







Systèmes de Référence Temps-Espace

Developments to improve the stability of optical lattice clocks

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Outline

- Progress in the development of a Hg lattice clock
- **Two developments to improve the stability of optical lattice clocks**
 - Non-destructive detection for Sr lattice clocks
 - Laser stabilization using spectral hole burning

Properties of neutral mercury

Isotope

196

198

199

200

201

202

Abundance (%)

0.15

10.1

17.0

23.1

13.2

29.6

Spin

0

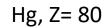
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1/2

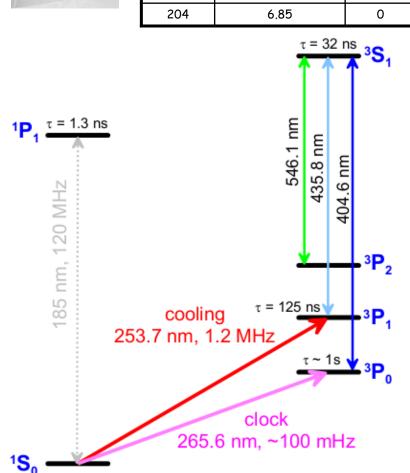
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3/2

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□ Low sensibility to BBR and static electric fields At 300 K: -1.6x10⁻¹⁶ (~Sr/30, Yb/15)

- 7 isotopes, 6 of them >6% abundant (2 fermions, 4 bosons)
 - ¹⁹⁹Hg is a fermion with nuclear spin ½
- High sensitivity to α variations

Phys. Rev. A **70**, 014102 (2004) $\frac{\delta \nu}{-} = 0.81 \frac{\delta \alpha}{-}$

High vapor pressure

No oven, 2D-MOT possible

Mostly unexplored in the ultra-cold regime

Challenge: deep UV laser sources

Status in 2015

□ Uncertainty of 1.7E-16 with 199Hg

- Limited by
 - Level of control of the lattice frequency
 - The maximum trap depth achievable (56 E_R)
 - The stability >1.2E-15 @1s
 - Reliability / measurement time cumulated

□ Frequency ratio against Cs, Rb and 87Sr

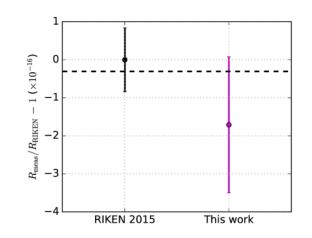
- From local comparisons
- Frac. uncertainty for Hg/Sr: 1.8E-16
- In agreement with RIKEN

Yamanaka et al., PRL114,230801 (2015)

199Hg in the CCTF-2015 LoR

And SRS since CCTF-2017

Ratio	Measured value
$rac{ u_{ m Hg}/ u_{ m Cs}}{ u_{ m Hg}/ u_{ m Rb}} u_{ m Hg}/ u_{ m Sr}$	122 769.55 729 311 011 (45) 165 124.754 879 997 258 (62) 2.629 314 209 898 909 15 (46)

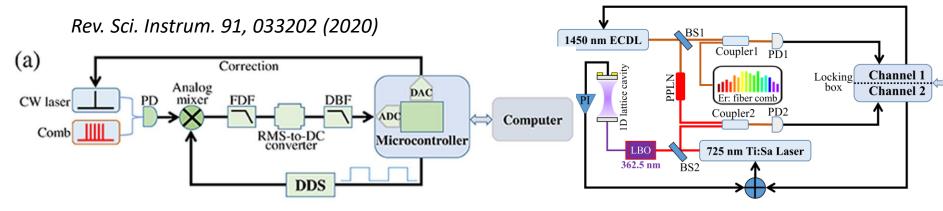


Tyumenev et al., New J. Phys. 18, 113002 (2016)

Some improvements: lattice frequency lock

Accurate control of the lattice frequency

- Scalar shift ~1.1E-17/MHz at a depth of 100 $E_R \rightarrow$ control at the 10 kHz needed
- \rightarrow a frequency doubled ECDL at 1450 nm to connect the comb to the TiSa of the lattice trap
- →an accurate locking scheme that can tolerate relatively broad laser linewidth and poor signal to noise ratio to avoid adding constrain to the existing comb



Lock accuracy in our application: 1 kHz

 \rightarrow Hg clock average nicely to 1E-17, where it used to flicker at a few 1E-17

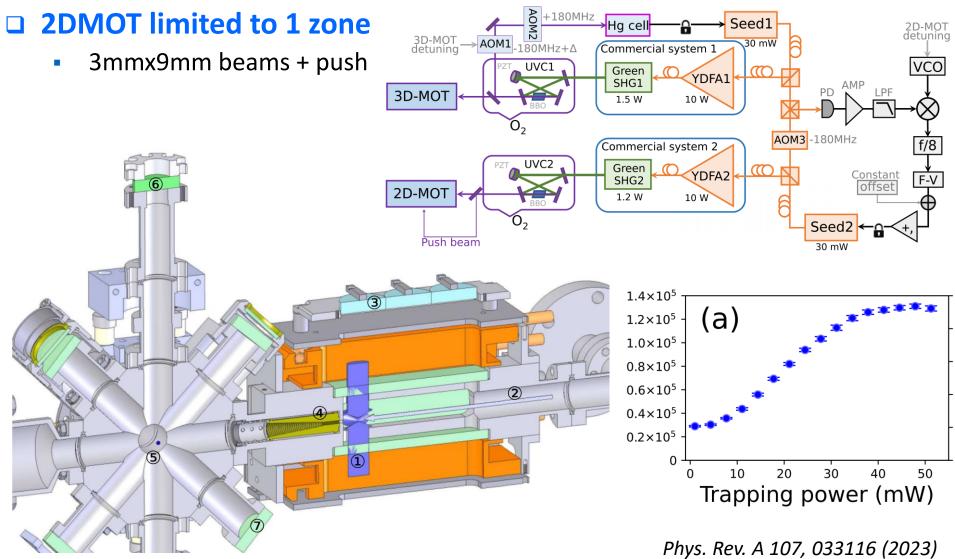
The scheme can find application on other cases

- Accurate frequency control of lasers with broad linewidth
- Tolerates extremely poor signal to noise ratio → can considerably extend the usable wavelength range of a given frequency comb in certain applications.

Some improvements: 2DMOT

□ Laser system based on a commercial YDFA + SHG system

1.2 W in the VIS, 50 mW in UV, linewidth ~100 kHz (natural linewidth 1.3 MHz)

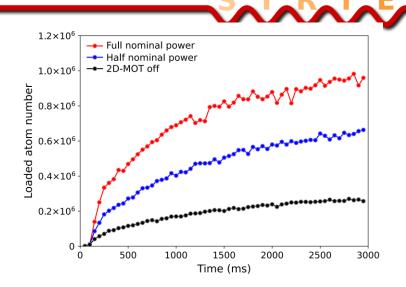


Stability with 2DMOT

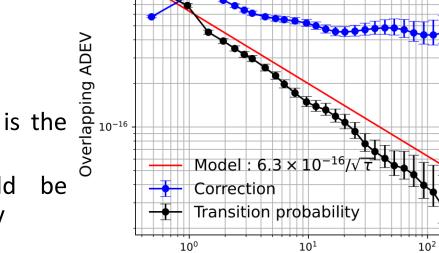
2DMOT improve loading rate by up to 4.5

□ Stability of 6.3E-16 at 1 s

- With a clock cycle of 482 ms with 160 ms probe (duty cycle of 0.33, instead of <0.1 before)
- With a 4E-16 laser (10cm cavity)



Time (s)



 10^{-15}

Limits

- Reliability: SHG 508 \rightarrow 254 nm is the weak point (3DMOT and 2DMOT)
- Significant improvement should possible with optimized geometry
- With more power too.

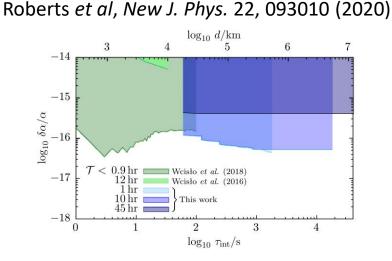
Comparisons

Improved reliability

- Capability to participate in coordinated fiber link comparisons, to cumulate many hours of measurements over a campaign (ex: 100h during 2023 campaign)
- Still requiring the presence of an operator

$\hfill\square$ Contribution to search for α variation and dark matter

 Search for topological dark matter exploits the spatial distribution of clock connected by links





Fiber link collaborations INRIM, PTB, NPL, SYRTE LPL, RENATER (REFIMEVE infrastructure) EURAMET projects



Comparisons

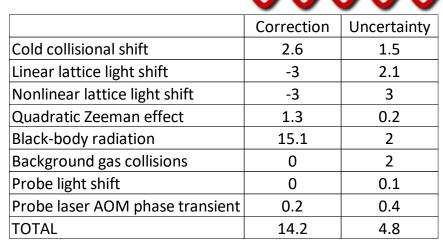
□ Uncertainties Hg (x1E-17)

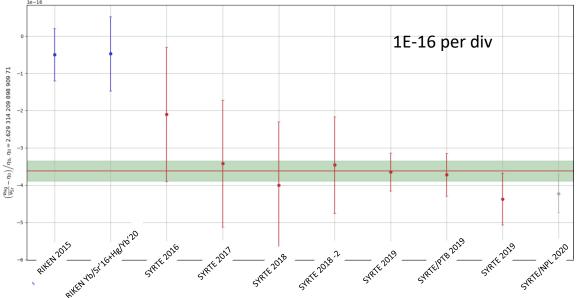
- Improved thanks to lattice control, improved stability
- Limited by stability/reliability, coupling between lattice depth and collision shift, lattice lifetime

199Hg/Sr comparisons

Uncertainties dominated

by Hg





□ Other comparisons with Yb+(E3), Yb+(E2), Yb

Including the March/April 2023 campaign still under analysis

Toward using bosonic isotopes

□ 199Hg 3P0 state has a relatively short lifetime

~1.5 s ightarrow will be a limit for the next generation of ultra-stable lasers

□ Bosonic isotope can be used with a quenching scheme

- Lifetime unlimited, probing time adaptable to laser characteristics
- Done successfully in 174Yb and 88Sr Taichenachev et al., Phys. Rev. Lett., 96, 3,(2006)

□ Modification to implement this scheme

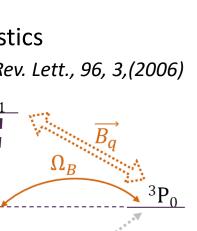
- Broadly tunable ultra-stable probe to address any of the isotopes while preserving 4E-16 stability
- Pulse magnetic field to quench the transition up to 15.5 mT

Used 199Hg to maximize coupling of light to trapped atoms. (Rabi frequency of 2 kHz)

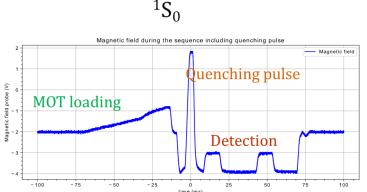
 \rightarrow inferred Rabi frequency for bosons >50 Hz

Search for the transition in 198Hg

- Current uncertainty is 10.3 MHz
- \rightarrow not easy but feasible



 $\alpha \sqrt{I} |B| \cos\theta$



 ^{3}P

 I, Ω_L

Non-destructive detection for Sr lattice clocks

- Quantum non-destructive measurements to improve clocks
 - uW transition, trapped Cs, dispersive measurement

New J. Phys. 12 065032 (2010)

 uW transition, trapped Rb, -18.5 dB below SQL, cavity enhanced-dispersive measurement, bichromatic cavity

Nature 529 505 (2016)

• Optical transition, Yb lattice, cavity in strong coupling

Nature 588, 414 (2020)

Optical transition, Sr lattice, transverse cavity in strong coupling arXiv:2211.08621

Goals of this work

- A practical implementation in optical lattice clock that goes beyond proof-ofprinciple
- Possible benefit for a transportable clock of both classical and quantum nondestructivity
 - →to relax the constrain on transportable ultra-stable cavities
 By reducing the impact of the Dick effect

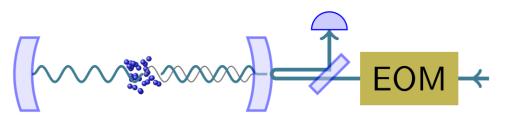
By extending the probe time beyond the laser coherence time using e.g. the atomic phase lock scheme

R. Kohlhaas, et al., Phys. Rev. X 5, 021011 (2015)

Detection method

PDH like detection immune to technical noise

2nd sidebands resonating



$$\phi^{\mathrm{at}} = rac{N}{S} rac{3\lambda^2}{2\pi} rac{\Delta/\Gamma}{s + 4\Delta^2/\Gamma^2}$$

Amplified by the cavity finesse \rightarrow

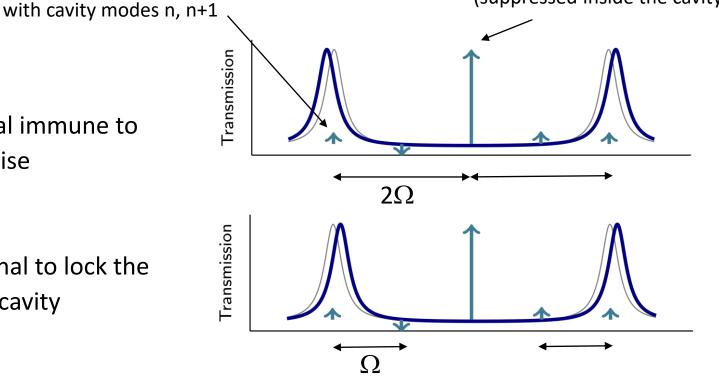
 $\mathcal{F}\phi_{\mathsf{at}}$

Probe carrier at the atomic resonance frequency (suppressed inside the cavity

At 2 Ω Atomic signal immune to

technical noise

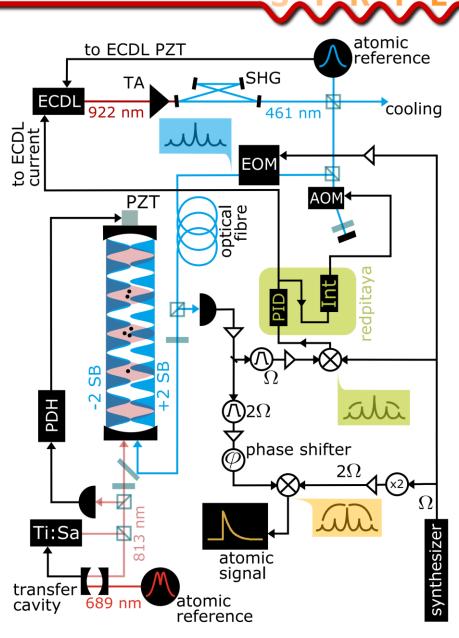
At Ω PDH-like signal to lock the laser to the cavity



1st generation experimental setup

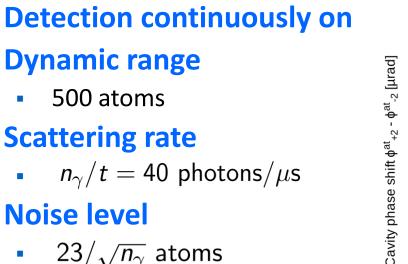
Features of the setup

- Bi-chromatic cavity (lattice 813 nm
 + detection 461 nm)
- \rightarrow self-aligned
- High-finesse at 461 nm (16000)
 →100 fold increase of the SNR
- Heterodyne dual-mode detection
 homogeneous atom-cavity coupling
- Fast digital servo system
 →possibility of pulsed operation



Vallet et al. New J. Phys. 19 083002 (2017)

Experimental results and limits



500 atoms

□ Scattering rate

 $n_{\gamma}/t = 40$ photons/ μ s

□ Noise level

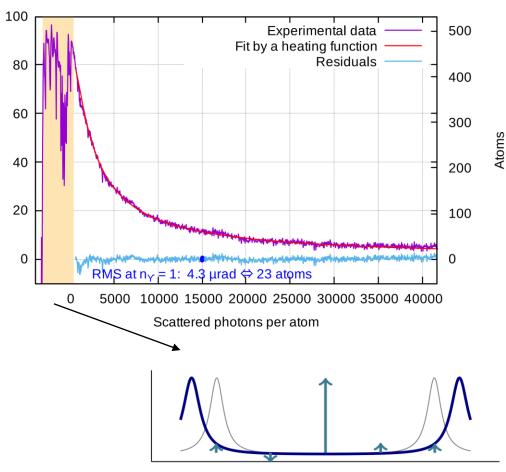
 $23/\sqrt{n_\gamma}$ atoms

□ 3 regimes of optimization

- Destructive (long, $n_{\gamma} >> 1$) $\delta N \ll 1$ atom
- Classically non-destructive $\delta N < 4$ atoms
- Quantum non-destructive: $n_{\gamma} < 1$

 $\delta N > 23$ atoms $\Rightarrow \delta N < \sqrt{N}$ for N > 500

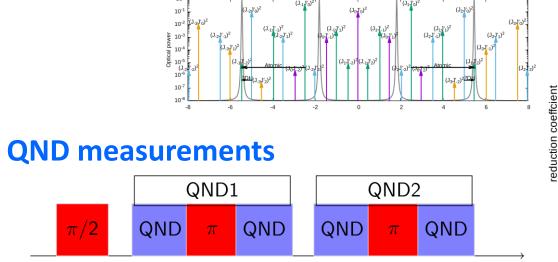
$\Box \rightarrow$ Need to improve dynamic range and noise/destructivity tradeoff

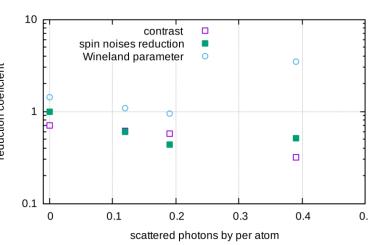


2nd experimental setup and prospects

- New system with two independent modulation frequencies
 - Large detuning (5.7 GHz) → low scattering rate
 - Tracking of the cavity resonances ightarrow improved dynamic range
 - Pulsed operation

As in Hobson et al., Optics Express 27 37099 (2019)

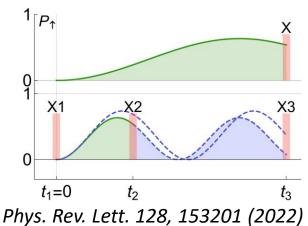




- Observation of correlations between QND meas.
- Wineland parameter $\xi = 0.95$

Protocol to exploit this detection (with ICFO)

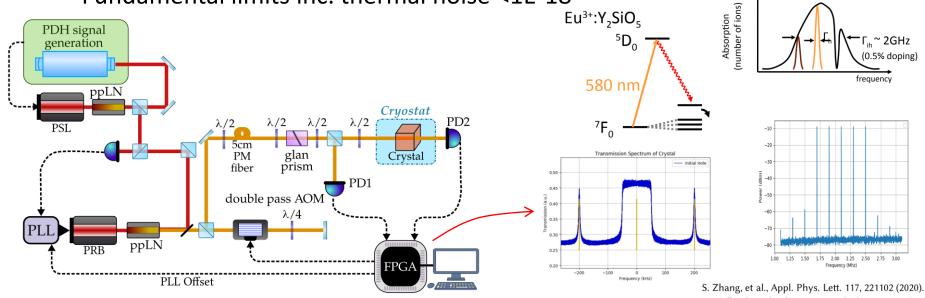
- Rabi interrogation with 3 QND measurements
- Optimization accounting imperfections
- \rightarrow significant gain possible (7.9 dB for 2E4 atoms)



Spectral hole burning for lasers <1E-17

- Laser stabilization to spectral features obtained by spectral hole burning (SHB)
 - Have extremely low fundamental limits.
 - One possibility to beat the thermal noise limit of room temperature cavities.
 - Was investigated at NIST, HHU Nat. Photon. 5, 688 (2011) Phys. Rev. Lett.107, 223202 (2011)
- □ Eu3+ ions doped into Y₂SiO₅ matrix at cryogenic temperature
 - Long-lived spectral holes can be burned in the optical absorption spectrum.
 - $T_2 = 2.6 \text{ ms} (1.4 \text{ K}) \rightarrow \Gamma_h = 122 \text{ Hz expected} \cdot 1 2 \text{ kHz routinely observed}$.





N. Galland, et al., Opt. Lett. 45, 1930-1933 (2020).S. Zhang, et al., Phys. Rev. A 107, 013518 (2023).

S. Zhang, et al., Appl. Phys. Lett. 117, 221102 (2020).
N. Galland, et al., Phys. Rev. Applied 13, 044022 (2020).
S. Zhang, et al., Phys. Rev. Research 2, 013306 (2020).

Spectral hole burning at 3.2K & dilution temperature

□ At 3.2K, several studies of sensitivities

- Stark, strain/acceleration, magnetic field, temperature
- →sensitivities and perturbations present in the environment compatible with <1E-17, except for the temperature
- Probing schemes setup and stability of 1.7E-15 @1s for 1 hole.

□ At <1K, measurement of the temperature sensitivity

Drastically reduced sensitivity and a turning point

