

# Quantum sensor applications to (low and high)-energy physics

M. Doser, CERN

# ① Some words on the landscape

Clarification of terms

Quantum sensors for low energy particle physics

Quantum sensors for low energy new particle physics experiments

# ② Quantum detectors for high energy particle physics

quantum sensors register a change of quantum state caused by the interaction with an external system:

- transition between superconducting and normal-conducting
- transition of an atom from one state to another
- change of resonant frequency of a system (quantized)

**Then, a "quantum sensor" is a device, the measurement (sensing) capabilities of which are enabled by our ability to manipulate and read out its quantum states.**

and because the commensurate energies are very low, unsurprisingly, quantum sensors are ideally matched to low energy (particle) physics;

→ focus on activities both in low energy and high energy particle physics

(I will not however be talking about entanglement and its potential applications)

# quantum sensors & particle physics: what are we talking about?

## quantum technologies

1 superconducting devices (TES, SNSPD, ...) / cryo-electronics

2 spin-based, NV-diamonds

3 optical clocks

4 ionic / atomic / molecular

5 optomechanical sensors

6 metamaterials, 0/1/2-D materials

## domains of physics

search for NP / BSM

Axions, ALP's, DM & non-DM  
UL-particle searches

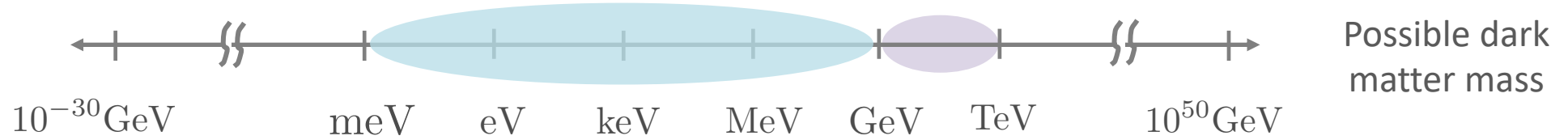
tests of QM

wavefunction collapse,  
decoherence

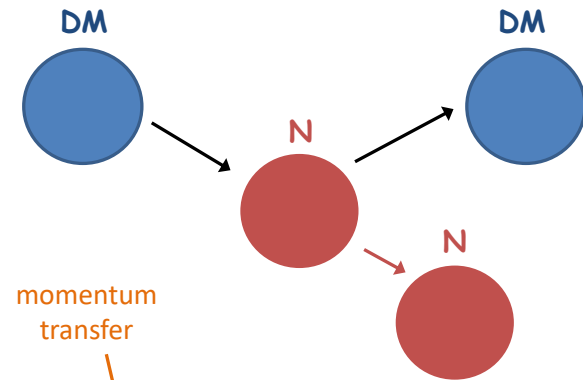
EDM searches & tests of  
fundamental symmetries

# An example: searching for dark matter

from: Yonit Hochberg, Ascona / 2023  
<https://indico.cern.ch/event/1198655>



Looking for nuclear recoils:  
think billiard balls

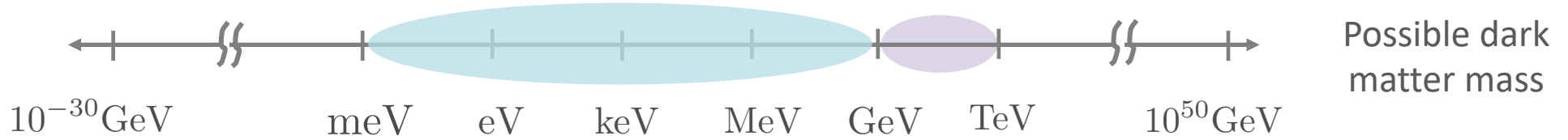


momentum transfer

$$E_{\text{NR}} = \frac{q^2}{2m_N} = \frac{(m_{\text{DM}}v)^2}{2m_N} \gtrsim E_{\text{threshold}} \sim \text{keV}$$

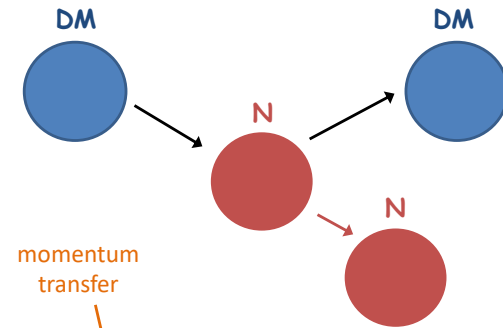
# An example: searching for dark matter

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Possible dark matter mass

Looking for nuclear recoils:  
think billiard balls



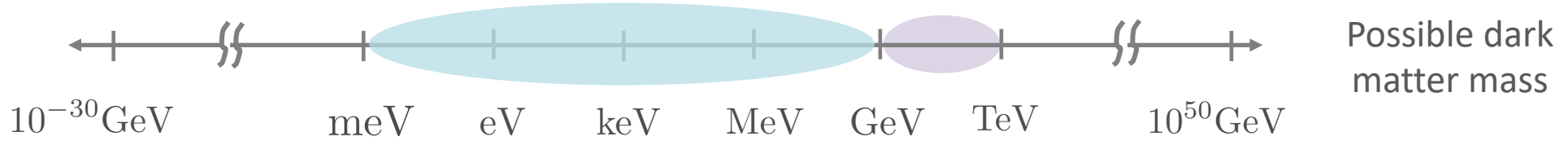
$$E_{\text{NR}} = \frac{q^2}{2m_N} = \frac{(m_{\text{DM}}v)^2}{2m_N} \gtrsim E_{\text{threshold}} \sim \text{keV}$$

light dark matter  
doesn't have enough punch  
to kick the heavy nuclei

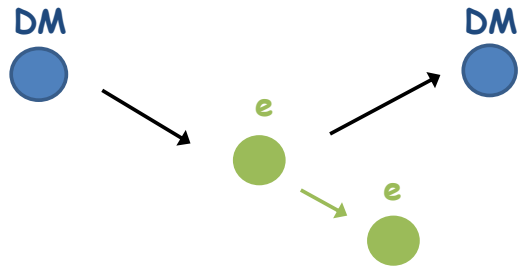
Lose sensitivity @ O(GeV) masses

# An example: searching for dark matter

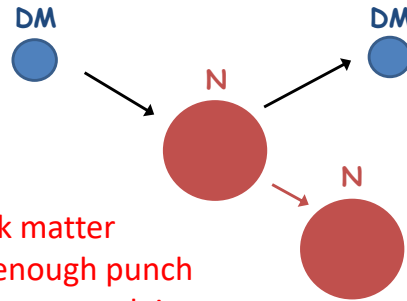
from: Yonit Hochberg, Ascona / 2023  
<https://indico.cern.ch/event/1198655>



Light dark matter: scatter off electrons!



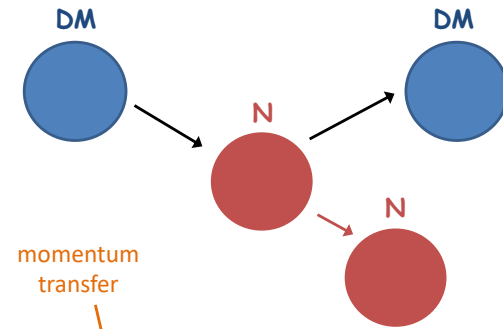
light dark matter can give enough punch to kick the light electrons



light dark matter doesn't have enough punch to kick the heavy nuclei

Lose sensitivity @ O(GeV) masses

Looking for nuclear recoils:  
think billiard balls

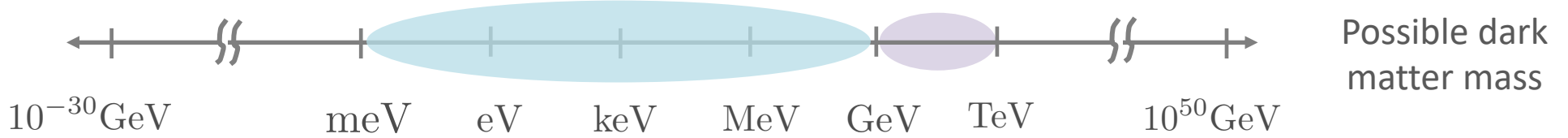


momentum transfer

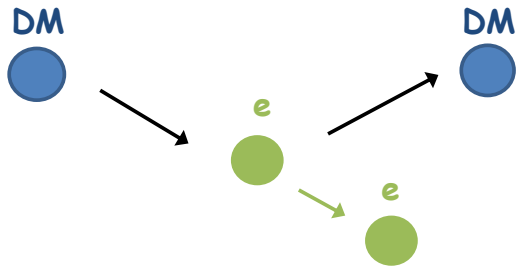
$$E_{\text{NR}} = \frac{q^2}{2m_N} = \frac{(m_{\text{DM}}v)^2}{2m_N} \gtrsim E_{\text{threshold}} \sim \text{keV}$$

# An example: searching for dark matter

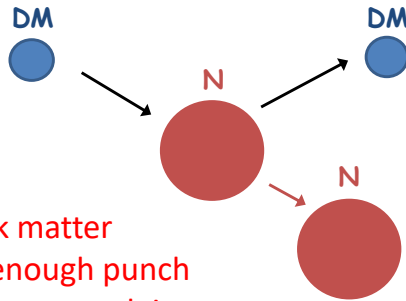
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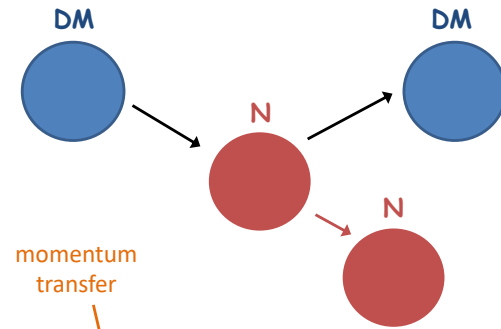


light dark matter can give enough punch to kick the light electrons



light dark matter doesn't have enough punch to kick the heavy nuclei

Looking for nuclear recoils: think billiard balls

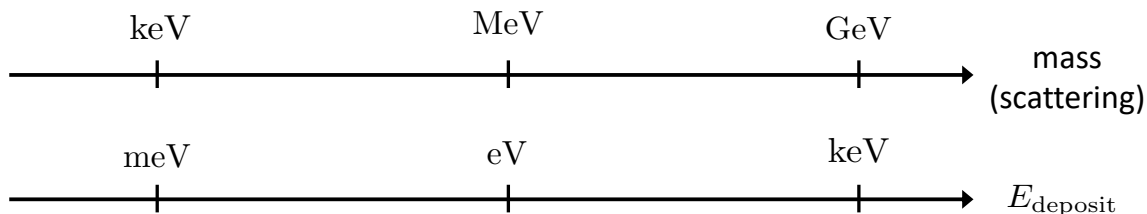


$$E_{NR} = \frac{q^2}{2m_N} = \frac{(m_{DM}v)^2}{2m_N} \gtrsim E_{\text{threshold}} \sim \text{keV}$$

Lose sensitivity @ O(GeV) masses

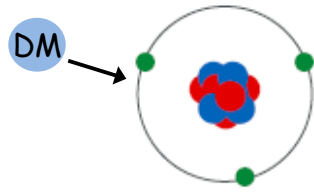
Dark matter scattering: kinetic energy

$$m_{DM}v^2 \sim 10^{-6}m_{DM}$$





## Atomic ionization

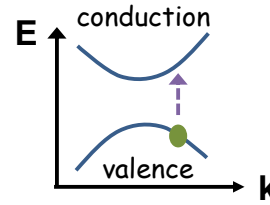


Xenon:  $\sim 12$  eV

$$m_{\text{DM}} \gtrsim 10 \text{ MeV}$$

Essig, Mardon, Volansky, 2012

## Semiconductors



Ge, Si, Diamond, SiC:  $\sim$  eV

$$m_{\text{DM}} \gtrsim \text{MeV}$$

Essig, Mardon, Volansky, 2012  
 Graham, Kaplan, Rajendran, Walters, 2012  
 Kurinsky, Yu, YH, Blas, 2019  
 Griffin, YH, et al, 2020

## Superconducting systems

- Ground state = Cooper pairs; Binding energy (gap)  $\sim$  meV  $\rightarrow$   $m_{\text{DM}} \sim$  keV
- The idea: DM scatters with Cooper pairs, deposits enough energy, breaks Cooper pairs  $\rightarrow$  detect

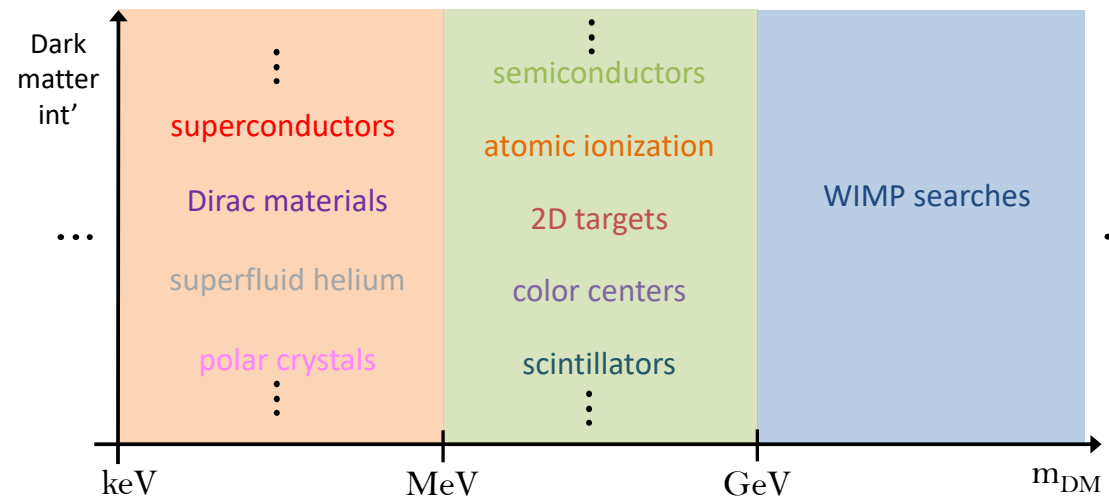
### Excitations

Excitation concentration philosophy

YH, Zhao, Zurek, PRL 2015  
 YH, Pyle, Zhao, Zurek, JHEP 2015

Sensor + target philosophy

YH, Charaev, Nam, Verma, Colangelo, Berggren, PRL 2019



Burgeoning field in recent years

1 superconducting devices (TES, SNSPD, ...) / cryo-electronics

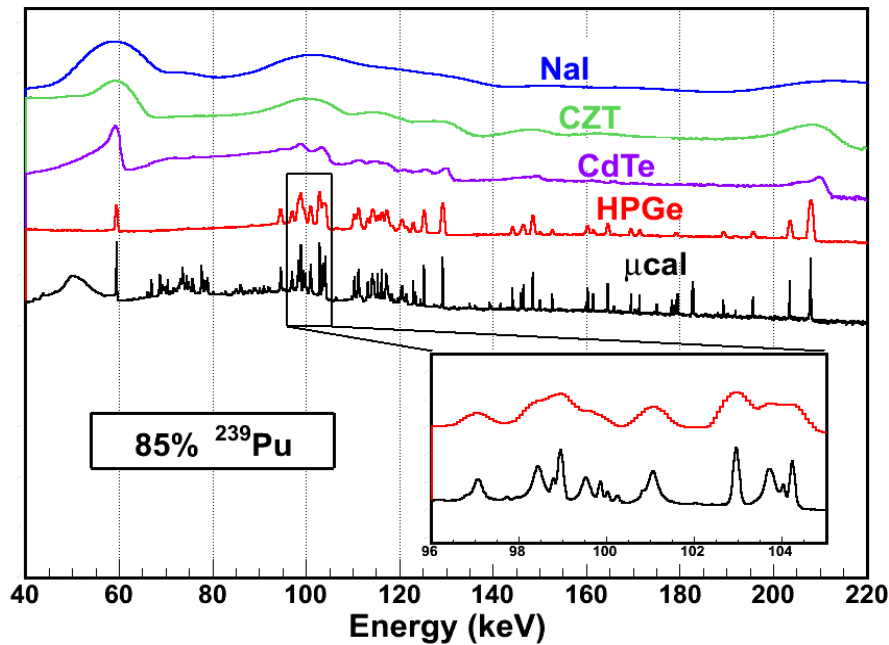
superconducting bolometers in the form of a transition edge sensor (TES) : 1990

invention of kinetic inductance detectors (KID): 2000's

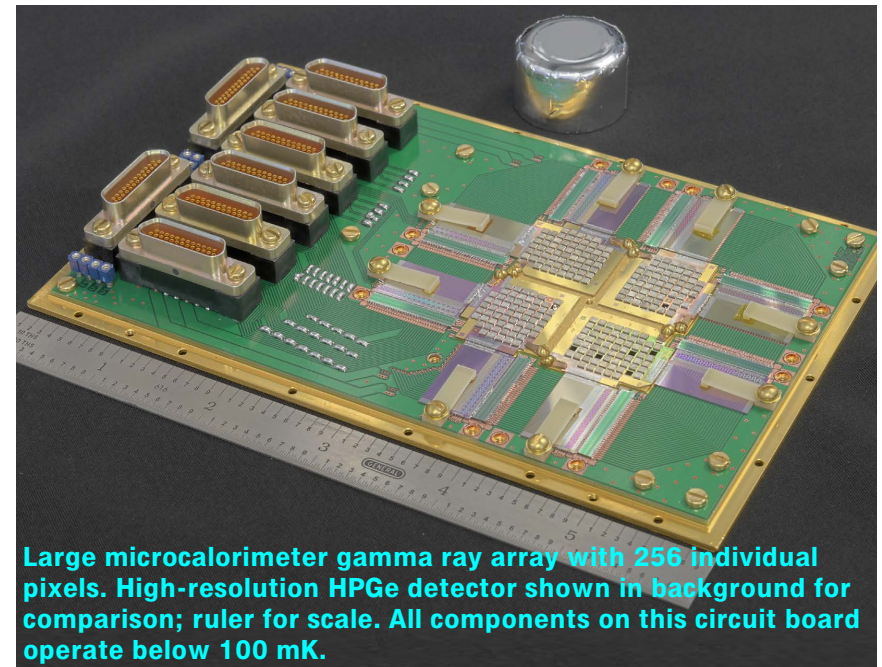
Photons incident on a strip of superconducting material break Cooper pairs and create excess quasiparticles. The kinetic inductance of the superconducting strip is inversely proportional to the density of Cooper pairs, and thus the kinetic inductance increases upon photon absorption. This inductance is combined with a capacitor to form a microwave resonator whose resonant frequency changes with the absorption of photons

P Guss · 2015 — Transition Edge Sensor Using High-Temperature Superconductivity

SDRD FY15 Final Review, Sept. 17, 2015



<https://www.osti.gov/servlets/purl/1506422>



Large microcalorimeter gamma ray array with 256 individual pixels. High-resolution HPGe detector shown in background for comparison; ruler for scale. All components on this circuit board operate below 100 mK.

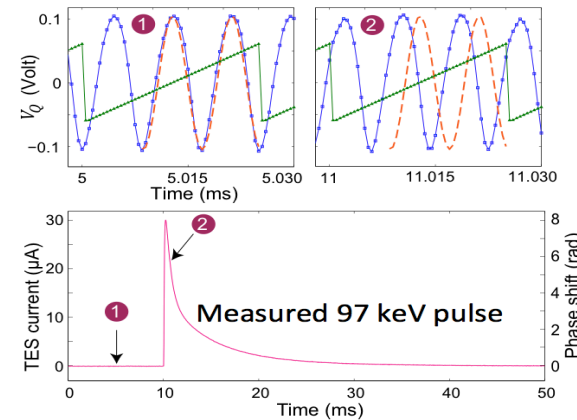
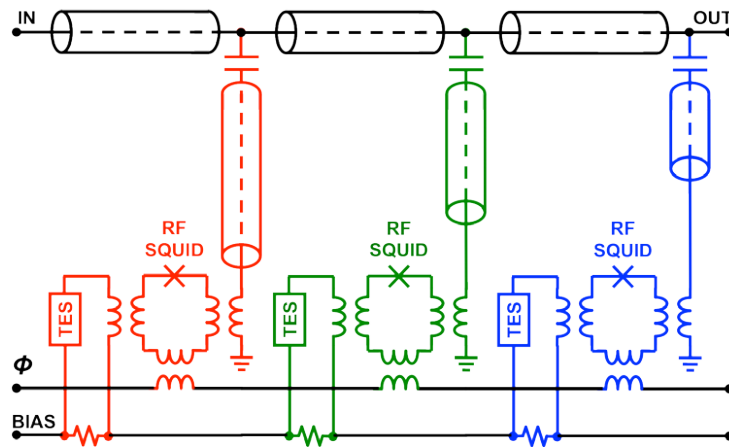
1 superconducting devices (TES, SNSPD, ...) / cryo-electronics

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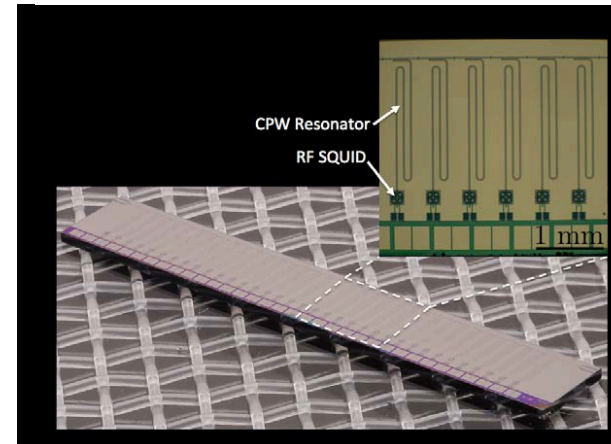
<https://www.osti.gov/servlets/purl/1506422>

Scalable readout is the key for very large arrays

Four innovations:  $\mu$ Cal, Microwave Multiplexing, RF-SQUIDS, Software Radio



- UCSB: 1,000 low-resolution sensors and softwaredefined radio (SDR) readout
- CU: 2 high-resolution sensors + RFSQUIDS without SDR
- CU + LANL: Demonstrated negligible readout noise from RF-SQUID
- *No one has combined all together at scale*



1 superconducting devices (TES, SNSPD, ...) / cryo-electronics

superconducting bolometers in the form of a transition edge sensor (TES) : 1990

invention of kinetic inductance detectors (KID): 2000's

Photons incident on a strip of superconducting material break Cooper pairs and create excess quasiparticles. The kinetic inductance of the superconducting strip is inversely proportional to the density of Cooper pairs, and thus the kinetic inductance increases upon photon absorption. This inductance is combined with a capacitor to form a microwave resonator whose resonant frequency changes with the absorption of photons

introduction of the travelling wave parametric amplifier (TWPA) for quantum noise limited coherent amplification of large bandwidth: 2010

Superconducting qubits for quantum computing, superconductor/spin-system quantum memory elements: 2020

Superconducting device physics is a whole technology, based on thin film deposition techniques and lithographic patterning and allowing large scale integration with high degree of functionality

Superconducting nanowire single photon detector

Superconducting quantum interference device

Josephson junction parametric amplifier

	Microwave	Submillimetre	Far infrared	Optical	High energy
	10 – 100 GHz 3 cm- 3 mm	100 GHz – 1 THz 3 mm – 300 μm	1 – 10 THz 300 – 30 μm	2 μm – 300 nm	UV, Yray and Xray
SIS mixers		•			
HEB			•		
CEB		•			
TES	•	•	•	•	•
KID	•	•	•	•	
SNSPD			•	•	
SQUID	•				
JJPA	•				
TWPA	•	•			

1 superconducting devices (TES, SNSPD, ...) / cryo-electronics

superconducting bolometers in the form of a transition edge sensor (TES) : x-ray imaging

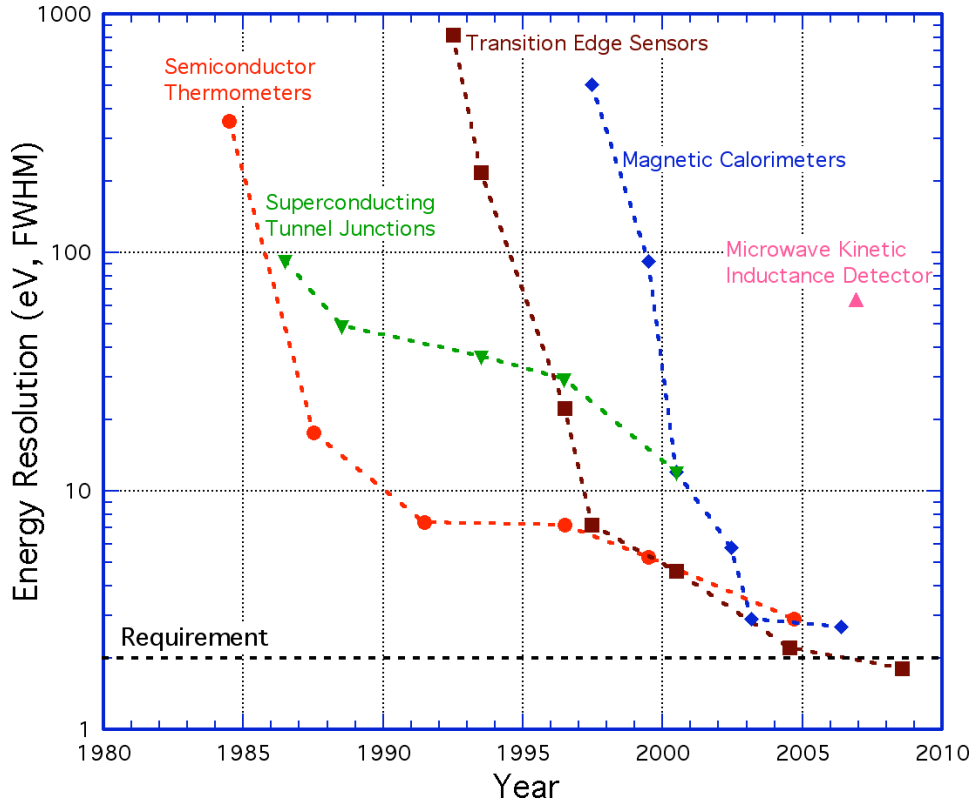
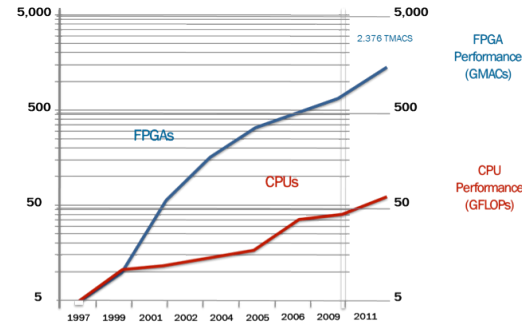
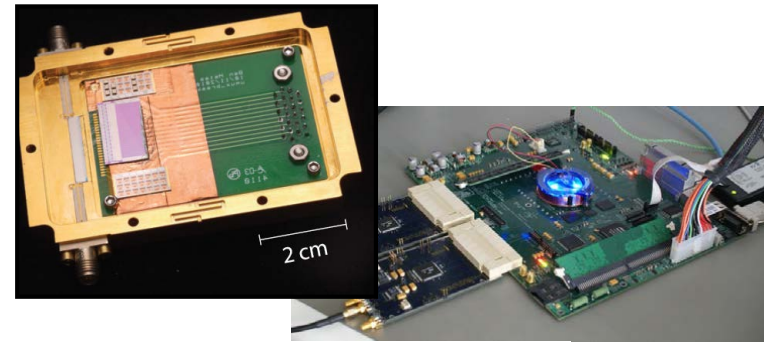


Fig. 1. Time evolution of energy resolution at 6 keV as a function of time for three microcalorimeter technologies.

Development of Low-Temperature Detectors For Generation-X and Other Missions Requiring High-Resolution, Large-Format, X-ray Detector Arrays  
 A White Paper submitted to *Electromagnetic Observations from Space (EOS)* Discipline Program, Simon Bandler *et al.*

- 1,000 resonators on two coaxial cables!
- Existing commercial electronics
- Functionality shared with telecommunications industry
- Benefit from steady improvement in FPGAs, ADCs, and DACs



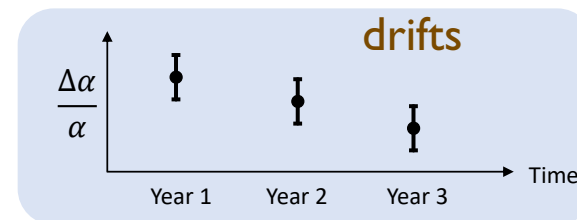
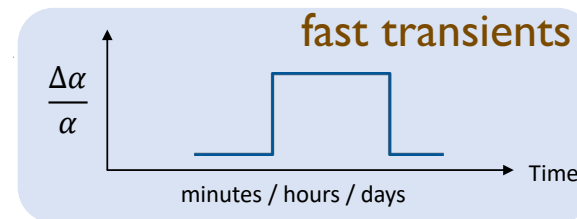
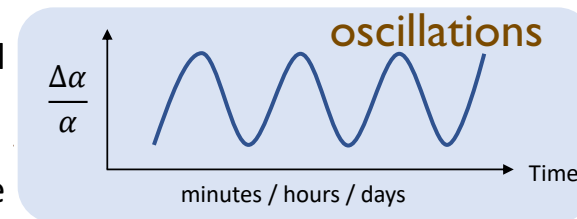
2 spin-based, NV-diamonds, ...

3 optical clocks

search for NP / BSM

Signal characteristics

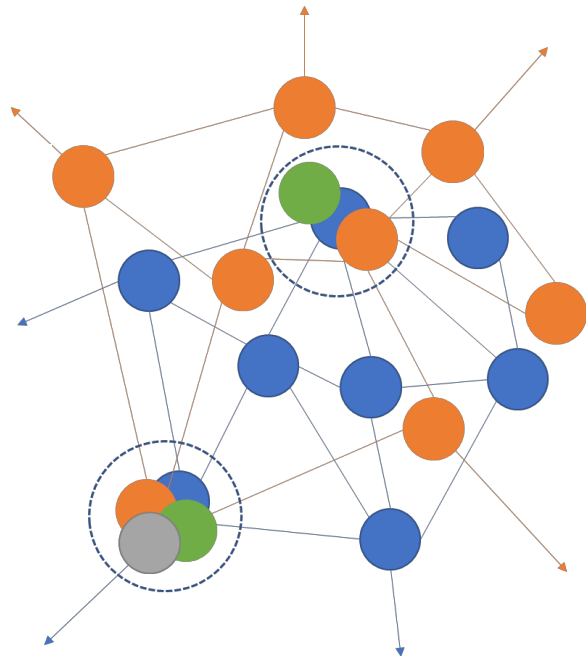
- The only possibility of detecting **transient** events such as topological defects, solitons, Q balls and dark stars
- **Oscillations** of dark matter fields at different locations as long as the distance is below the coherence length (100 km: mass  $\sim 10^{-9}$  eV)
- Sensors with **similar sensitivities and different systematics** are necessary to confirm any measurements and reject false positives
- Using multiple sensors increases the detection confidence and sensitivity
- **Multimessenger** detection, discriminating between different couplings



# particle physics: what are we talking about? Networked experiments!

search for NP / BSM

networks of sensors



magnetometers

2

Afach et al, arXiv:2102.13379v2

atomic clocks

nuclear, HCl,  
molecules

3

Wcislo et al, Sci. Adv. 4, 4869 (2018)

optical fiber networks

3

Roberts et al, New J. Phys. 22, 093010 (2020)

Investigate very light scalar and pseudo-scalar DM candidates over  $\sim 10$  orders of magnitude in mass and different couplings

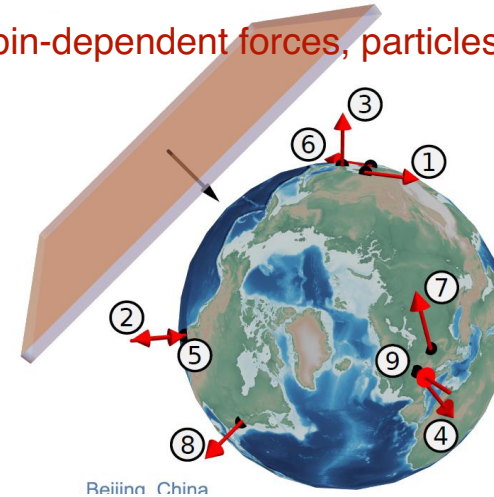
2 spin-based, NV-diamonds, ...



**GNOME:** arXiv:2102.13379 (2021)

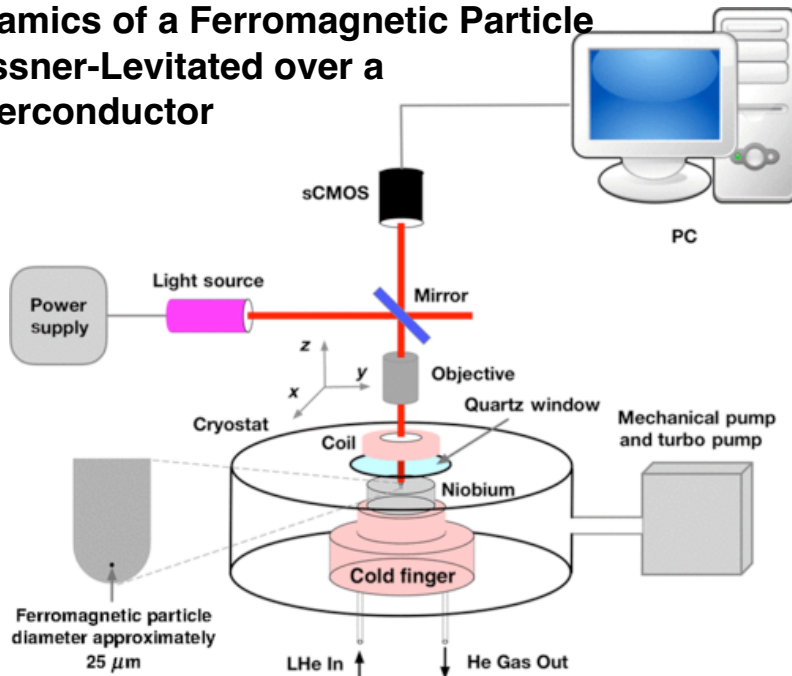
Global Network of **Optical (atomic) Magnetometers** for Exotic searches

→ Spin-dependent forces, particles coupling to spin

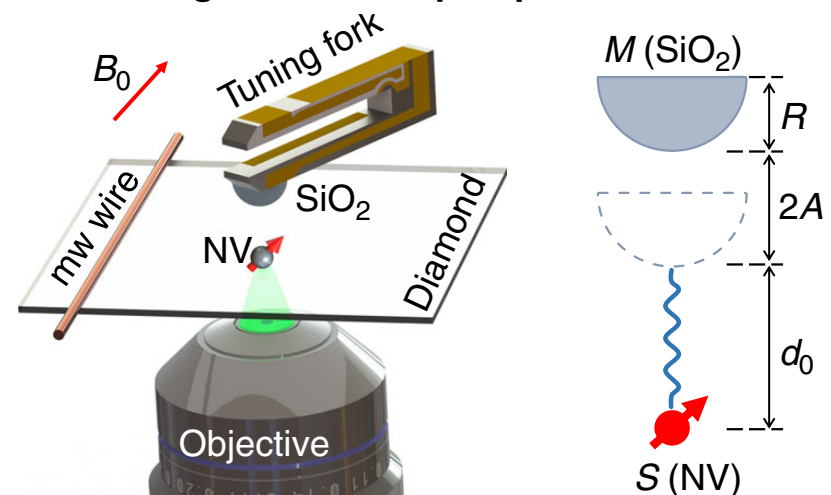


ALP's can interact with atomic spins; passage of an ALP domain wall → simultaneous / correlated signals ~ **transient magnetic field pulse** (down to  $O(pT)$ )

Dynamics of a Ferromagnetic Particle Meissner-Levitated over a Superconductor



Searching for an exotic spin-dependent interaction with a single electron-spin quantum sensor



Rong, X., Wang, M., Geng, J. *et al.*, *Nat Commun* **9**, 739 (2018).  
<https://doi.org/10.1038/s41467-018-03152-9>



### 3 optical clocks

- **Atomic clocks** measure with extreme precision atomic and molecular spectra

- Spectroscopy lends itself to measure variations of:

$$\alpha = \frac{1}{4\pi\epsilon_0} \frac{e^2}{\hbar c}$$

$$\mu = \frac{m_p}{m_e}$$

- Different clocks have different sensitivities to variations of  $\alpha$  and  $\mu$

$$\frac{\delta\omega}{\omega} = K_\alpha \frac{\delta\alpha}{\alpha} + K_\mu \frac{\delta\mu}{\mu}$$

- Clocks are “naturally” networked, need to compare at least 2

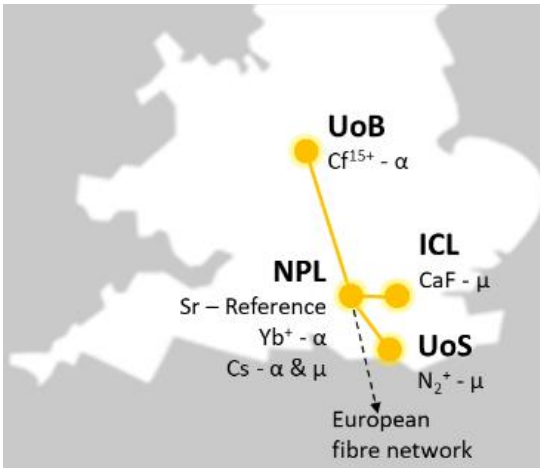
A common, stable, and insensitive frequency reference (a Sr clock), against which all the clocks of the network can measure variations

Clock		$K_\alpha$	$K_\mu$
Highly-charged ion clock	Cf <sup>15+</sup> (775 nm)	59	0
Atomic clock	Yb <sup>+</sup> (467 nm)	-5.95	0
Molecular ion clock	N <sub>2</sub> <sup>+</sup> (2.31 μm)	0	0.5
Molecular clock	CaF (17 μm)	0	0.5
Atomic clock	Sr (698 nm)	0.06	0
	Cs (32.6 mm)	2.83	1

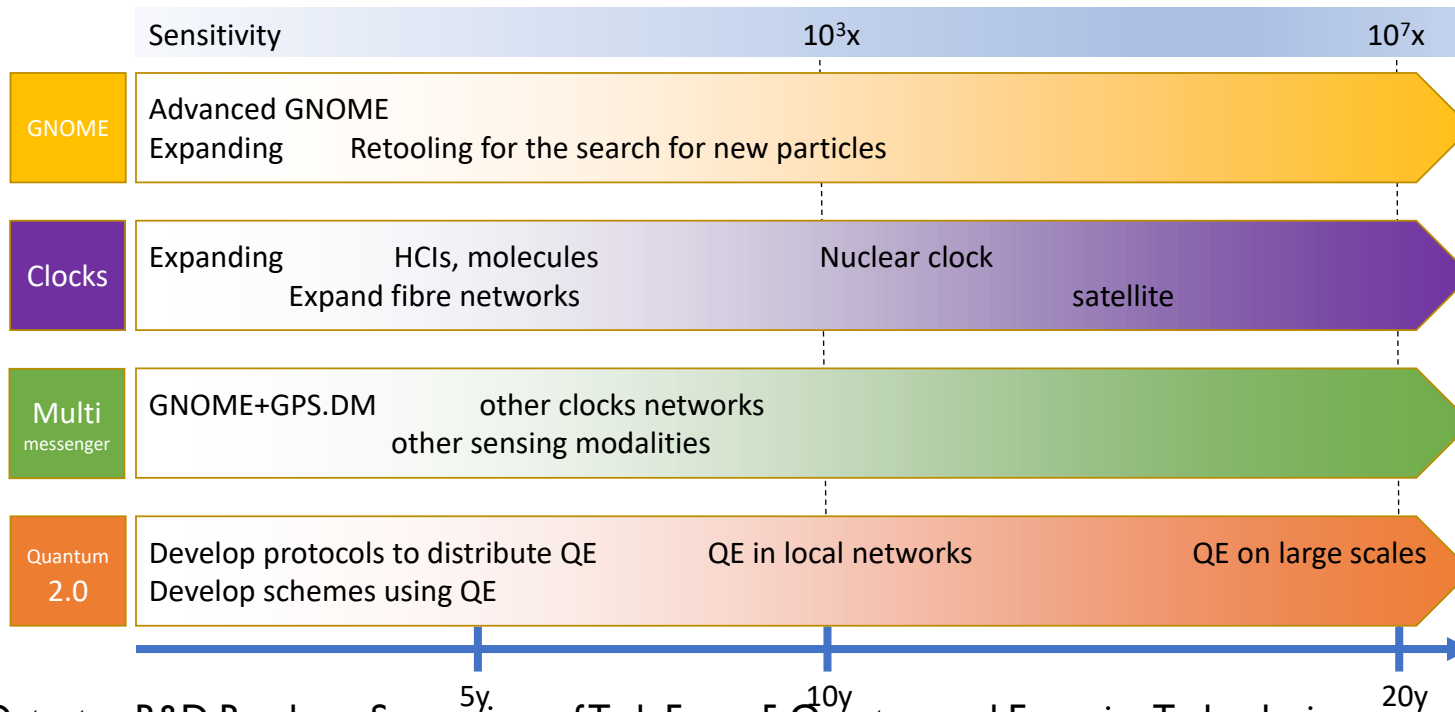
#### Quantum 2.0 (~ 5-10 years from now !)

- Entanglement between sensors can give an advantage when measuring multiple non-commuting parameters of dealing with “nuisance” parameters [PRA 95, 012326 (2017)]
- Creating a super-stable global network of clocks synchronized with entanglement [Nat. Phys. 10, 582 (2014)]
- Need more measurement schemes

3 optical fiber networks



QSNET



ECFA Detector R&D Roadmap Symposium of Task Force 5 Quantum and Emerging Technologies

<https://indico.cern.ch/event/999818/>

Giovanni Barontini (Birmingham)

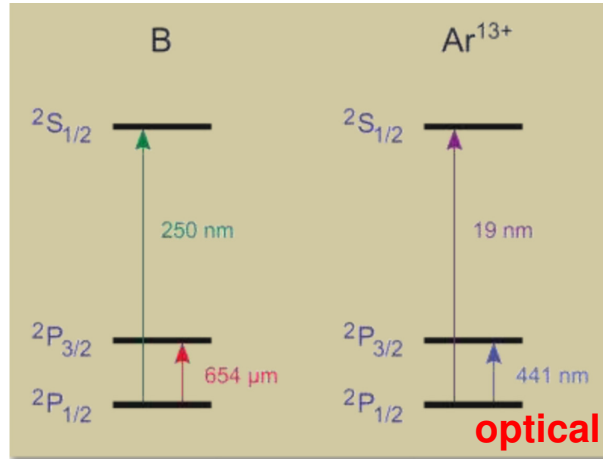
# 4 ionic / atomic / molecular

novel systems with greatly enhanced sensitivity wrt standard atomic systems

## HCI's

e.g.  $^{18}\text{Ar} \rightarrow \text{Ar}^{13+}$   
(5 e<sup>-</sup> remain)

Lorentz violation searches  
variation of fund. constants  
DM searches  
tests of QED  
5th force searches



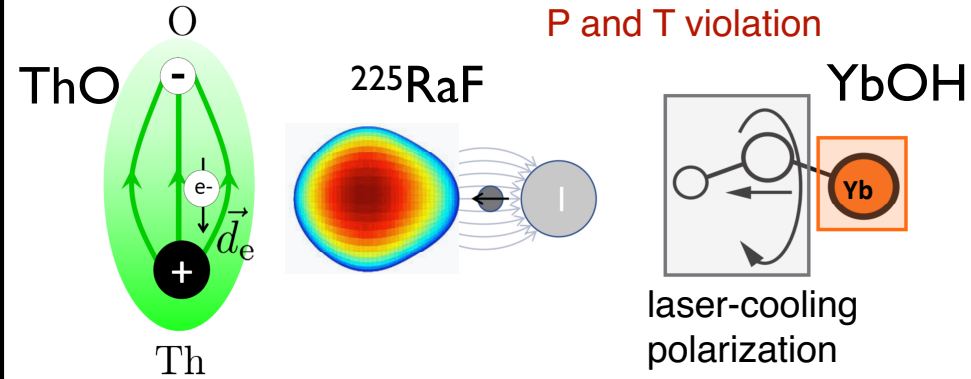
Quantum Sensors for New-Physics Discoveries  
<https://iopscience.iop.org/journal/2058-9565/page/Focus-on-Quantum-Sensors-for-New-Physics-Discoveries>

Search for new physics with atoms and molecules  
Rev. Mod. Phys. 90, 025008 (2018)

Barontini et al. EPJ Quantum Technology (2022) 9:12  
<https://doi.org/10.1140/epjqt/s40507-022-00130-5>

## eEDM's in molecules

P and T violation



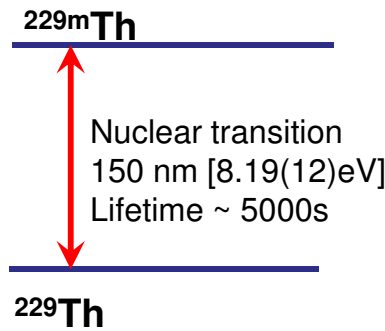
## nuclear clock ( $^{229}\text{Th}$ )

typical nuclear energy levels are in MeV. Six orders of magnitude from ~few eV we can access by lasers!

only ONE exception:

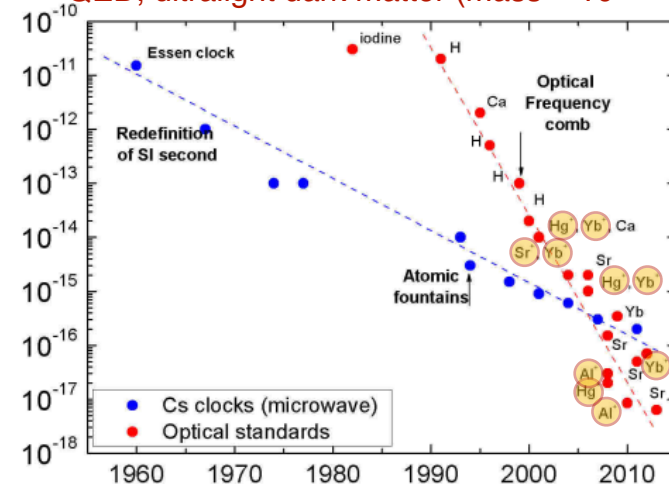
sensitivity to **dark matter**, variation of  $\alpha$   
~ 8.19 eV / MeV ~ 10<sup>5</sup> wrt current clocks

Review: E. Peik, et al., arXiv:2012.09304, in press, Quantum Science and Technology (2021)



## molecular / ion clocks

QED, ultralight dark matter (mass ~ 10<sup>-22</sup> – 10<sup>-15</sup> eV)



Hg<sup>+</sup>  
Al<sup>+</sup>  
Yb<sup>+</sup>  
CaF  
N<sub>2</sub><sup>+</sup>

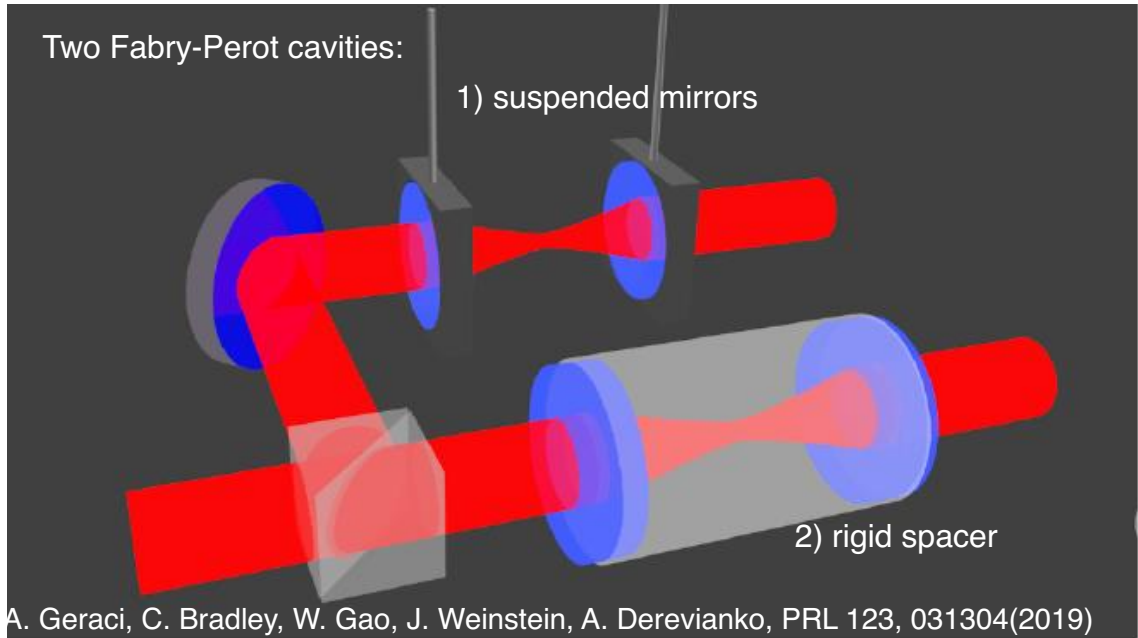
N<sub>2</sub><sup>+</sup> : 10<sup>-18</sup> <https://doi.org/10.1103/PhysRevA.89.032509>

ECFA Detector R&D Roadmap Symposium of Task Force 5 Quantum and Emerging Technologies

<https://indico.cern.ch/event/999818/> Marianna Safronova (University of Delaware) David Hume (NIST)

# 5 optomechanical sensors

- Wavelike DM
  - Scalar-like: **Optical-cavity-based detectors**
  - Vector-like: **Accelerometer detectors**
- Particle/extended-object DM
  - Gravitational Wave detectors for DM (also Axions)
  - Levitated microspheres
  - Windchime experiment



Isotropic strain in material objects due to variation of atomic size from interaction with ultra-light DM Manley et al. PRL 124, 151301 (2020)

$$\phi(t, \mathbf{r}) \approx \frac{\hbar}{m_\phi c} \sqrt{2\rho_{\text{DM}}} \cos [2\pi f_\phi t - \mathbf{k}_\phi \cdot \mathbf{r} + \dots]$$

$k_\phi = m_\phi v / \hbar$   
 $m_\phi$  is the mass of the DM field,  
 $v$  is the relative speed of the DM

$$\frac{\delta m_e(t, \mathbf{r})}{m_{e,0}} = d_{m_e} \sqrt{4\pi \hbar c} E_P^{-1} \phi(t, \mathbf{r})$$

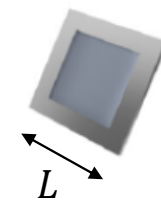
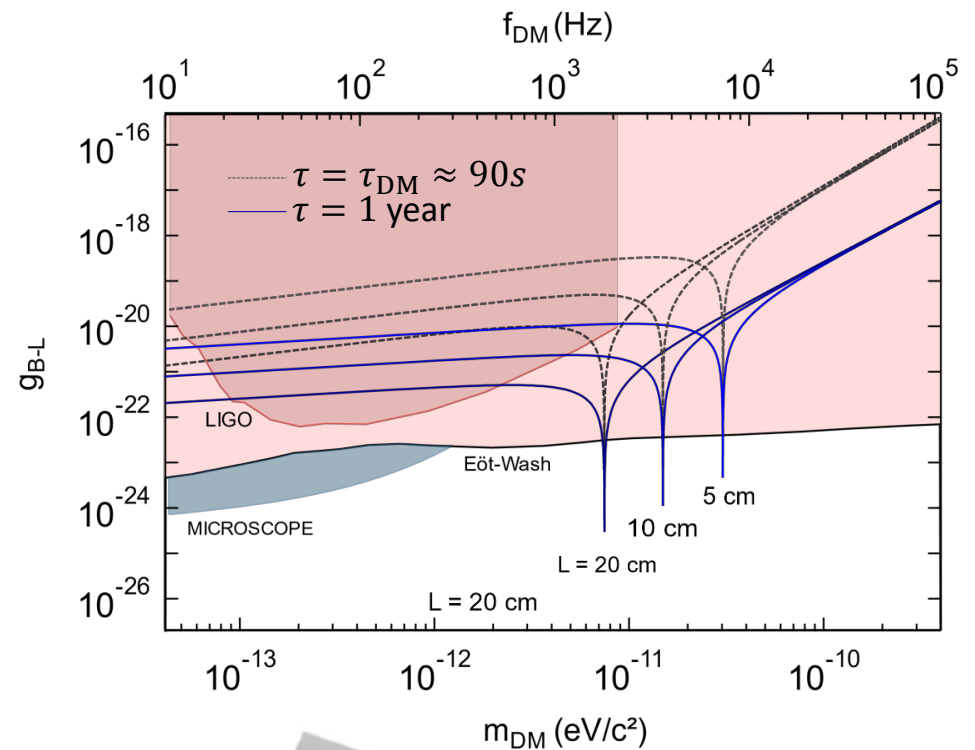
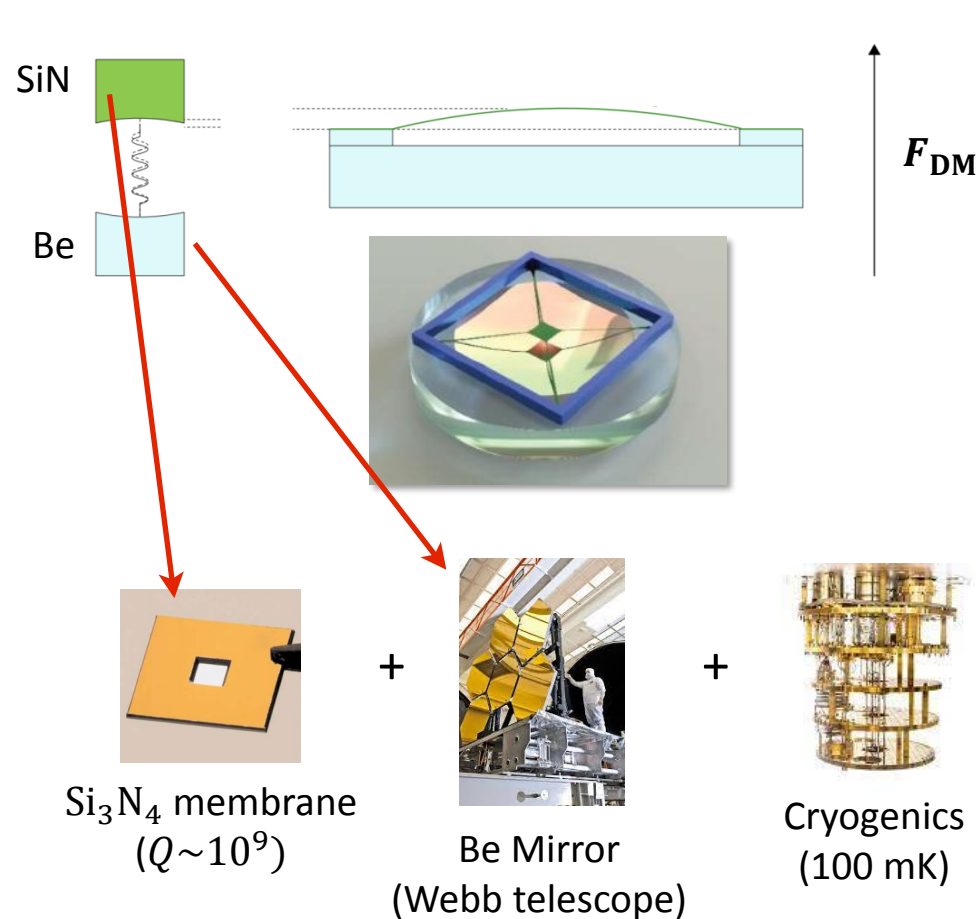
$$\frac{\delta \alpha(t, \mathbf{r})}{\alpha_0} = d_e \sqrt{4\pi \hbar c} E_P^{-1} \phi(t, \mathbf{r})$$

strain :

$$h = -\frac{\delta \alpha}{\alpha_0} - \frac{\delta m_e}{m_{e,0}}$$

**5** accelerometers

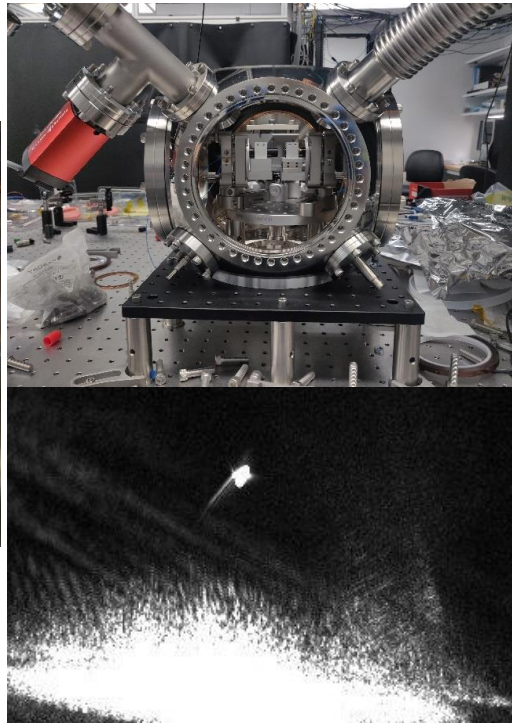
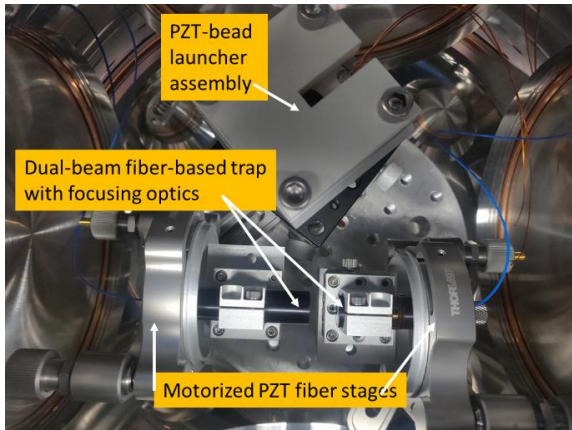
differential acceleration between materials of different composition (probe of B-L couplings)



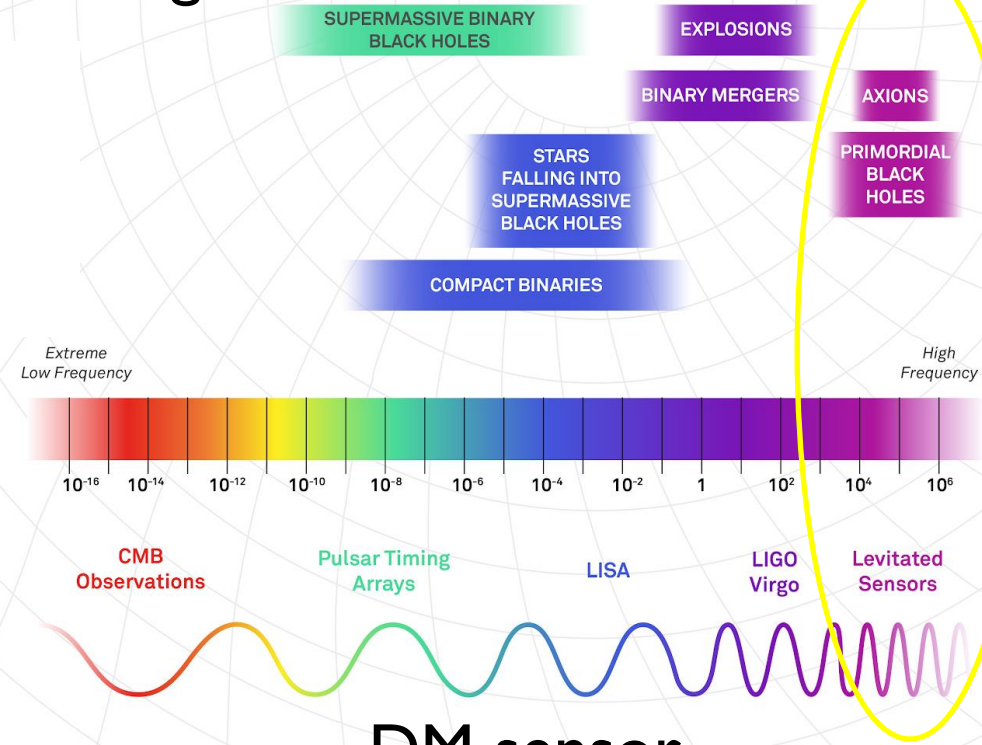
Manley et. al. PRL 126 061301 (2021)

5 levitated microspheres

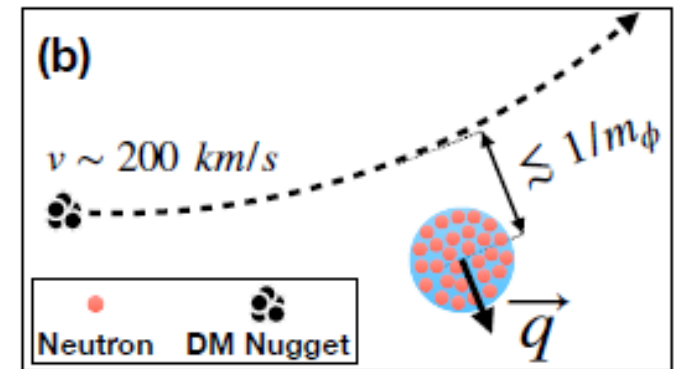
nanoparticle standing-wave trap (optical lattice)



gravitational wave sensor



DM sensor



Search for Composite Dark Matter with Optically Levitated Sensors

Fernando Monteiro, Gadi Afek, Daniel Carney, Gordan Krnjaic, Jiayang Wang, and David C. Moore  
Phys. Rev. Lett. **125**, 181102 – Published 28 October 2020

G. Ranjit, et.al. , *Phys. Rev. A*, **93**, 053801 (2016)  
C. Montoya et. al. arXiv:2103.03420 (2021)

= zeptonewton sensing

ECFA Detector R&D Roadmap Symposium of Task Force 5 Quantum and Emerging Technologies

<https://indico.cern.ch/event/999818/>

Andrew Geraci (Northwestern)

## 5 “Wind chime”

Planck-scale DM: measure the gravitational effect of flying-by DM on an array of accelerometers

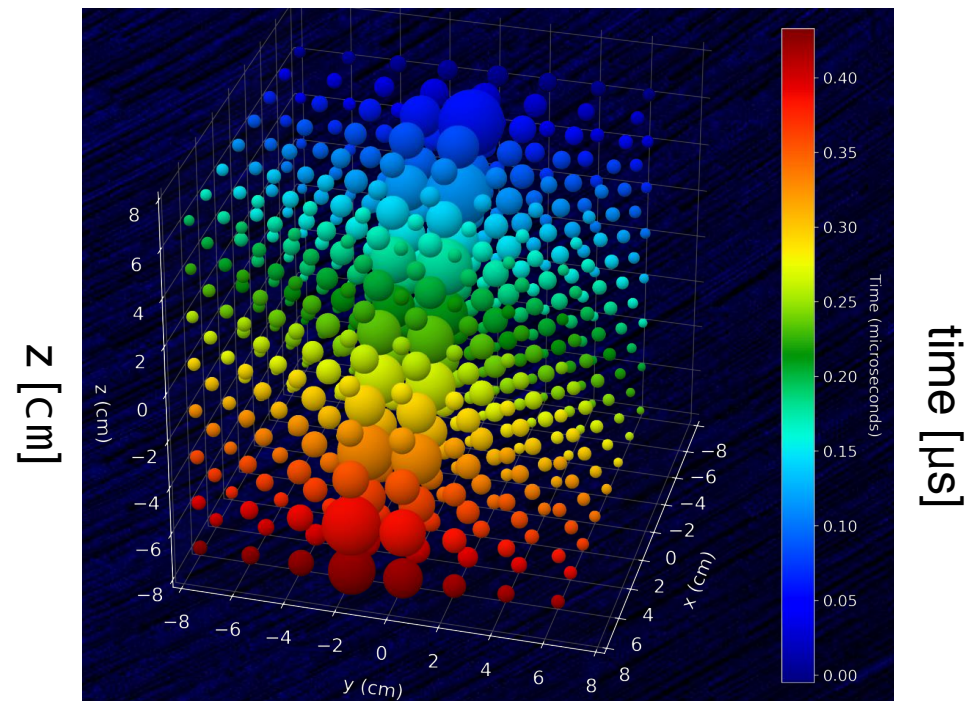
vision: “Our proposed strategy is to build a **three-dimensional array of force sensors**. A heavy DM particle passing through the array will exert a **small but correlated force** on the sensors nearest its trajectory. Much **like tracking a particle in a bubble chamber**, we can then pick out this correlated force signal along the DM “track” through the array.”

<https://arxiv.org/abs/1903.00492>

flux:  $1/\text{m}^2/\text{year}$

“With a **billion** detectors at the gram scale, Planck-scale gravitational DM detection is achievable.”

sensors: “cryogenic opto-mechanical devices”, e.g.: atoms in a lattice that are continuously optically probed



6 metamaterials, 0/1/2-D materia

metamaterial : materials that obtain their properties from their *structure* rather than the material of which they are *composed*

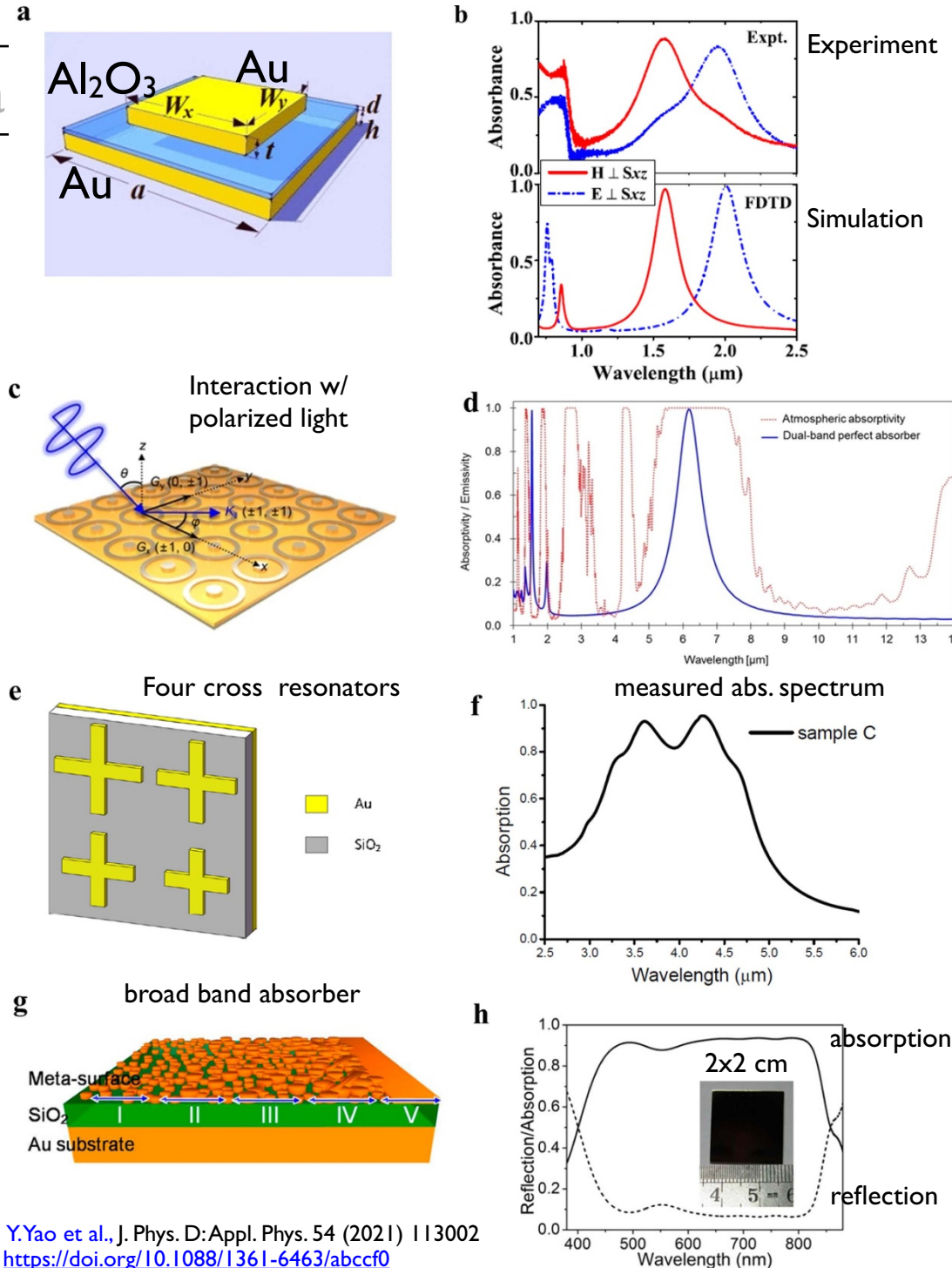
These are *engineered composite materials* mainly consisting of artificially designed periodic sub-wavelength structures.

One particular application revolves around the absorption / reflectance characteristics in the IR in very compact devices.

(relatively) *broadband* absorption:

Importantly: the properties can be designed!

Quantum sensors for low energy particle physics

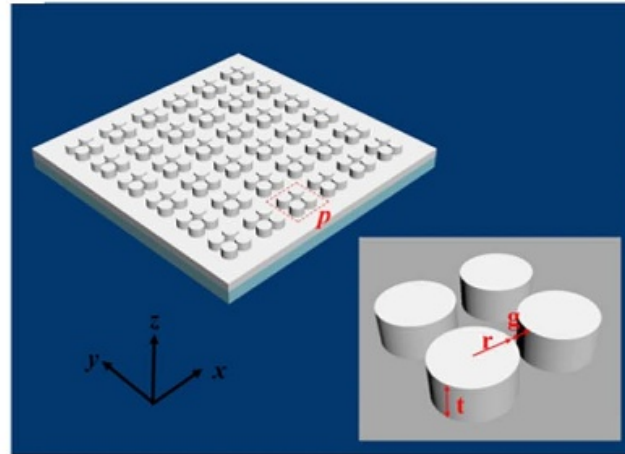


Y.Yao et al., J. Phys. D: Appl. Phys. 54 (2021) 113002  
<https://doi.org/10.1088/1361-6463/abccf0>

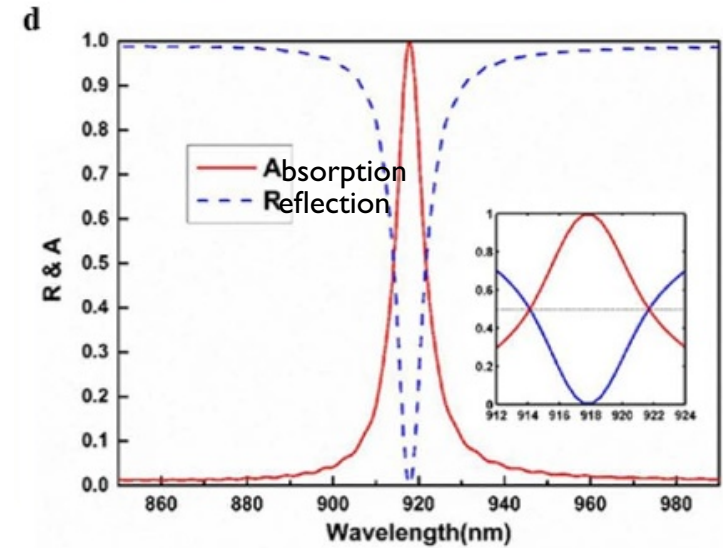


6 metamaterials, 0/1/2-D materials

Schematic of the MMA  
 $r = 80\text{nm}$   
 $t = 100\text{ nm}$   
 $g = 20\text{ nm}$



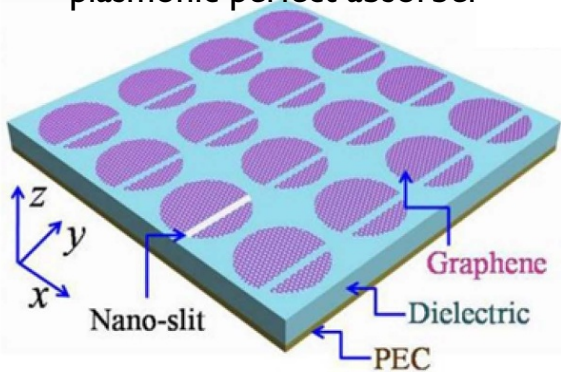
(relatively) *narrowband* absorption:



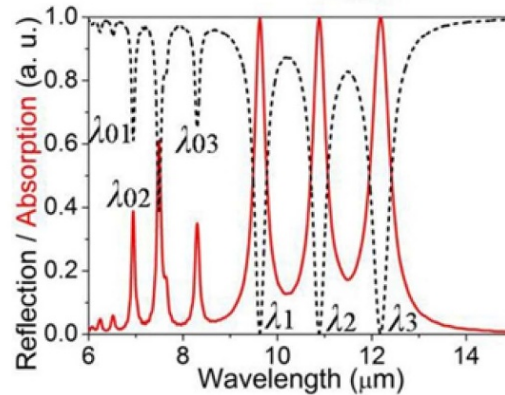
Y.Yao et al., J. Phys. D: Appl. Phys. 54 (2021) 113002 <https://doi.org/10.1088/1361-6463/abccf0>

2D materials based MMAs: strong tunability by tuning gate voltage

Schematic of the graphene plasmonic perfect absorber



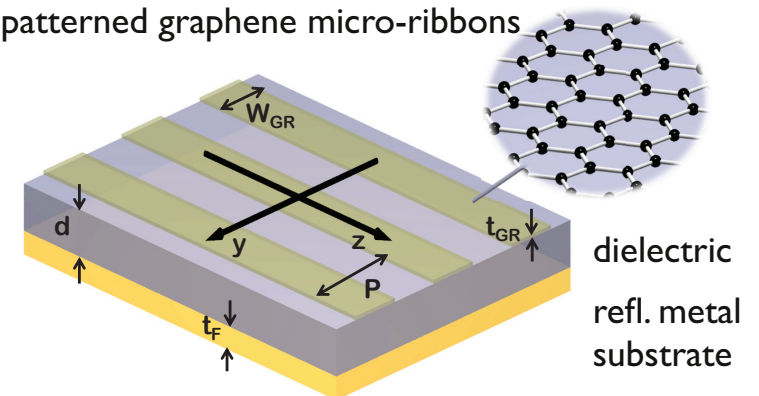
Absorption and reflection of the absorber.



Y.Yao et al., J. Phys. D: Appl. Phys. 54 (2021) 113002 <https://doi.org/10.1088/1361-6463/abccf0>

plasmonic metamaterials: perfect absorption of light

patterned graphene micro-ribbons



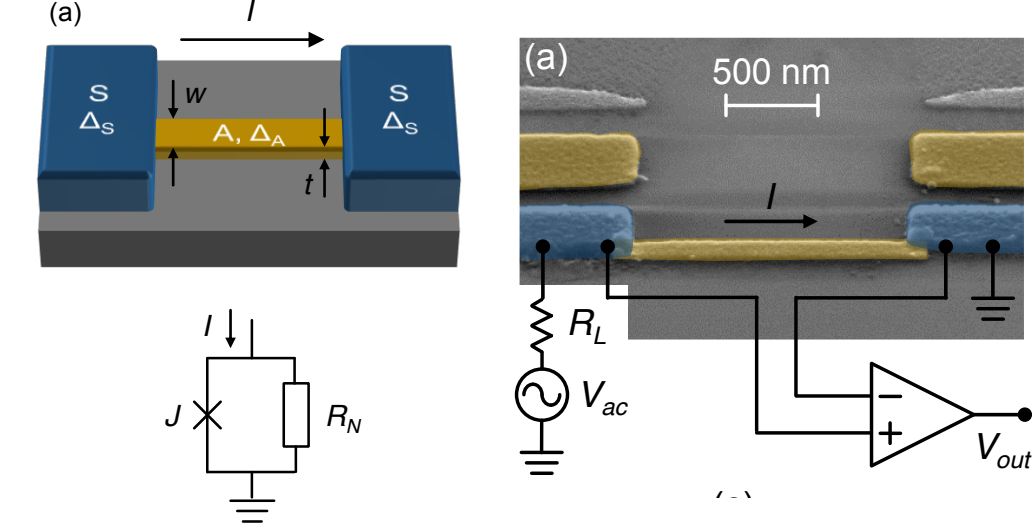
R.Alaei et al., Opt. Express. 20 (2012) 28017 <https://doi.org/10.1364/OE.20.028017>

6 metamaterials, 0/1/2-D materials

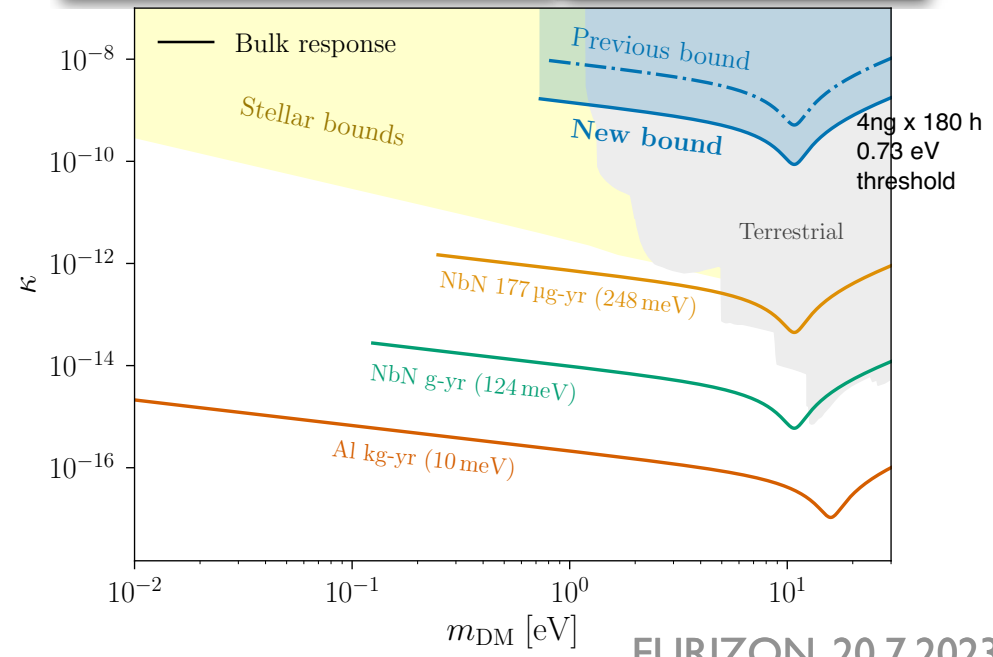
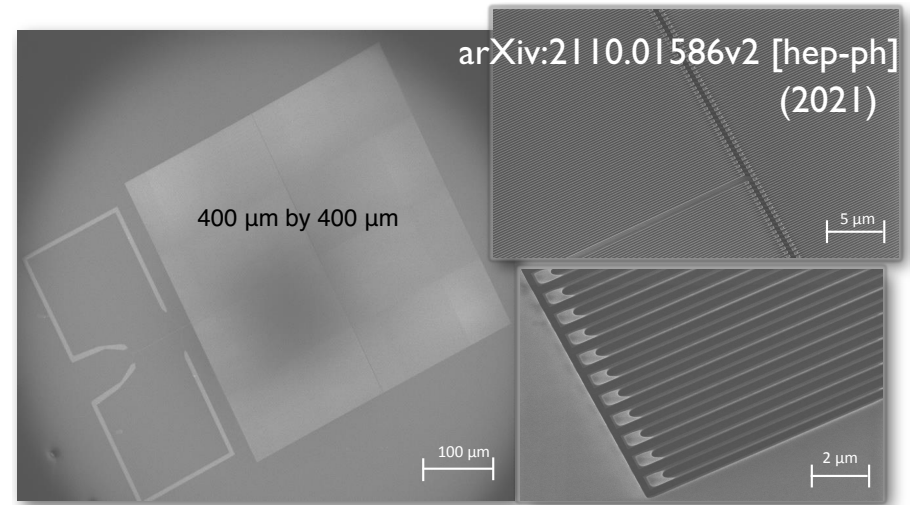
I-D junction: JES: Josephson escape sensor, nanoscale transition edge sensor (nano-TES)

well suited to detecting dark matter / axion interactions with electrons

superconducting single-photon detection arXiv:2101.08558v2



Superconducting nanowire single-photon detectors (SNSPDs): sub-eV energy deposition thresholds



directional detection of light dark matter

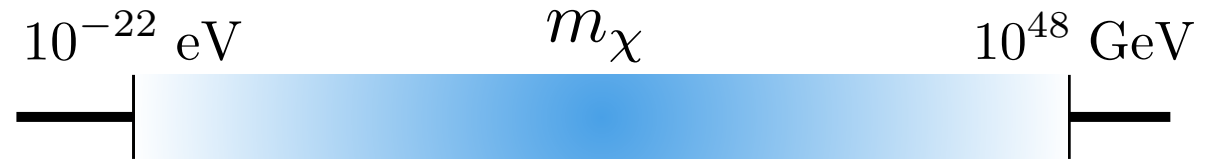
dark matter scattering produces initial excitations with an anisotropic distribution which is preserved during de-excitation

<https://arxiv.org/abs/2109.04473>

Quantum sensors for new  
particle physics experiments

# quantum sensors & particle physics: what are we talking about?

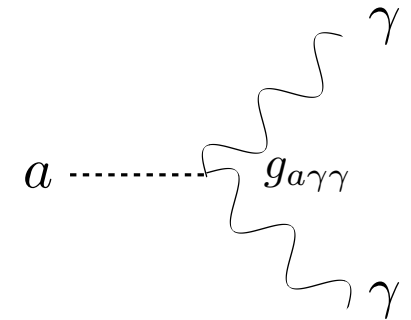
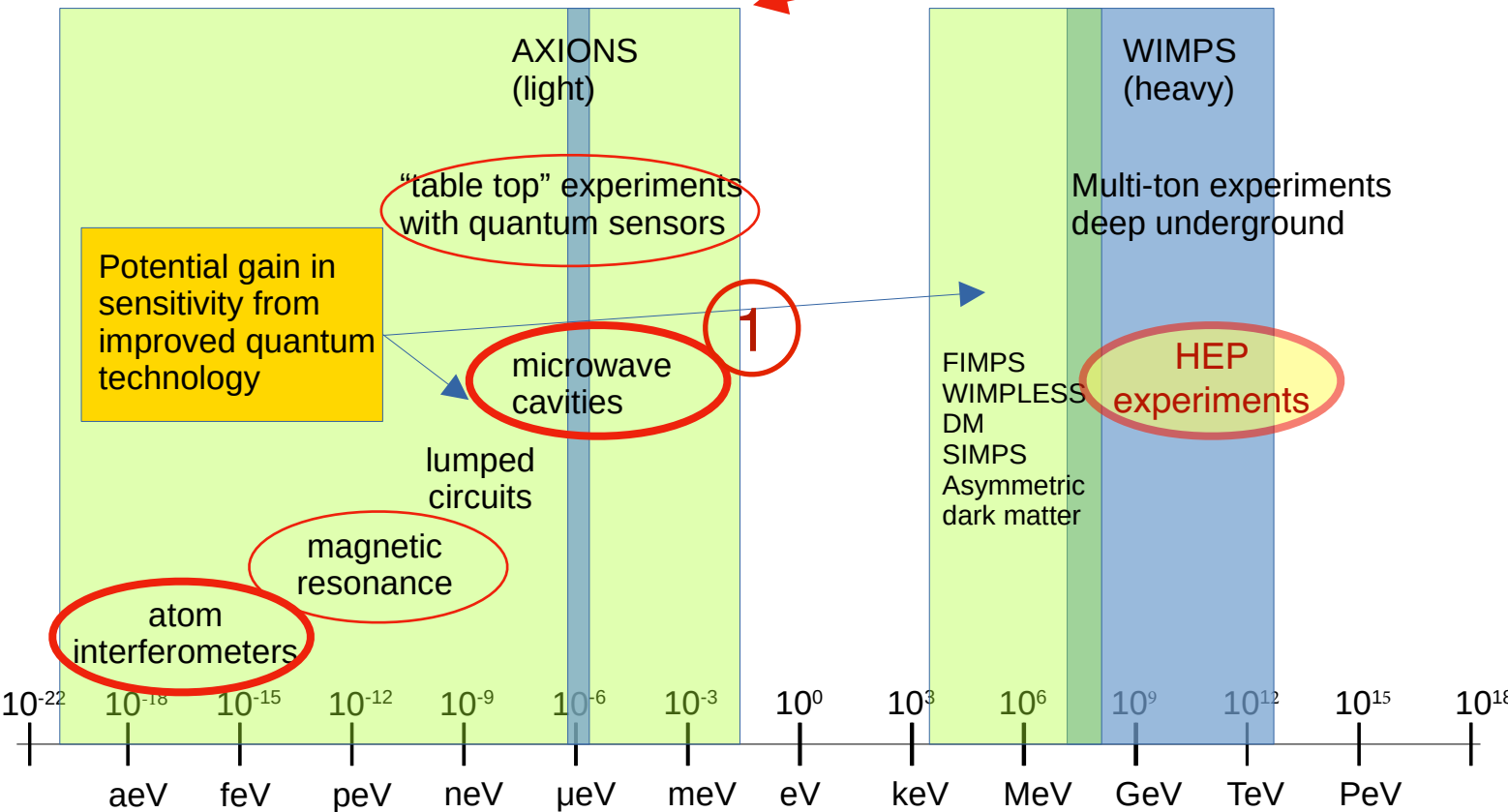
**Axions, ALP's, DM & non-DM UL-particle searches**



cavity size = axion size  
axion mass = unknown

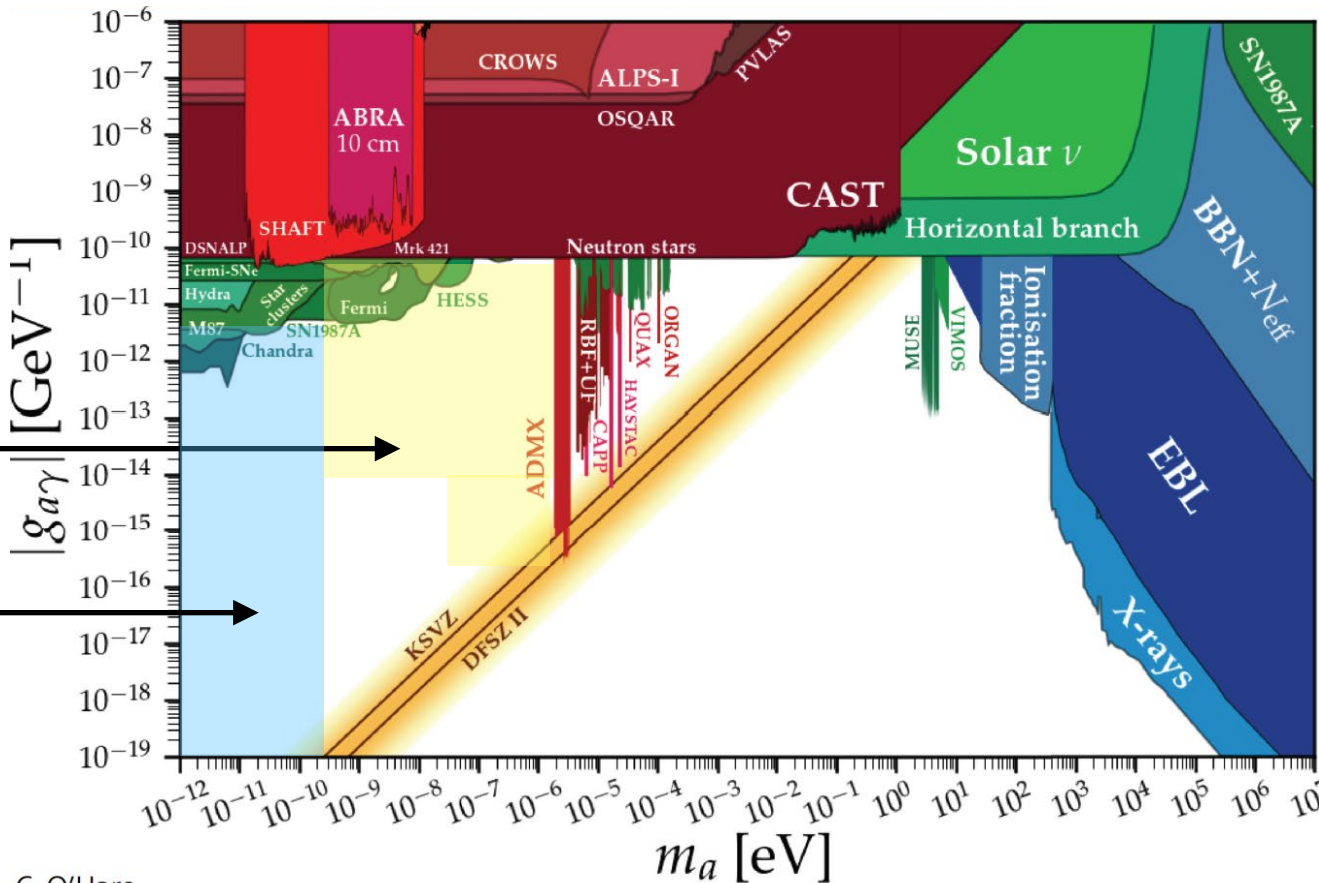
$$F \sim g_{a\gamma}^2 m_A^2 B^4 V^2 T_{sys}^{-2} G^4 Q,$$

system noise temperature  
cryo-amplifiers JPA



(but not only...)

particle physics: what are we talking about?



C. O'Hare

DMRadio



CASPER  
electric  
NMR



spin-based, NMR (2)

microwave cavities (1)

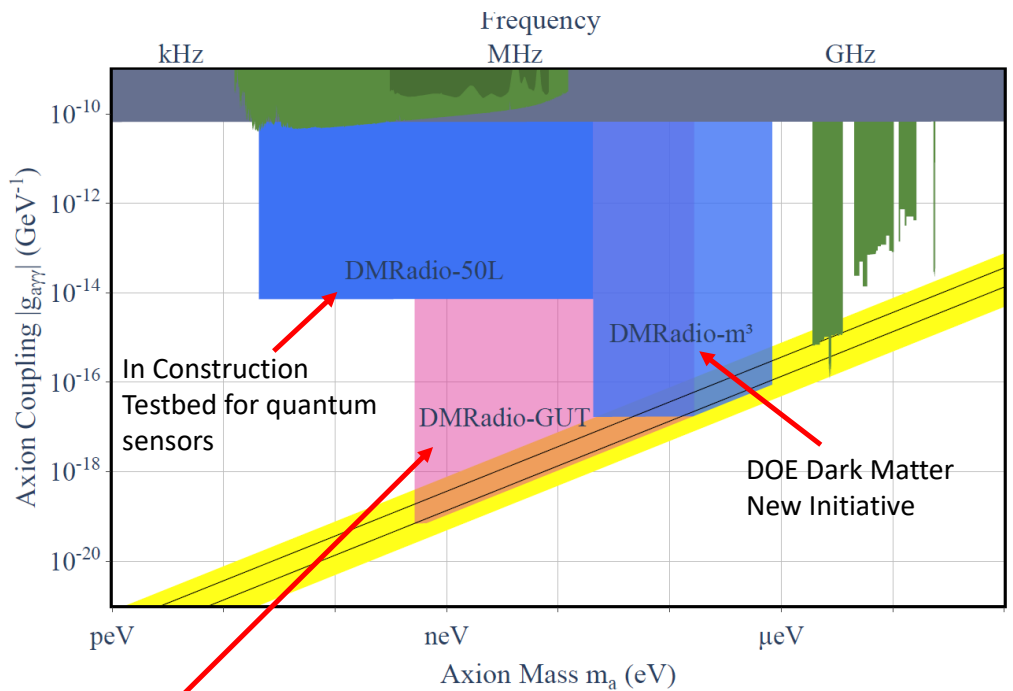
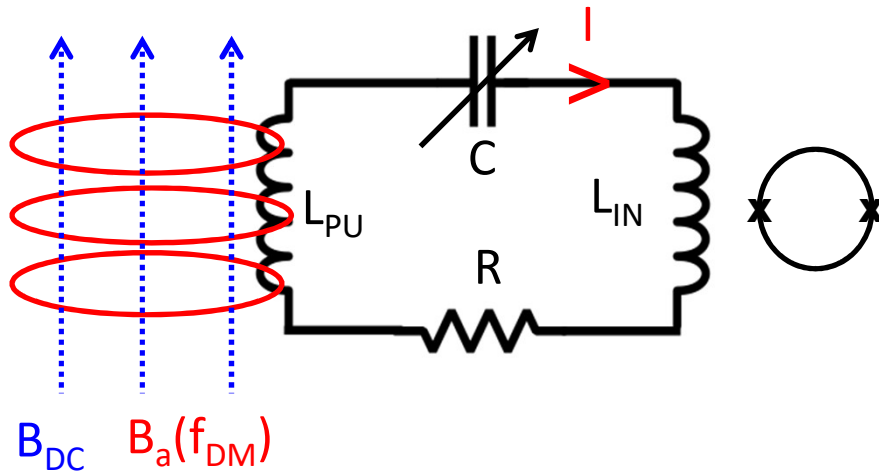
bolometers, TES (1)

electromagn. resonators

# DMRadio

Focus on electromagnetic interaction: axions and photons mix in the presence of a strong magnetic field

Focus on detecting not a particle (a photon), but a field



- Axion field converts to oscillating EM signal in background DC magnetic field
- Detect using tunable resonator
- Signal enhancement when resonance frequency matches rest-mass frequency  $\nu_{DM}=mc^2/h$
- SQUID's, RF Quantum upconverters, cryoamplifiers (e.g. JJPA)

# CASPER electric NMR

Focus on different interactions: the **electric dipole moment (EDM) interaction** and the **gradient interaction with nuclear spin I**. The EDM interaction arises from the coupling of the axion to the gluon field.

→ spin  $\sigma$  to axion coupling:

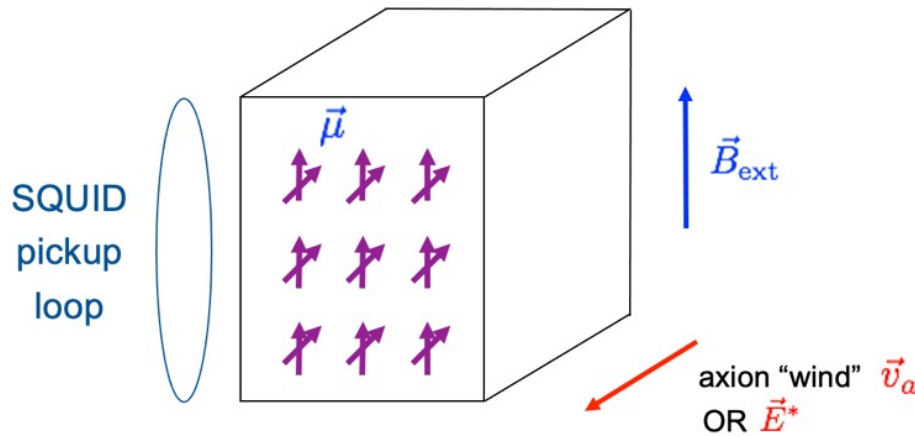
$$H_e \propto a \sigma \cdot E^*$$

**CASPER-electric**

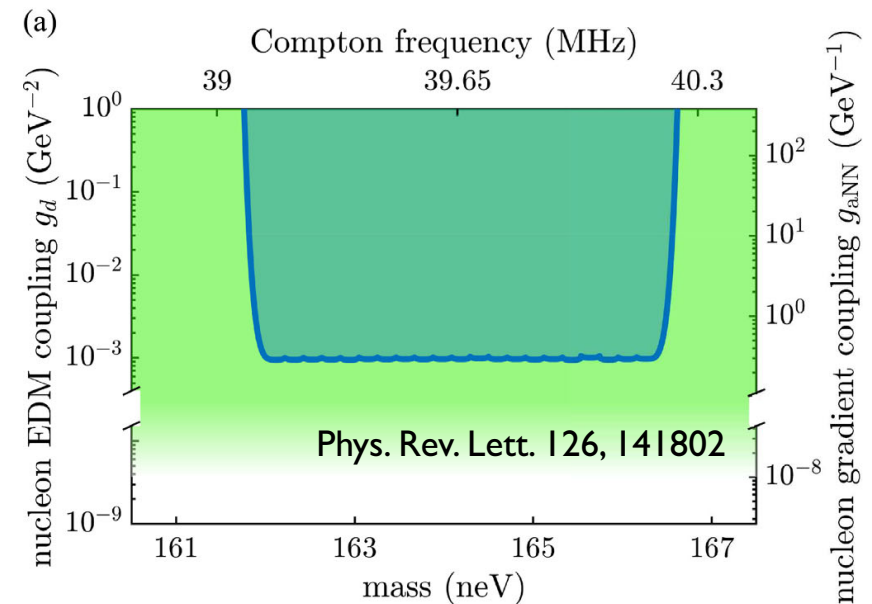
→ spin  $\sigma$  to axion gradient coupling:

$$H_g \propto \sigma \cdot \nabla a$$

**CASPER-gradient**



**Cosmic Axion Spin Precession Experiment** is based on a precision measurement of  $^{207}\text{Pb}$  solid-state nuclear magnetic resonance in a polarized ferroelectric crystal. Axion-like dark matter can exert an oscillating torque on  $^{207}\text{Pb}$  nuclear spins via the electric dipole moment coupling  $g_d$  or via the gradient coupling  $g_{aNN}$ .



numerous improvements possible → many orders of magnitude in mass and sensitivity range

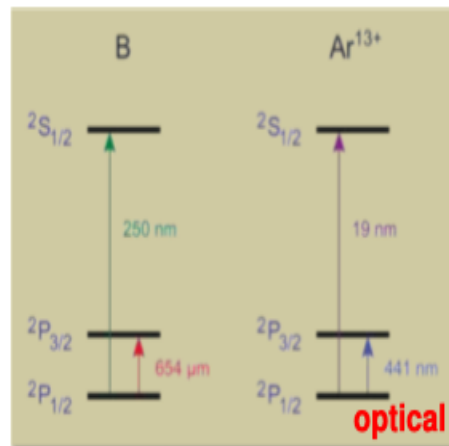
# particle physics @ CERN : what are we talking about?

tests of QED, T-violation, P, Lorentz-violation, DM searches

## HCI's in Penning traps

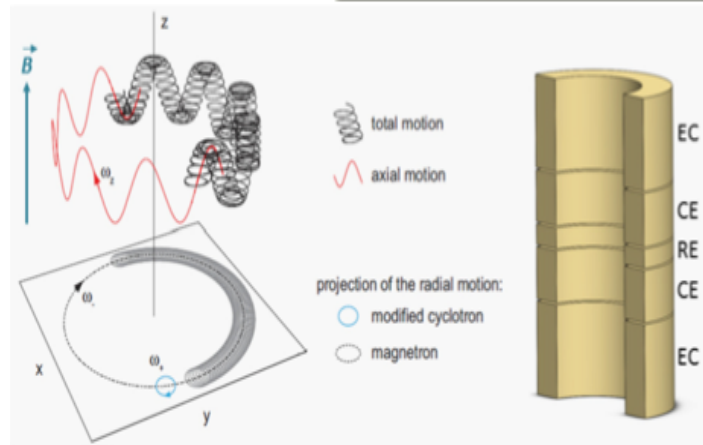
Scaling with a nuclear charge  $Z$

- Binding energy  $\sim Z^2$
- Hyperfine splitting  $\sim Z^3$
- QED effects  $\sim Z^4$
- Stark shifts  $\sim Z^{-6}$



@ CERN:

eEDM's in molecules  
 nuclear clock ( $^{229}\text{Th}$ )  
 molecular / ion clocks



K. Blaum et al., Quantum Sci. Technol. 6 014002 (2021)

Quantum Sensors for New-Physics Discoveries  
<https://iopscience.iop.org/journal/2058-9565/page/Focus-on-Quantum-Sensors-for-New-Physics-Discoveries>

**Present status:  $5 \times 10^{-12}$  relative mass precision**

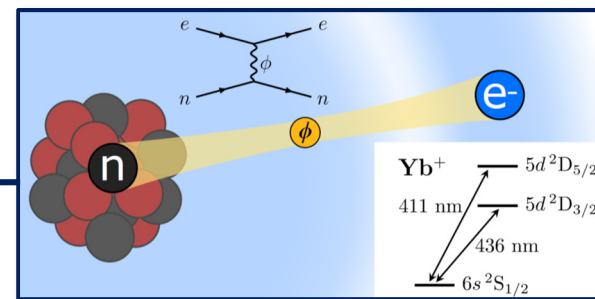
R. X. Schüssler et al., Nature 581, 42 (2020)



HCI: **much larger** sensitivity to variation of  $\alpha$  and dark matter searches than current clocks

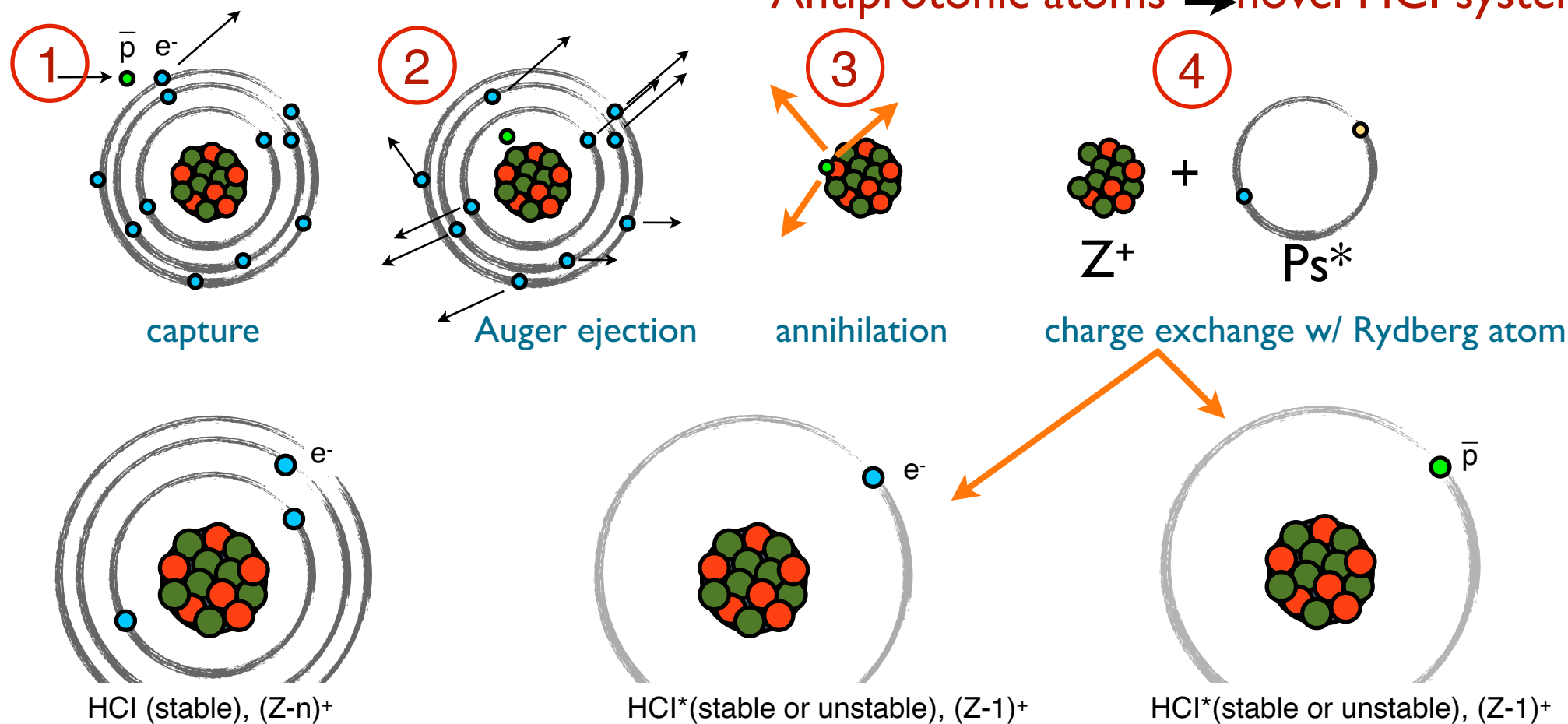
- Searches for the variation of fundamental constants
- Tests of QED: precision spectroscopy
- Fifth force searches: precision measurements of isotope shifts with HCIs to study non-linearity of the King plot

linear relation in isotope shifts between two transitions caused by e.g. new force mediated by weakly interacting boson  
Mikami, K. Et al.. Probing new intra-atomic force with isotope shifts. The European Physical Journal C. 77. 10.1140/epjc/s10052-017-5467-4.

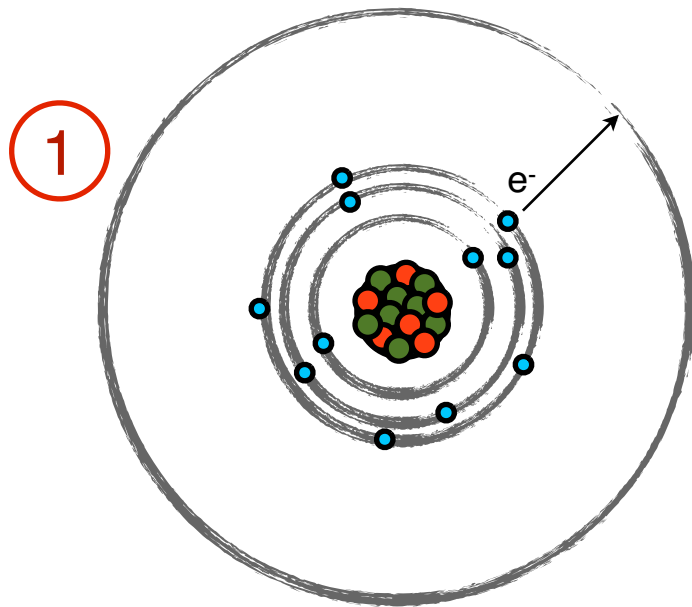


Review on HCIs for optical clocks: Kozlov *et al.*, Rev. Mod. Phys. **90**, 045005 (2018)

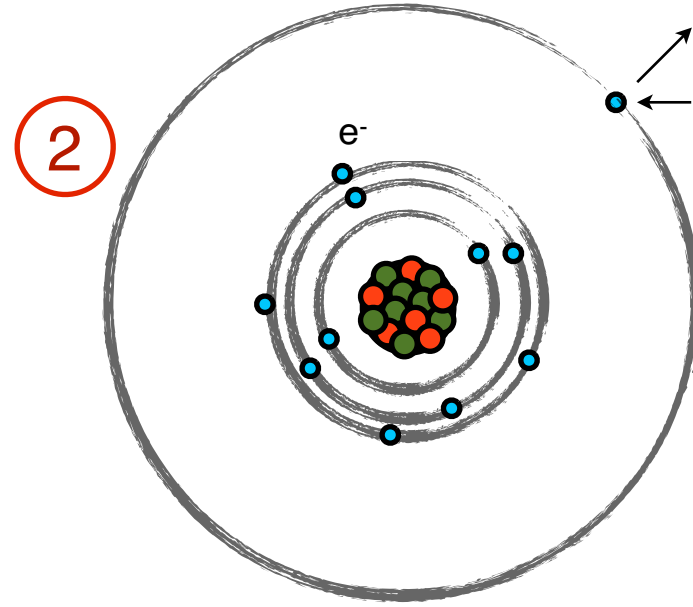
## Antiprotonic atoms → novel HCI systems



Antiprotonic Rydberg atom: exotic couplings, similar approach as spectroscopy of muonic atoms, CPT tests



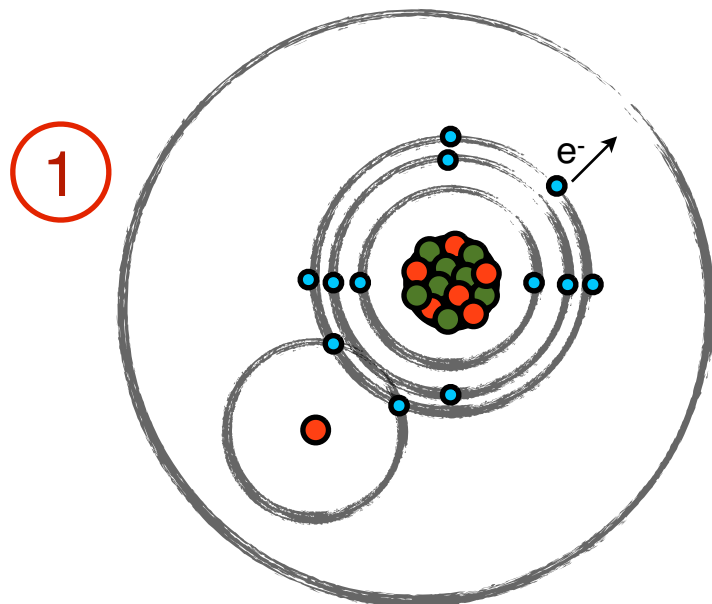
Rydberg excitation



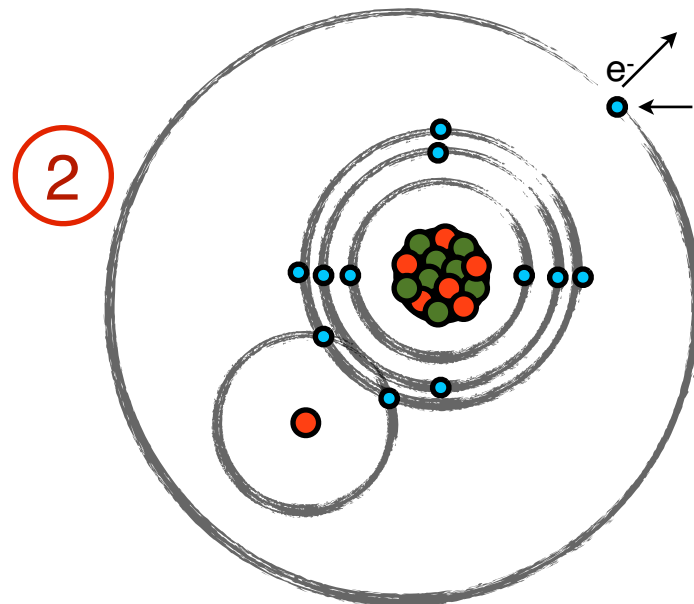
charge exchange

at end of cascade, p is very close to nucleus... investigate long-range behavior of strong interaction?

Antiprotonic Rydberg molecule:  $\bar{p}$ EDM?



Rydberg excitation



charge exchange

similar approach as eEDM in molecules

end of part I...

# Some words on signal read out ...

superconducting devices (TES, SNSPD, ...) / cryo-electronics

spin-based, NV-diamonds

optical clocks

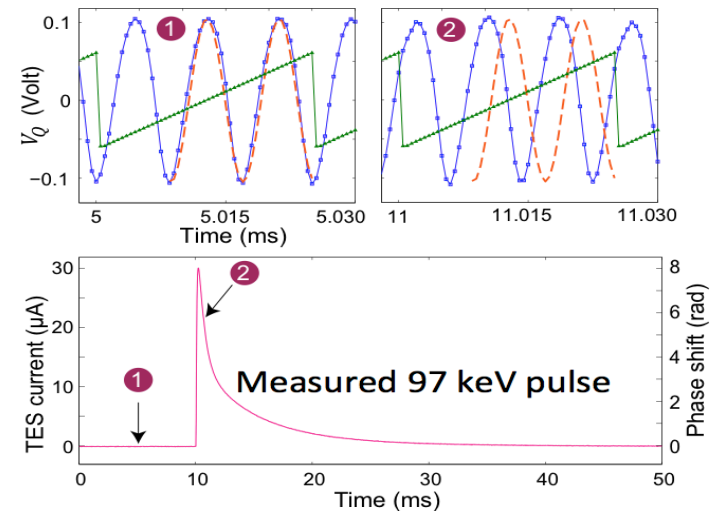
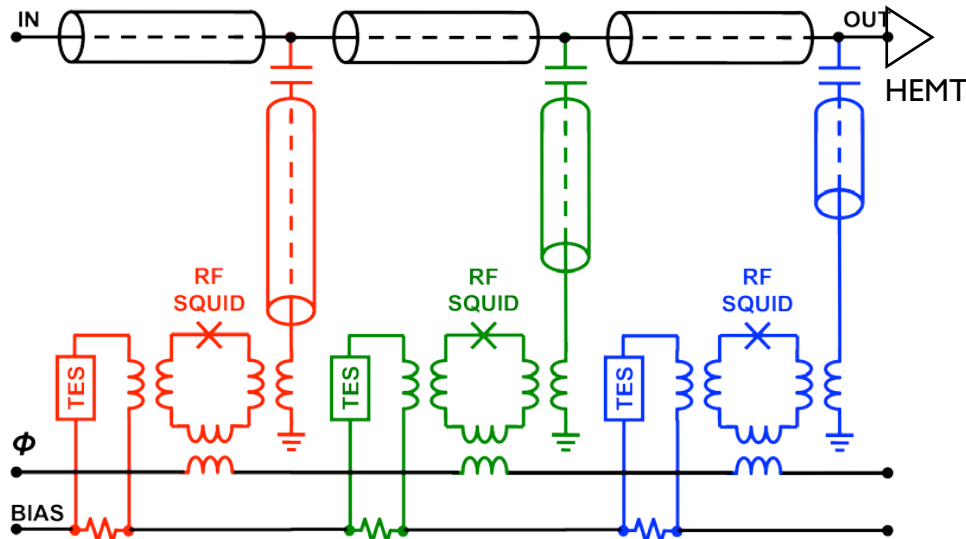
ionic / atomic / molecular

# Some words on signal read out ...

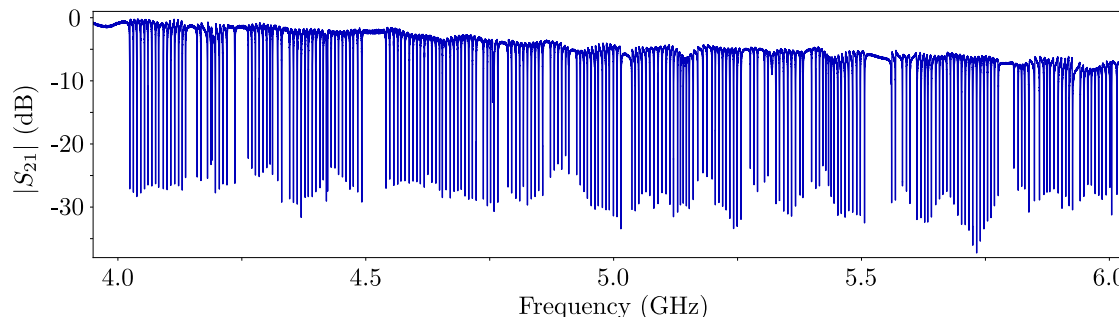
<https://arxiv.org/abs/2209.05621>

superconducting devices (TES, SNSPD, ...) / cryo-electronics

Microwave SQUID multiplexing ( $\mu$ mux) is a form of frequency-domain multiplexing that allocates this bandwidth between input channels by the use of distinct, high-Q microwave resonances, each coupled to its own rf-SQUID and reading out the current signal from its own detector. As a superposition of microwave tones passes through the circuit, each tone is modulated by its own SQUID/resonator circuit before being amplified by the HEMT and brought to room temperature on a single coaxial cable.

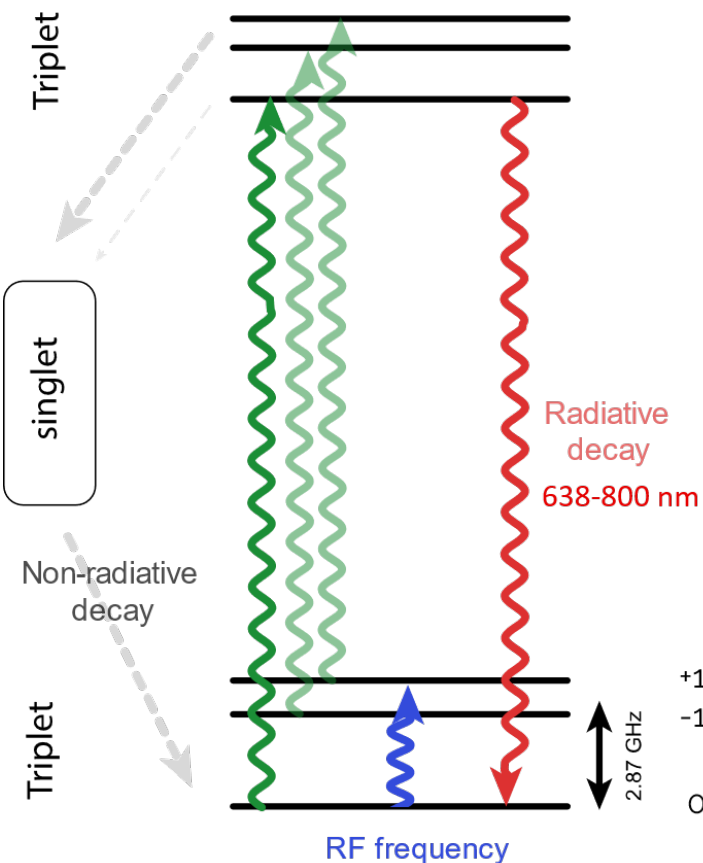
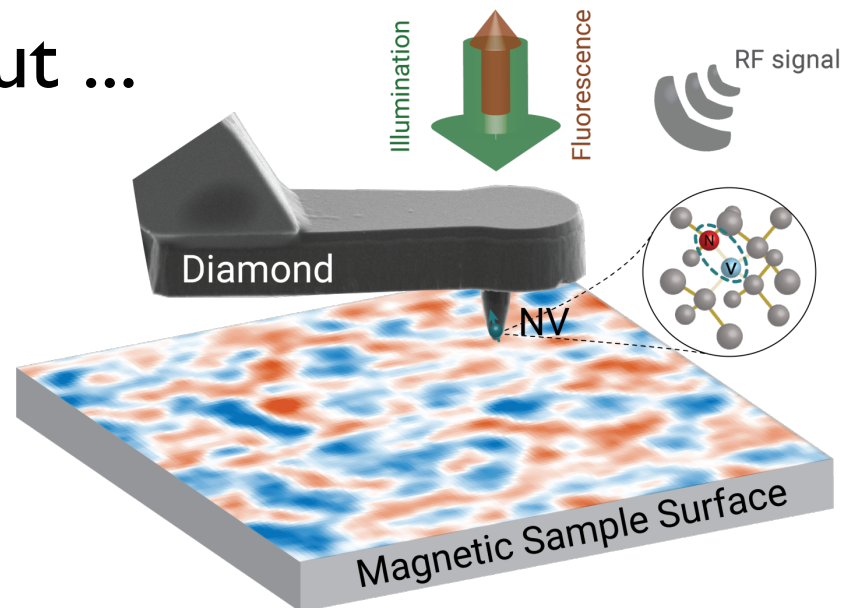


Critical current brings TES close to critical transition: Detection-induced transition leads to change in resonant frequency of circuit; slightly different resonant frequencies for each SQUID circuit (through voltage biasing): unique identification of each pixel. Injected microwave signal is then shifted in phase and amplitude



# Some words on signal read out ...

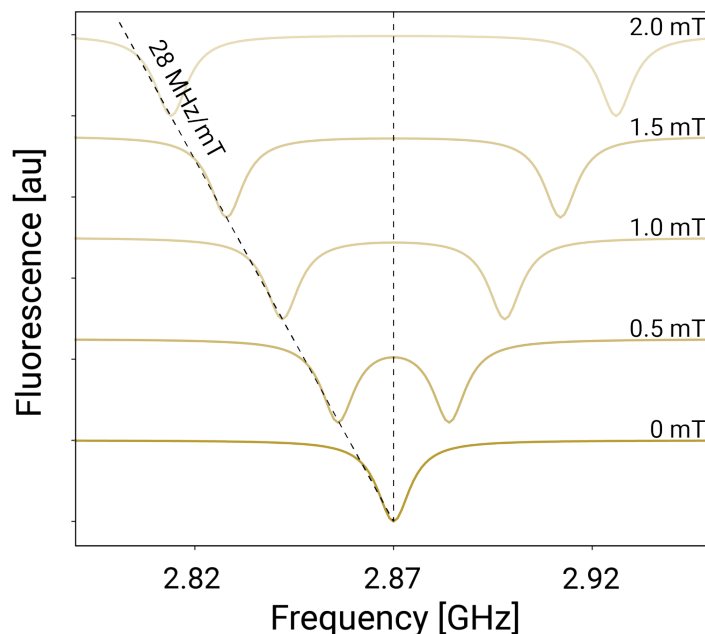
spin-based, NV-diamonds



two level system, with an energy spacing in the microwave range

State readout is based on the fact that the excited state couples to a "dark state", reducing the fluorescence rate compared to the ground state

Probe ground state population by detecting red fluorescence light



Magnetic field sensitivity: Zeeman shift: microwave resonance frequency changes

# Some words on signal read out ...

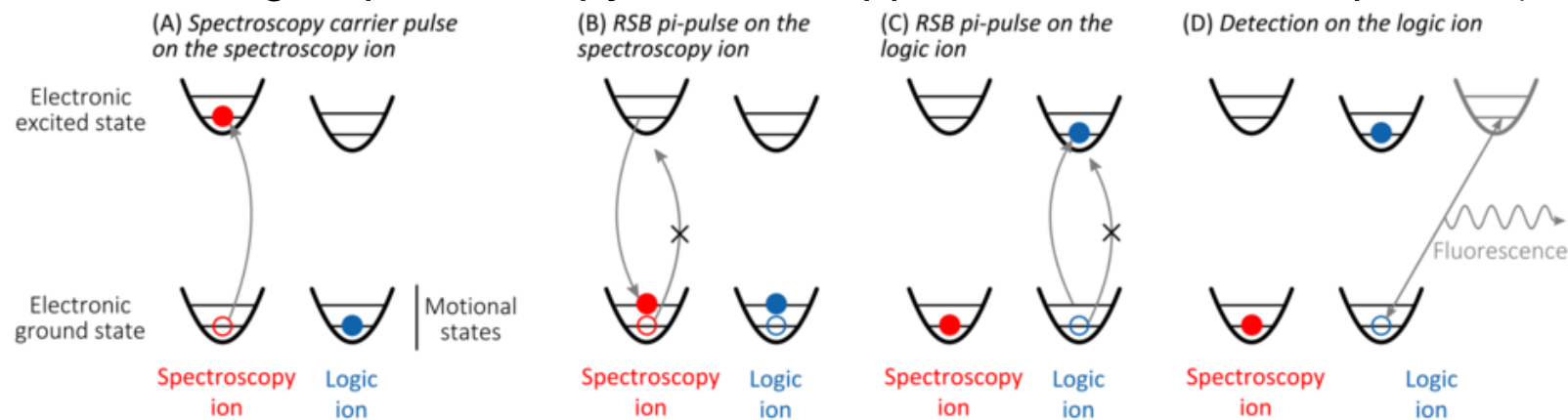
ionic / atomic / molecular

optical clocks

Various schemes: e.g. antiprotonic atoms (very low statistics): resonant transitions result in annihilation

Generally: fluorescence from resonant state back to ground state

Quantum logic spectroscopy: via co-trapped ions of another species (e.g. for  $\text{Al}^+$ )



<https://www.quantummetrology.de/eqm/research/iqloc/quantum-logic-spectroscopy/>

(Quantum non-demolition measurement)

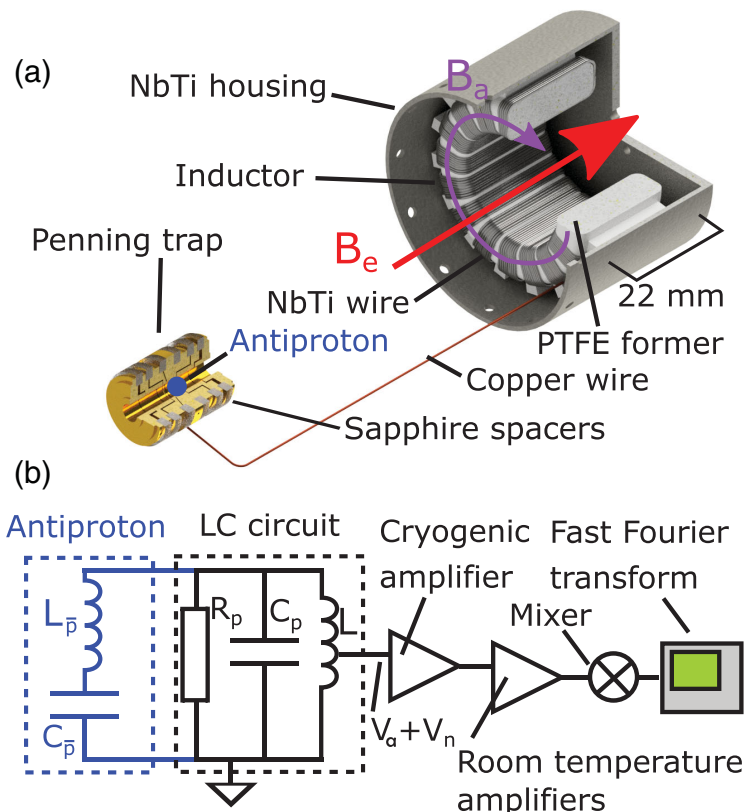
# Quantum sensors for new particle physics experiments: Penning traps

search the noise spectrum of fixed-frequency resonant circuit for peaks caused by dark matter ALPs converting into photons in the strong magnetic field of the Penning-trap magnet

Resolving single antiproton spin flips requires the highest Q and lowest temperature LC resonant detectors ever built: **BASE-CERN** is the state of the art

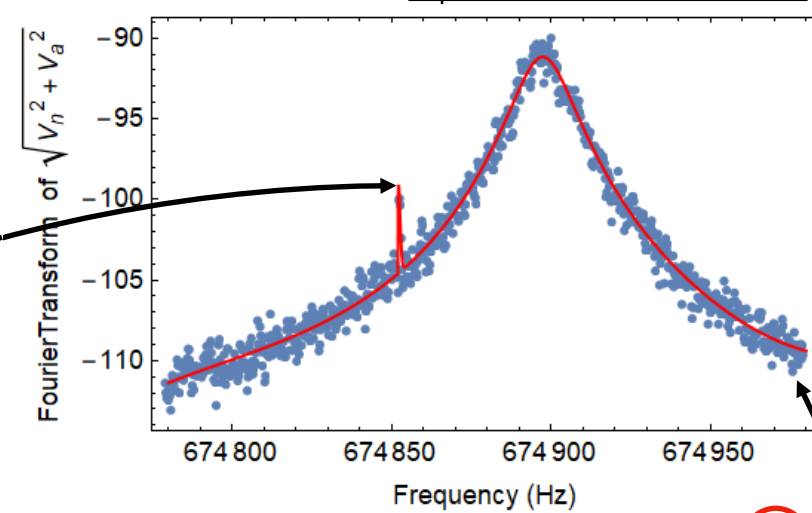
Constraints on the Coupling between Axion-like Dark Matter and Photons Using an Antiproton Superconducting Tuned Detection Circuit in a Cryogenic Penning Trap

J. Devlin et al., **BASE** collaboration, Physical Review Letters 126, 041301 (2021)



H. Nagahama et al., Rev. Sci. Instrum. 87, 113305 (2016)

<https://indico.cern.ch/event/1002356/>



resonator background  $\propto \sqrt{T_Z}$   
from antiproton spin-flip

The axion signal

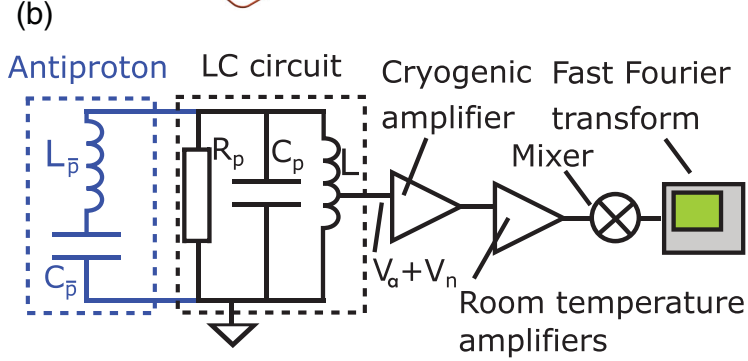
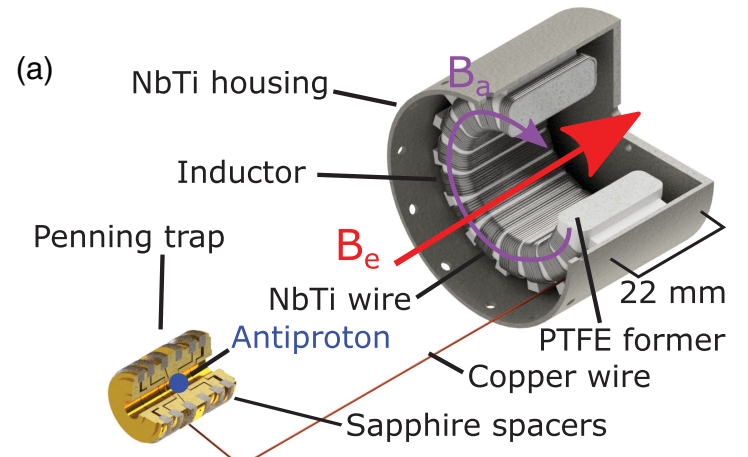
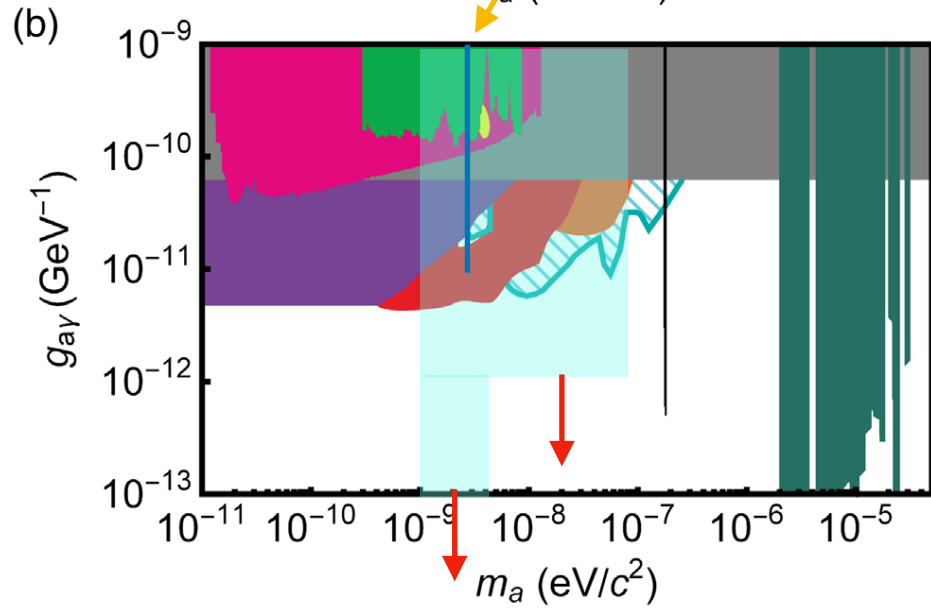
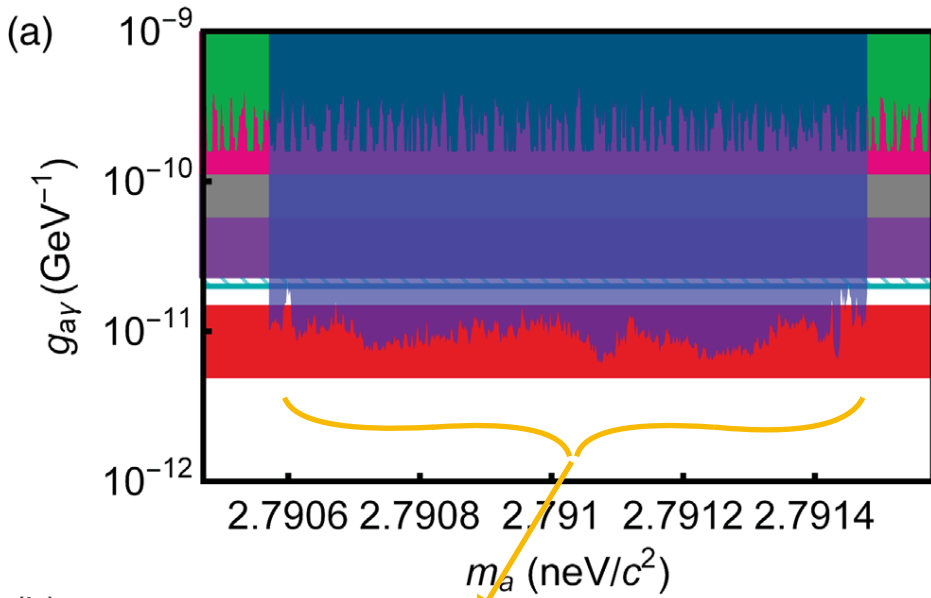
$$V_a = \frac{\pi}{2} Q \sqrt{f(\nu, Q, \mathbf{q})} \kappa \nu_a l N_T (r_2^2 - r_1^2) g_{a\gamma} |\mathbf{B}_e| \sqrt{\rho_a \hbar c}.$$

- $f(\nu, Q, \mathbf{q})$  is a lorentzian line-shape function proportional to  $\text{Re}\{Z\}$
- $e_n$  is the equivalent input noise of the amplifier
- $\kappa$  is the coupling constant
- $Q$  is the resonator Q-factor
- $N_T$  is the number of turns
- $l$  is the length of the toroid along the magnet B field
- $r_1$  is the inner radius of the toroid
- $r_2$  is the outer radius
- $g_{a\gamma}$  is the coupling constant
- $B$  is the static magnetic field
- $\rho_a$  is the dark matter density



# Tunability!

# Quantum sensors for new particle physics experiments: Penning traps



currently developing **superconducting tunable capacitors & laser-cooled resonators**

7 T magnet + broader FFT span: one month  $\longrightarrow$   
 2 and 5 neV to an upper limit of  $1.5 \times 10^{-11} \text{ GeV}^{-1}$

Limits	Hints
<div style="display: inline-block; width: 15px; height: 15px; background-color: purple; border: 1px solid black; margin-right: 5px;"></div> SN-1987A <div style="display: inline-block; width: 15px; height: 15px; background-color: gray; border: 1px solid black; margin-left: 10px; margin-right: 5px;"></div> CAST <div style="display: inline-block; width: 15px; height: 15px; background-color: black; border: 1px solid black; margin-left: 10px; margin-right: 5px;"></div> ADMX-SLIC	<div style="display: inline-block; width: 15px; height: 15px; background-color: lightblue; border: 1px solid black; margin-right: 5px;"></div> Excess
<div style="display: inline-block; width: 15px; height: 15px; background-color: orange; border: 1px solid black; margin-right: 5px;"></div> H.E.S.S. <div style="display: inline-block; width: 15px; height: 15px; background-color: blue; border: 1px solid black; margin-left: 10px; margin-right: 5px;"></div> BASE <div style="display: inline-block; width: 15px; height: 15px; background-color: green; border: 1px solid black; margin-left: 10px; margin-right: 5px;"></div> ABRACADABRA	<div style="display: inline-block; width: 15px; height: 15px; border: 1px solid black; margin-right: 5px;"></div> $\gamma$ rays
<div style="display: inline-block; width: 15px; height: 15px; background-color: darkgreen; border: 1px solid black; margin-right: 5px;"></div> Cavities <div style="display: inline-block; width: 15px; height: 15px; background-color: magenta; border: 1px solid black; margin-left: 10px; margin-right: 5px;"></div> SHAFT <div style="display: inline-block; width: 15px; height: 15px; background-color: red; border: 1px solid black; margin-left: 10px; margin-right: 5px;"></div> FERMI-LAT	<div style="display: inline-block; width: 15px; height: 15px; background-color: yellow; border: 1px solid black; margin-right: 5px;"></div> Pulsars

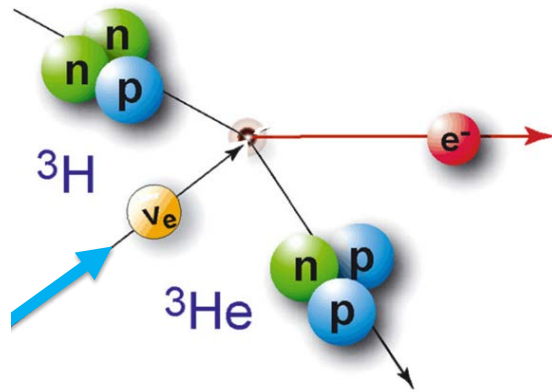
# Challenge: Detection of primordial neutrino flux: $E_\nu \sim \mu\text{eV}$

neutrino capture on  $\beta$ -decay nuclei as a detection method for the Cosmic Neutrino Background (CNB) was laid out in the original paper by Steven Weinberg in 1962

S. Weinberg, Universal Neutrino Degeneracy, Phys. Rev. 128, 1457 (1962)

$\frac{d\Gamma}{dE_e} [\text{yr}^{-1}\text{eV}^{-1}]$

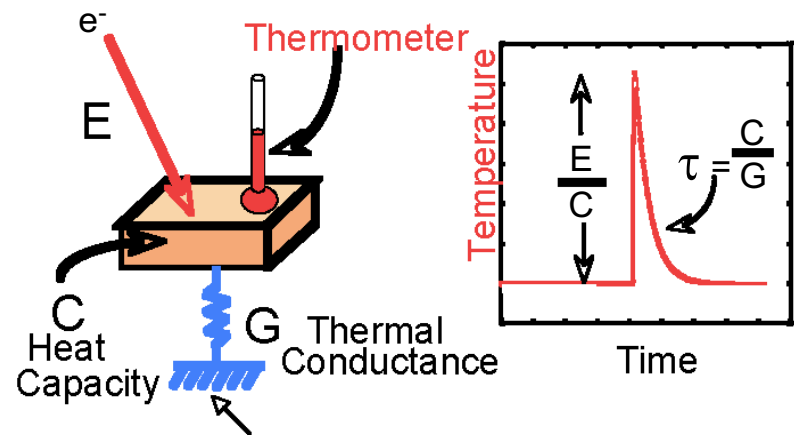
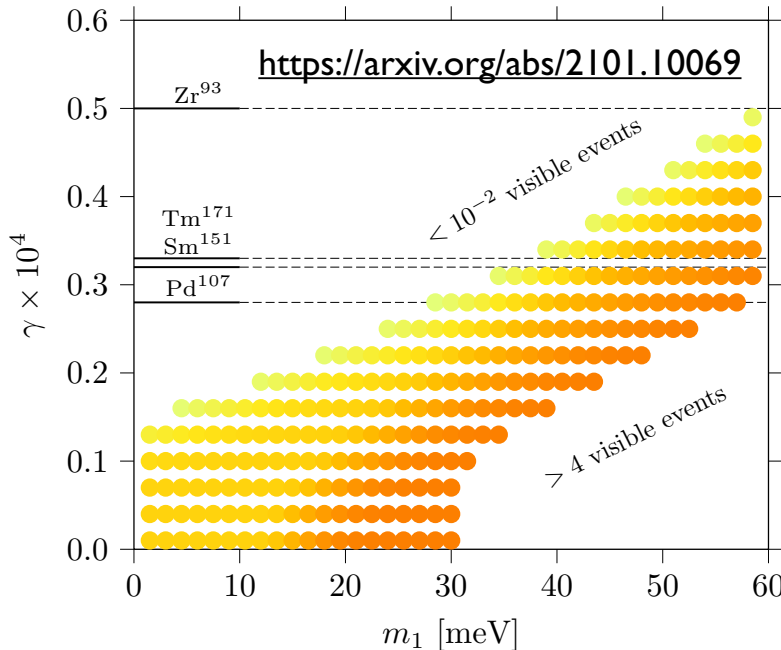
detection!  
directionality!  
energy!



**PTOLEMY** experimental quest  
relic neutrino capture  
RF tracker (EM radiation in B field)  
microcalorimeter (TES: sub-eV for  $\sim\text{eV}$ )

<https://indico.cern.ch/event/1090250/>

event rate low!  
Better use  
heavier  
isotopes than  
tritium as  
targets!



$\sim 100 \text{ mK}$  cold bath (refrigerator)

**MUCH R&D needed!**

Quantum detectors for  
high energy particle physics

typically not obvious, given that most detectors rely on detecting the product of many interactions between a particle and the detector (ionization, scintillation, Cerenkov photons, ...)

handful of ideas that rely on quantum devices, or are inspired by them, but do not necessarily use them as quantum detectors per se, but rather their properties to enhance / permit measurements that are more difficult to achieve otherwise

these are not developed concepts, but rather the kind of approaches one might contemplate working towards



very speculative!

main focus on tracking / calorimetry / timing      closely related: nanostructured materials

→ Frontiers of Physics, M. Doser et al., 2022

## what are the challenges?

- tracking: hit positions, material budget, vertexing
- timing: TOF  $\sim$  sub-ns (ideally ps) for PID
- calorimetry: shower shape, timing, granularity, particle flow
- redundancy / independent modes of measurement
- novel observables: helicity / polarization
- power budget, event rate, PU (timing)
- sensitivity!

## Metamaterials, 0 / 1 / 2-dimensional materials (quantum dots, nanolayers)

ultra-fast scintillators based on perovskites

chromatic calorimetry (QDs)

active scintillators (QCL, QWs, QDs)

GEMs (graphene)

## Atoms, molecules, ions

Rydberg TPC's

## Spin-based sensors

helicity detectors

## Metamaterials, 0 / 1 / 2-dimensional materials (quantum dots, nanolayers)

ultra-fast scintillators based on perovskites

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GEMs (graphene)

## Atoms, molecules, ions

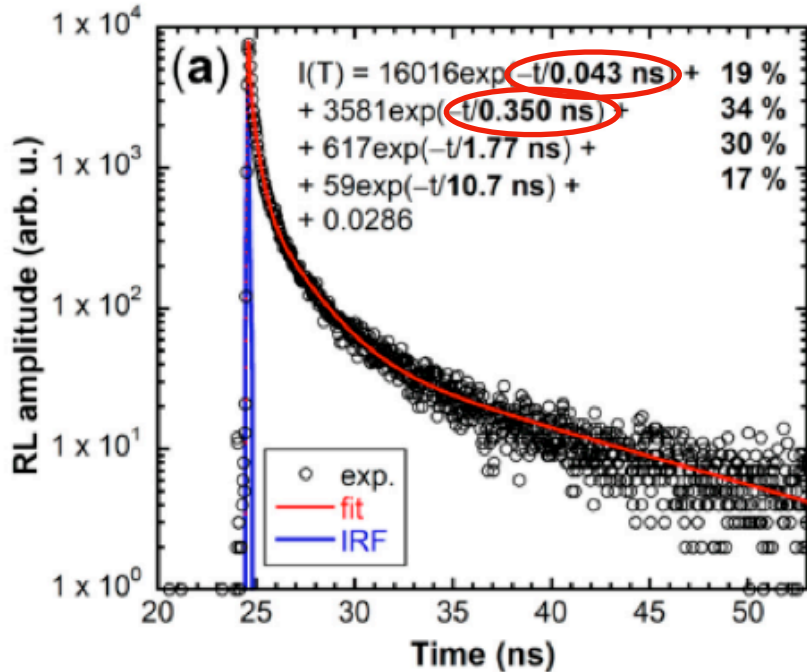
Rydberg TPC's

## Spin-based sensors

helicity detectors

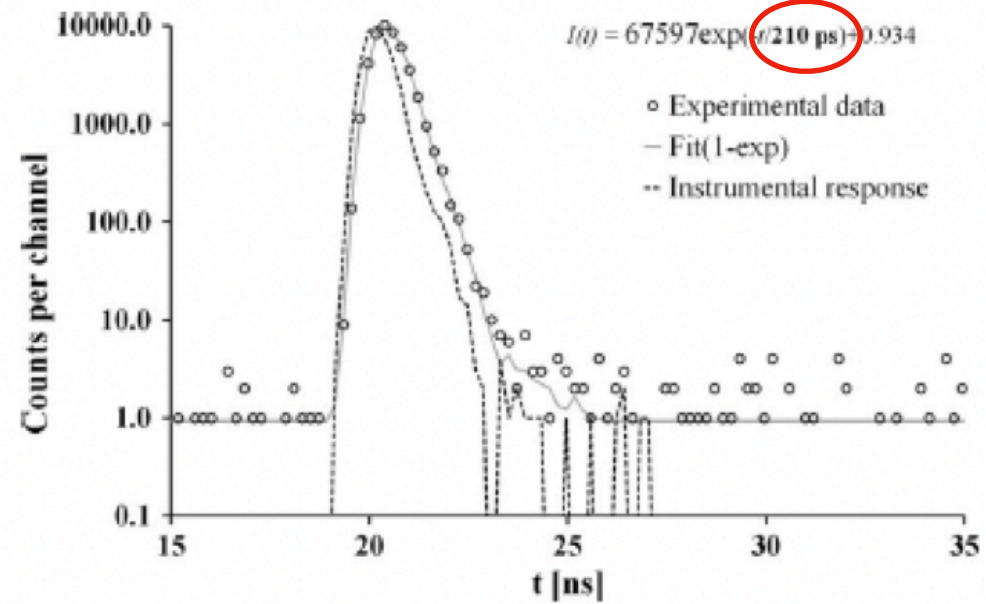
# Quantum dots: timing

Etiennette Auffray-Hillemans / CERN



Scintillation decay time spectra from CsPbBr<sub>3</sub> nanocrystal deposited on glass

K. Decka et al., Scintillation Response Enhancement in Nanocrystalline Lead Halide Perovskite Thin Films on Scintillating Wafers. *Nanomaterials* 2022, 12, 14. <https://doi.org/10.3390/nano12010014>



**Fig. 9.** Photoluminescence decay of ZnO:Ga sample at room temperature. Excitation nanoLED 339 nm, emission wavelength set at 390 nm. Decay curve is approximated by the convolution of instrumental response (also in figure) and single exponential function  $I(t)$  provided in the figure.

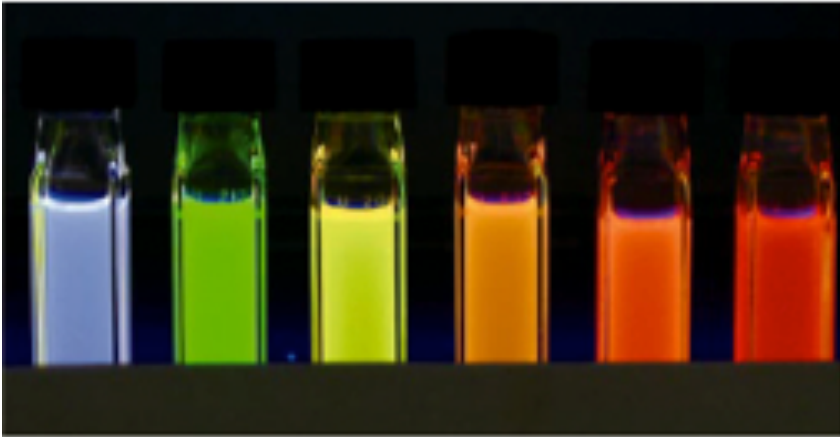
Lenka Prochazkova et al., *Optical Materials* 47 (2015) 67–71

Concerns: integrated light yield (need many photons to benefit from rapid rise time)

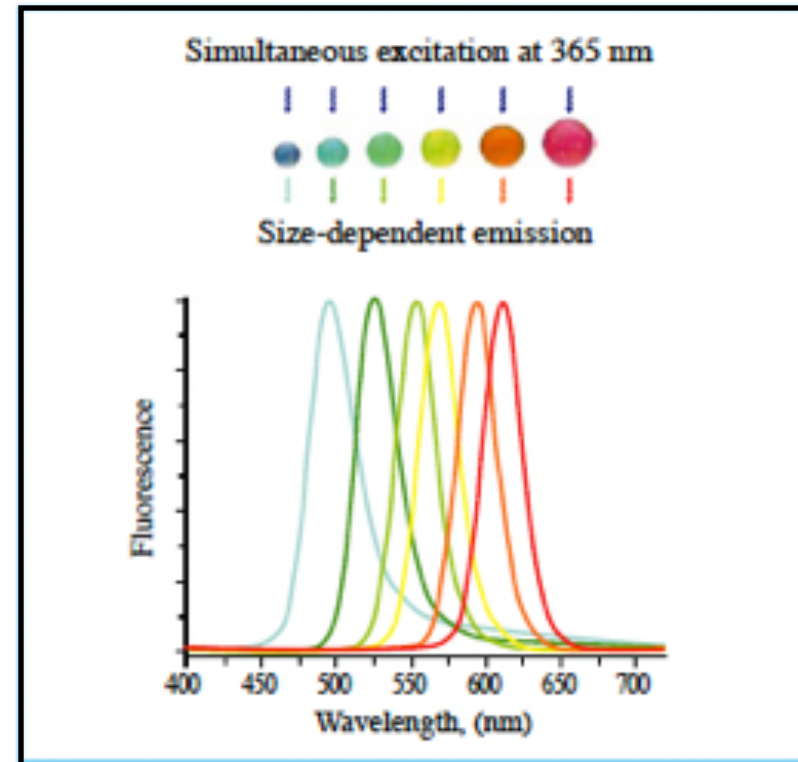


# Quantum dots: timing

Etiennette Auffray-Hillemans / CERN



Hideki Ooba, "Synthesis of Unique High Quality Fluorescence Quantum Dots for the Biochemical Measurements," AIST TODAY Vol.6 , No.6 (2006) p.26- 27

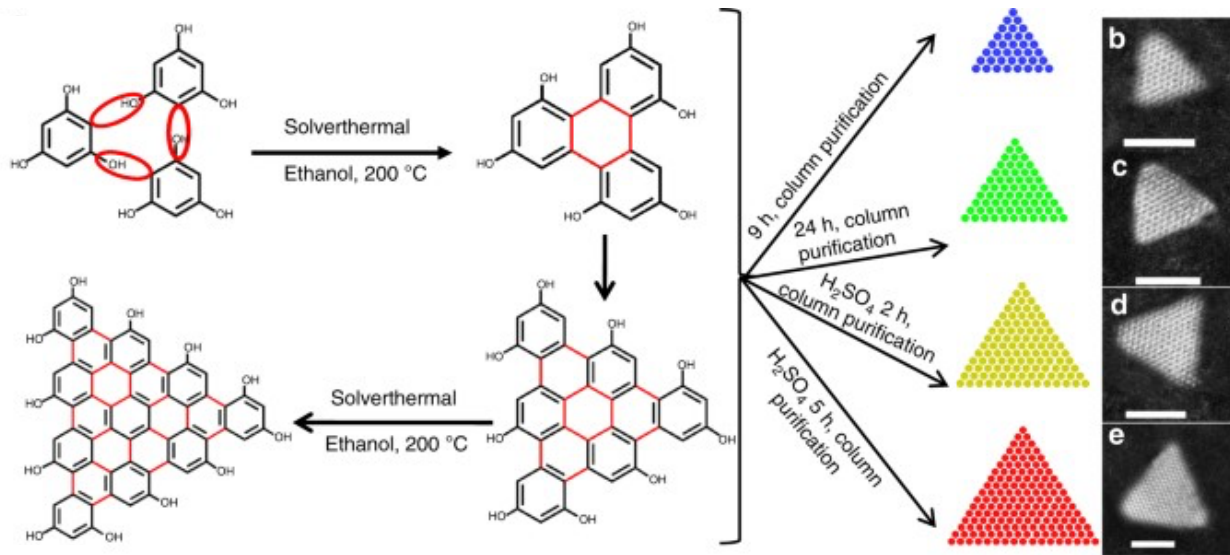


chromatic tunability → optimize for quantum efficiency of PD (fast, optimizable WLS)

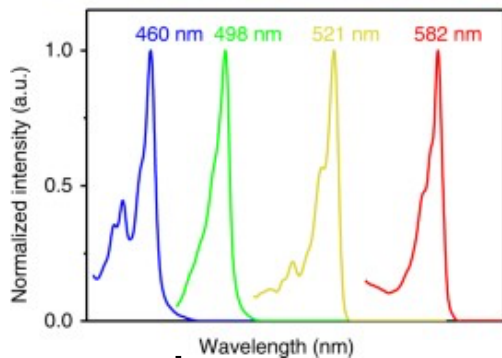
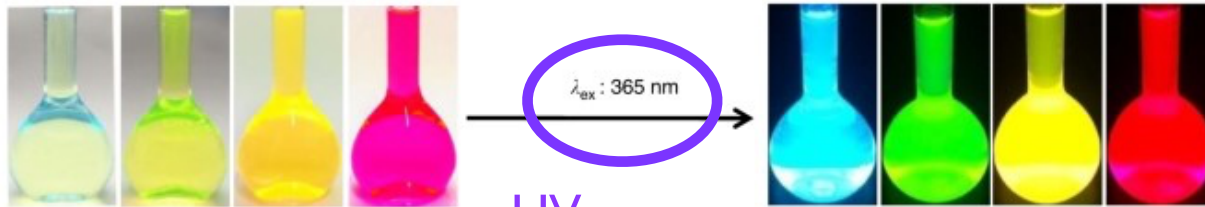
deposit on surface of high-Z material → thin layers of UV → VIS WLS

embed in high-Z material ? two-species (nanodots + microcrystals) embedded in polymer matrix?  
 → quasi continuous VIS-light emitter (but what about re-absorbtion?)

# Quantum dots: chromatic calorimetry



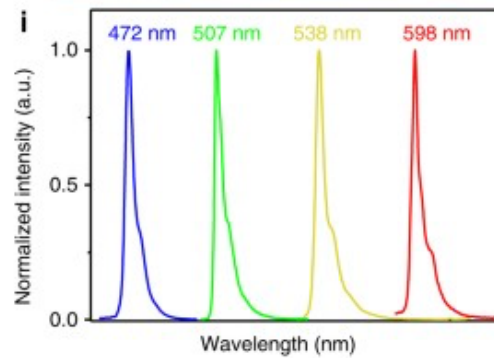
idea: seed different parts of a “crystal” with nanodots emitting at different wavelengths, such that the wavelength of a stimulated fluorescence photon is uniquely assignable to a specific nanodot position



absorption

UV illumination

e.m. shower



emission

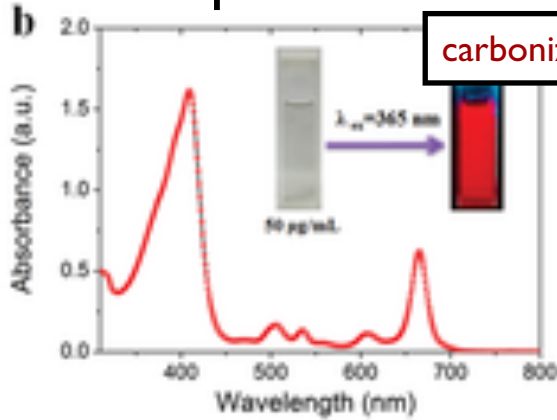
requires:

- narrowband emission (~20nm)
- only absorption at longer wavelengths
- short rise / decay times

select appropriate nanodots

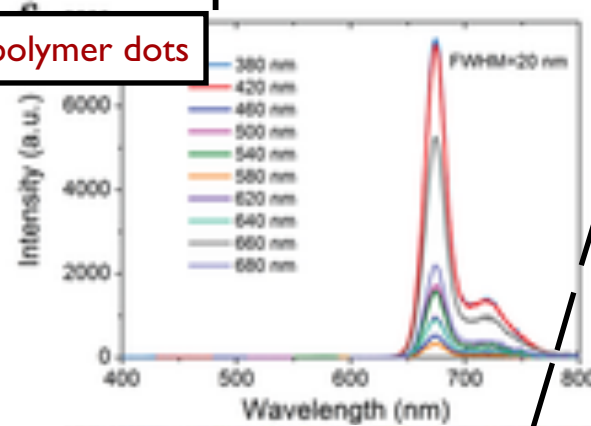
e.g. **triangular carbon nanodots**

absorption spectrum



carbonized polymer dots

emission spectrum



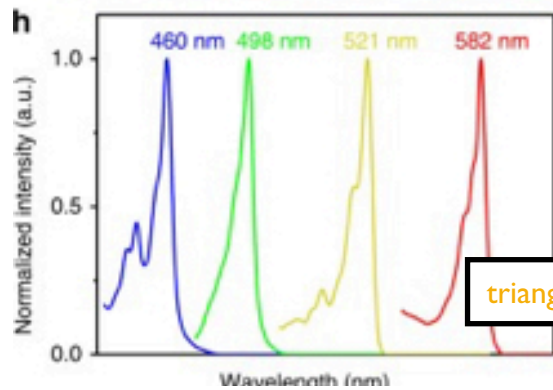
leftmost nanodots:  
absorb wavelengths  $< 650 \text{ nm}$   
emit at  $> 680 \text{ nm}$

next band:  
absorb wavelengths  $< 590 \text{ nm}$   
emit at  $> 590 \text{ nm}$

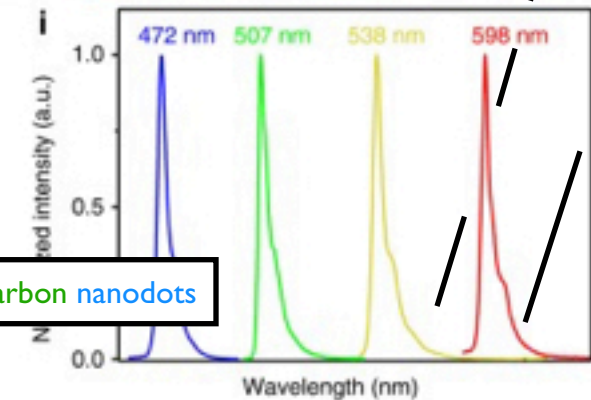
...

rightmost nanodots:  
absorb wavelengths  $< 410 \text{ nm}$   
emit at  $> 420 \text{ nm}$

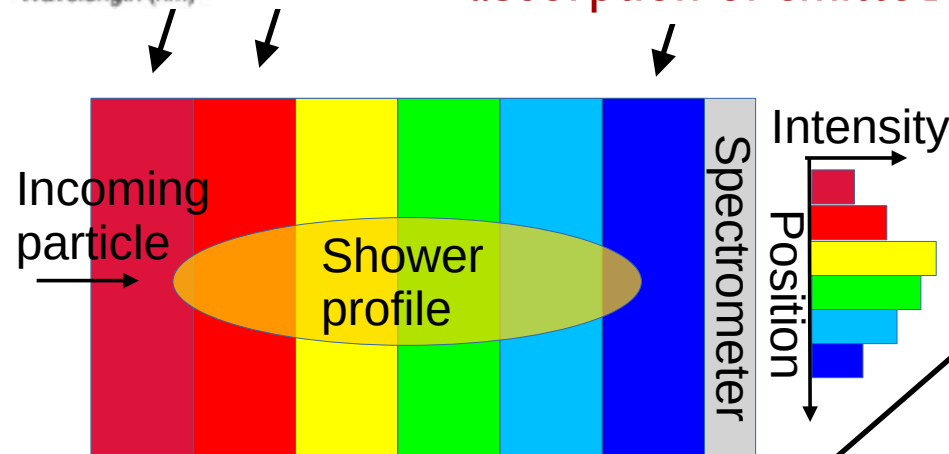
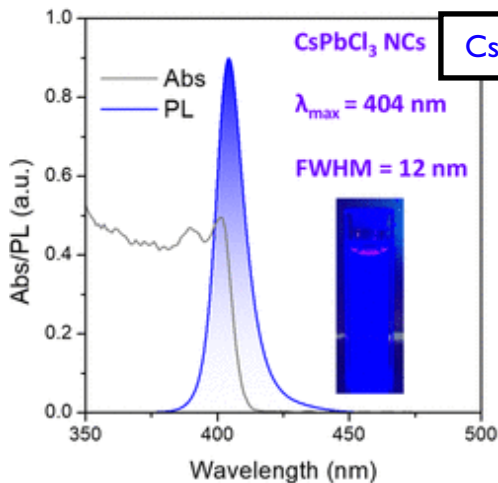
if high-Z substrate transparent  
in 400-700 nm, then no re-  
absorption of emitted light



triangular carbon nanodots



$\text{CsPbCl}_3$  nanocrystals

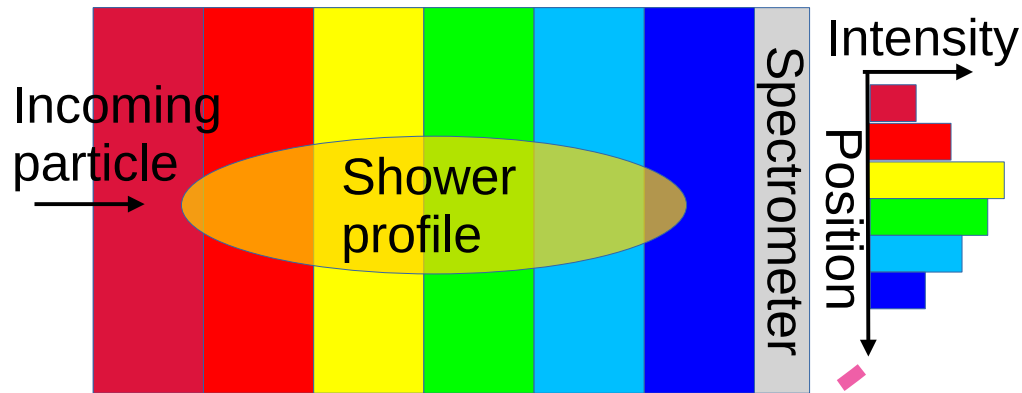


Metalenses?

M. Khorasaninejad & F. Capasso, Science 358, 6367 (2017)

(shower profile via **spectrometry**)

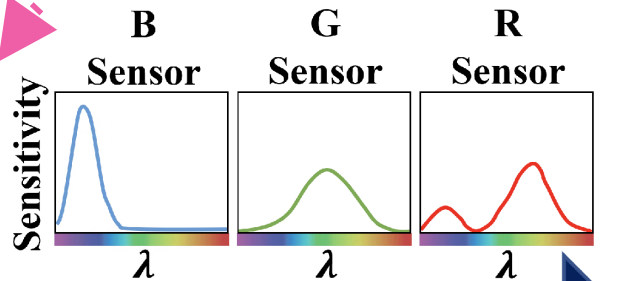
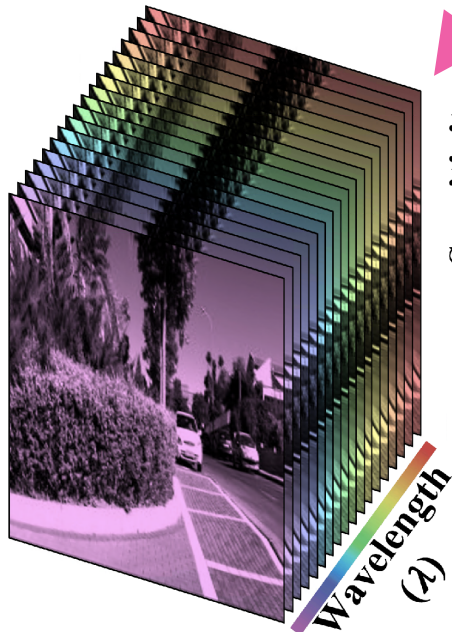
# Quantum dots: chromatic calorimetry (shower profile via spectrometry)



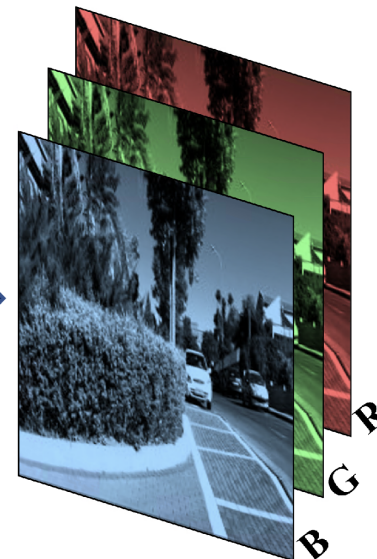
- different options for spectrometry:
- monochromators + PD
  - light guiding fiber / each layer
  - light guiding fiber to spectrometer

A Rehabilitation of Pixel-Based Spectral Reconstruction from RGB Images  
 Yi-Tun Lin and Graham D. Finlayson  
*Sensors* **2023**, 23(8), 4155; <https://doi.org/10.3390/s23084155>

Hyperspectral Image  
(Ground-truth Spectra)



RGB Image



RGB Imaging

Spectral Reconstruction (SR)

# Active scintillators (QWs, QDs, QWDs, QCLs)

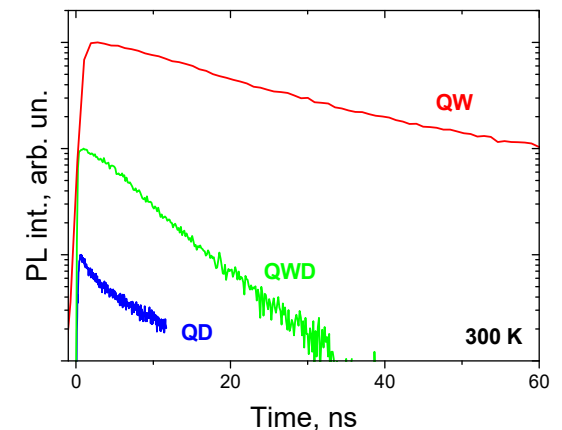
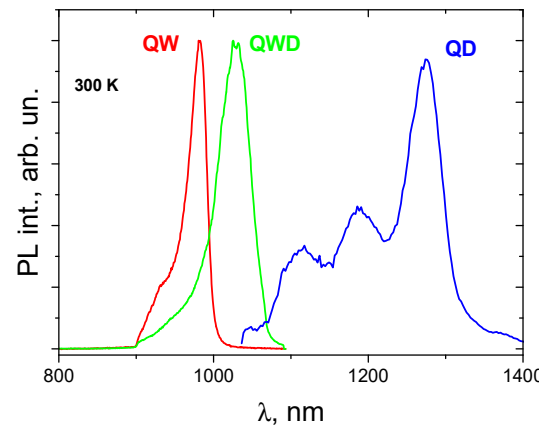
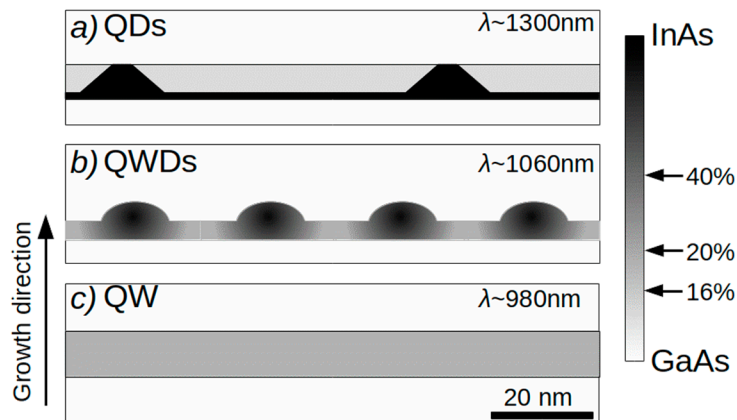
standard scintillating materials are **passive**

- can not be amplified
- can not be turned on/off
- can not be modified once they are in place

is it possible to produce **active** scintillating materials?

- electronically amplified / modulable
- pulsed / primed
- gain adapted in situ

existing QD's, QWD's are elements of optoelectronic devices, typically running at 10 GHz, quite insensitive to temperature



Light Emitting Devices Based on Quantum Well-Dots, Appl. Sci. 2020, 10, 1038; doi:10.3390/app10031038

Emission in **IR!** Silicon is transparent at these wavelengths...  
 Can this IR light be transported *through* a tracker to outside PDs?

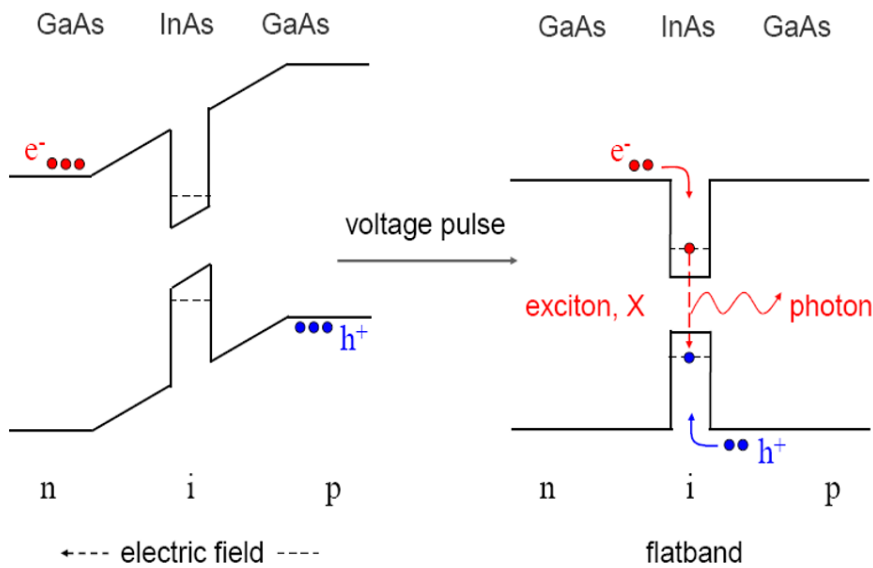
QD's are radiation resistant

R. Leon et al., "Effects of proton irradiation on luminescence emission and carrier dynamics of self-assembled III-V quantum dots," in IEEE Transactions on Nuclear Science, 49, 6, 2844-2851 (2002), doi: 10.1109/TNS.2002.806018.

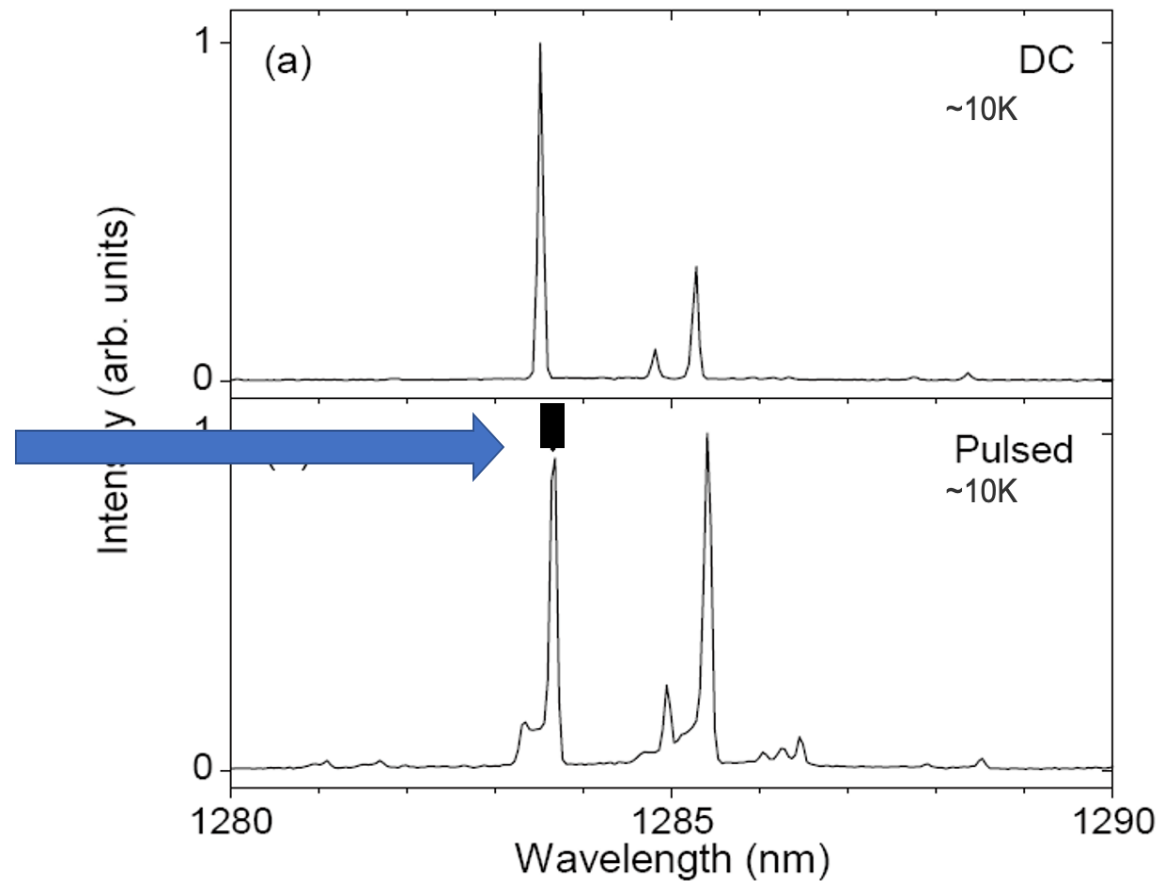
# Active scintillators (QWs, QDs, QWDs, QCLs)

QD's produce sharp atom-like emission peaks

generate photons by optical pumping or electrical injection of electrons into the QD

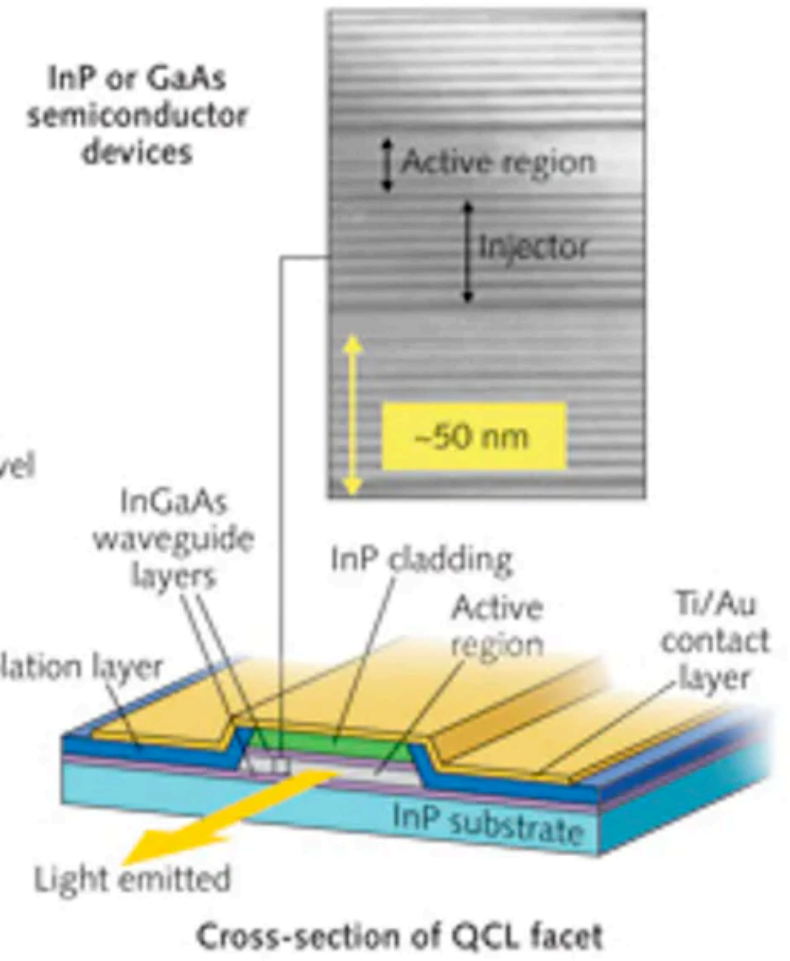
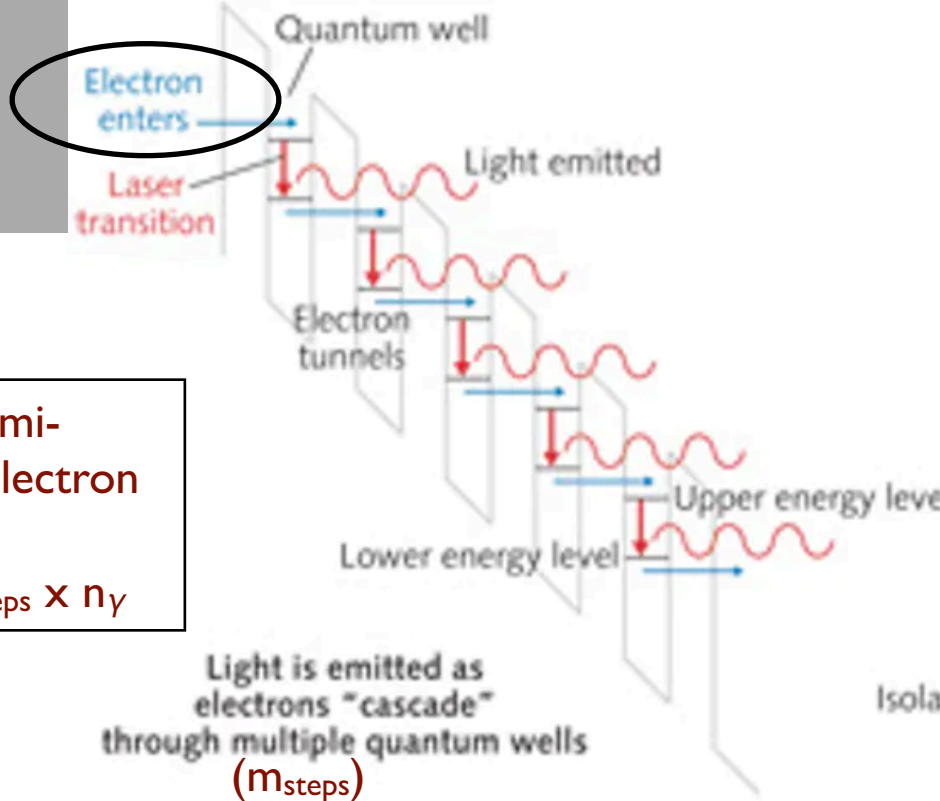


## Electroluminescence (DC and pulsed)



# Active scintillators (QCLs, QWs, QDs, QWDs)

<https://www.laserfocusworld.com/test-measurement/spectroscopy/article/16556856/quantumcascade-lasers-qcls-enable-applications-in-ir-spectroscopy>



Couple bulk semiconductor to electron injection layer:  
 $n_e \rightarrow m_{steps} \times n_y$

Light is emitted as electrons "cascade" through multiple quantum wells ( $m_{steps}$ )

Emitted light is IR~THz, normally mono-chromatic but tunable from  $3 \mu m \sim 12 \mu m$

Radiation resistant ([Radiation Physics and Chemistry 174](#), 2020, 108983)

Quantum dots and wells: <https://arxiv.org/abs/2202.11828>

submicron pixels

DoTPIx

= single n-channel MOS transistor, in which a buried quantum well gate performs two functions:

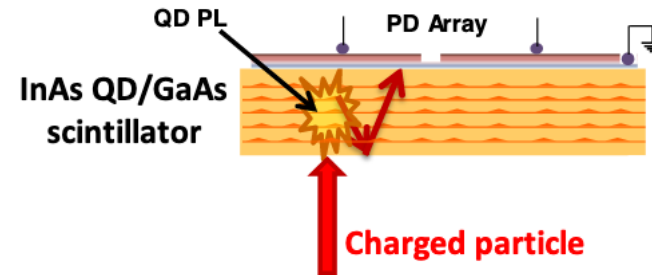
- as a hole-collecting electrode and
- as a channel current modulation gate

Novel Sensors for Particle Tracking: a Contribution to the Snowmass Community Planning Exercise of 2021

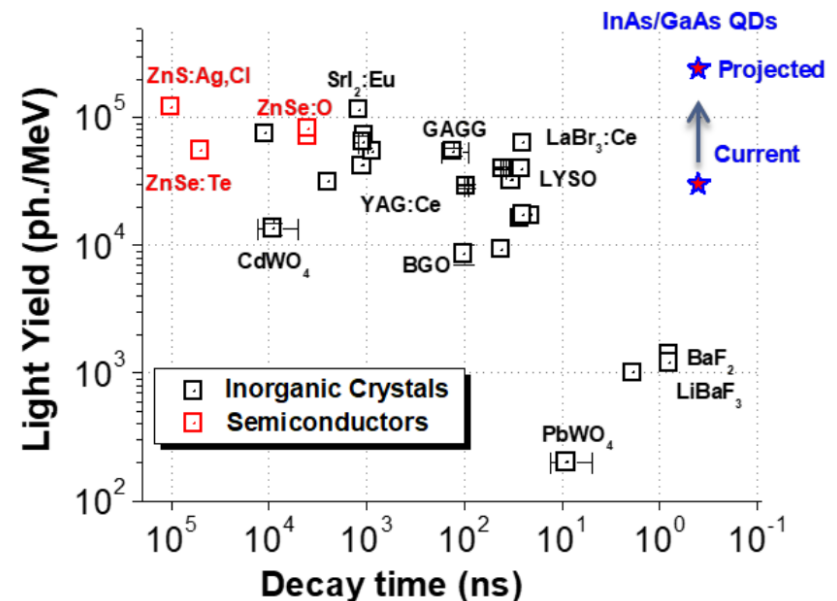
[M.R. Hoferkamp, S. Seidel, S. Kim, J. Metcalfe, A. Sumant, H. Kagan, W. Trischuk, M. Boscardin, G.-F. Dalla Betta, D.M.S. Sultan, N.T. Fourches, C. Renard, A. Barbier, T. Mahajan, A. Minns, V. Tokranov, M. Yakimov, S. Oktyabrsky, C. Gingu, P. Murat, M.T. Hedges](#)

<https://arxiv.org/abs/2202.11828>

scintillating (chromatic) tracker



IR emission from InAs QD's integrated PD's (1-2 μm thick)





# 2-D materials for MPGDs

Florian Brunbauer / CERN

State-of-the-art MPGDs:

- high spatial resolution
- good energy resolution
- timing resolution <25ps (PICOSEC Micromegas)

use of 2-D materials to improve:

- tailor the primary charge production process,
- protect sensitive photocathodes in harsh environments
- improve the performance of the amplification stage

tunable work function

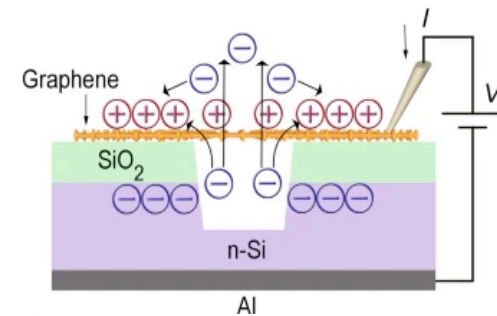
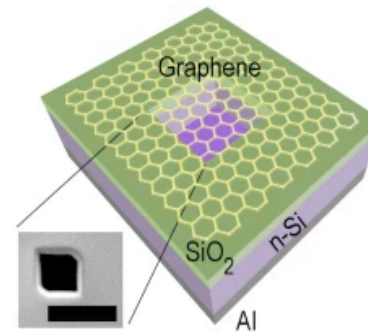
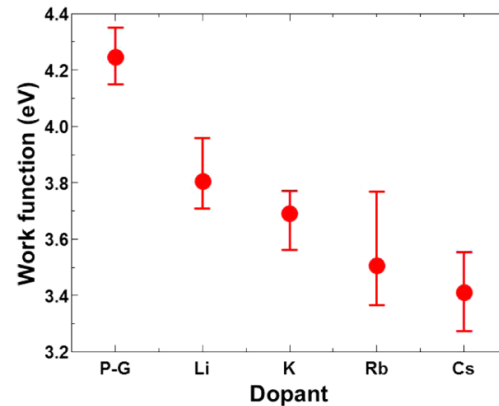
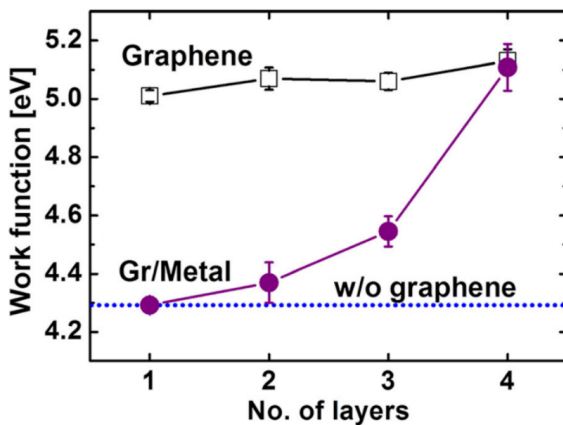
amplification

efficiency of the photocathode → timing resolution; QE  
tune via resonant processes in low dimensional coating structures

(additionally, encapsulation of semiconductive as well as metallic (i.e. Cu) photocathodes increases operational lifetime)

back flow of positive ions created during charge amplification to the drift region can lead to significant distortions of electric fields

Graphene has been proposed as selective filter to **suppress** ion back flow while **permitting** electrons to pass:  
Good transparency (up to ~99.9%) to very low energy (<3 eV) electrons (?)



Tuning the work function of graphene toward application as anode and cathode, Samira Naghdi, Gonzalo Sanchez-Arriaga, Kyong Yop Rhee, <https://arxiv.org/abs/1905.06594>

Space charge neutralization by electron-transparent suspended graphene, Siwapon Srisophonpan, Myungji Kim & Hong Koo Kim, [Scientific Reports 4, 3764 \(2014\)](https://doi.org/10.1038/s41598-014-03764-4)

## Metamaterials, 0 / 1 / 2-dimensional materials (quantum dots, nanolayers)

ultra-fast scintillators based on perovskites

chromatic calorimetry (QDs)

active scintillators (QCL, QWs, QDs)

GEMs (graphene)

## Atoms, molecules, ions

Rydberg TPC's

## Spin-based sensors

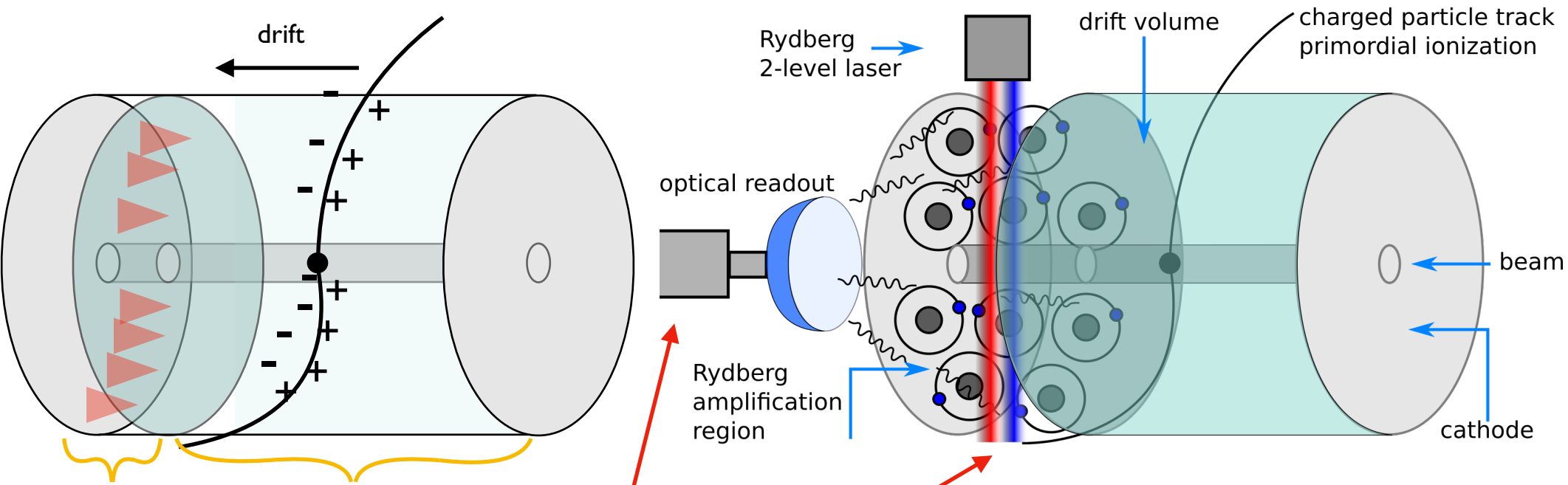
helicity detectors

## Superconducting sensors

# Rydberg atom TPC's

Georgy Kornakov / WUT

Act on the **amplification** region



enhanced electron signal through “priming” of gas in **amplification** region:  $\longrightarrow$  effective reduction of ionization threshold of gas in amplification region  
 $\longrightarrow$  higher electron yield

Rydberg atoms can serve to up-convert THz / GHz radiation into the optical regime  $\longrightarrow$  optical R/O of avalanche intensities

# Rydberg atom TPC's

Georgy Kornakov / WUT

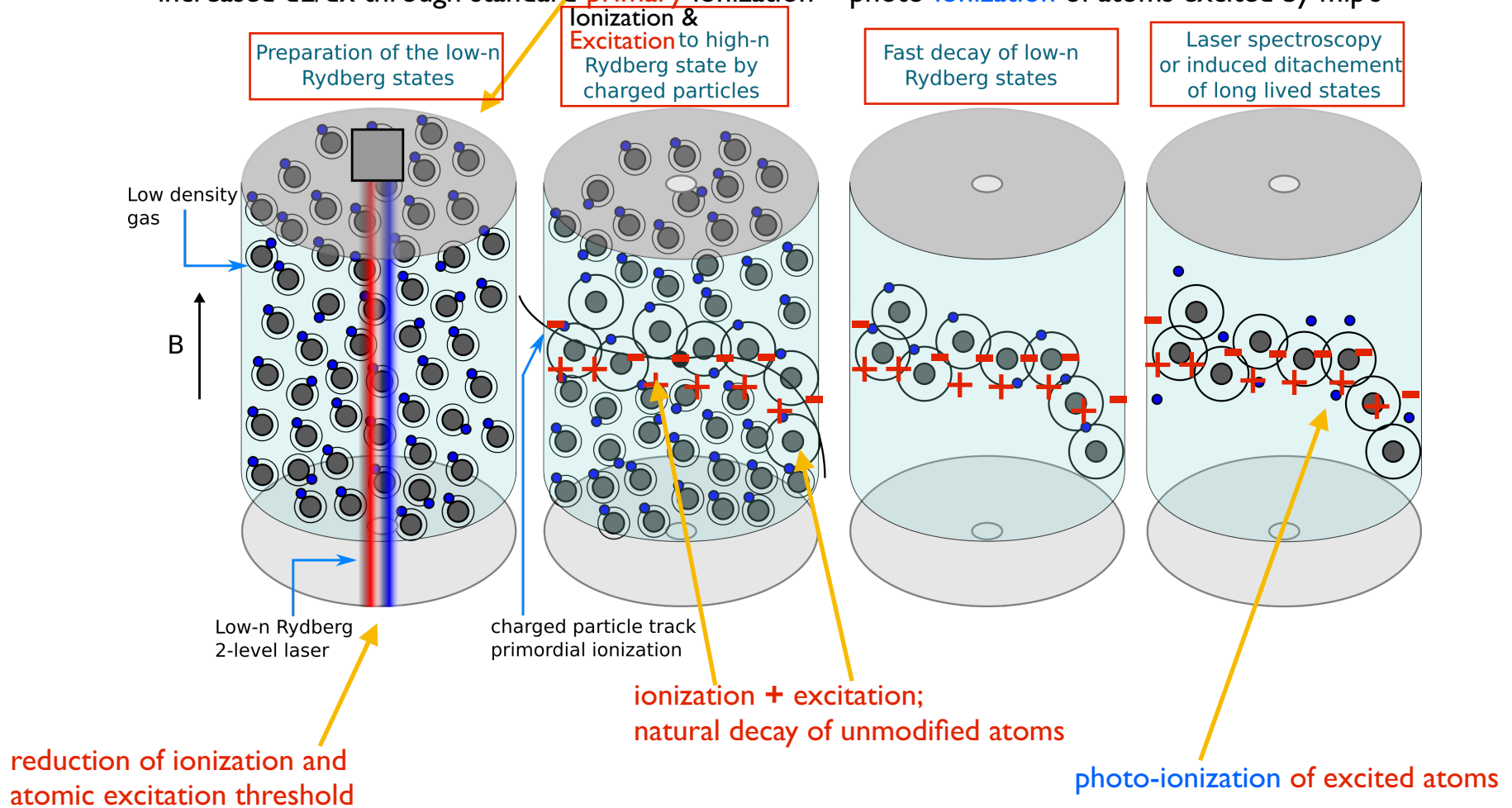
Act on the drift region

principle carries over to drift region:

enhanced electron signal through “priming” of gas in drift region:

effective reduction of ionization threshold of gas in drift region

increased  $dE/dx$  through standard primary ionization + photo-ionization of atoms excited by mip's



Metamaterials, 0 / 1 / 2-dimensional materials (quantum dots, nanolayers)

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active scintillators (QCL, QWs, QDs)

GEMs (graphene)

Atoms, molecules, ions

Rydberg TPC's

Spin-based sensors

helicity detectors

Superconducting sensors

# optically polarizable elements: Nitrogen-vacancy diamonds (NVD)

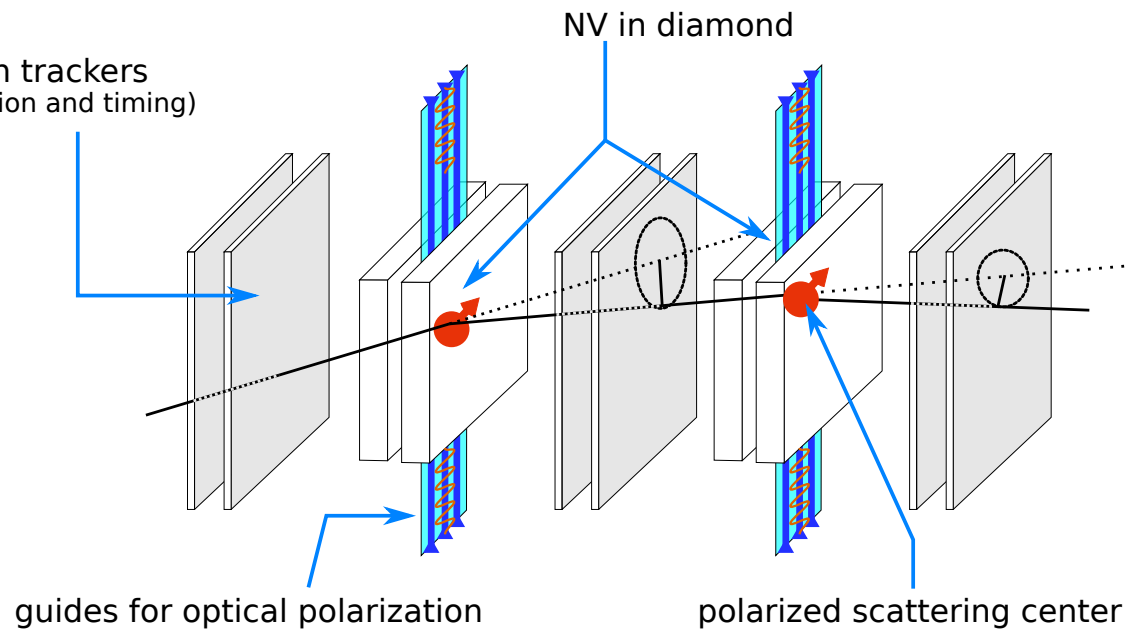
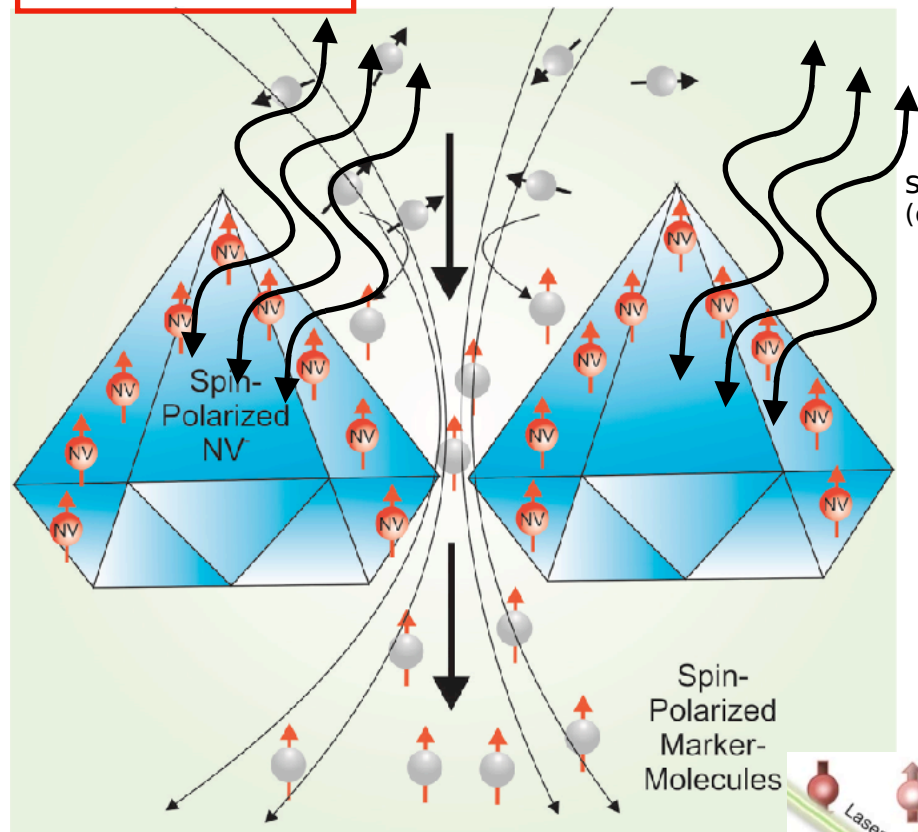
Georgy Kornakov / WUT

spin-spin scattering for helicity determination: usually with polarized beams and/or polarized targets

introduce polarized scattering planes to extract track-by-track particle helicity

$10^{16} \sim 10^{18} / \text{cm}^3$

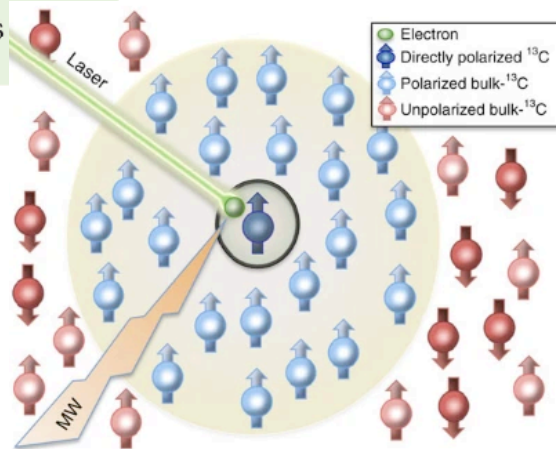
use NV-diamonds to scatter mips (spin-dependent scattering)



© Dr. Christoph Nebel, Fraunhofer IAF

[https://www.metaboliqs.eu/en/news-events/MetaboliQs\\_PM\\_first\\_year.html](https://www.metaboliqs.eu/en/news-events/MetaboliQs_PM_first_year.html)

Diamond plates of up to  $8 \times 8 \text{ mm}^2$  in size, fabricated by Element Six



Local and bulk  $^{13}\text{C}$  hyperpolarization in nitrogen-vacancy centred diamonds at variable fields and orientations, G. Alvarez et al., *Nature Communications* **6**, 8456 (2015)

<https://www.nature.com/articles/ncomms9456>

$\times 10^2$

## Metamaterials, 0 / 1 / 2-dimensional materials (quantum dots, nanolayers)

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active scintillators (QCL, QWs, QDs)

GEMs (graphene)

## Atoms, molecules, ions

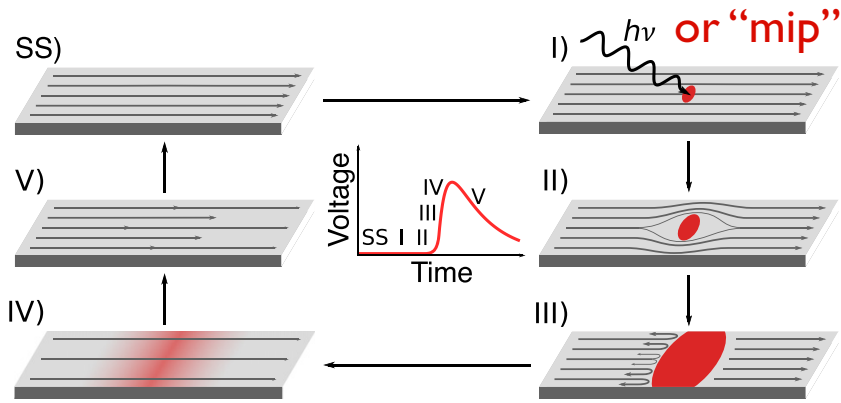
Rydberg TPC's

## Spin-based sensors

helicity detectors

## Superconducting sensors

# Extremely low energy threshold detectors: SNSPD



## SNSPD's Near term future

Parameter	SOA 2020	Goal by 2025
Efficiency	98% @ 1550nm	>80% @ 10 $\mu$ m
Energy Threshold	0.125 eV (10 $\mu$ m)	12.5 meV (100 $\mu$ m)
Timing Jitter	2.7 ps	< 1ps
Active Area	1 mm <sup>2</sup>	100 cm <sup>2</sup>
Max Count Rate	1.2 Gcps	100 Gcps
Pixel Count	1 kilopixel	16 megapixel
Operating Temperature	4.3K	25 K

Snowmass2021 - Letter of Interest

### Superconducting Nanowire Single-Photon Detectors

Moving to SC strips conventional lithography  $\rightarrow$  scale up  
Development towards SC SSPM

QT4HEP22-- I. Shipsey

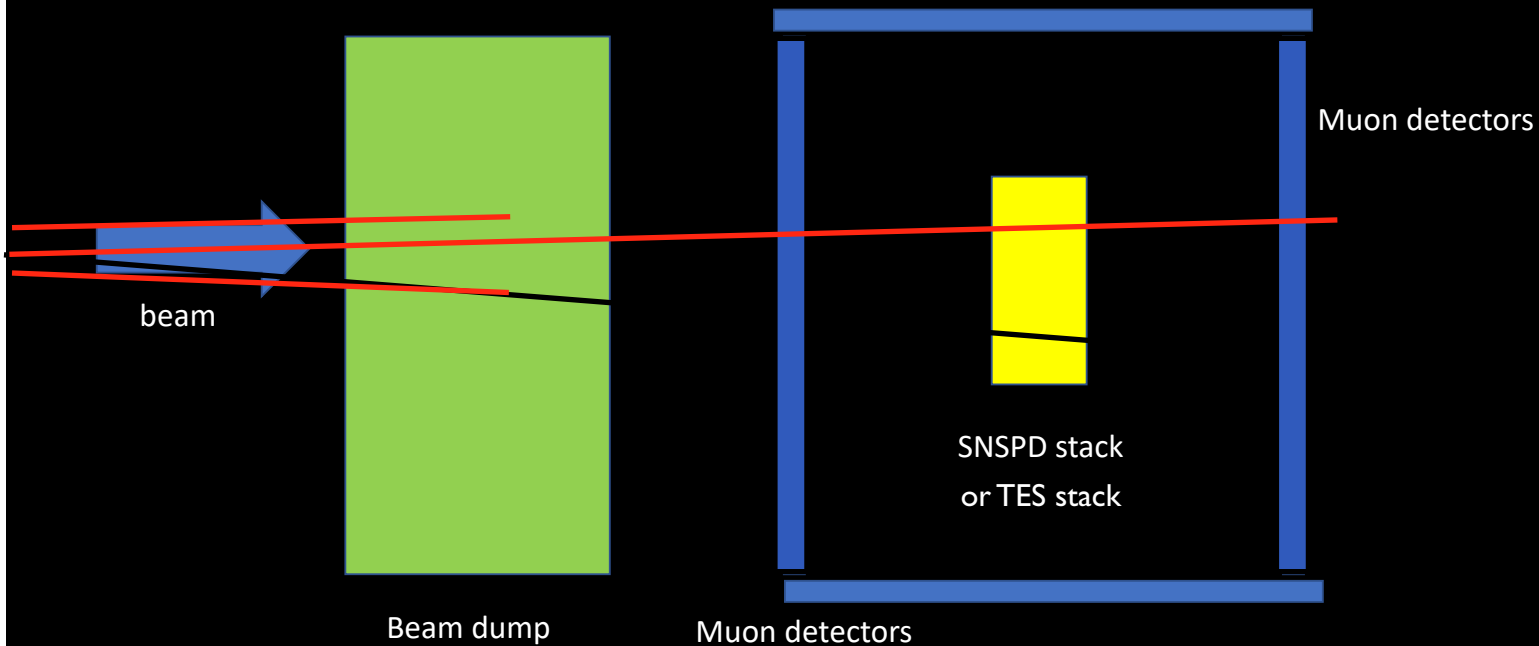
125

#### Contact Information:

Karl Berggren, berggren@mit.edu  
Ilya Charaev, charaev@mit.edu  
Jeff Chiles, jeffrey.chiles@nist.gov  
Sae Woo Nam, saewoo.nam@nist.gov  
Valentine Novosad, novosad@anl.gov  
Boris Korzh, bkorzh@jpl.nasa.gov  
Matt Shaw, mattshaw@jpl.nasa.gov

Search for Beyond Standard Model **milli-charged particles?**

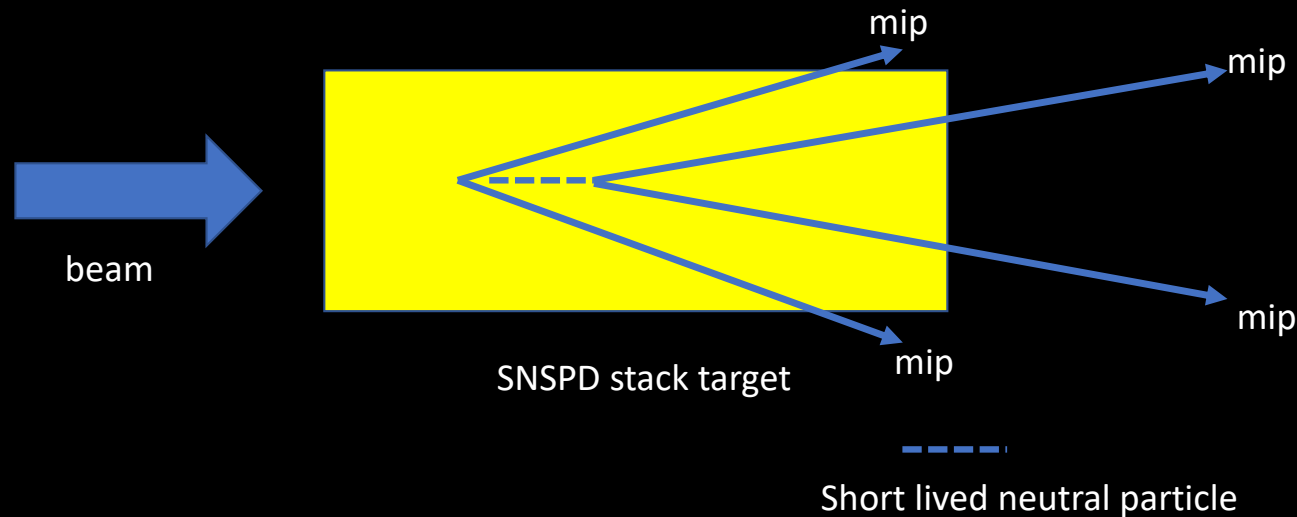
**mip: ~20 keV/100  $\mu$ m**





# Extremely low energy threshold detectors: SNSPD

A way to measure the lifetime of very short-lived particles?



a fixed target experiment with a very thinly layered ( $\sim 10$  nm layers) SNSPDs as target and make a thick stack perhaps a mm thick: very short-lived neutral particles would appear as a  $n \times 10$  nm gap in the signal plane stack between where the mip projectile interacts and the short-lived particle decays into mips. Addition of a B-field helpful

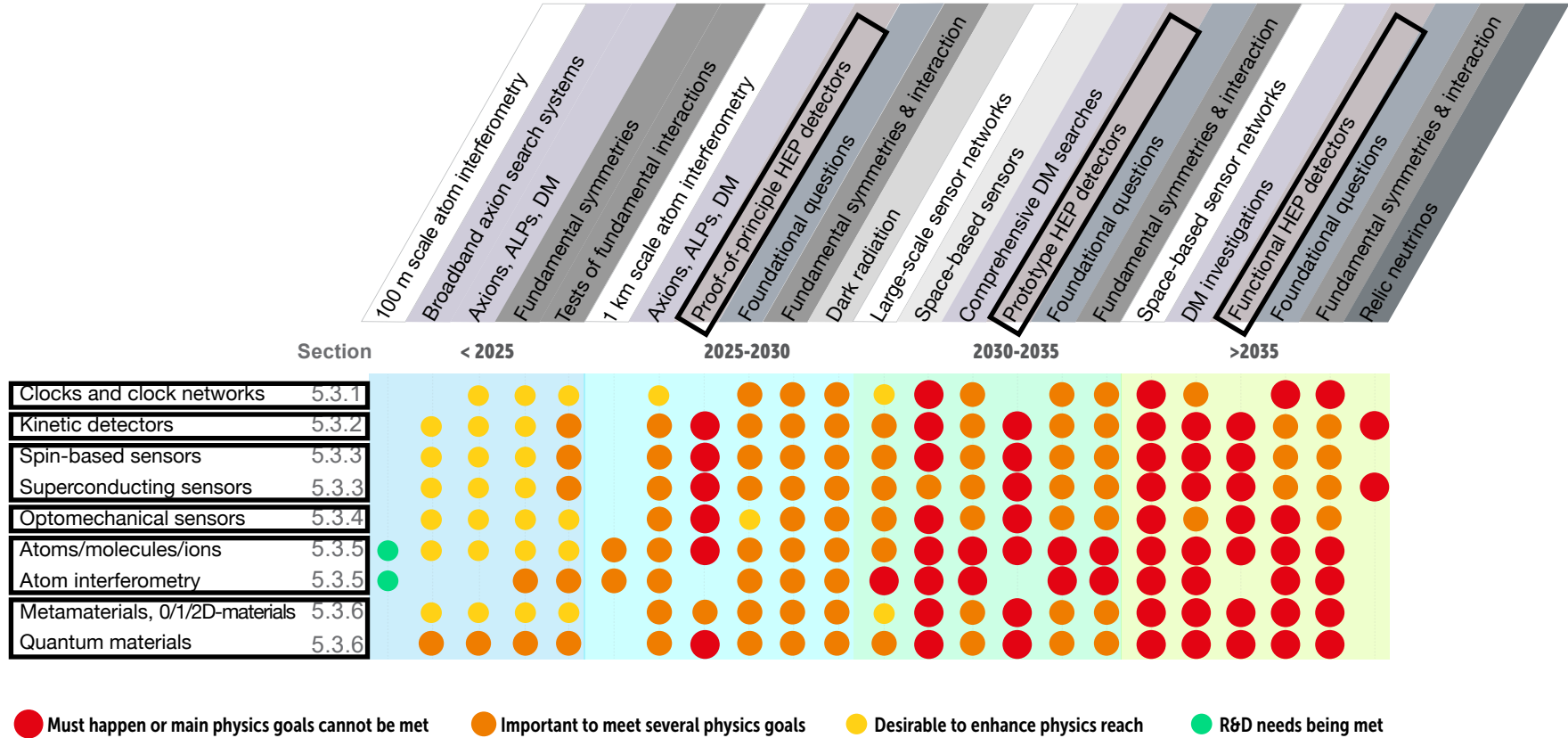
QT4HEP22-- I. Shipsey

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# RECFA Detector R&D roadmap 2021

<https://cds.cern.ch/record/2784893>

## Chapter 5: Quantum and Emerging Technologies Detectors



## Chapter 4: Particle Identification and Photon Detectors

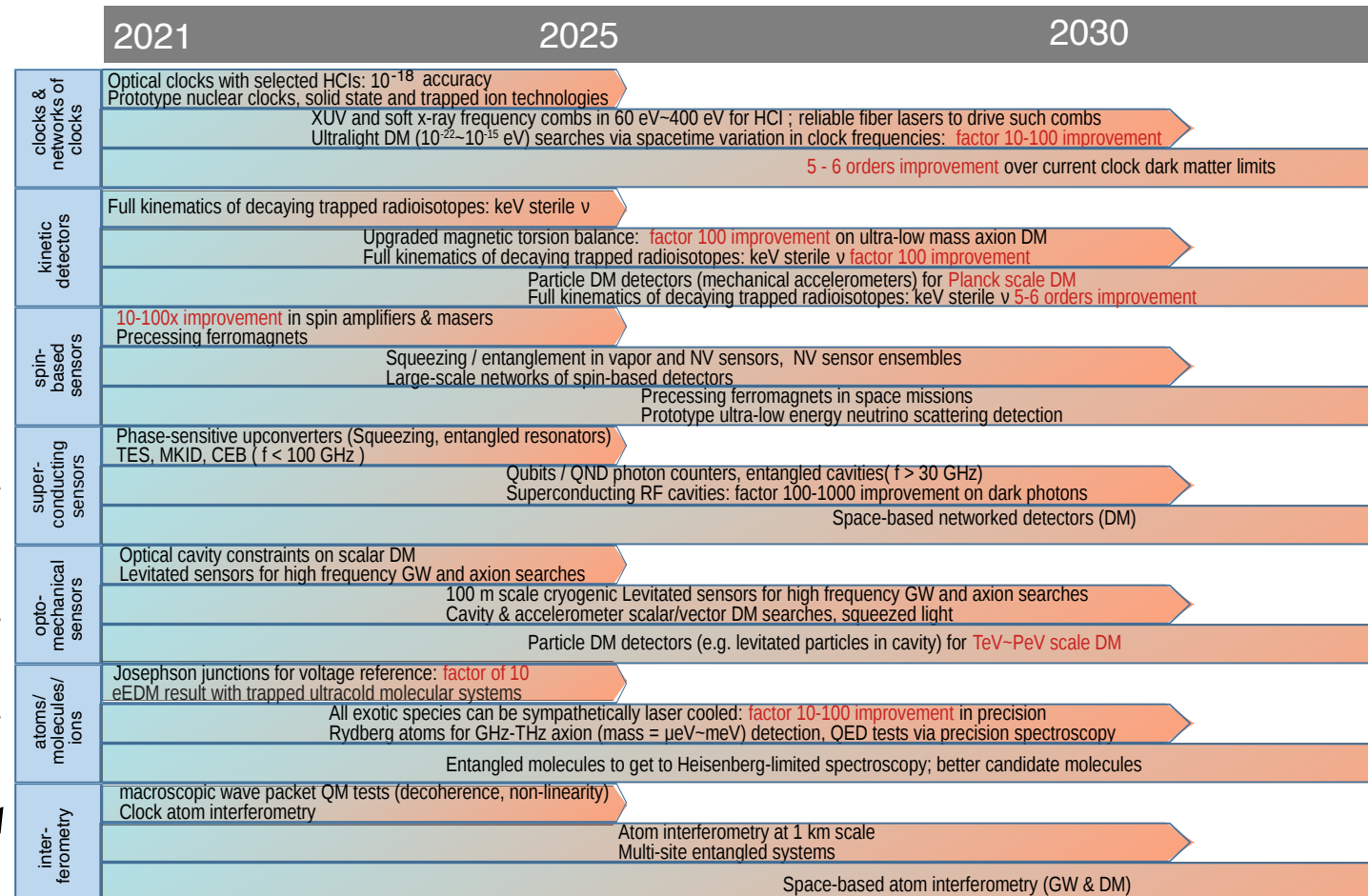
It is recommended that several “blue-sky” R&D activities be pursued. The development of solid state photon detectors from novel materials is an important future line of research, as is the development of cryogenic superconducting photosensors for accelerator-based experiments. Regarding advances in PID techniques, gaseous photon detectors for visible light should be advanced. Meta-materials such as photonic crystals should be developed, giving tune-able refractive indices for PID at high momentum. Finally, for TRD imaging detectors, the detection of transition radiation with silicon sensors is an important line of future research.

## What's next?

These potential applications of quantum sensors also in HEP require dedicated R&D to evaluate their potential and feasibility.

In line with the RECFA R&D roadmap, it makes sense to consider a quantum-sensing R&D program that brings together the following strands:

- Clocks and clock networks 5.3.1
- Kinetic detectors 5.3.2
- Spin-based sensors 5.3.3
- Superconducting sensors 5.3.3
- Optomechanical sensors 5.3.4
- Atoms/molecules/ions 5.3.5
- Atom interferometry 5.3.5
- Metamaterials, 0/1/2D-materials
- Quantum materials 5.3.6



also for HEP!

thank you!

# quantum sensors (an electromagnetic perspective)

	Microwave	Submillimetre	Far infrared	Optical	High energy
	10 – 100 GHz 3 cm- 3 mm	100 GHz – 1 THz 3 mm – 300 μm	1 – 10 THz 300 – 30 μm	2 μm – 300 nm	UV, Yray and Xray
SIS mixers		●			
HEB			●		
CEB		●			
TES	●	●	●	●	●
KID	●	●	●	●	
SNSPD			●	●	
SQUID	●				
JJPA	●				
TWPA	●	●			

Stafford Withington (Cambridge)

## ECFA Detector R&D Roadmap Symposium of Task Force 5 Quantum and Emerging Technologies

Symposium: April 12, 2022

<https://indico.cern.ch/event/999818/>

14 presentations  
first block covering physics landscape  
following blocks focusing on technologies  
discussion of three important points

### ECFA Detector R&D Roadmap Symposium of Task Force 5: Quantum and emerging technologies

Monday 12 Apr 2021, 09:00 → 18:30 Europe/Zurich

09:00 → 09:15 Introduction

09:15 → 11:00 science targets – Overview and Landscape

9:15 EDM searches & tests of fundamental symmetries Peter Fierlinger / TU Munich

9:45 Tests of QM [wavefunction collapse, size effects, temporal separation, decoherence] Angelo Bassi

10:15 Multimessenger detection [including atom interferometer or magnetometer networks] Giovanni Barontoni / Birmingham

10:45 Axion and other DM (as well as non-DM Ultra-light) particle searches Mina Arvanitaki / Perimeter Institute

11:15 → 11:30 Coffee break

11:30 → 12:30 Experimental methods and techniques - Overview and Landscape

11:30 Precision spectroscopy and clocks, networks of sensors and of entangled systems [optical atomic clocks] David Hume / NIST

12:00 Novel ionic, atomic and molecular systems [RaF, multiatomic molecules, exotic atoms] Marianna Safranova / U. Delaware

12:30 → 13:30 Lunch break

13:30 → 16:00 Experimental and technological challenges, New Developments

13:30 Superconducting platforms [detectors: TES, SNSPD, Haloscopes, including single photon detection] Alexander Romanenko (FNAL)

14:00 High sensitivity superconducting cryogenic electronics, low noise amplifiers Stafford Withington / Cambridge

14:30 Broadband axion detection Kent Irwin / Stanford

15:00 Mechanical / optomechanical detectors Andrew Geraci / Northwestern

15:30 Spin-based techniques, NV-diamonds, Magnetometry Dima Budker / Mainz

16:00 → 16:15 Coffee break

16:15 → 18:30 Experimental and technological challenges, New Developments

16:15 Calorimetric techniques for neutrinos and axions potential speaker identified

16:35 Quantum techniques for scintillators potential speaker identified

16:55 Atom interferometry at large scales (ground based, space based) Jason Hogan / Stanford

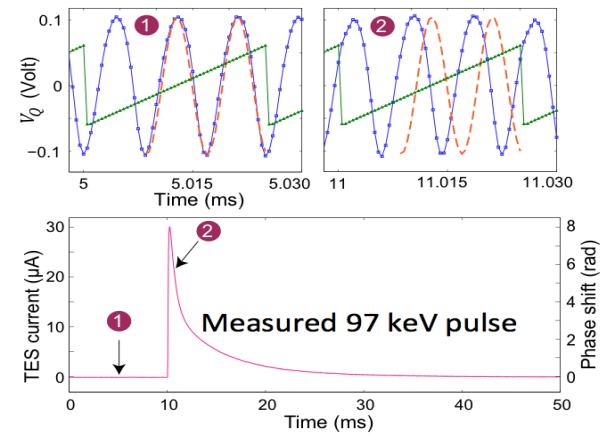
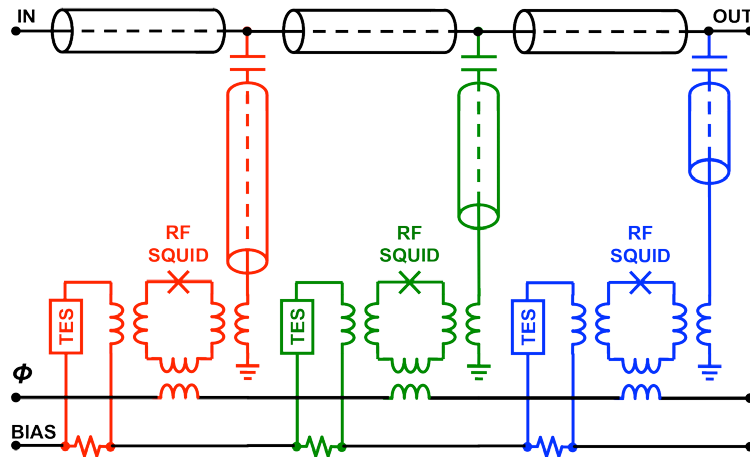
17:25 → 18:15 Discussion session : discussion points

- Scaling up from table-top systems
- Networking – identifying commonalities with neighboring communities
- Applying quantum technologies to high energy detectors

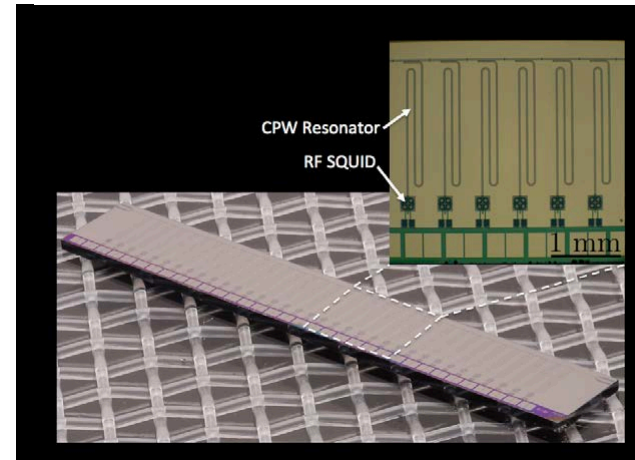
18:15 → 18:30 Wrap-up

## Scalable readout is the key for very large arrays

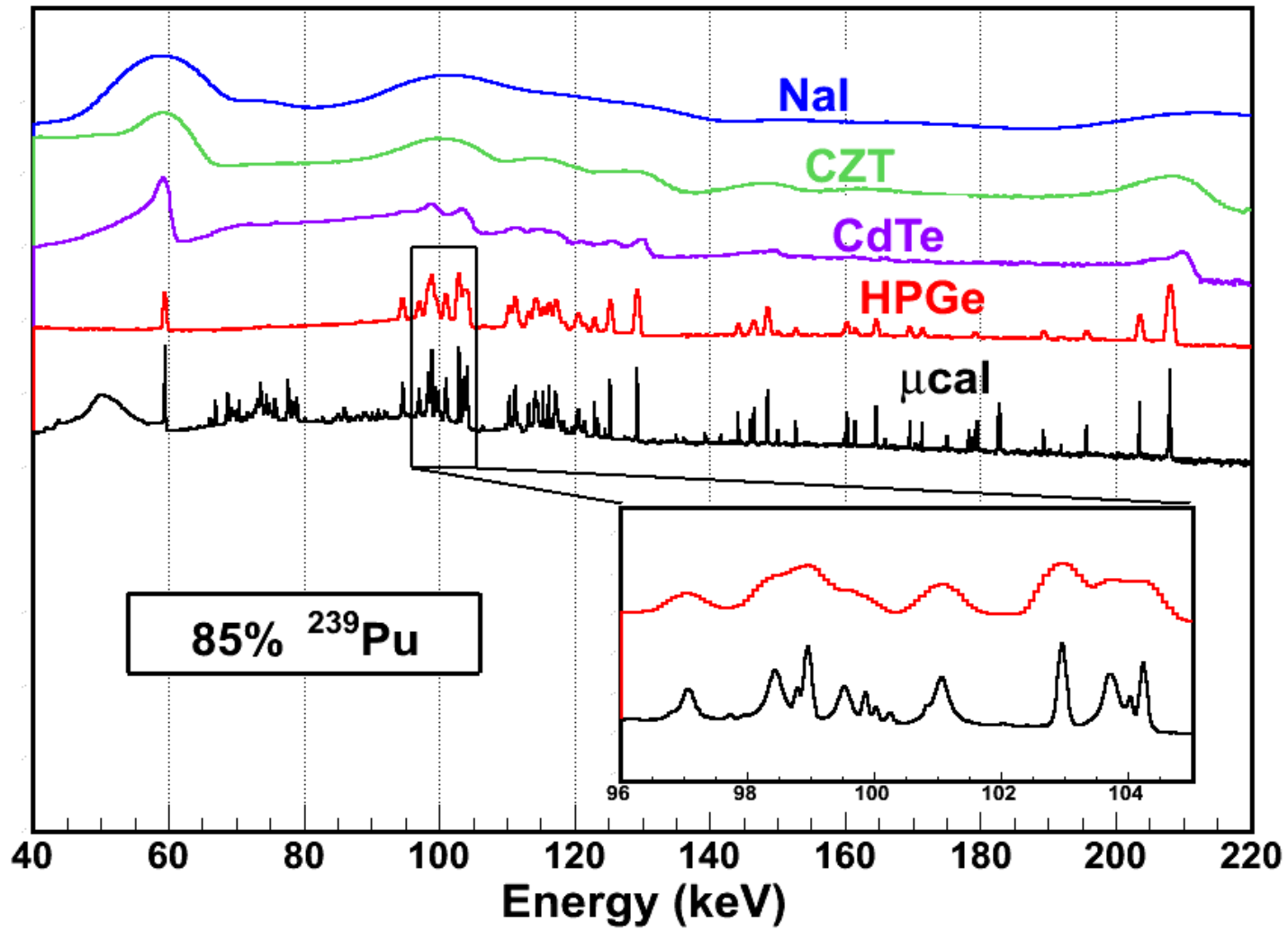
*Four innovations:  $\mu$ Cal, Microwave Multiplexing, RF-SQUIDS, Software Radio*



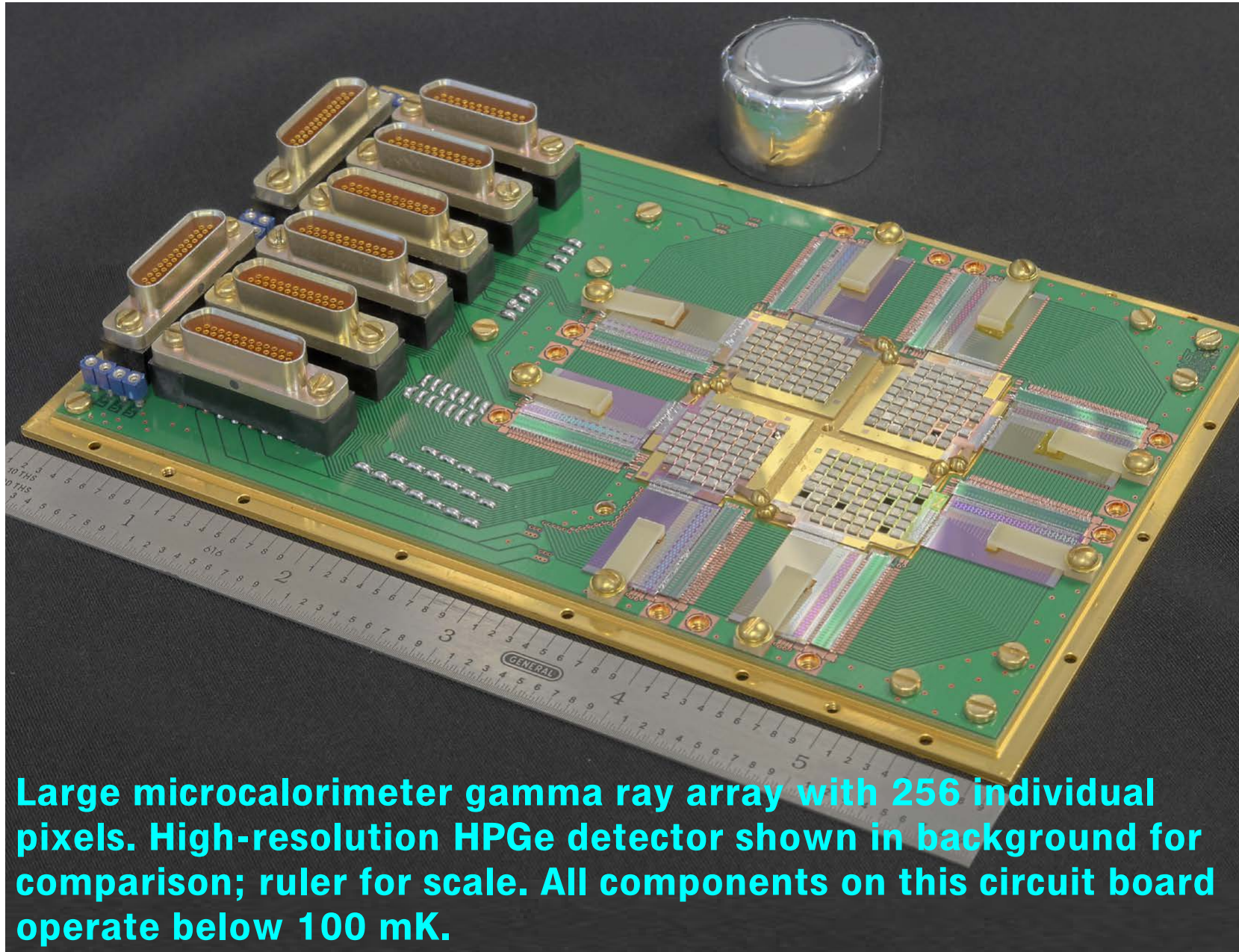
- UCSB: 1,000 low-resolution sensors and softwaredefined radio (SDR) readout
- CU: 2 high-resolution sensors + RFSQUIDS without SDR
- CU + LANL: Demonstrated negligible readout noise from RF-SQUID
- *No one has combined all together at scale*



What is your technical approach? Describe background and your techniques.





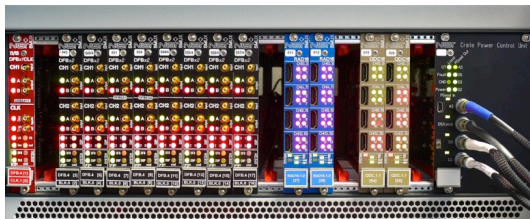
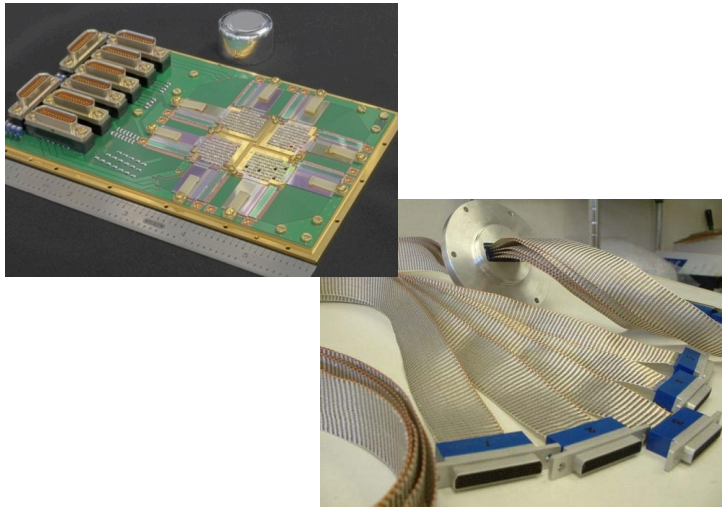


**Large microcalorimeter gamma ray array with 256 individual pixels. High-resolution HPGe detector shown in background for comparison; ruler for scale. All components on this circuit board operate below 100 mK.**

## Microwave readout reduces complexity and cost

### Present Techniques

- 500 wires per 1,000 sensors
- Elaborate wiring harnesses and connectors
- Unique control electronics



### Microwave Techniques

- 1,000 resonators on two coaxial cables!
- Existing commercial electronics
- Functionality shared with telecommunications industry
- Benefit from steady improvement in FPGAs, ADCs, and DACs

