

Frustrated Dark Matter

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Next Generation Models

- (I) the model should preferably be a theoretically consistent extension of one of the DM simplified models already used by the LHC Collaborations;
- (II) the model should still be generic enough to be used in the context of broader, more complete theoretical frameworks;
- (III) the model should have a sufficiently varied phenomenology to encourage comparison of different experimental signals and to search for DM in new, unexplored channels;
- (IV) the model should be of interest beyond the DM community, to the point that other direct and indirect constraints can be identified.

Frustrated Dark Matter

All mediator fields coupling both to χ and to SM fields carry SM gauge charges that preclude renormalizable gauge-invariant interactions between the dark matter and any SM fermion.

Interactions of the dark matter are frustrated in the sense that the specific mediator assignments preclude its tree level interaction with the SM

Bipartite Mediator Sector

$$\text{SM} \longleftrightarrow \text{mediators} \left\{ \begin{array}{l} \varphi \text{ (scalar)} \\ \psi \text{ (Dirac)} \end{array} \right\} \longleftrightarrow \text{DM } \chi.$$

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{med}} + \mathcal{L}_{\chi},$$

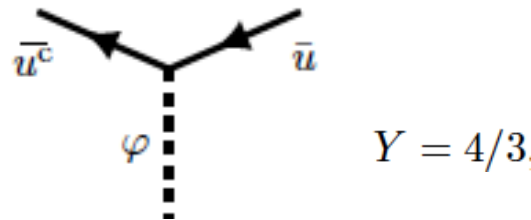
$$\mathcal{L}_{\text{med}} = (D_{\mu}\varphi)^{\dagger s}(D^{\mu}\varphi)_s - m_{\varphi}^2\varphi^{\dagger s}\varphi_s + \bar{\psi}^s(i\not{D} - m_{\psi})\psi_s + \mathcal{L}_{\text{decay}}$$

$$\mathcal{L}_{\chi} = \bar{\chi}(i\not{D} - m_{\chi})\chi + y_{\chi}(\varphi^{\dagger s}\bar{\chi}\psi_s + \text{H.c.})$$

Sextet Mediators

Field	Description	$SU(3)_c \times SU(2)_L \times U(1)_Y$ representation	Couples to SM?
χ	Dark matter	$(1, 1, 0)$	
φ	Scalar mediator	$(\mathbf{6}, 1, \frac{4}{3})$	✓
ψ	Dirac mediator		

$$\mathcal{L}_{\text{decay}} = \lambda_{IJ} K_s^{ij} \varphi^{\dagger s} \bar{q}_{RI}^c q_{RJ} + \text{H.c.} \quad \text{with} \quad q \in \{u, d\},$$



Why Sextets ?

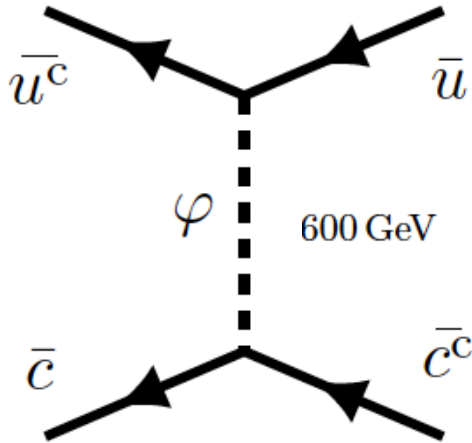
One of few BSM extensions that allows renormalizable couplings of BSM state to quarks (interesting BSM extension just for this reason)

Allows Frustrated DM scenario through renormalizable coupling of scalar messenger only

Less explored and less constrained but perfectly allowed extension of SU(3)

Interesting phenomenological feature of 'SUSY-like' signatures but with heavier fermionic particles

$D^0-\bar{D}^0$ mixing



$$(\lambda_{11}\lambda_{22})^2 \leq 9.3 \times 10^{-7} \left(\frac{m_\varphi}{\text{TeV}}\right)^2$$

★ $\lambda_{22} = 0$

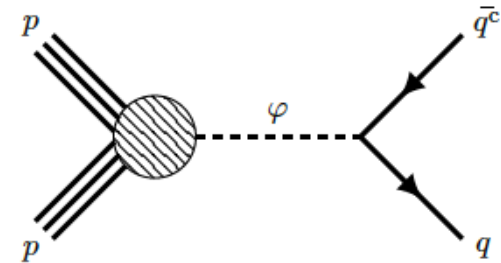
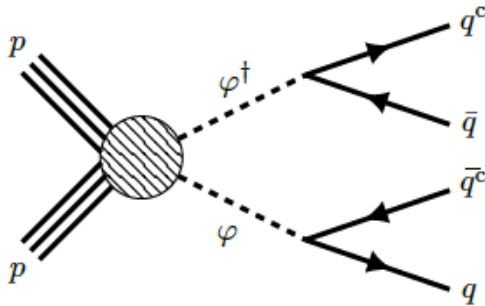
Choose MFV framework to kill mixing

Scalar Mediators at LHC

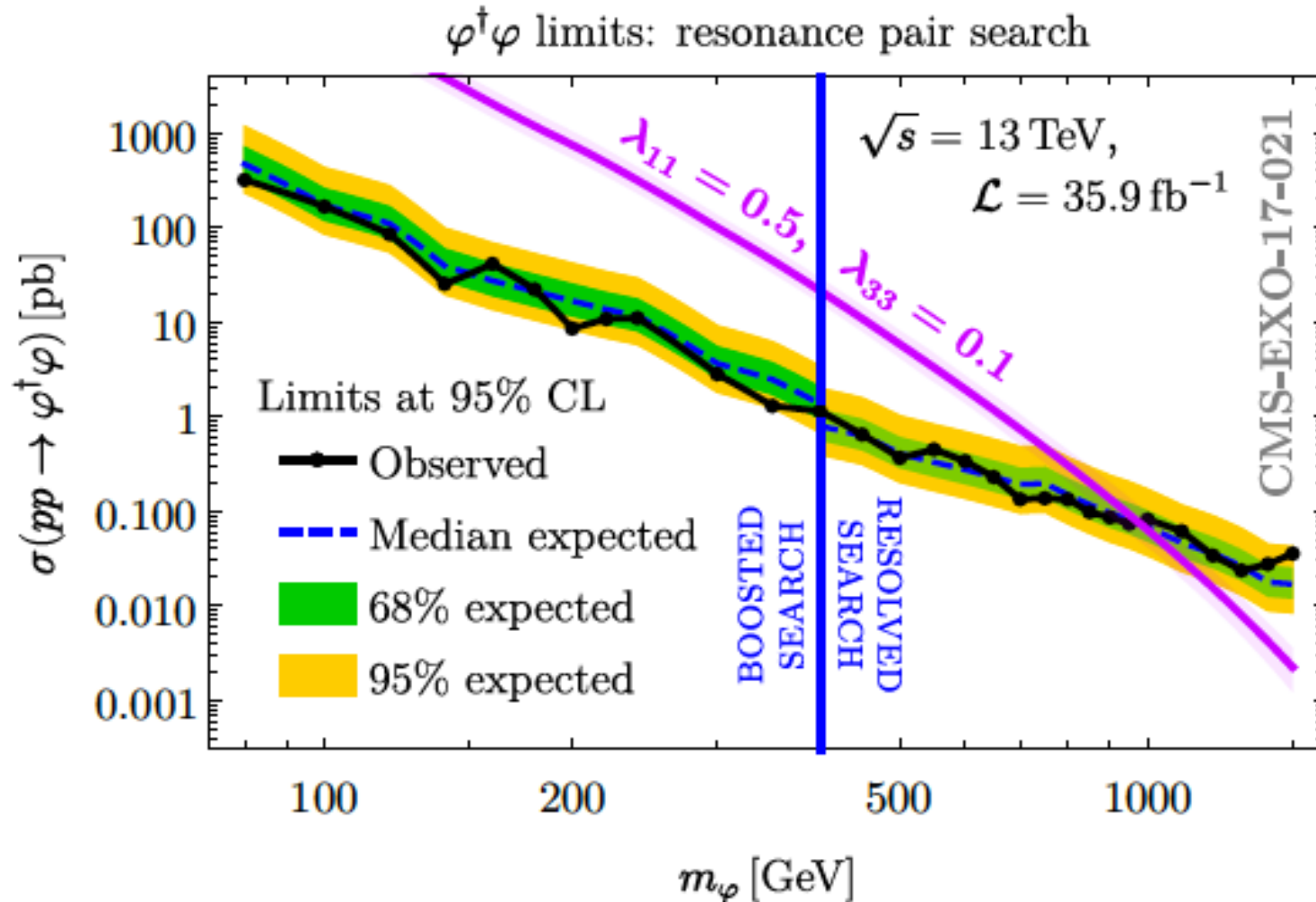
$$gg \rightarrow \varphi^\dagger \varphi$$

$$pp \rightarrow \varphi^\dagger \varphi \rightarrow uu\bar{u}\bar{u} \ (tt\bar{t}\bar{t})^{\star} \quad pp \rightarrow \varphi \rightarrow uu \ (tt)^{\star}$$

$$pp \rightarrow \varphi^\dagger \varphi \rightarrow uutt \ ?$$



Scalar Mediator 4 quark Searches

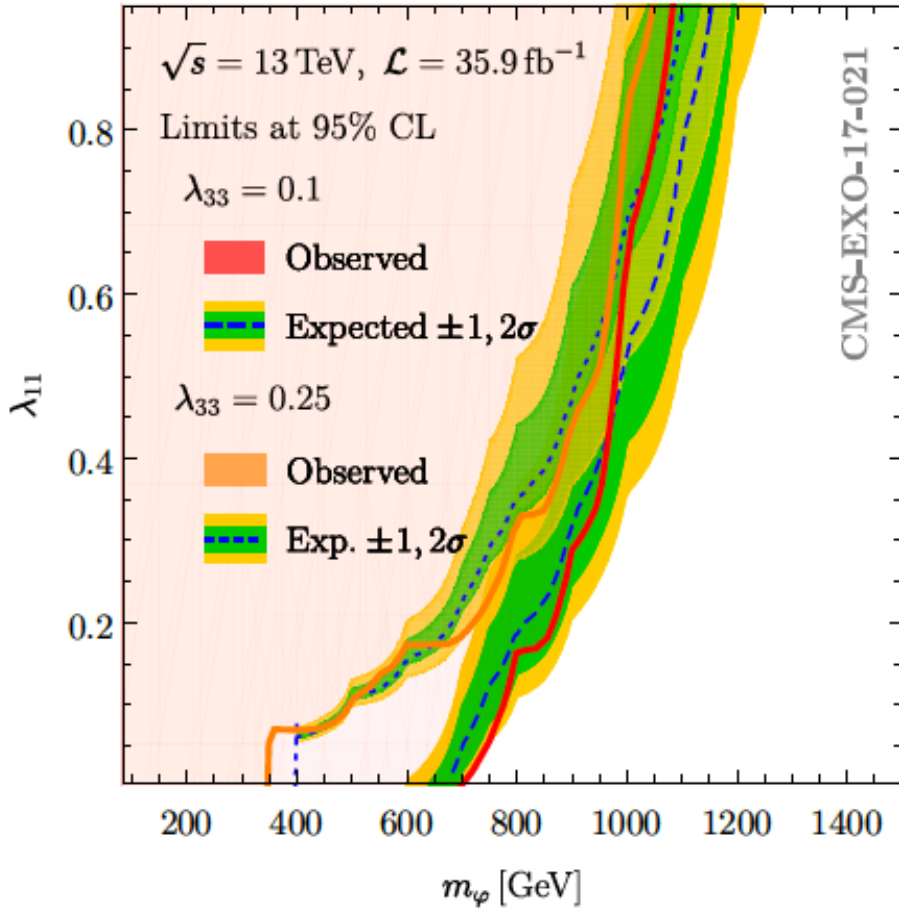


CMS-EXO-17-021

$\sigma(pp \rightarrow \varphi^\dagger\varphi \rightarrow uu\bar{u}\bar{u})$ for $\lambda_{11} = 0.5, \lambda_{33} = 0.1$

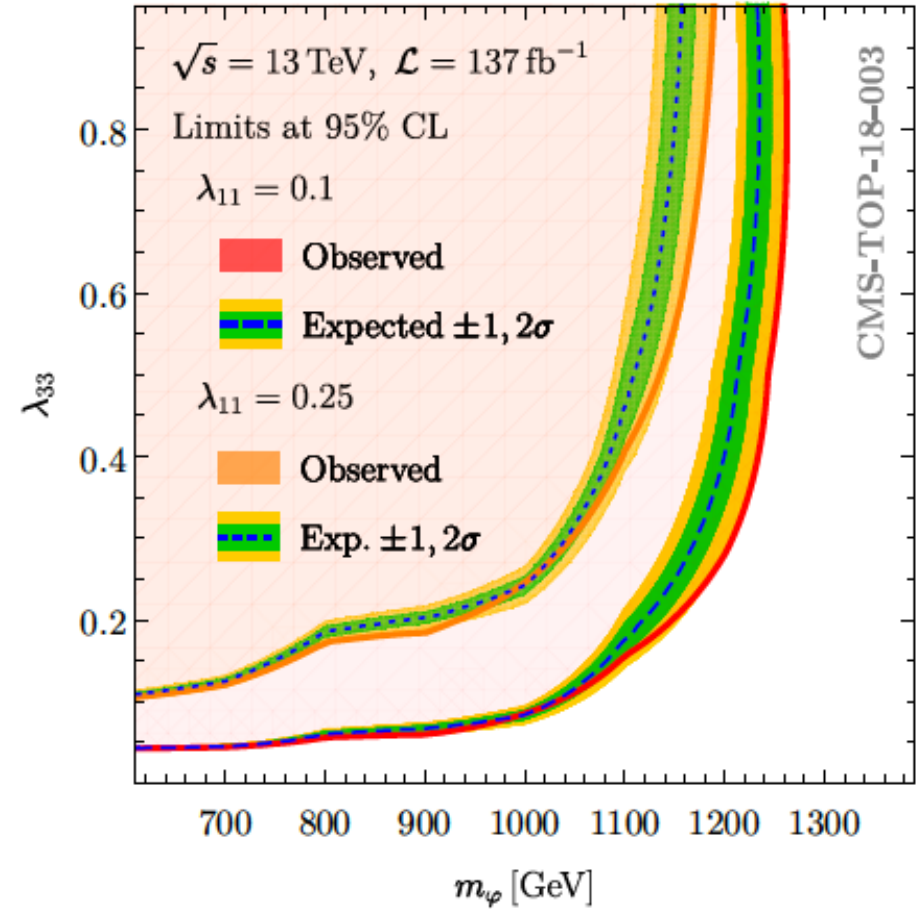
Scalar Mediator 4 quark Searches

$\varphi^\dagger\varphi$ limits: resonance pair search



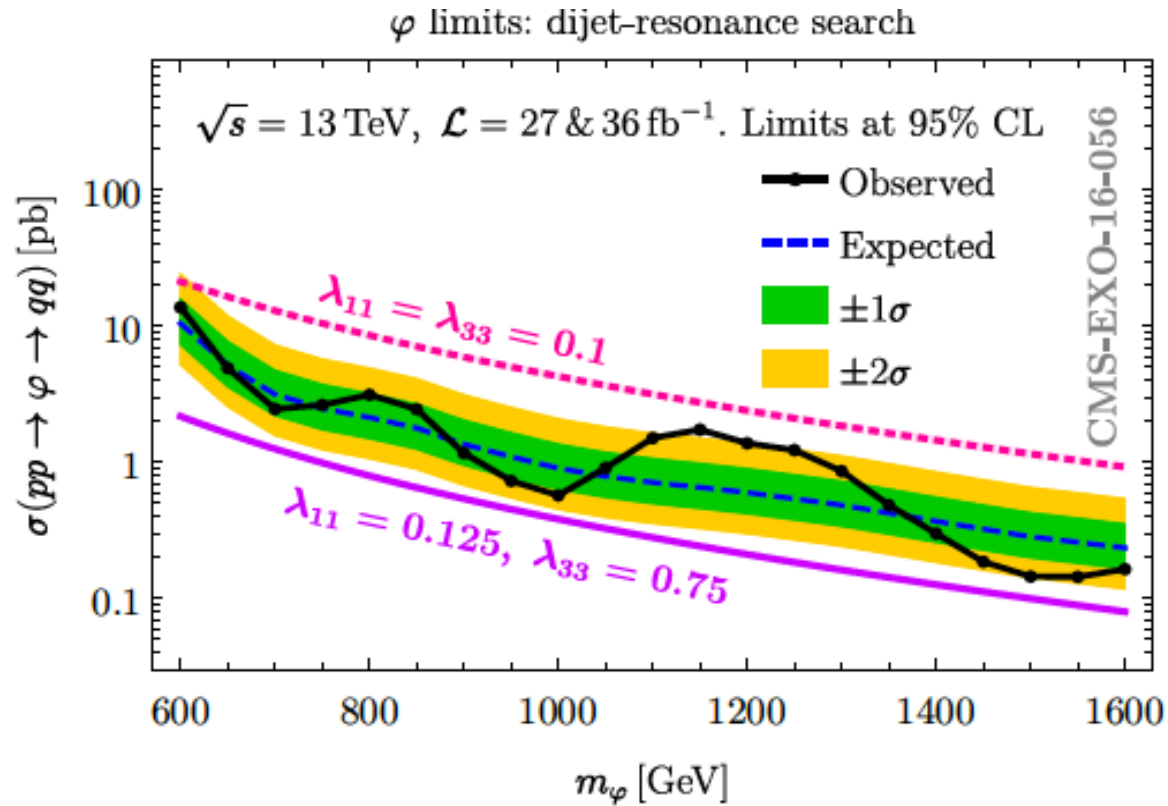
$$\sigma(pp \rightarrow \varphi^\dagger\varphi \rightarrow uu\bar{u}\bar{u})$$

$\varphi^\dagger\varphi$ limits: $\sigma(t\bar{t}t\bar{t})$ measurement

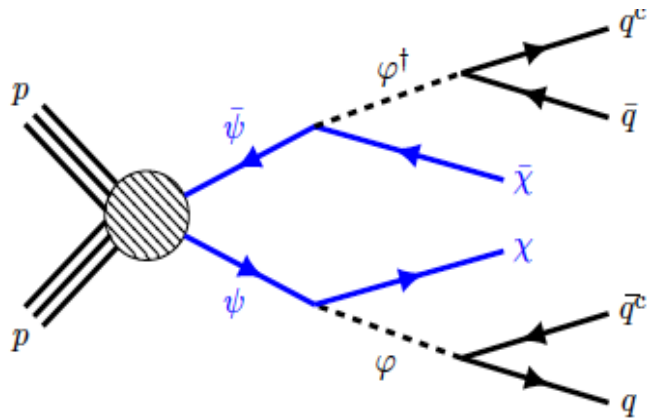


$$\sigma(pp \rightarrow \varphi^\dagger\varphi \rightarrow t\bar{t}t\bar{t})$$

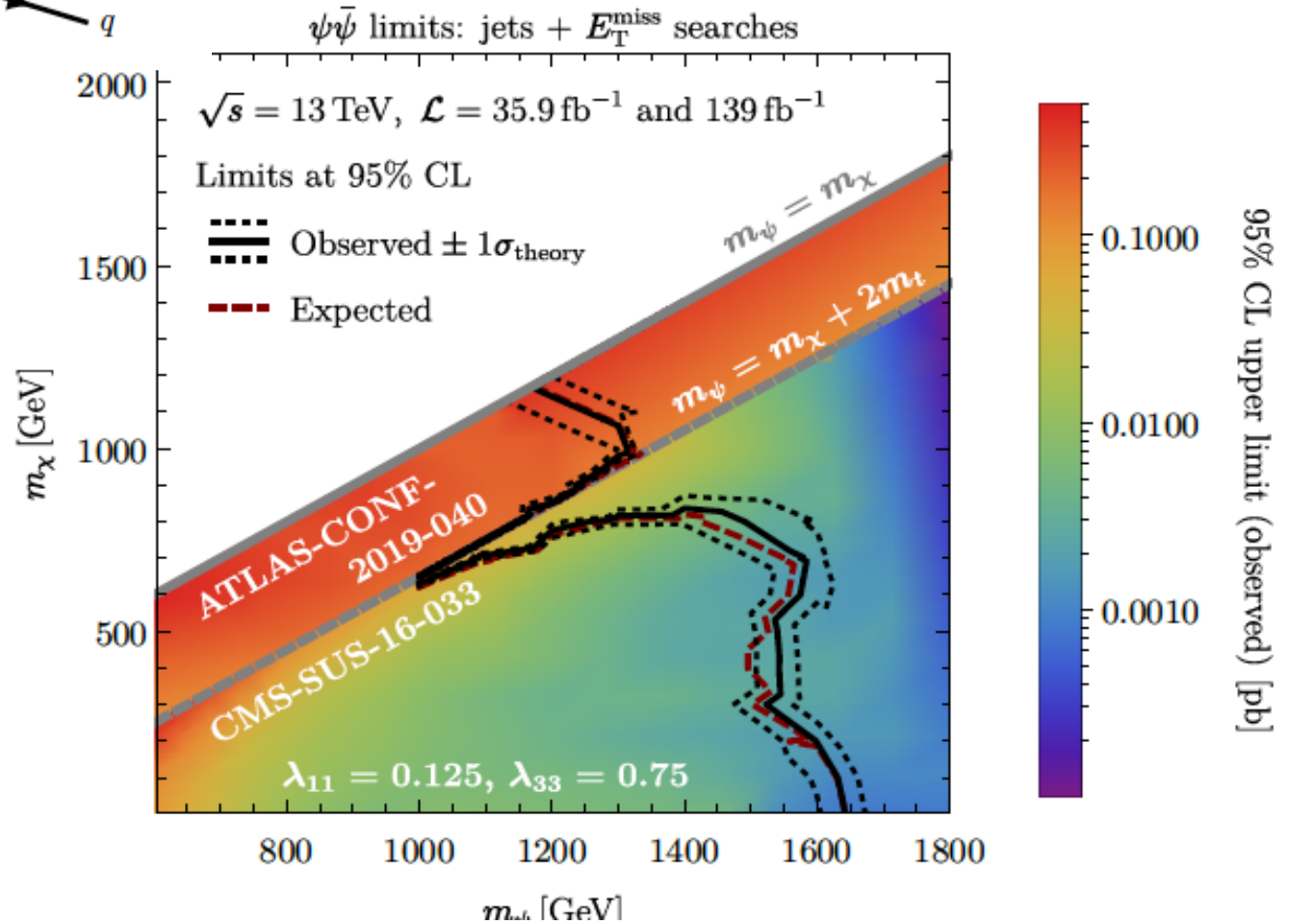
Scalar Mediator di-jet searches



Fermionic Mediators at LHC

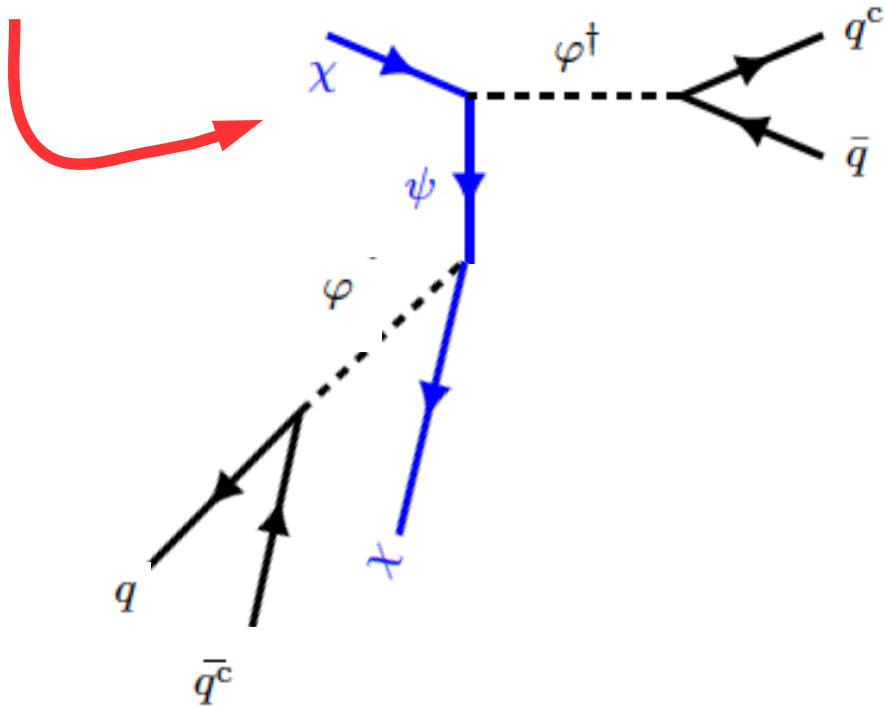


Limits computed from
 ATLAS-CONF-2019-040
 CMS-SUS-16-033
 SUSY jets + MET channels



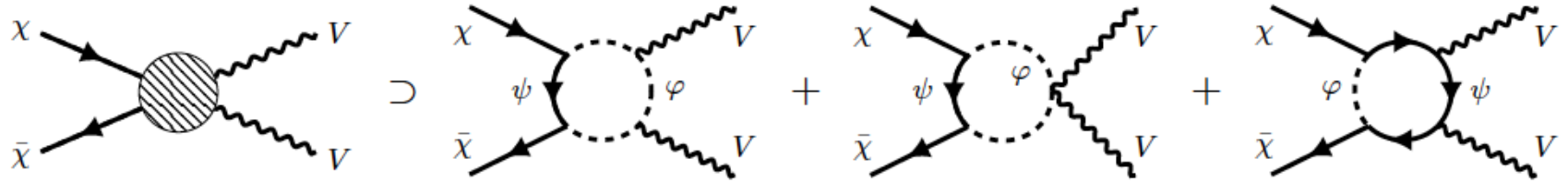
More Channels?

1) Large MET



2) $pp \rightarrow \varphi^\dagger \varphi \rightarrow u\bar{u}t\bar{t}$

Dark Matter Loop Interactions



$$\mathcal{L}_{\text{eff}}^{\text{R}} = \lambda_{\text{s}} \bar{\chi} \chi B_{\mu\nu} B^{\mu\nu} + i\lambda_{\text{p}} \bar{\chi} \gamma^5 \chi B_{\mu\nu} \tilde{B}^{\mu\nu} + \kappa_{\text{s}} \bar{\chi} \chi \text{tr} G_{\mu\nu} G^{\mu\nu} + i\kappa_{\text{p}} \bar{\chi} \gamma^5 \chi \text{tr} G_{\mu\nu} \tilde{G}^{\mu\nu}$$

$$\mathcal{L}_{\text{eff}}^{\text{T}} = \varrho_1 \bar{\chi} i\partial^{\{\mu} \gamma^{\nu\}} \chi \mathcal{G}_{\mu\nu}^{(2)} + \varrho_2 \bar{\chi} i\partial^\mu i\partial^\nu \chi \mathcal{G}_{\mu\nu}^{(2)}$$

$$\mathcal{G}_{\mu\nu}^{(2)} \equiv \text{tr} \left[G_\mu{}^\rho G_{\rho\nu} + \frac{1}{4} \eta_{\mu\nu} G_{\alpha\beta} G^{\alpha\beta} \right],$$

For our model
Coefficients are
Computed as

$$\lambda_{\text{s}} = -\frac{\alpha_1}{27\pi} \frac{y_\chi^2}{m_\psi m_\phi^2},$$

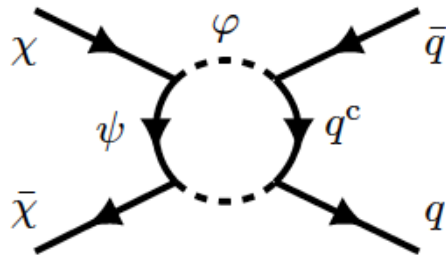
$$\kappa_{\text{s}} = -\frac{5\alpha_3}{96\pi} \frac{y_\chi^2}{m_\psi m_\phi^2},$$

$$\lambda_{\text{p}} = \tilde{\kappa}_{\text{p}} = 0,$$

$$\varrho_1 = \frac{5\alpha_3}{48\pi} y_\chi^2 \frac{1}{m_\phi^2 (m_\phi^2 - m_\psi^2)^2} \left[m_\phi^2 - m_\psi^2 + m_\phi^2 \ln \frac{m_\psi^2}{m_\phi^2} \right],$$

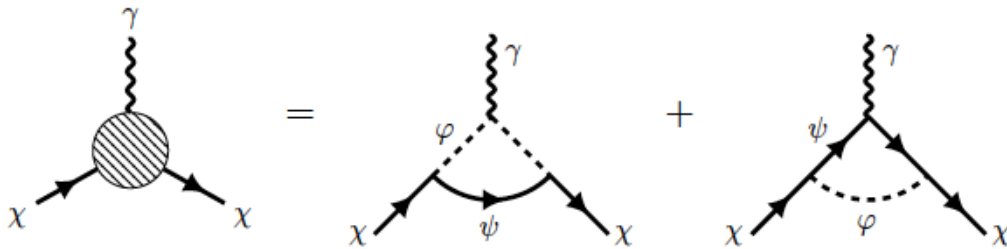
$$\varrho_2 = \frac{5\alpha_3}{12\pi} y_\chi^2 \frac{m_\psi}{(m_\phi - m_\psi)^5} \left[3(m_\phi^4 - m_\psi^4) + (m_\psi^4 + 4m_\phi^2 m_\psi^2 + m_\psi^4) \ln \frac{m_\psi^2}{m_\phi^2} \right].$$

Loop coupling to photon and quarks



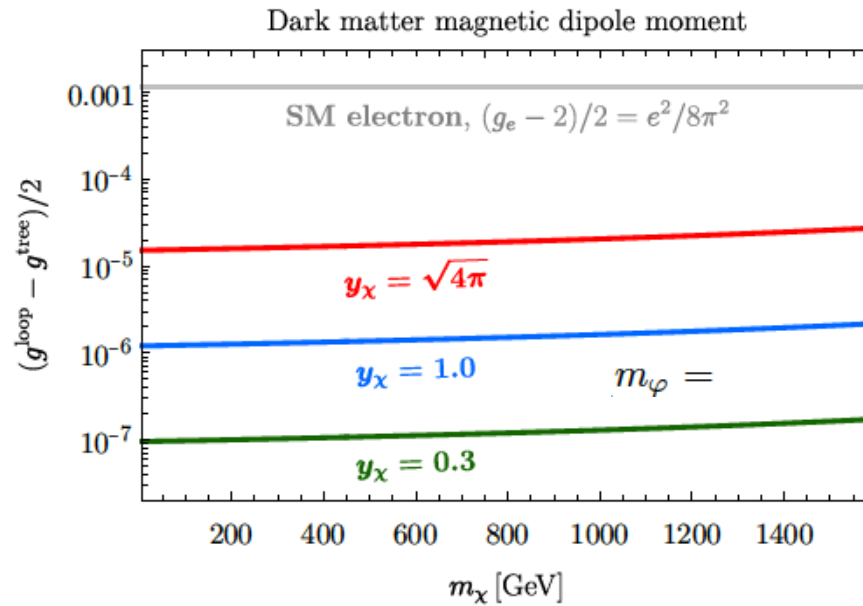
$$\mathcal{L}_{\text{eff}}^{\text{Q}} = \nu_{II} \left[(\bar{\chi}\gamma^\mu\chi)(\bar{u}_I\gamma_\mu u_I) + (\bar{\chi}\gamma^\mu\chi)(\bar{u}_I\gamma_\mu\gamma^5 u_I) \right]$$

$$\nu_{II} = \frac{1}{2} \frac{\lambda_{II}^2}{(4\pi)^2} \frac{y_\chi^2}{m_\psi m_\phi}$$



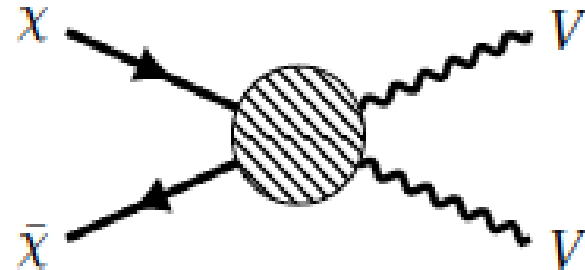
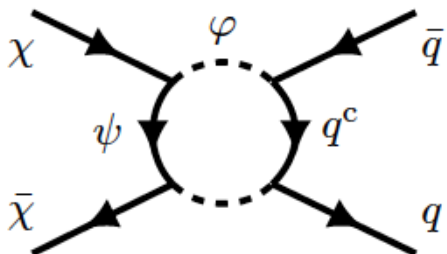
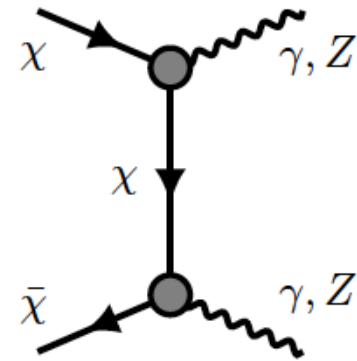
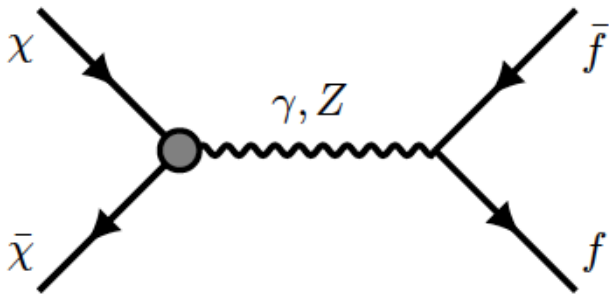
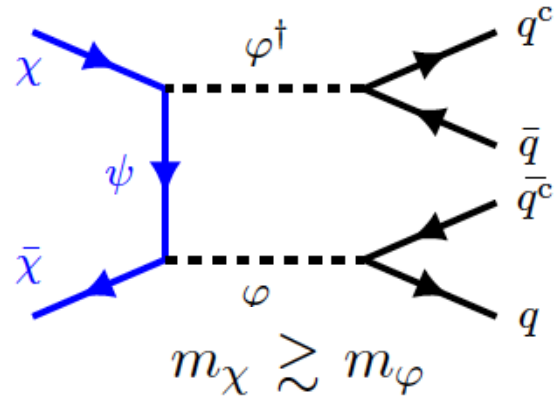
$$\mathcal{L}_{\text{eff}}^{\text{EM}} = A_1 \bar{\chi}\gamma^\mu\chi \partial^\nu B_{\mu\nu} + \frac{1}{4} A_2 \bar{\chi}\sigma^{\mu\nu}\chi B_{\mu\nu} + A_3 \bar{\chi}\gamma^\mu\gamma^5\chi \partial^\nu B_{\mu\nu} + \frac{1}{4} A_4 i\bar{\chi}\sigma^{\mu\nu}\gamma^5\chi B_{\mu\nu},$$

Dark Matter EDM



$$m_\varphi = 1.15 \text{ TeV} \text{ and } m_\psi = 1.6 \text{ TeV}$$

DM annihilation channels

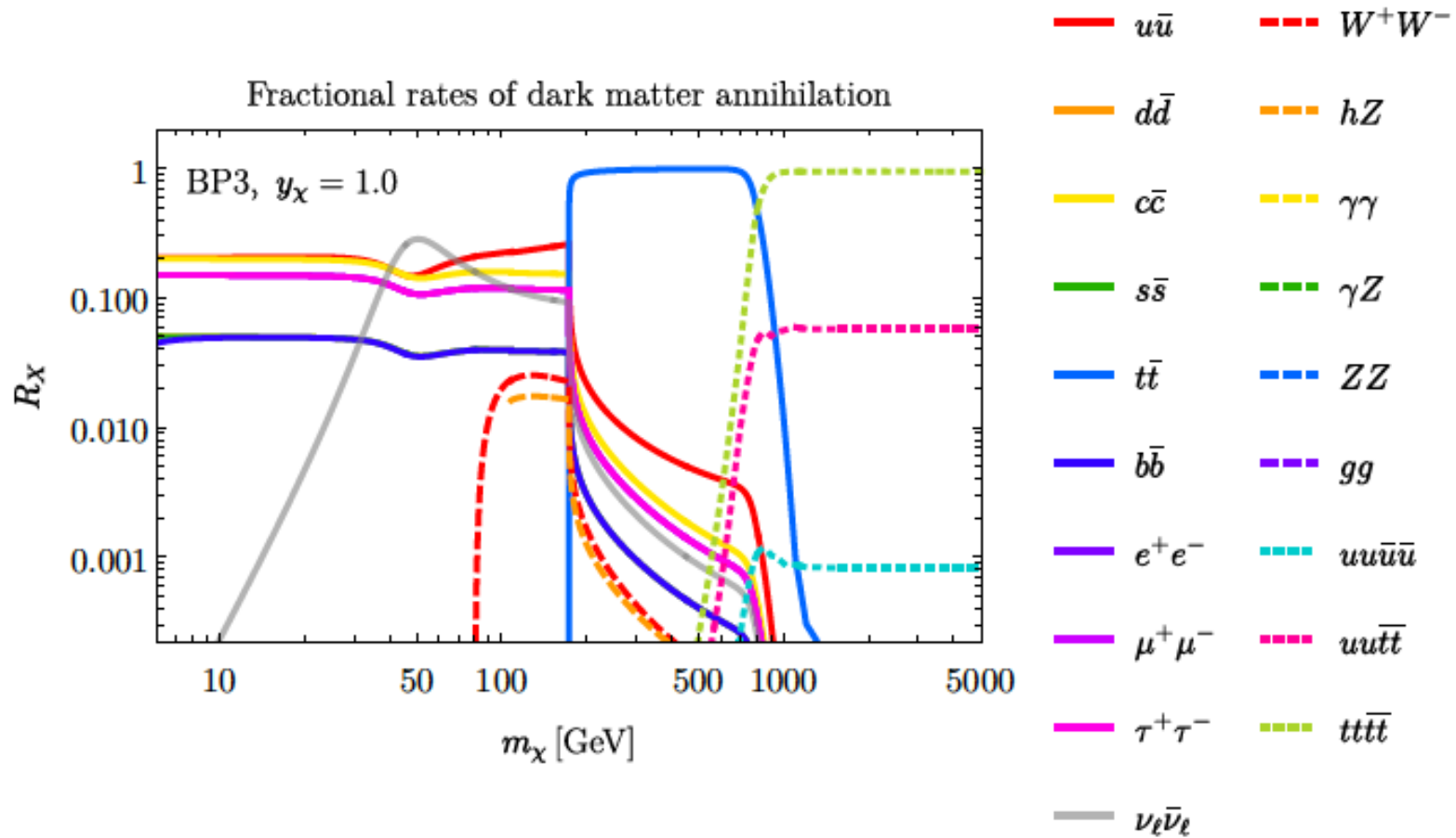


Channel X	$\langle \sigma v_\chi \rangle (\chi\bar{\chi} \rightarrow X)$
$f\bar{f}, f \in \{\ell, q \neq \{u, t\}\}$	$\frac{17\alpha_Q}{16c_w^2} N_c Q_f^2 [\mathcal{F}(A_1^s, A_2^s)]^2$
$q_l \bar{q}_l, q_l \in \{u, t\}$	$\frac{17\alpha_Q}{16c_w^2} N_c Q_f^2 [\mathcal{F}(A_1^s, A_2^s)]^2 + \frac{1}{4\pi c_w} N_c [\mathcal{H}(A_1^s, A_2^s, \iota_{ll})]^2$
$\nu_\ell \bar{\nu}_\ell$	$\frac{\alpha_Q}{16c_w^2} [\mathcal{F}(A_1^s, A_2^s)]^2$
W^+W^-, hZ	$2\langle \sigma v_\chi \rangle (\chi\bar{\chi} \rightarrow \nu_\ell \bar{\nu}_\ell)$
$\gamma\gamma$	$\frac{2}{\pi} c_w^4 m_\chi^2 \left(\frac{A_2^t}{4}\right)^4$
γZ	$2\left(\frac{s_w}{c_w}\right)^2 \langle \sigma v_\chi \rangle (\chi\bar{\chi} \rightarrow \gamma\gamma)$
ZZ	$\left(\frac{s_w}{c_w}\right)^4 \langle \sigma v_\chi \rangle (\chi\bar{\chi} \rightarrow \gamma\gamma)$
gg	$\frac{64}{\pi} m_\chi^4 \kappa_s^2 v_\chi^2$

$$[\mathcal{F}(A_1, A_2)]^2 = \left[\frac{A_2}{4} + m_\chi A_1 \right]^2$$

$$[\mathcal{H}(A_1, A_2, \iota_{ll})]^2 = 2eQ_f m_\chi \frac{A_2 \iota_{ll}}{4} + 2eQ_f m_\chi^2 A_1 \iota_{ll} + m_\chi^2 \iota_{ll}^2$$

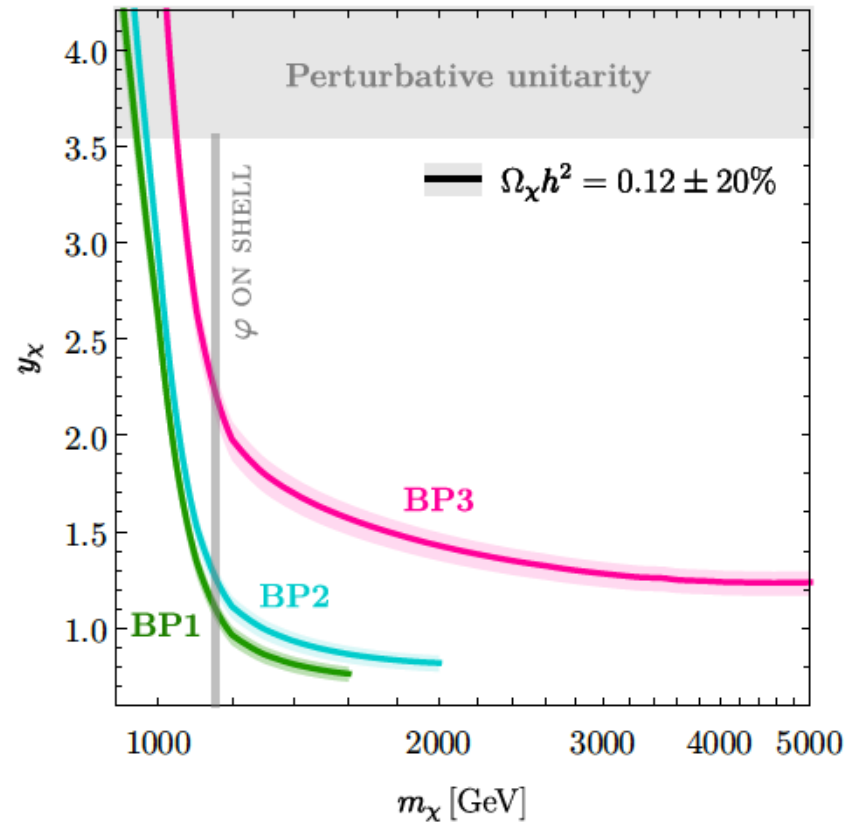
$$R_X = \frac{\langle \sigma v_\chi \rangle (\chi\bar{\chi} \rightarrow X)}{\langle \sigma v_\chi \rangle} \quad \text{with} \quad \langle \sigma v_\chi \rangle = \sum_X \langle \sigma v_\chi \rangle (\chi\bar{\chi} \rightarrow X)$$



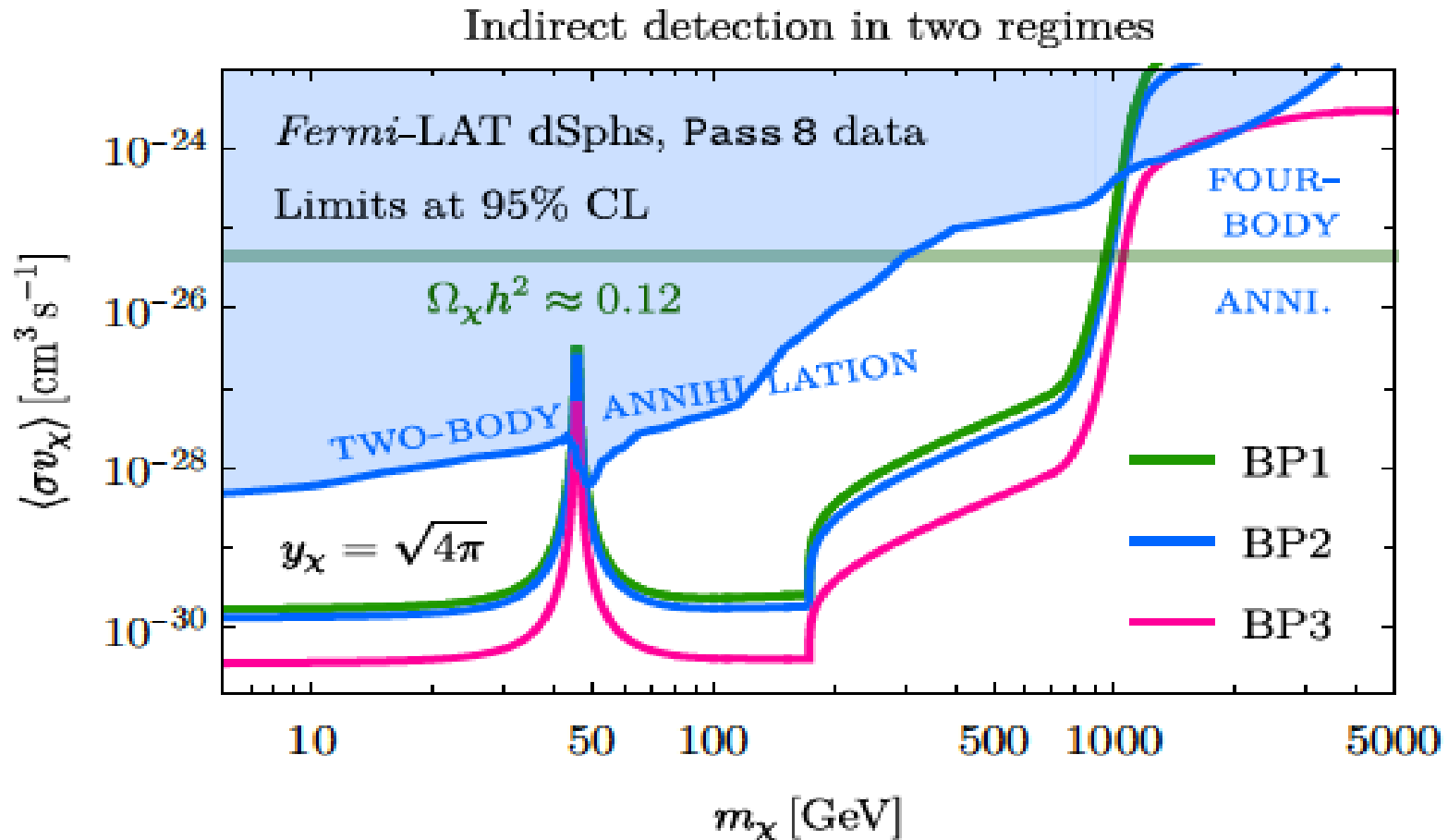
	m_χ range [GeV]	m_φ [GeV]	m_ψ [GeV]	λ_{11}	λ_{33}
BP1	(0, 1600)	1150	1600	0.125	0.75
BP2	(0, 2000)		2000		
BP3	(0, 5000)		5000		



Achieving the observed relic density with y_x



Indirect Detection Bounds



Computed with **DARKFLUX** <https://github.com/carpenterphysics/DarkFlux>.

arXiv:2202.03419v2

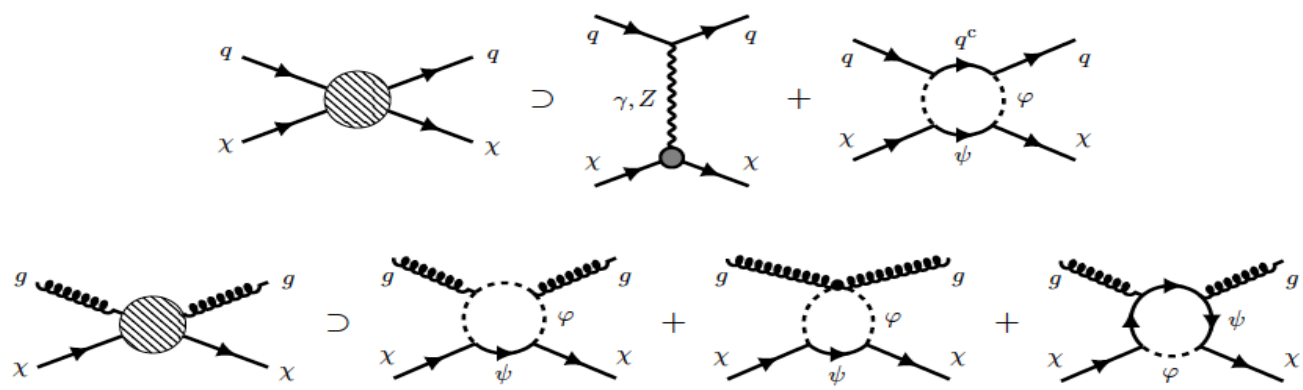
New Directions

Explore New Frustrated Models

Explore MFV scenario, is it less constrained?

More Collider signatures, e.g. large MET events

Extend DARKFLUX use



Limits on Yukawa coupling [BP3, $m_\psi = 5.0$ TeV]

