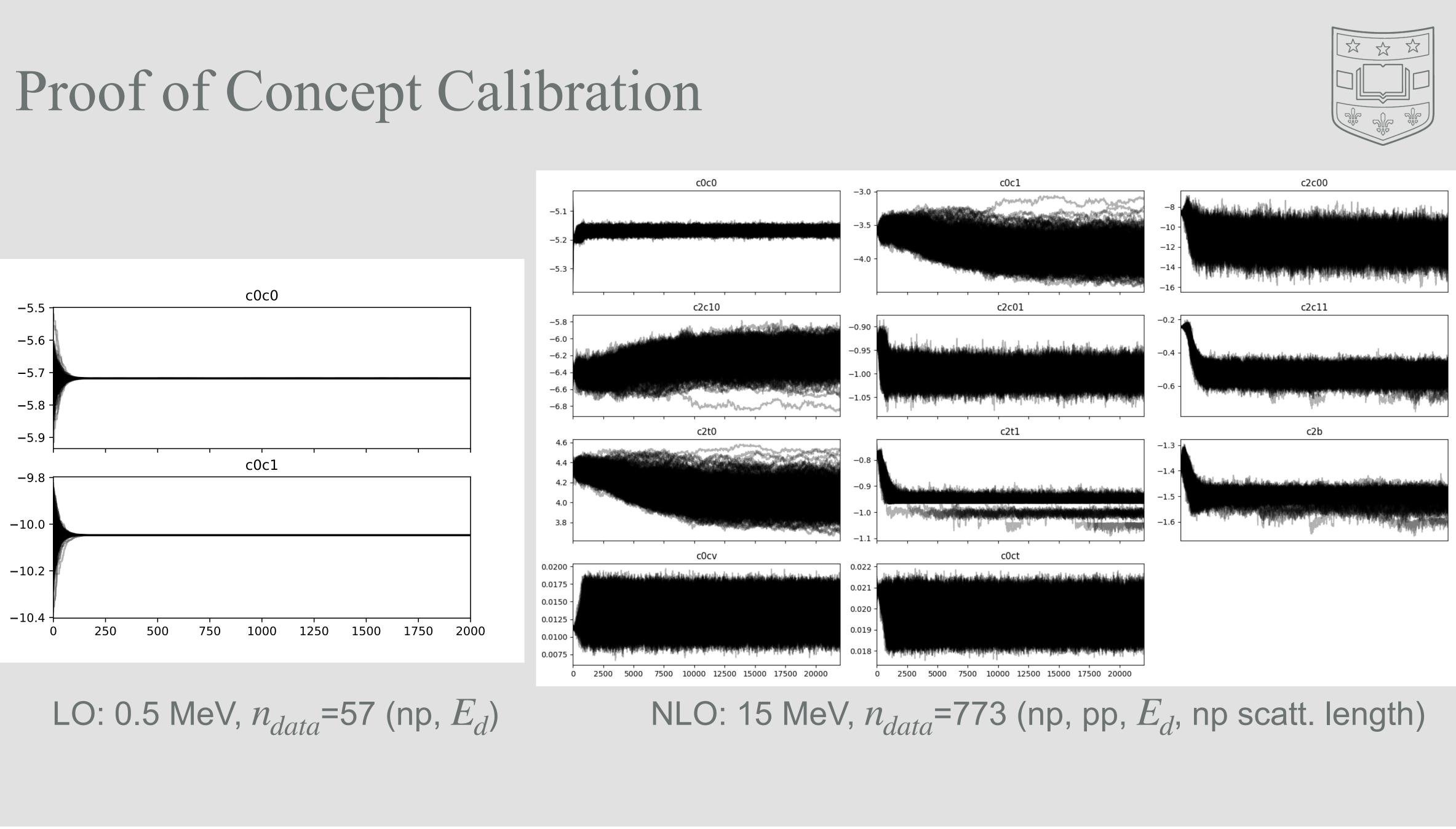
The posterior gives the distribution for a certain set of parameters, a (LECs), given data, y (scattering data), and any other information, I, which we maximize.

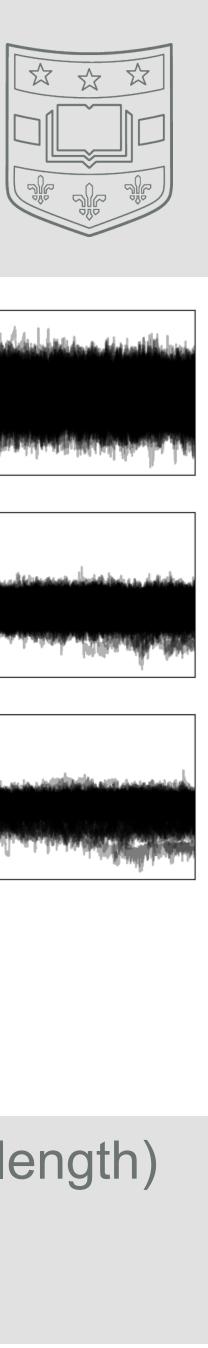
The likelihood gives the probability of scattering data for a given parameter set •Find region of minimum χ^2

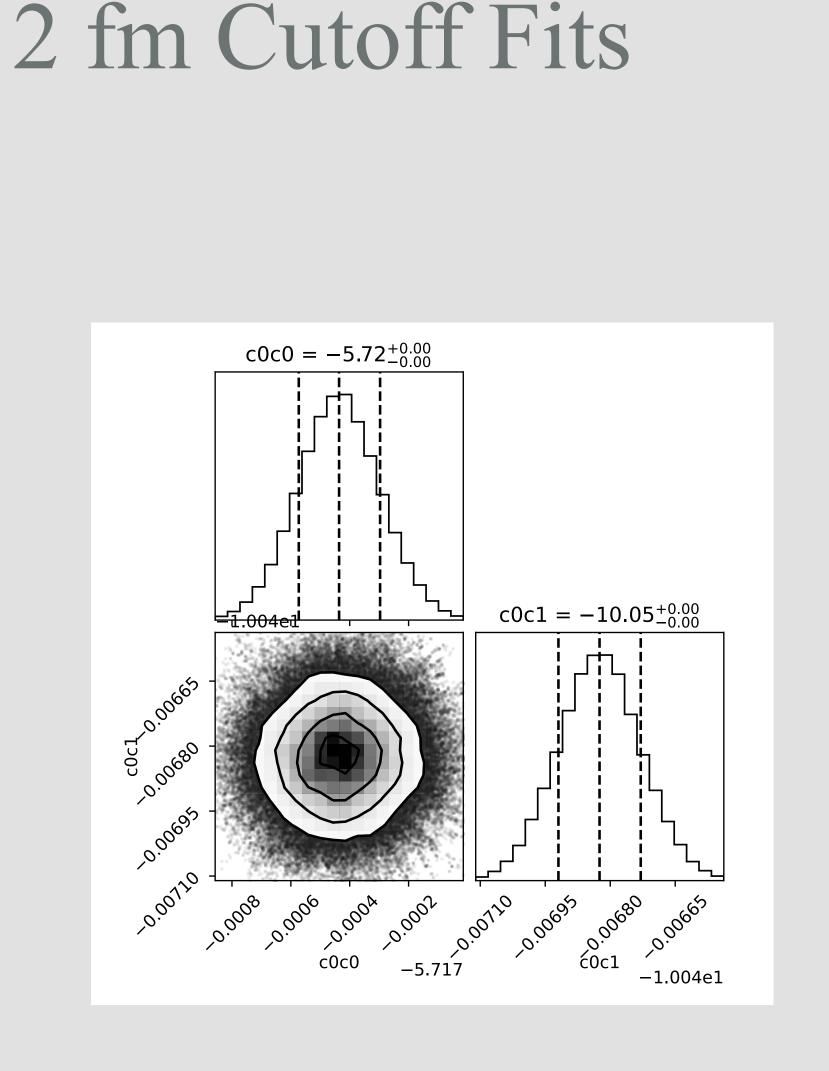
The prior encodes any external information we know about the parameters. • Initial point/naturalness \rightarrow ours is relatively uninformative

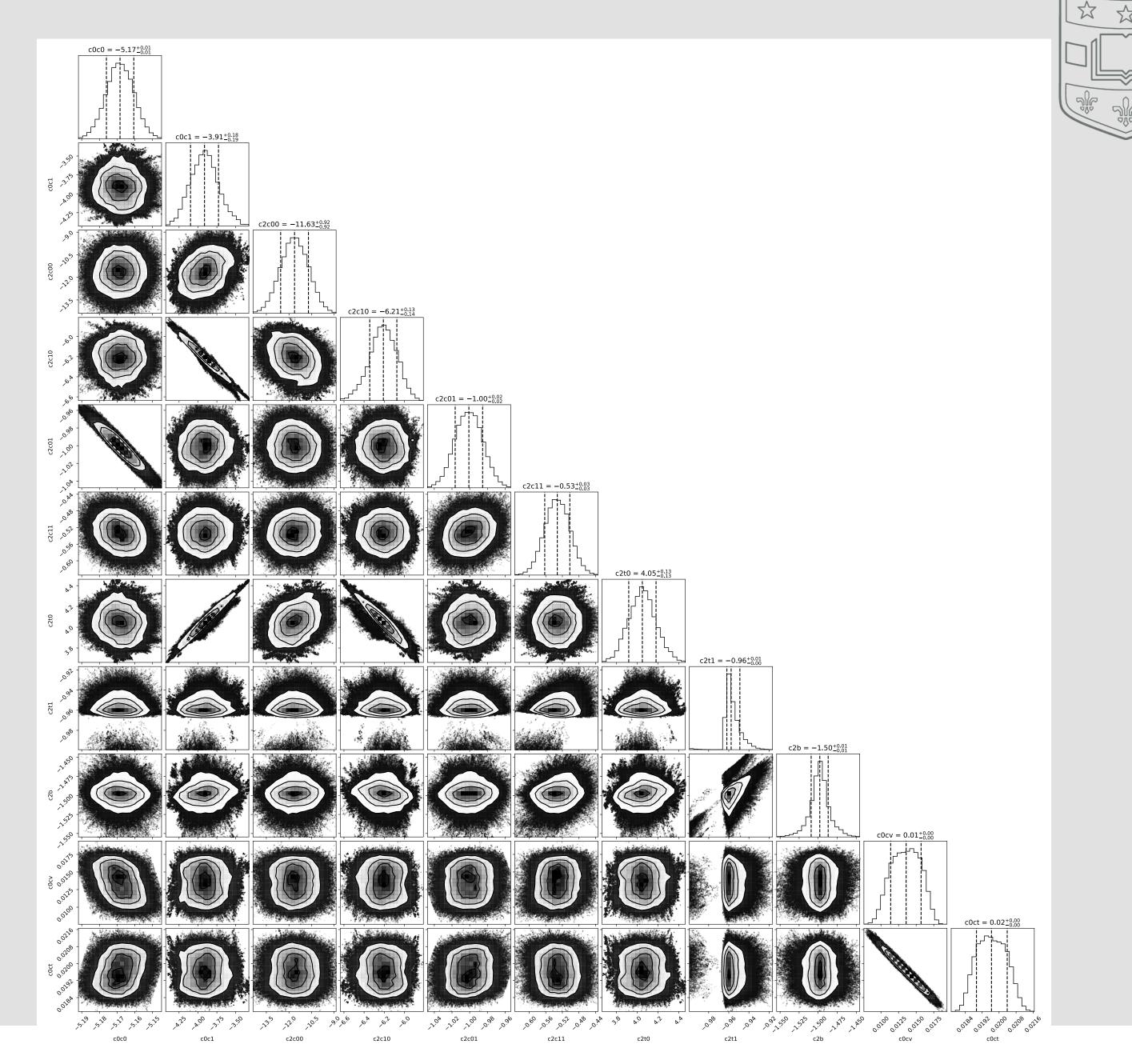














Full Bayesian model calibration

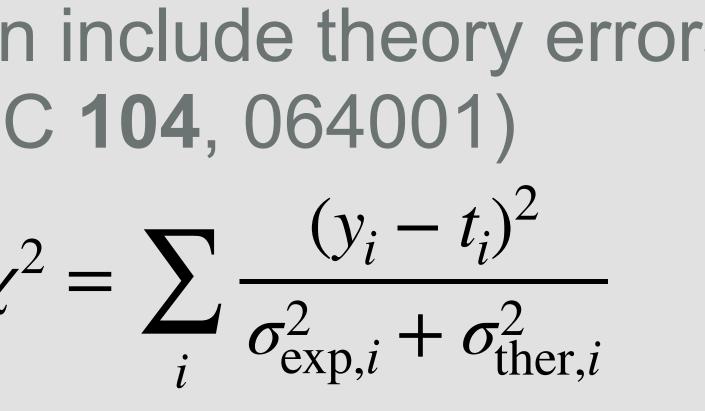
In our model calibration, we can include theory errors (Wesolowski et. al. Phys. Rev. C 104, 064001)

$$\chi^2 = \sum_{i} \frac{(y_i - t_i)^2}{\sigma_{\exp,i}^2} \to \chi^2$$

Where $\sigma_{\text{ther},i}^{2} = \frac{(y_{\text{ref},i}\,\bar{c}\,Q_{i}^{n+1})^{2}}{1-O_{i}^{2}}, \quad Q_{i} = \frac{\max[p_{\text{soft}},p_{i}]}{\Lambda_{b} \sim m_{\pi}}$

and $y_{ref,i}$ sets the scale of the correction for observable y_i , and \bar{c} sets the magnitude of the correction.



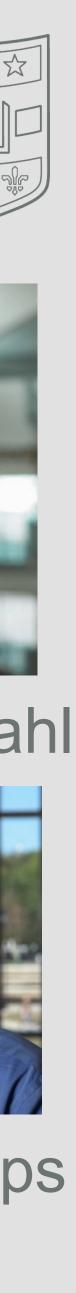




Dick Furnstahl



Daniel Phillips



Correlated theory errors

We can correlate our theory errors as well:

$$\sigma_{\text{ther},i}^{2} = \frac{(y_{\text{ref},i}\,\bar{c}\,Q_{i}^{n+1})^{2}}{1-Q_{i}^{2}} \to \sigma_{\text{ther},ij}^{2} = \frac{y_{\text{ref},i}\,y_{\text{ref},j}\,\bar{c}^{2}\,Q_{i}^{n+1}Q_{j}^{n+1}}{1-Q_{i}\,Q_{j}}e^{-|p_{i}-p_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}}e^{-|\theta_{i}-\theta_{j}|/2l_{p}$$

with the goodness of fit determined by the Mahalanobis distance ("modified" χ^2)

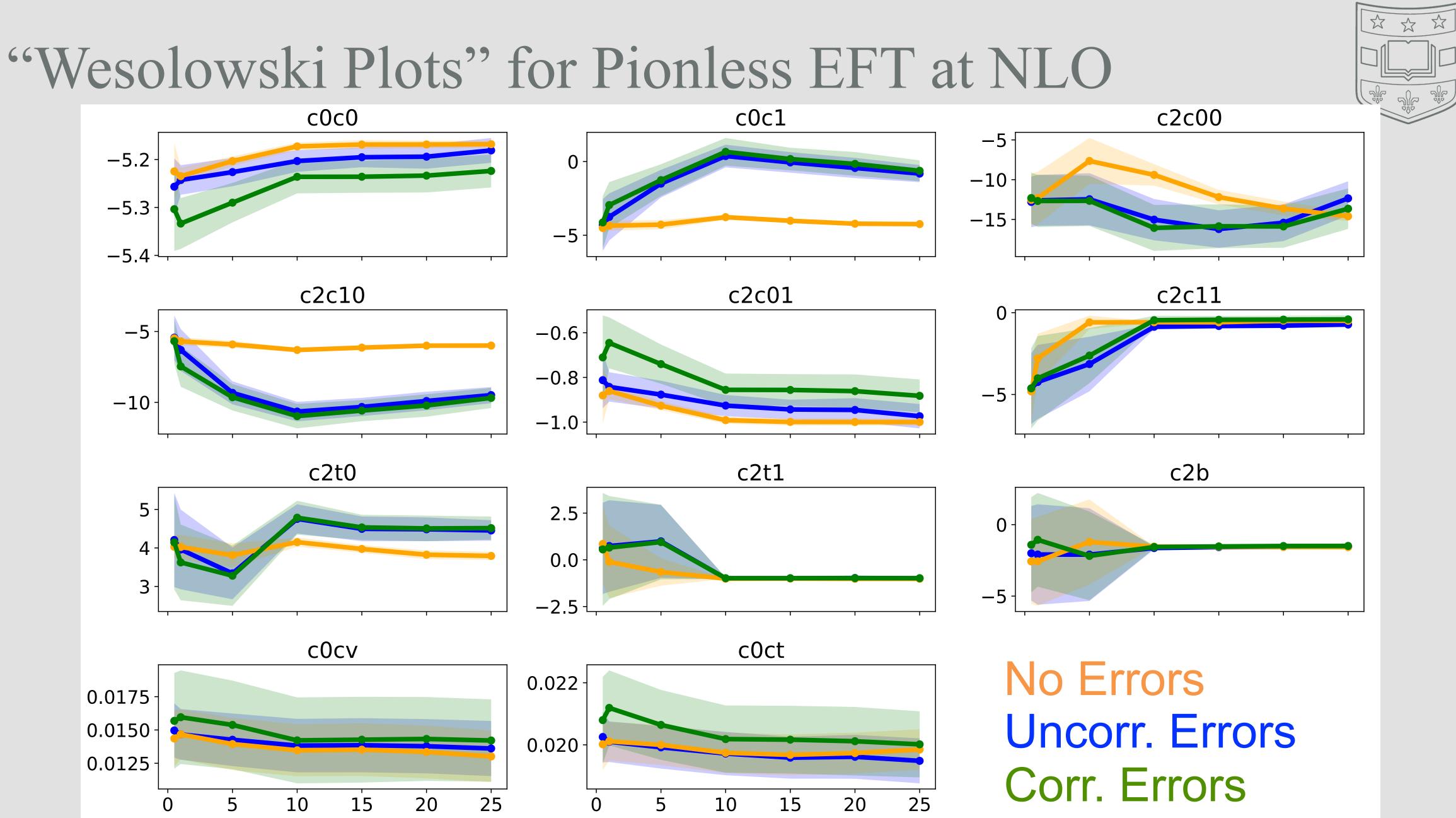
$$d_M(\vec{a}) = \left(\vec{y} - \vec{t}(\vec{a})\right)^T (\sigma_{\exp}^2 + \sigma_{\text{ther},ij}^2)^{-1} \left(\vec{y} - \vec{t}(\vec{a})\right)$$

Correlations on data introduces strong degeneracies in the covariance matrix, so we use Gaussian processes to smooth the correlations.

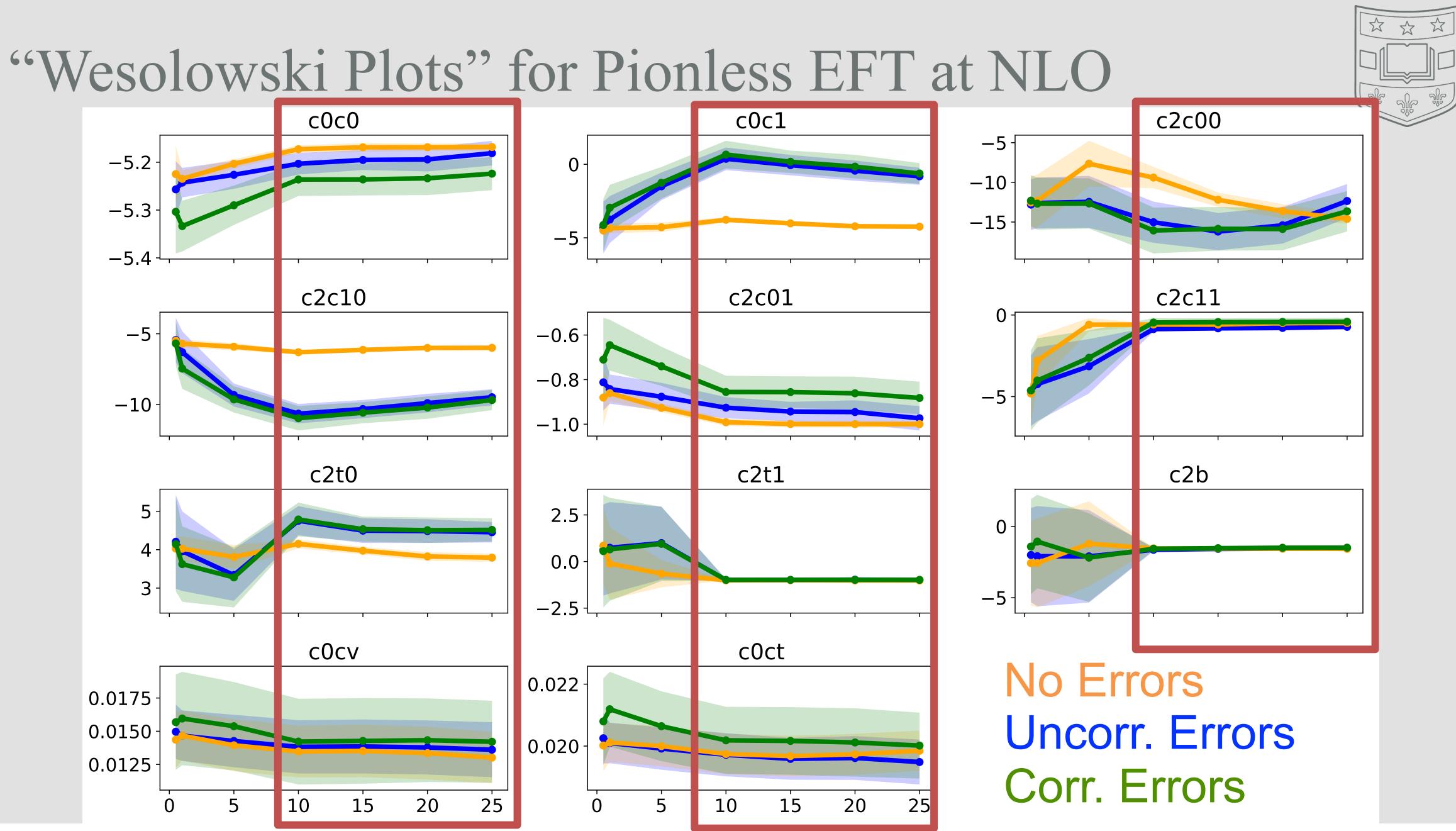














Full posterior

scale).

Full posterior for our EFT:

Prior for \overrightarrow{a} Posterior for \overline{c}^2 Posterior for $\Lambda_{\rm h}$ Likelihood for \overrightarrow{a} Total posterior The posterior is found via sampling of $\{\vec{a}, \vec{c}^2, \Lambda_h\}$.



We have two more parameters to estimate in a full Bayesian model calibration: \bar{c}^2 (scale of truncation error) and Λ_h (EFT breakdown

$P(\overrightarrow{a}, \overrightarrow{c}^2, \Lambda_b | \overrightarrow{y}_{exp}, I) \propto P(\overrightarrow{y}_{exp} | \overrightarrow{a}, \Sigma, I) P(\overrightarrow{a} | I) P(\overrightarrow{c}^2 | \Lambda_b, \overrightarrow{a}, I) P(\Lambda_b | \overrightarrow{a}, I)$



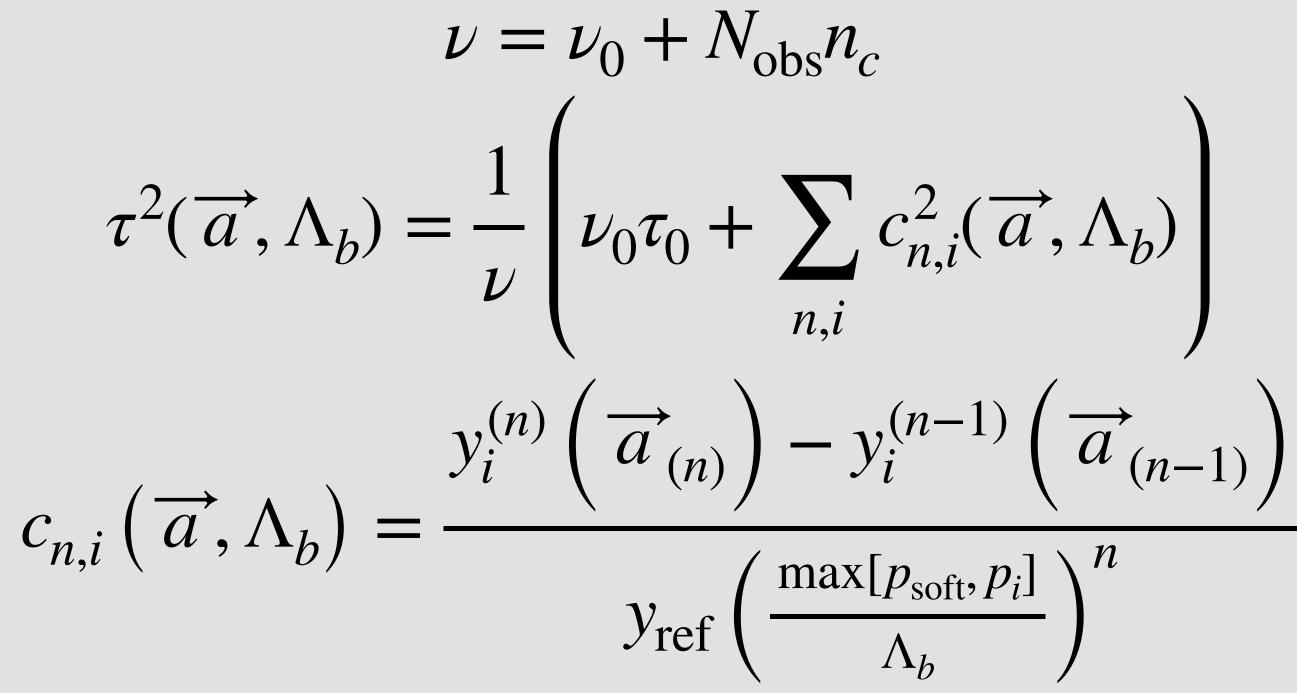
Posterior for \bar{c}^2

For \overline{c}^2 we choose (following Melendez et. al. Phys. Rev. C 100, 044001)

With hyper parameters:



$P(\overline{c}^2 \mid \overrightarrow{a}, \Lambda_b, I) \sim \chi^{-2} \left(\nu, \tau^2(\overrightarrow{a}, \Lambda_b) \right)$





Posterior for Λ_h

Our posterior for the breakdown scale also uses these hyper parameters:

$P(\Lambda_b | \overrightarrow{a}, I) \propto -$

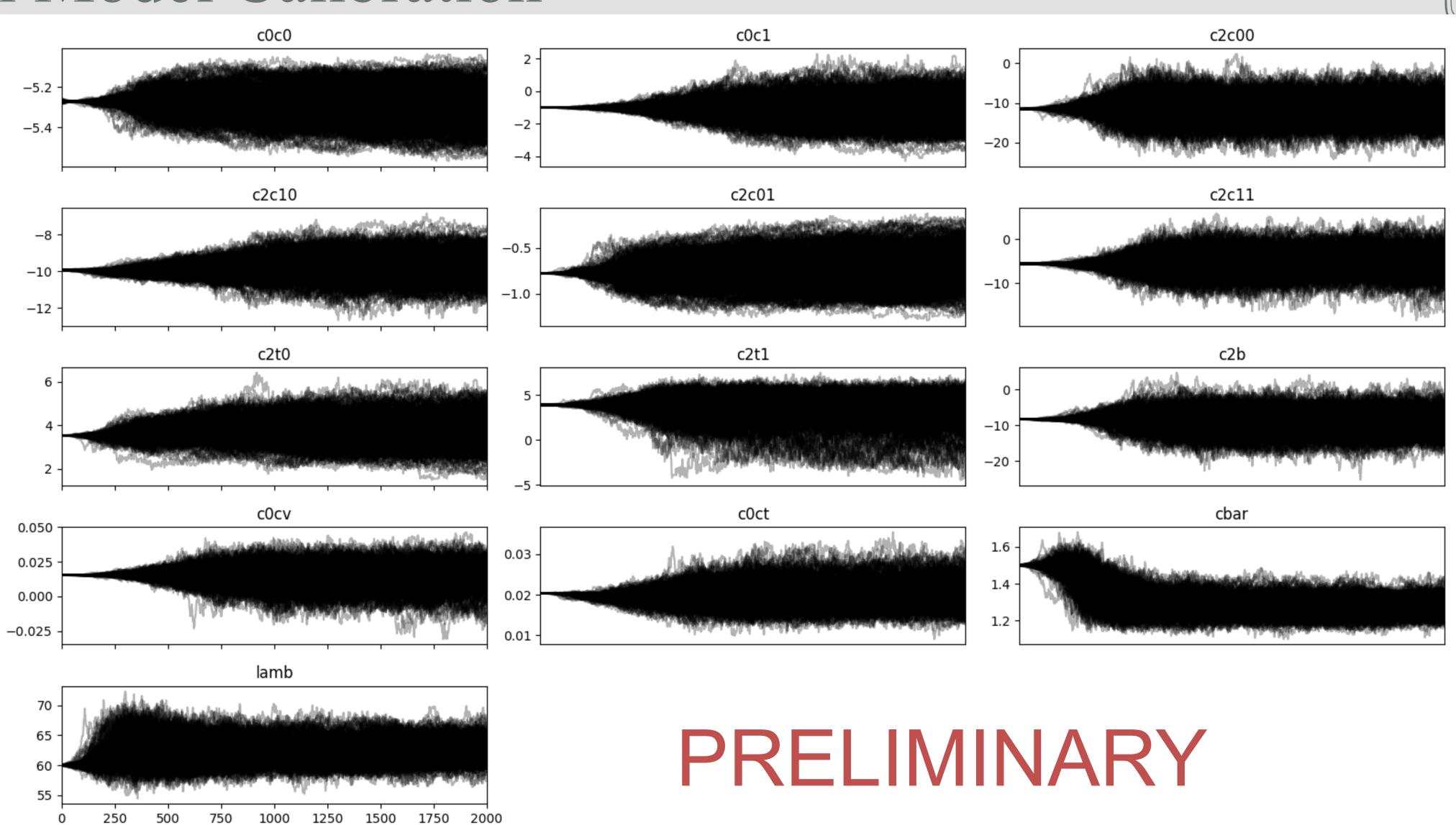
This posterior needs to be numerically normalized as the normalization constant is dependent on \vec{a} .



$$\frac{P(\Lambda_b | I)}{P(\Lambda_b | I)} = \frac{P(\Lambda_b | I)}{\prod_{n,i} \left(\frac{\max[p_{\text{soft}}, p_i]}{\Lambda_b}\right)^n}$$



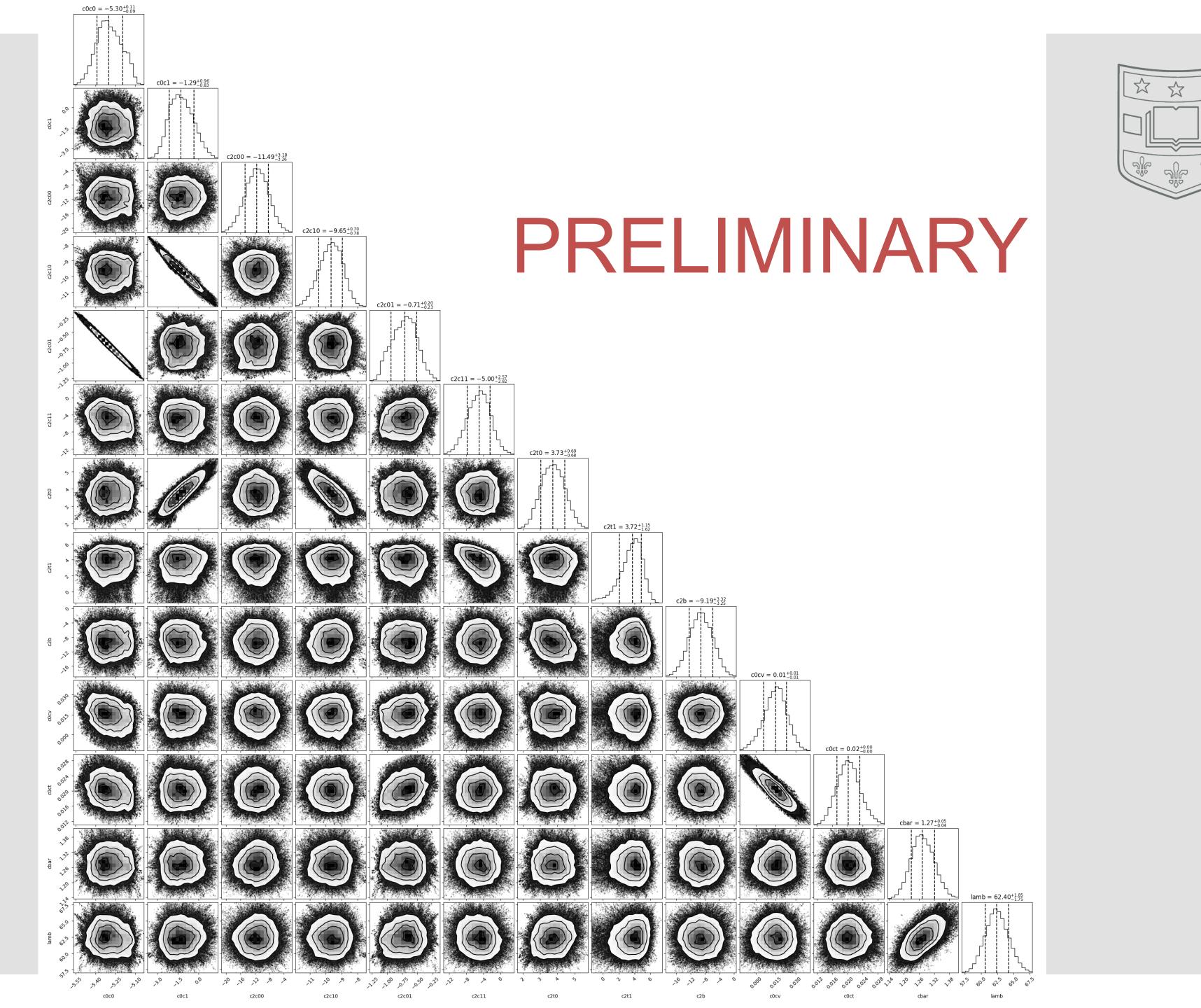
Full Model Calibration





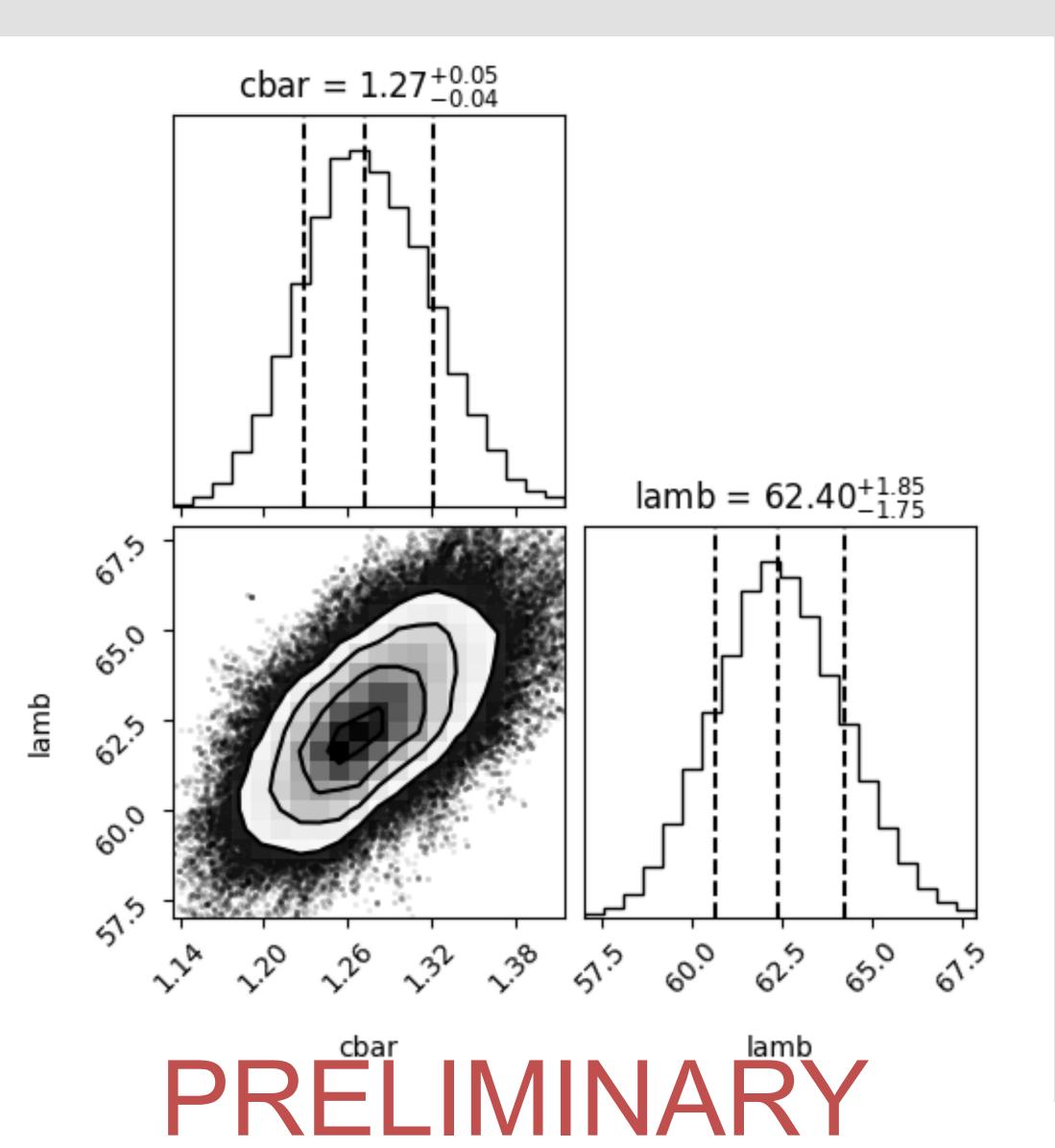
 $\overrightarrow{\Delta}$

Marginal Distributions





Truncation Error Parameters







Future Steps

- Investigate fits at N3LO in the pionless regime
- Emulation for calculation of scattering observables





Ozge Surer Miami University with uncertainty quantification in $\chi EFTs$





Matt Plumlee **Stefan Wild** LBNL Northwestern Include other degrees of freedom: π 's and Δ 's \rightarrow A new set of interactions



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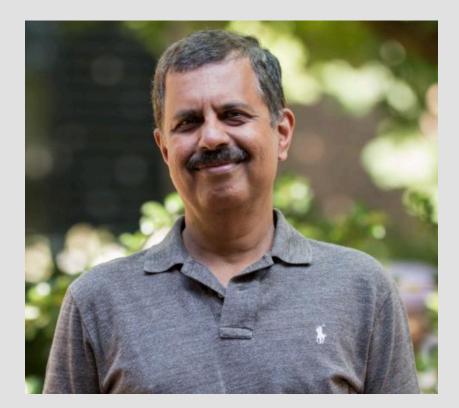


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Acknowledgements

<u>QMC@WashU</u> Piarulli (PI), Pastore (PI), Andreoli (PD), McCoy (PD), King (GS), Chambers-Wall (GS)

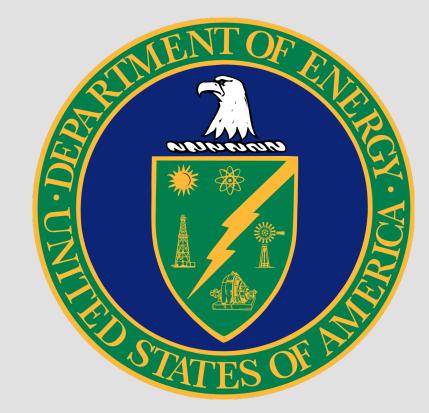


Sai Iyer WashU

Computational Resources







Fellowship/Travel

