

Machine & People Protection Issues

CAS Introduction to Accelerator Physics

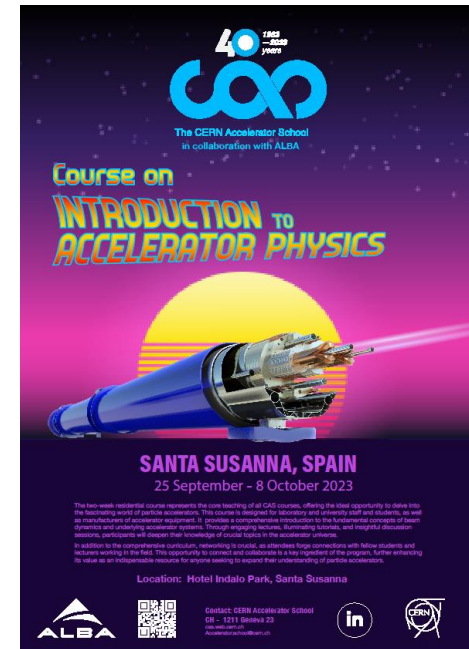
Santa Susanna, 1st of October 2023

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Lecture based on previous CAS & JUAS contributions by Daniela Kiselev, Xavier Queralt, Rüdiger Schmidt, Ivan Strasik, Markus Zerlauth...



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Introduction and Outline

Reasons for machine protection:

- **Protection of the environment:** Only necessary activation inside & outside of the facility should be produced
- **Protection of the accelerator:** Prevent for destruction of component, prevent for down-time, destruction & cost
- **Enable save operation:** Threshold values for reliable operation
- **Protection of people:** Important for workers and general public, following laws

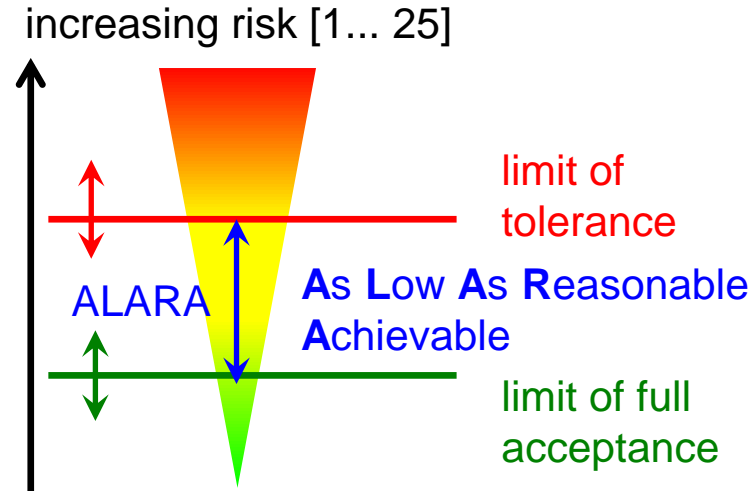
Outline of this talk:

1. Introduction to risk & destruction potential
2. Relevant atomic and nuclear physics
3. Collimation as passive protection
4. Measurements by Beam Loss Monitors
5. Design of Machine Protection System
6. Overview of personal safety

What Risk is acceptable?

Risk is a factor to prepare for decisions, it is not a physical quantity:

| | | | | | |
|-----------------------|---------------------|---------------------|---------------------|-------------------|-------------------|
| 5 Catastrophic | 5 | 10 | 15 | 20 | 25 |
| 4 Major | 4 | 8 | 12 | 16 | 20 |
| 3 Severe | 3 | 6 | 9 | 12 | 15 |
| 2 Minor | 2 | 4 | 6 | 8 | 10 |
| 1 Slight | 1 | 2 | 3 | 4 | 5 |
| Consequences | 1 Negligible | 2 Improbable | 3 Occasional | 4 Probable | 5 Frequent |
| Probability | | | | | |



Risk = probability of an accident x **consequences**

measured in terms of e.g. money, manpower, accelerator downtime, radiation pollution

- Intolerable or acceptable depends on e.g. maintenance access, destruction level, operation
 - Different accelerator facilities allows different risks e.g. medical ↔ research facilities
- ⇒ Risk must be weighted to foreseen usage, goals and possible achievements

What is the Risk for Accelerators?

Categories of destruction, consequences and risk:

➤ **Heating:** Lost beam heats the surrounding by its energy loss (by *atomic physics*)

⇒ **Consequence:** Material is melted and deformed ⇒ proper functionality hindered

⇒ **Type of risk:** Stop of operation

Example: Destroyed insertions, leak in vacuum chamber, quench of superconducting magnet

➤ **Activation:** Nuclear reaction by beam particles (*nuclear physics*)

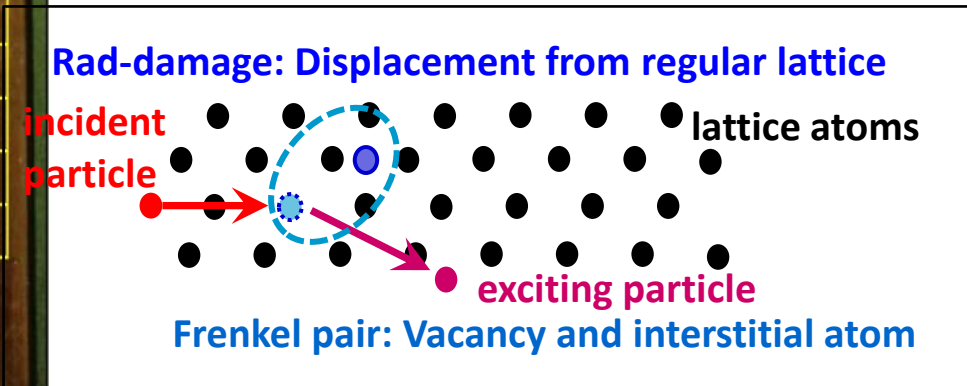
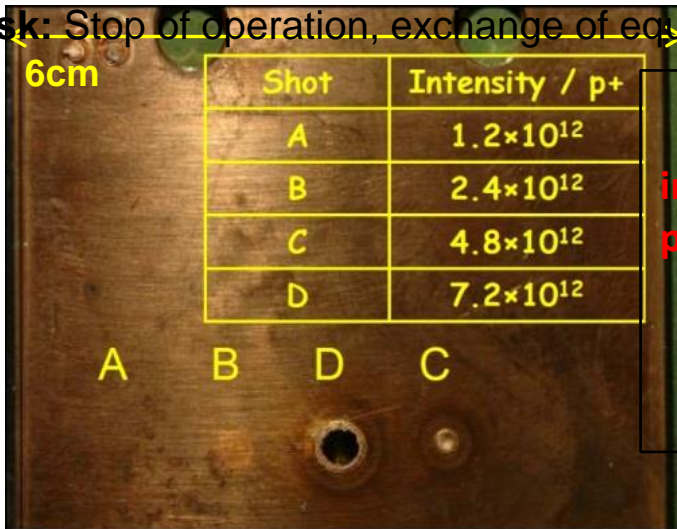
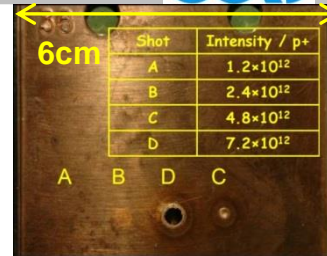
⇒ **Consequence:** Permanent activation ⇒ pollution, human access hindered

⇒ **Type of risk:** Maintenance impossible, expensive disposal

➤ **Radiation damage:** Displacement of lattice atoms, destruction of molecules (*atomic physics*)

⇒ **Consequence:** Degradation of material properties, faulty electronics

⇒ **Type of risk:** Stop of operation, exchange of equipment



What is the Risk for an Accelerators?

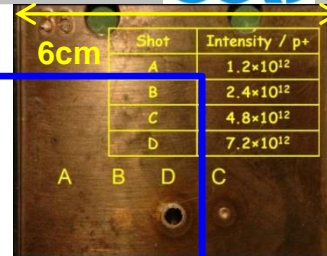
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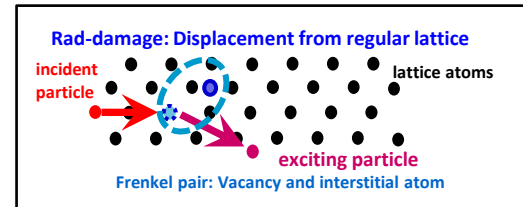
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⇒ **Type of risk:** Stop of operation, exchange of equipment



➤ **Financial aspects:** High cost of additional radiation shield

⇒ **Consequence:** Reconstruction of buildings

⇒ **Type of risk:** Insufficient budget, loss of operation permit



➤ **User requirements:** Less beam available for users

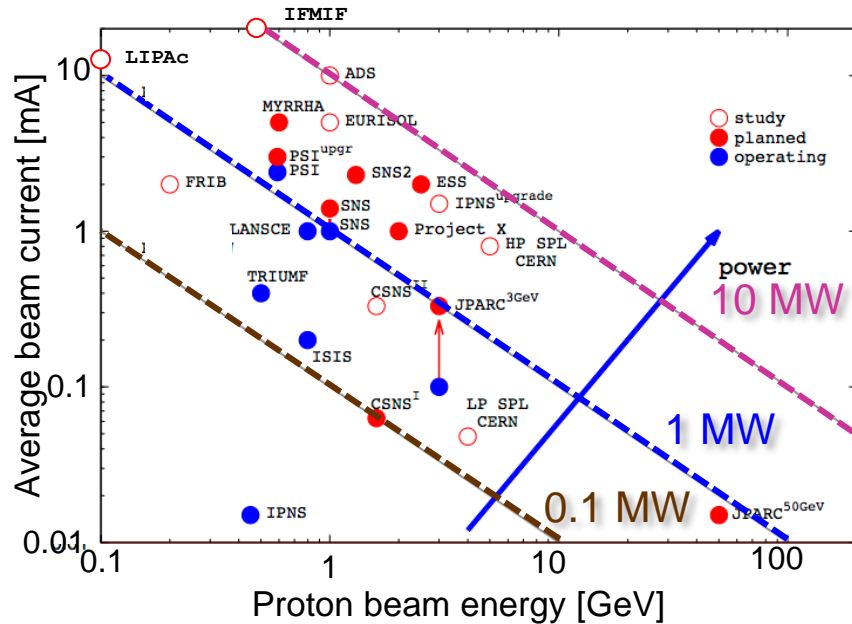
⇒ **Consequence:** Angry or disappointed users

⇒ **Type of risk:** Cancel financial support for accelerator facility

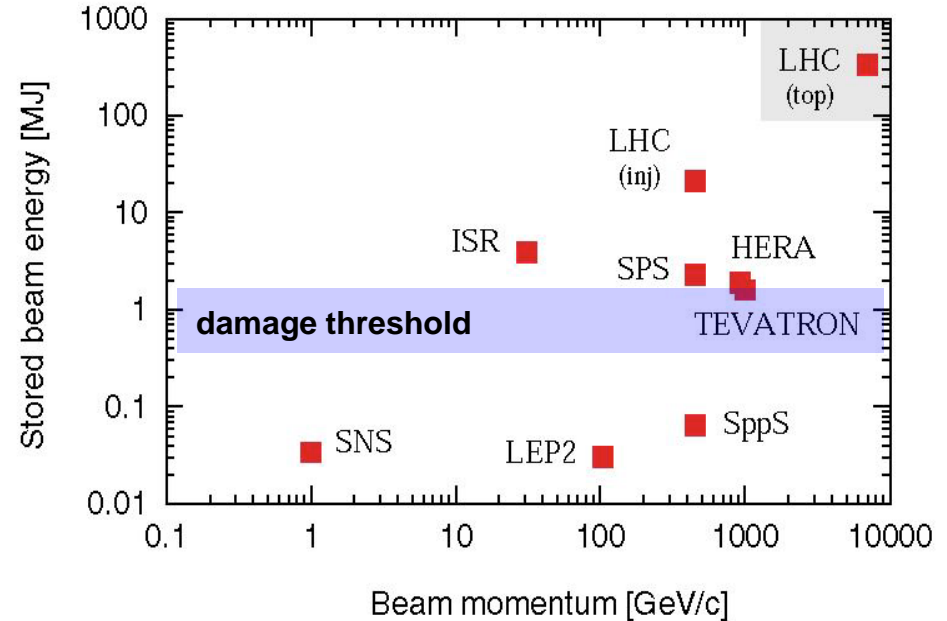


Stored Beam Energy at Accelerators

Beam power on fixed target proton accelerator:
LINACs, cyclotrons or extraction from synchrotrons

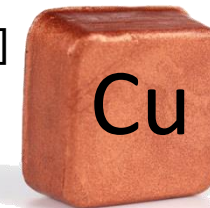


Stored beam energy within a synchrotron:
Mainly large circular collider



Examples: Energy of 1MJ correspondence:

- 1 MJ is the kinetic energy of 2600 kg with an velocity of 100 km/h
 - 1 MJ can heat and melt 1.5 kg of copper [equals cube (5.5 cm)³]
 - 1 MJ is liberated by the explosion of 0.25 kg TNT
- LINAC: 1 MW delivered within 1 s equals to 1MJ



$T_{melt} = 1080 \text{ }^{\circ}\text{C}$
 $\rho = 8.9 \text{ g/cm}^3$



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Overview: Interaction of Particles and Photons with Matter

Interaction with matter

General:

- Charged particles interacts with electrons
⇒ shorter range
- Neutrons ionizes only indirectly
⇒ longer range
- Atomic processes have larger cross section than nuclear processes

'Geometrical' cross section:

Very probable reactions have $\approx \sigma_{geo}$

Cross section σ_{geo} comparable to size:

- Size of **atom**: $r_{Bohr} = 0.053 \text{ nm}$

$$\sigma_{geo}^{atom} = \pi (r_{Bohr})^2 = 8,8 \cdot 10^{-17} \text{ cm}^2$$

$$\approx 10^{-16} \text{ cm}^2$$

- Size of **nucleus**: $r_{nucl} \approx 3 \text{ fm}$

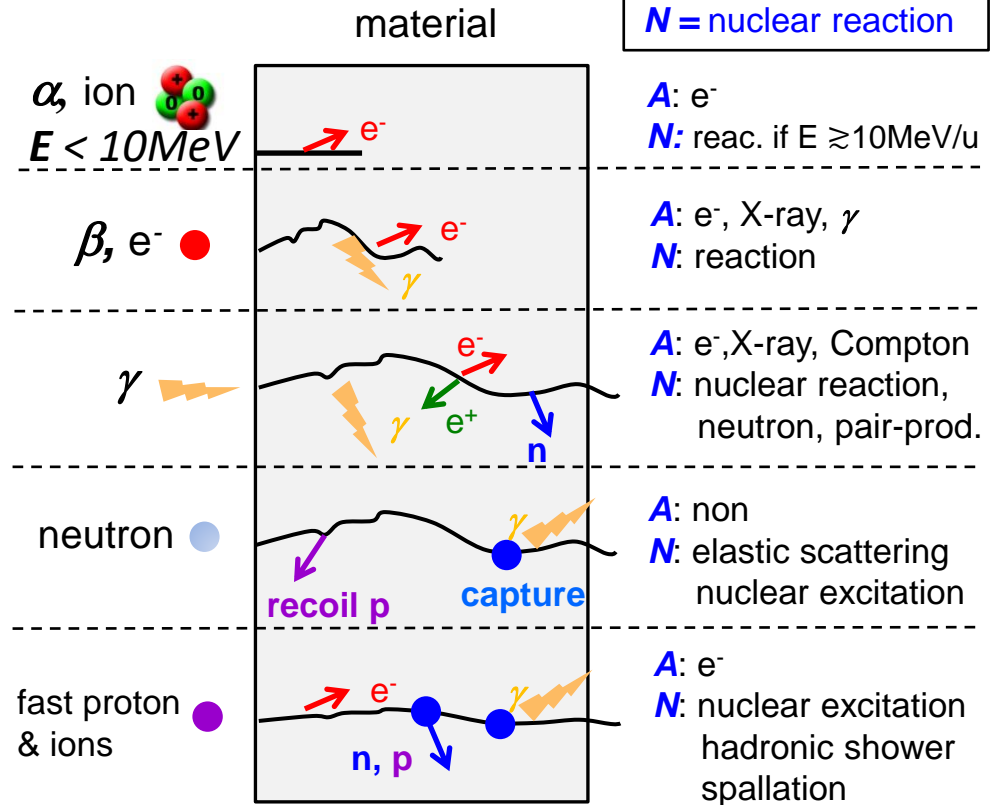
$$\sigma_{geo}^{nucl} = \pi (2 \cdot r_{nucl})^2$$

$$\approx 10^{-24} \text{ cm}^2 \equiv \mathbf{1 \text{ barn}}$$

Mean free path: $\lambda = \frac{1}{n \cdot \sigma} = \frac{M}{\rho N_A \cdot \sigma}$

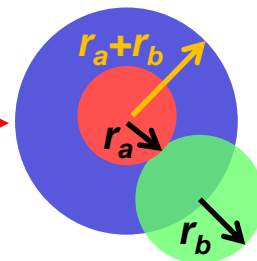
n target atom density [cm^{-3}], M molar mass, ρ density, N_A Avogadro number

A = atomic reaction
N = nuclear reaction



factor 10⁸

'beam'



Hard balls' 'geometrical' cross section:

$$\sigma_{geo} = \pi (r_a + r_b)^2 \text{ for any 'reaction'}$$

Energy Loss of Ions in Copper

Bethe-Bloch formula: $-\frac{dE}{dx} = 4\pi N_A r_e^2 m_e c^2 \cdot \frac{Z_t}{A_t} \rho_t \cdot Z_p^2 \cdot \frac{1}{\beta^2} \left(\frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 \cdot W_{max}}{I^2} - \beta^2 \right)$
 (simplest formulation)

Range: $R = \int_0^{E_{max}} \left(\frac{dE}{dx} \right)^{-1} dE$

with approx. scaling $R \propto E_{max}^{1.75}$

Numerical calculation for **ions**

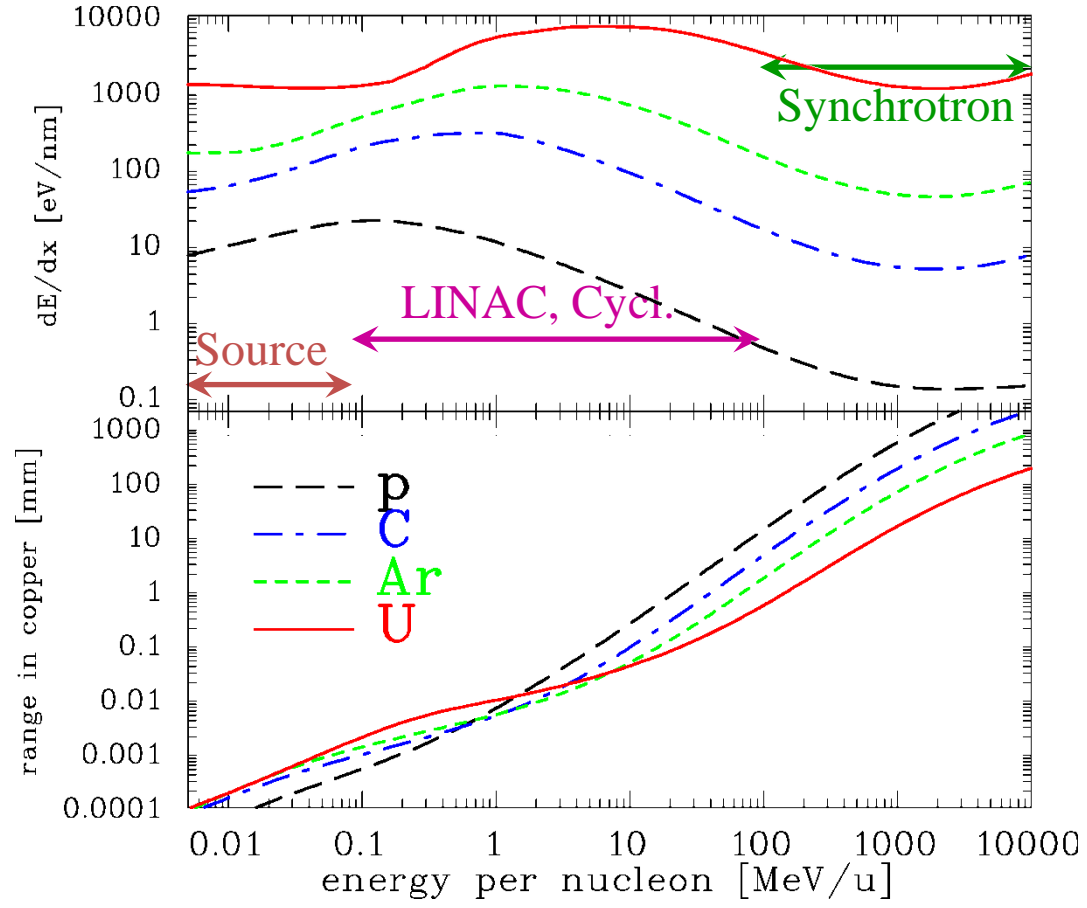
with semi-empirical model e.g. SRIM

Main modification $Z_P \rightarrow Z_p^{eff}(E_{kin})$

This is an atomic physics process:

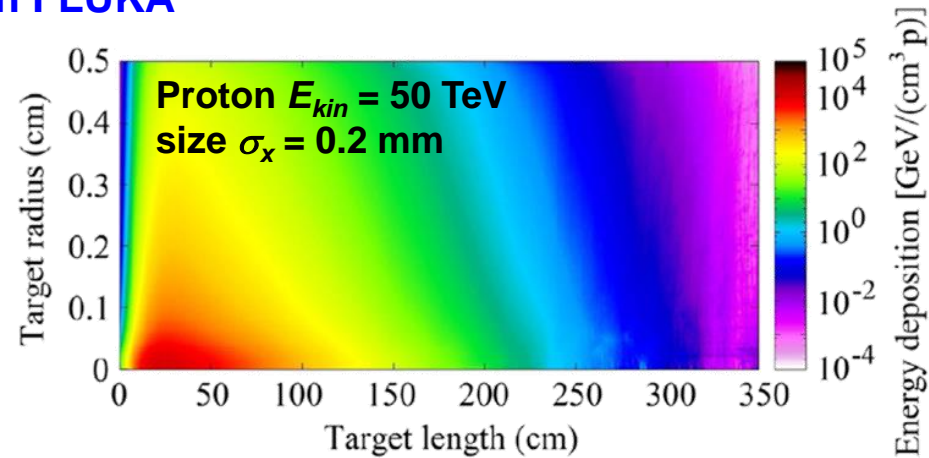
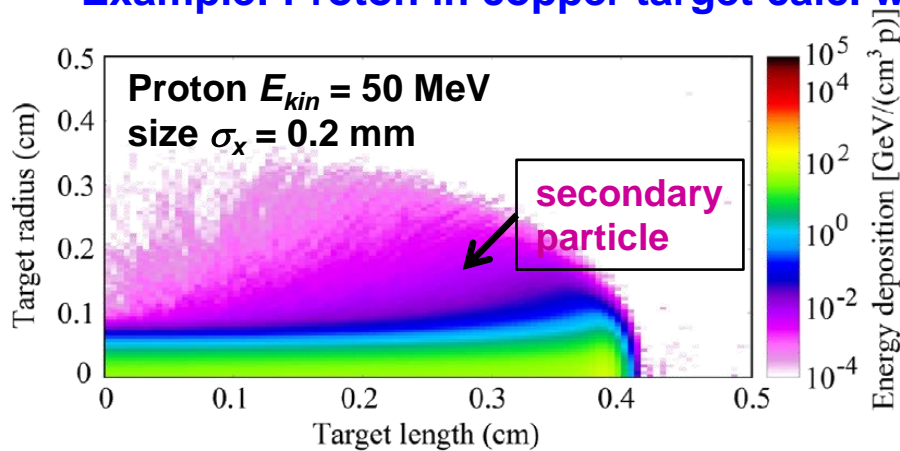
1. Projectile ions liberates fast electrons
2. Thermalization by collisions with further electrons
3. Transfer of energy to lattice (phonon)

⇒ **Heating of target**



Energy Loss and Heating: Calculations

Example: Proton in copper target calc. with FLUKA



Y. Nie et al., Phys Rev AB 20, 081001 (2017)



General method of calculation (simplified):

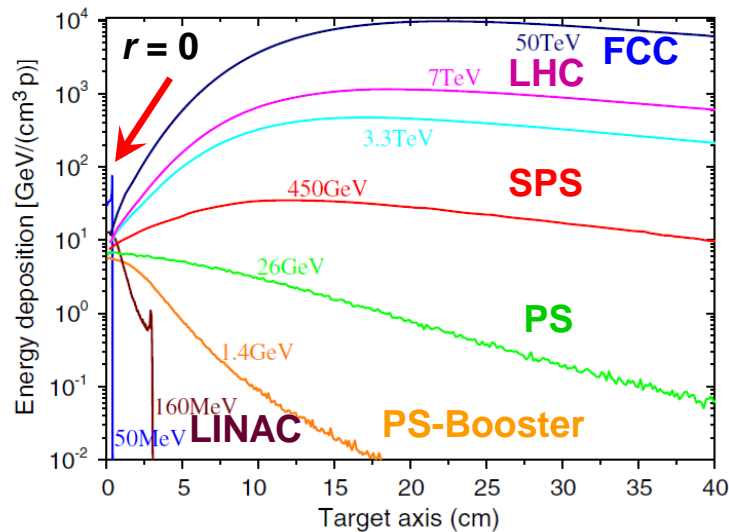
1. **Differential energy loss:** by Bethe-Bloch $\frac{dE}{dx}(x)$ via codes like SRIM, LISE, FLUKA, MARS...
2. **Energy deposition:** $\frac{dE}{dV} = -\frac{dE}{dx} \cdot \frac{N}{A}$ $\left[\frac{J}{cm^3} \right]$ with N : number of particles, A : beam cross section
3. **Temperature rise:** $\Delta T = \frac{dE}{dV} \cdot \frac{1}{\rho c_p}$ [K] for short bunches; ρ : mat. density, c_p specific heat
4. **Further material response:** Melting, evaporation, pressure and stress via e.g. ANSYS
5. **Secondary particles:** Nuclear reactions, fragmentation, spallation, shower.... → discussed later

Energy Loss and Heating: Calculations

Example: Proton in copper target calc. with FLUKA

Proton beam at the CERN accelerators

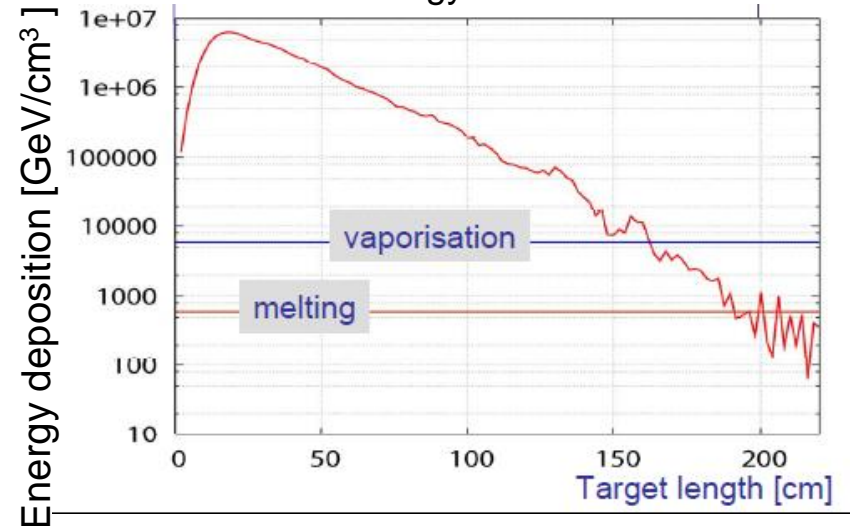
Energy deposition per protons



Nominal LHC parameters for protons:

$E_{kin} = 7 \text{ TeV}$, 2808 bunch

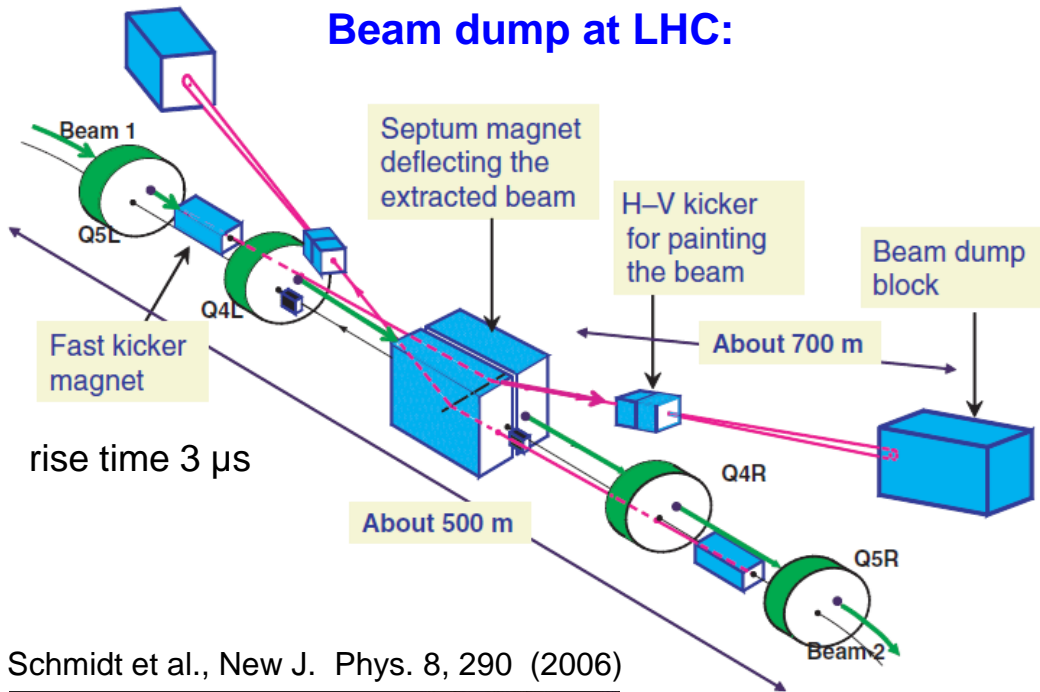
⇒ 380 MJ energy at center $r = 0$



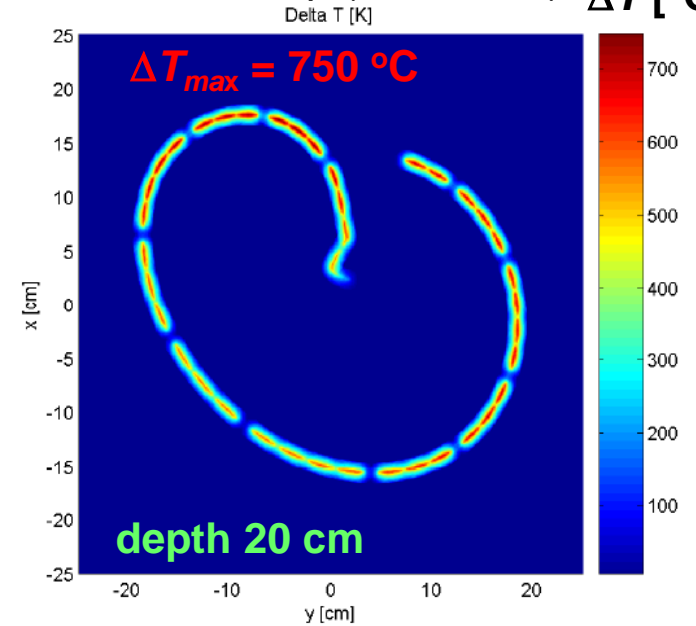
Remark: Low energetic proton have large energy deposition at short range e.g. $E_{kin} = 50 \text{ MeV}$

Beam Dump for high Intensity Beams

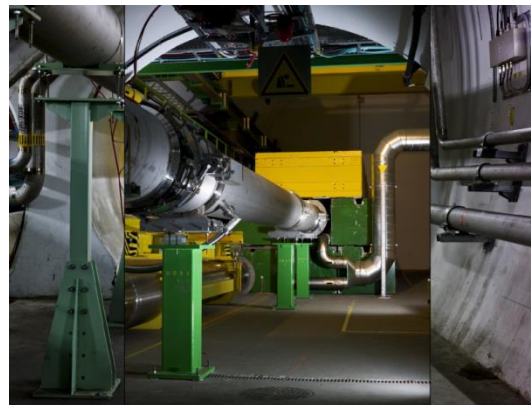
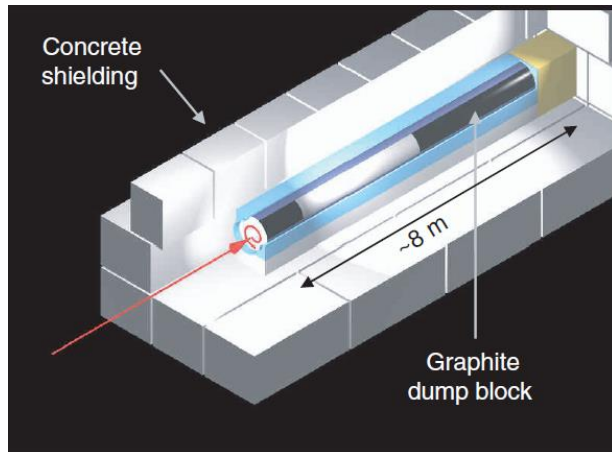
Beam dump at LHC:



Extraction of LHC within **one** turn 86 μ s on the beam dump (simulation): ΔT [°C]



R. Schmidt et al., New J. Phys. 8, 290 (2006)



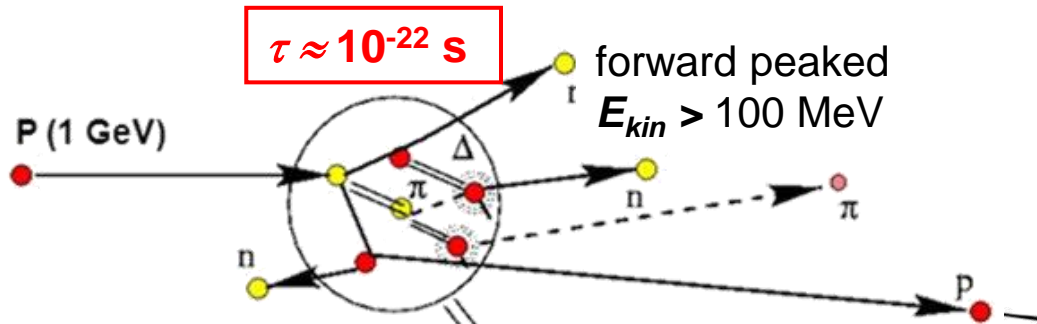
Beam dump at LHC:

7m long, \varnothing 0.7 m, graphite
900 tons of concrete shielding

Nuclear Physics Processes for Protons

Nuclear reactions via spallation for protons with $E_{kin} \geq 1$ GeV (simplified):

- Pre-equilibrium phases: π -exchange within $\approx 10^{-22}$ s with $E_{kin} > 20$ MeV \Rightarrow hadronic shower



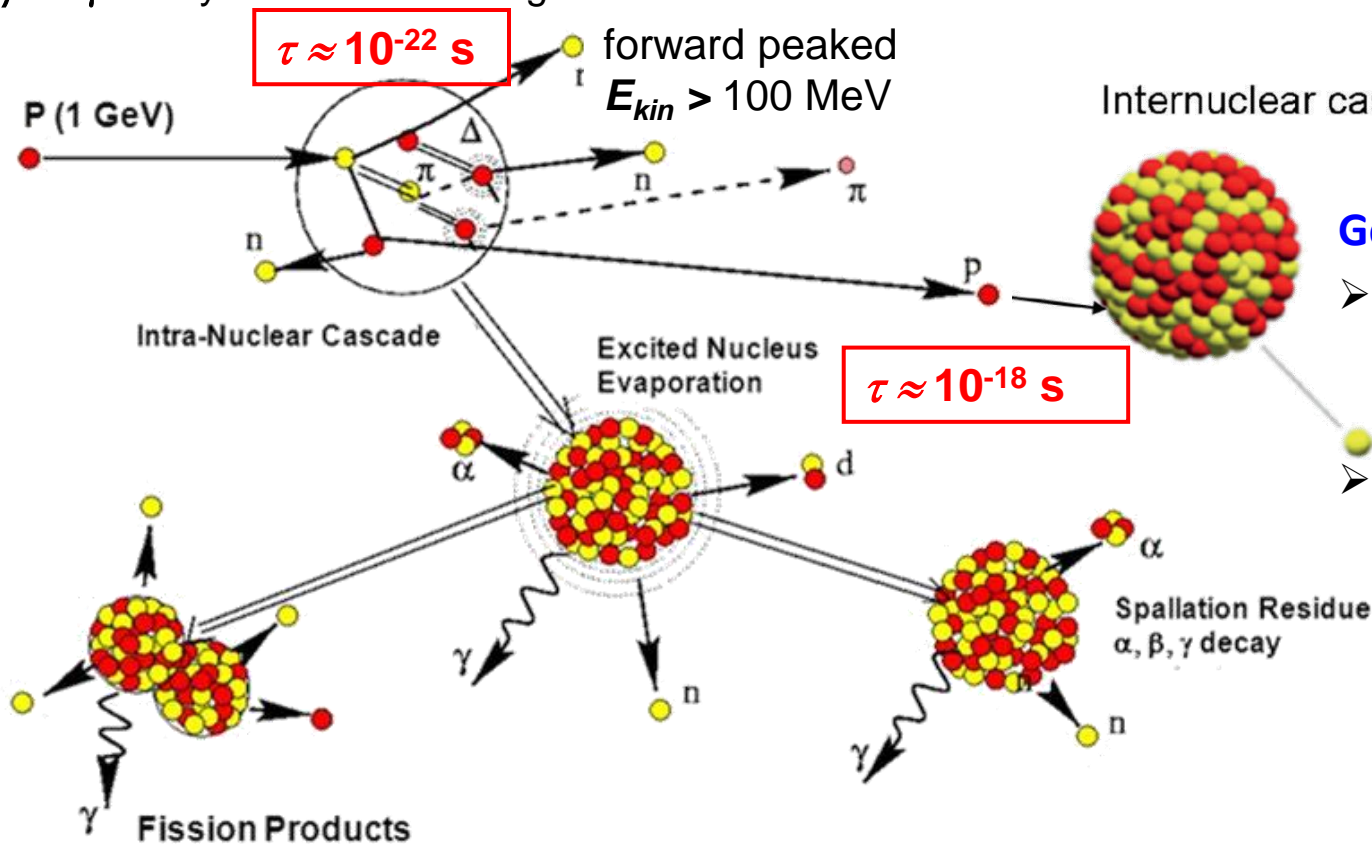
General properties:

- Binding energy:
 - ≈ 5 MeV outer nucleons
 - ≈ 50 MeV inner nucleons
- for $E_{kin} \gg 100$ MeV
comparable σ for n & p

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- Inter-nuclear cascade: Evaporation of n, p, d, α within $\approx 10^{-18}$ s with $E_{kin} \approx 1 - 10$ MeV
- Fission for heavy nuclei
- β & γ decay of nuclei with long lifetime $\tau \gg 10^{-9}$ s



General properties:

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 - ≈ 5 MeV outer nucleons
 - ≈ 50 MeV inner nucleons
- for $E_{kin} \gg 100$ MeV comparable σ for n & p

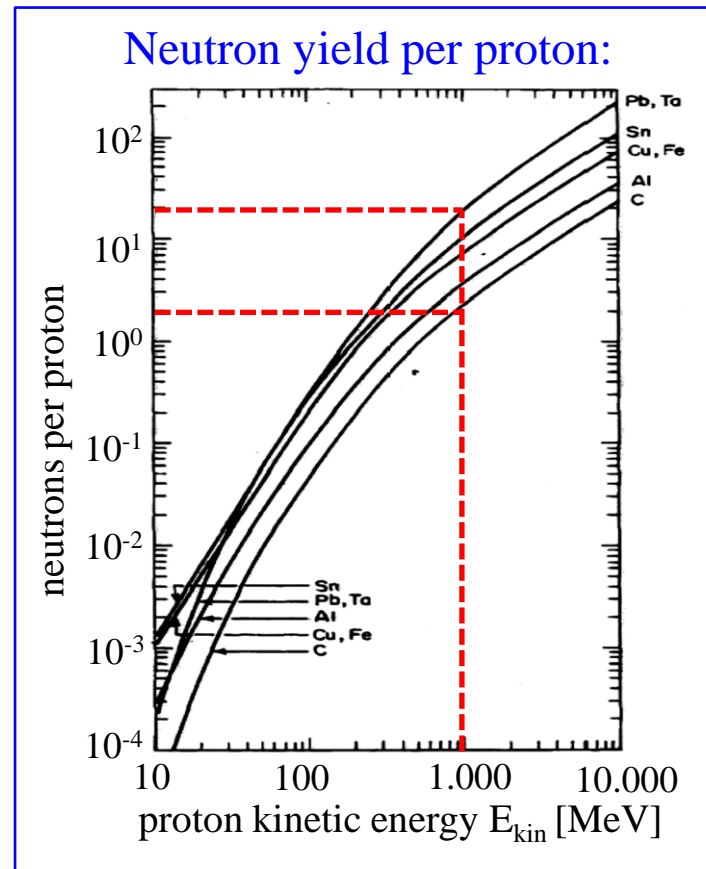
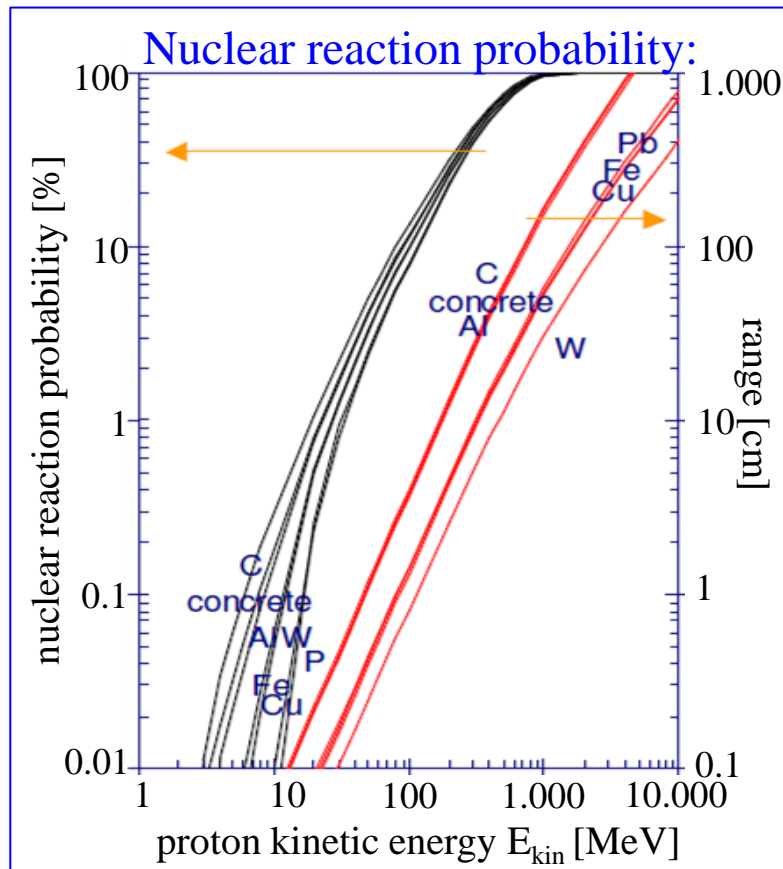
Result on long term $t > 1$ ms: Radioactive nuclei = activation

D. Kiselev, CAS 2011

Nuclear Physics Processes for Protons

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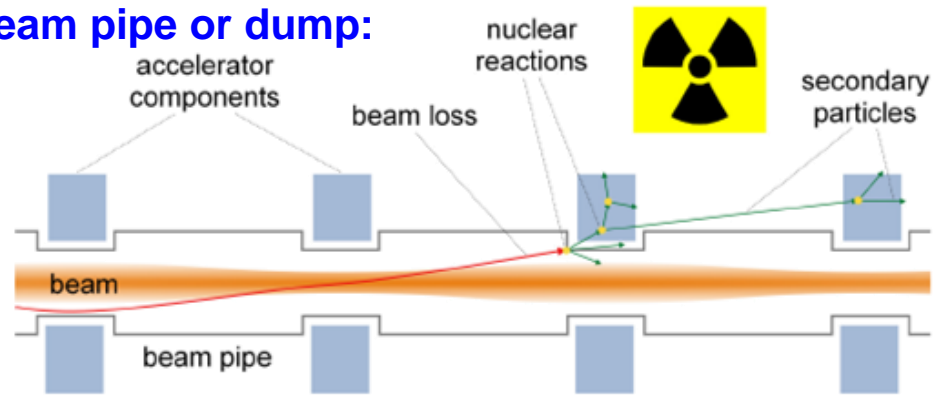
Thick target: Penetration depth comparable to range

R.H. Thomas, in Handbook on Acc. Phy. & Eng.

Impact of protons with $E_{kin} > 100$ MeV at beam pipe or dump:

- Emission of n, p, d, α , nucl. fragment ...
- β & γ decay of target nuclei on long time scale

Gracing incident: vacuum pipe is 'thick target'



Courtesy I. Strasik

Example of cross section for protons on steel beam pipe:

- Reaction: $\text{Fe} + p \rightarrow {}^{54}\text{Mn} + \text{something}$
- [$100 \text{ mb} \approx \frac{1}{10} \sigma_{\text{geo}}$]
- ${}^{54}\text{Mn}$ lifetime $t_{1/2} = 312$ days
- Electron capture $E = 1.3$ MeV to ${}^{54}\text{Cr}$ (excited) with X-ray emission of $E_\gamma = 0.54$ MeV
- ${}^{54}\text{Cr}$ decay via γ emission $E_\gamma = 0.83$ MeV

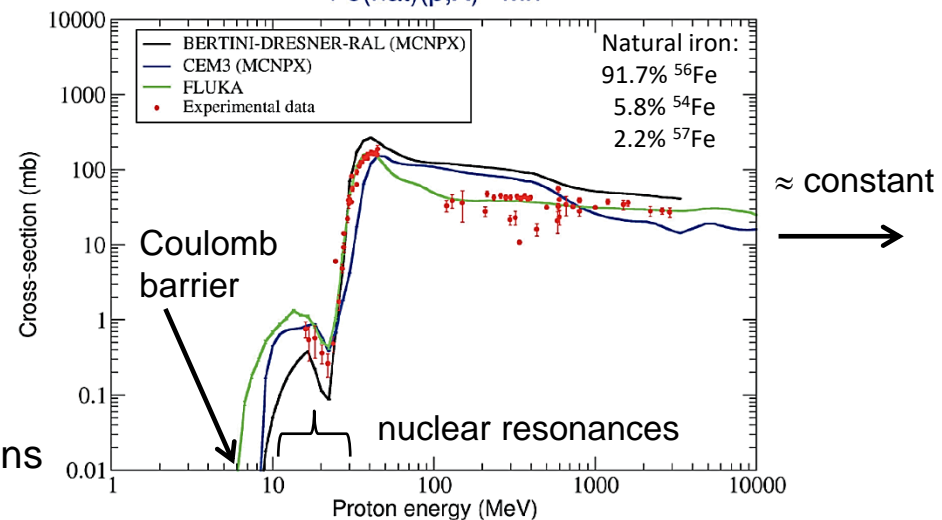
⇒ **activation of beam pipe**

Remark: Comparable cross section for fast neutrons

Coulomb barrier:

Kinetic energy required to overcome the electric potential to reach a distance for nuclear force ≈ 5 fm

$\text{Fe}(\text{nat})(p,X){}^{54}\text{Mn}$



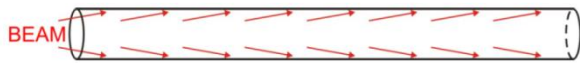
D. Kiselev, CAS 2011

Tolerable Beam Losses

Rule of thumb for proton beam with $E_{kin} > 100$ MeV:

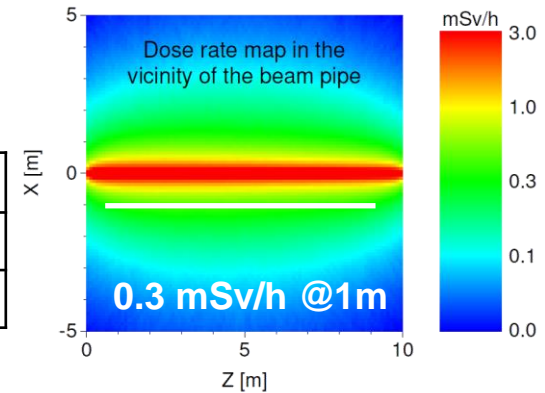
‘Beam loss below 1 W/m enables hands-on maintenance’

- **Example:** 1 W/m $\approx 6 \times 10^9$ protons/(m·s) at 1 GeV
 - **Care:** Most energy is lost by atomic process, while activation depends on nuclear physics
- \Rightarrow dependence on projectile and target



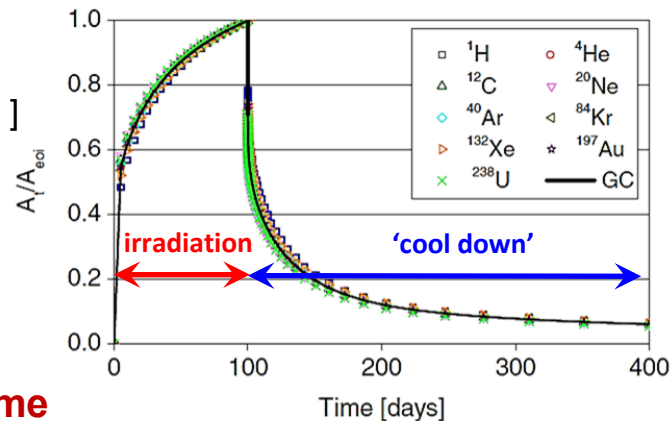
| | |
|-----------------------|-------------------|
| Natural background | ≈ 1 mSv/a |
| Medical X-ray CT | ≈ 3 mSv |
| Max. for rad. workers | 20 mSv/a |

Simulation for 1 GeV proton irradiation:
Stainless steel beam pipe after 1 W/m beam loss for 100 days & 4 h ‘cool down’



Simulation for 1 W/m losses for 1 GeV/u impact:

- 100 days irradiation of stainless steel No. 304 [Fe(70%), Cr(18%), Ni(10%), Mn(2%)]
 - Decrease of activation: $\approx 10\%$ after ≈ 1 year
 - Isotope mixture same for all ions
- \Rightarrow **highly activated material needs significant ‘cool down’ time**



I. Strasik et al., Phys Rev AB 13, 071004 (2010)

Rule of thumb: Light targets (C, Al ...) have lower activation for impact of same # particles

Secondary Particle Production for Electron Beams

Processes for interaction of electrons

For $E_{kin} < 10$ MeV:

Mainly electronic stopping \Rightarrow X-rays, slow e^-

For $E_{kin} > 10$ MeV:

Bremsstrahlungs- γ , forward peaked $E_\gamma = 5-50$ MeV

$\Rightarrow \gamma \rightarrow e^+ + e^-$ or $\mu^\pm \dots \rightarrow$ electro-mag. showers

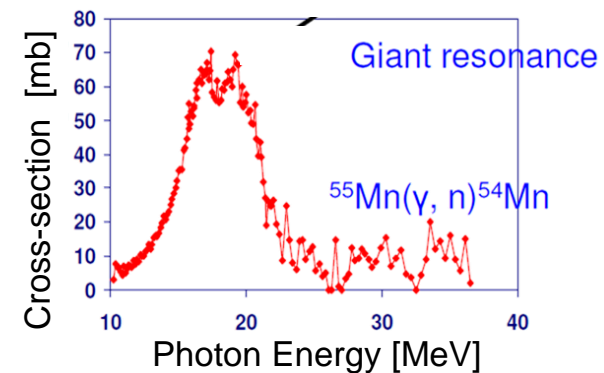
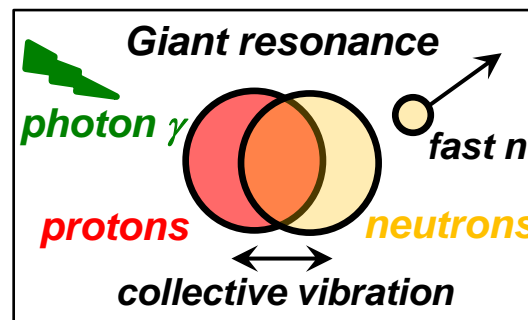
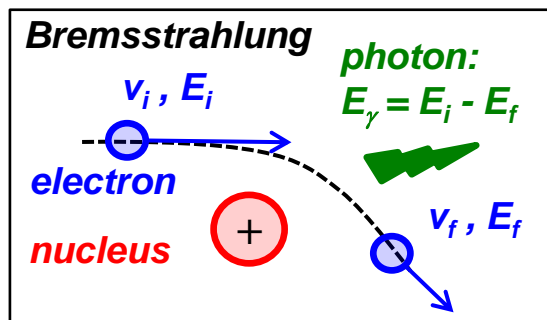
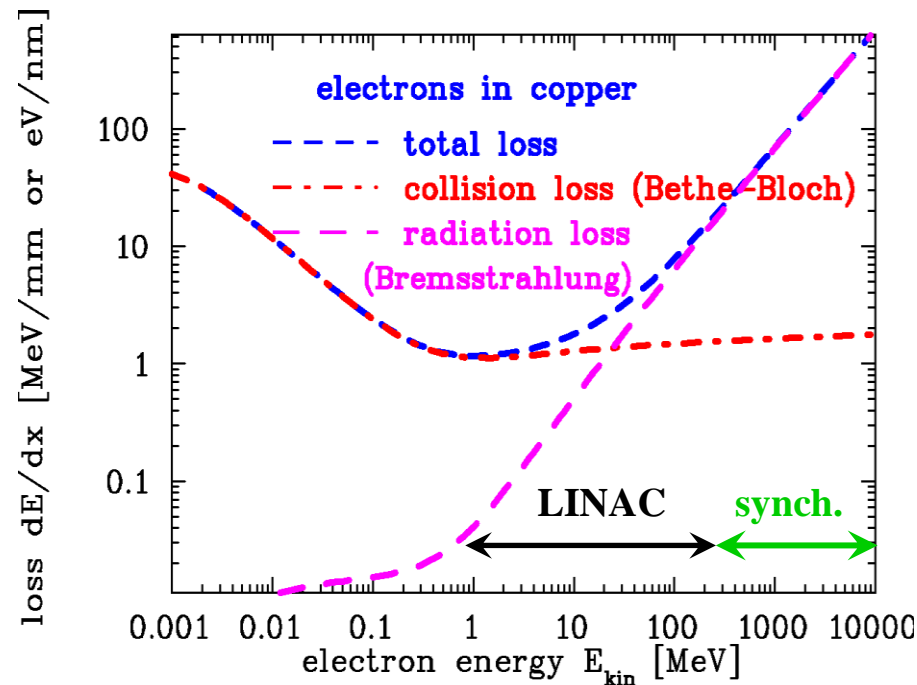
\Rightarrow Excitation of giant resonances $E_{res} \approx 10-30$ MeV

via (γ, n) , (γ, p) or (γ, np) with $\sigma_{giant} \approx \frac{1}{10} \sigma_{geo}$

\rightarrow Fast neutrons emitted (by nuclear interaction)

\rightarrow Neutrons: Long ranges in matter

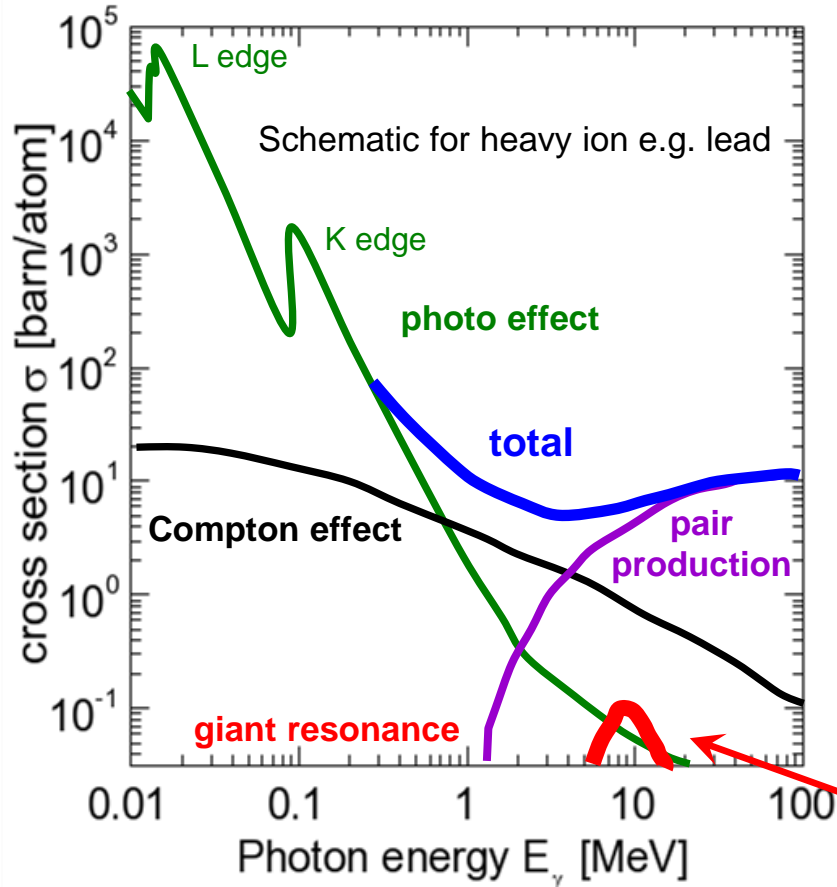
\Rightarrow **Neutrons at electron accelerators!**



Interaction of high Energy γ

At accelerators the γ are originated from nuclear reactions or Bremsstrahlung for e^- .

Example: Absorption in lead



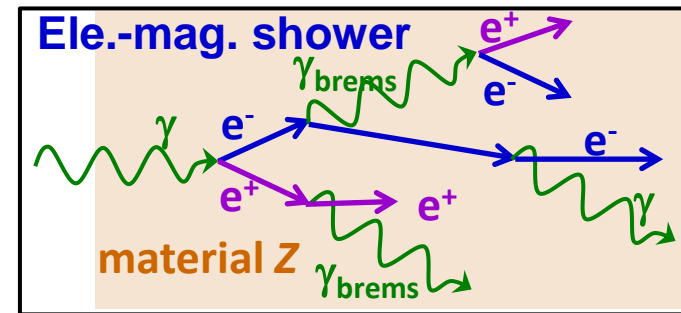
Atomic physics (Z = target nuclear charge):

Photo-effect: $\gamma + \text{atom} \rightarrow e^- + \text{atom}^+$
approx. material scaling $\sigma_{\text{photo}} \propto Z^4$

Compton-effect: $\gamma + \text{atom} \rightarrow \gamma' + e^- + \text{atom}^+$
approx. material scaling $\sigma_{\text{Comp}} \propto Z$

Pair production: $\gamma + \text{nucleus} \rightarrow e^- + e^+ + \text{nucleus}$
approx. material scaling $\sigma_{\text{pair}} \propto Z^2$

Ele.-mag. shower: for high E_γ
 $\gamma \rightarrow (e^-e^+) \rightarrow \gamma'_{\text{brems}} \rightarrow (e^-e^+) \rightarrow \gamma''_{\text{Brems}} \rightarrow \dots$



Nuclear physics:

Giant resonance: $\gamma + \text{nucleus} \rightarrow n + \text{nucleus}'$
small cross section but create free neutrons

Mass absorption coefficient $\mu = \frac{\rho N_A}{A} \cdot \sigma$
 ρ density, N_A Avogadro const., A atomic mass

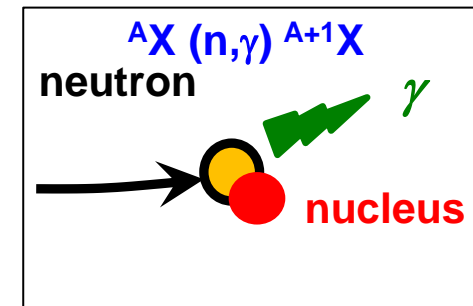
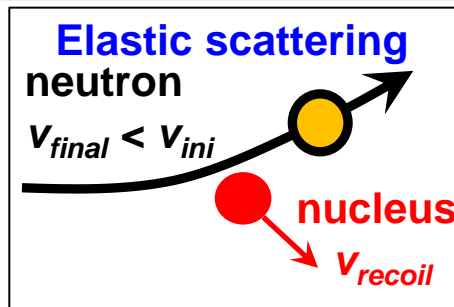
Courtesy C. Grupen, Xavier Queralt, JUAS

Interaction of Neutrons

Neutrons don't interact with electrons

Only nuclear physics processes:

- Elastic scattering: $X(n,n)X$
with X receiving recoil momentum
- Radiative capture with γ emission: ${}^AX(n,\gamma){}^{A+1}X$

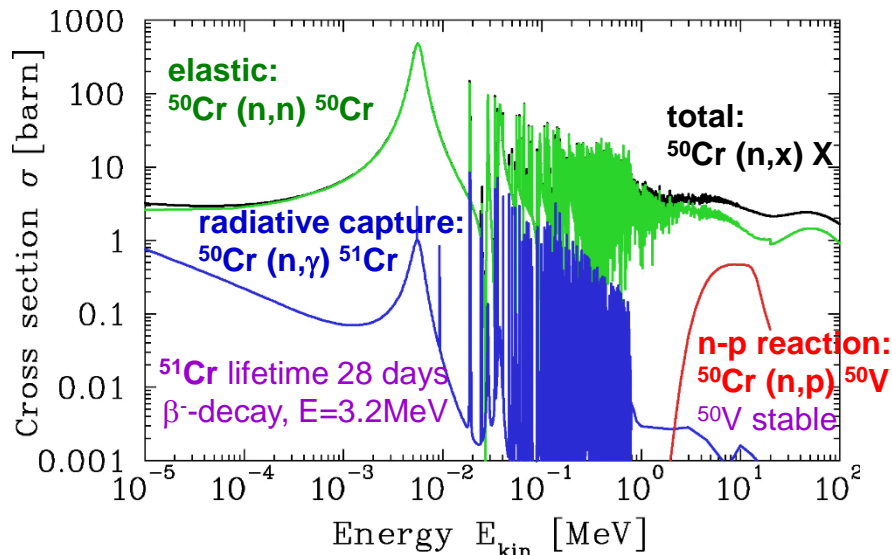


Example: Neutron on chromium ${}^{50}\text{Cr}$ (4.3% natural abundance, in stainless steel)

Elastic scattering: Cross section $\approx \sigma_{geo}$

Absorption: Large cross section at resonances
 γ emission and activation

For $E \gg 100$ MeV comparable cross section as proton

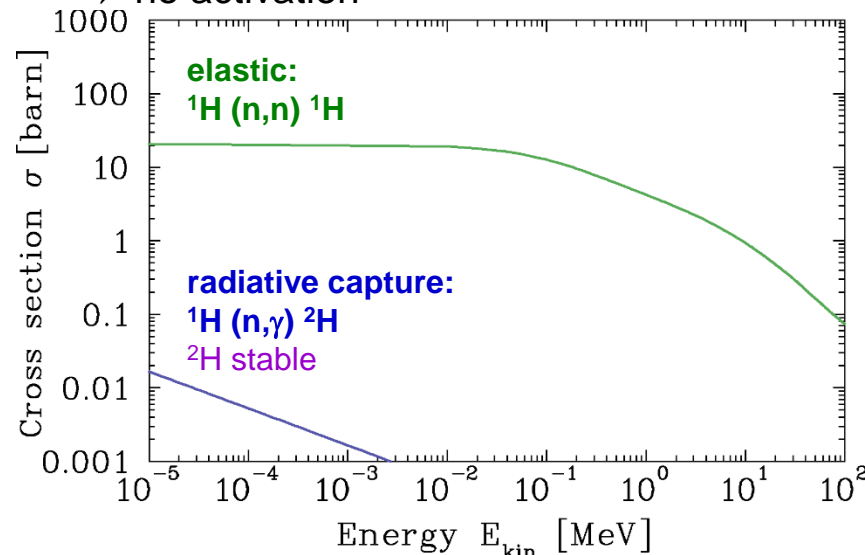


Example: Neutrons on H with $\approx 10 \sigma_{geo}$

e.g. organic materials, H_2O

→ effective moderator due to equal masses

→ no activation



Remark: Shielding of n by plastic ('paraffin') or concrete

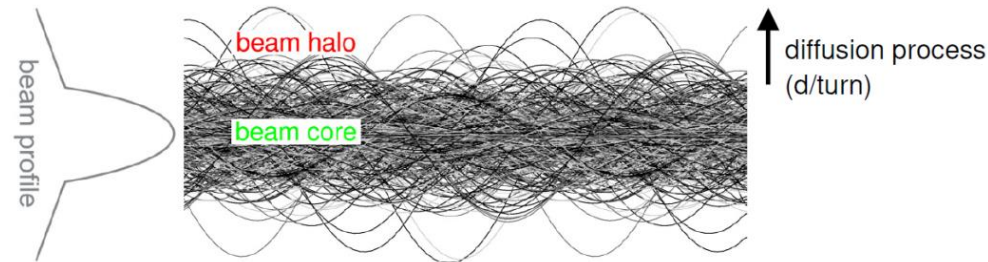
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Regular Losses from Halo

Halo formation at synchrotrons:

- Definition of halo: low density of particle with large betatron amplitude
- Caused by collective effect (e.g. space charge), resonances or machine errors
- Diffusion process (e.g. 1 μm per turn)



⇒ **unstable particles are lost**

Beam loss terminology: 'uncontrolled regular loss'

⇒ Beam halo collimation system at a synchrotron

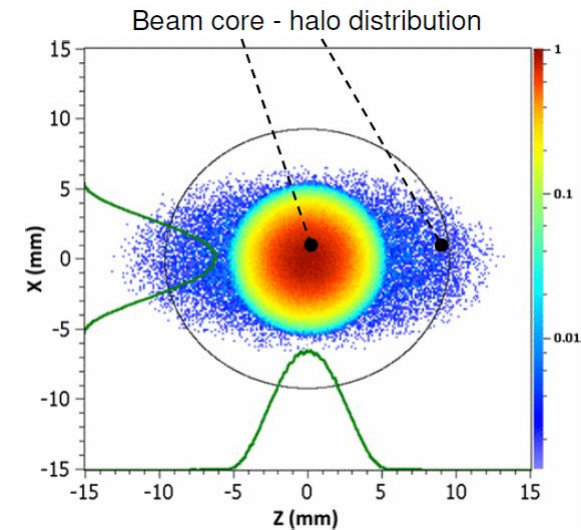
Goal: Low impurity beam

- **Warm synchrotron:** Protection of sensitive insertions (e.g. septum)
Concentration of loss at few locations
- **Super-conduction synch:** + quench protection of sc magnets
- **Collider:** + well defined condition for detector at IP
⇔ min. exp. background
Cleaning of collisional halo particles

⇒ Concentration of loss at dedicated locations i.e. 'controlled losses'

LINAC: Halo generation by long. and trans. mismatch

Goal: Quench protection of sc civilities



Remark:

- Halo might have other phase space distribution than core
- Halo formation and its mitigation is an actual topic

Quench Protection for superconducting Magnets

Superconducting magnets:

Beam particles energy loss
 ⇒ heat wires due to energy loss

Quench: Transition to normal-conducting phase

Goal: Beam dump prior to quench !!!

Simulation of temperature increase ΔT :

Energy deposition: $\frac{dE}{dV} = -\frac{dE}{dx} \cdot \frac{N}{A} \left[\frac{J}{cm^3} \right]$

N : number of particles, A : cross section

Temperature rise: $\Delta T = \frac{dE}{dV} \cdot \frac{1}{\rho c_p(T)} \text{ [K]}$

ρ : mat. density, c_p specific heat

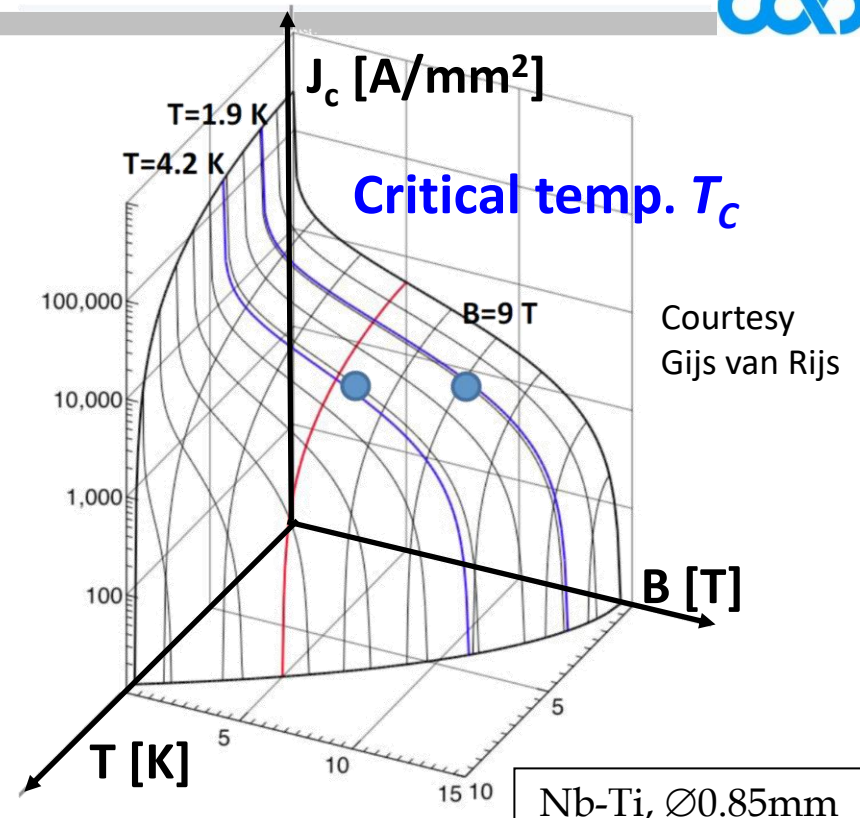
Temperature dependent specific heat:

Insulator: $c_{phonon}(T) \propto T^3$

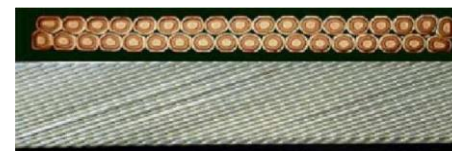
Normal conductor: $c_{NC}(T) \propto \alpha T + \beta T^3$

Superconductor: $c_{SC}(T) \propto T_c e^{-\gamma T/T_c}$

α, β, γ material constants



'Rutherford' cable strand



Nb-Ti, Ø0.85mm



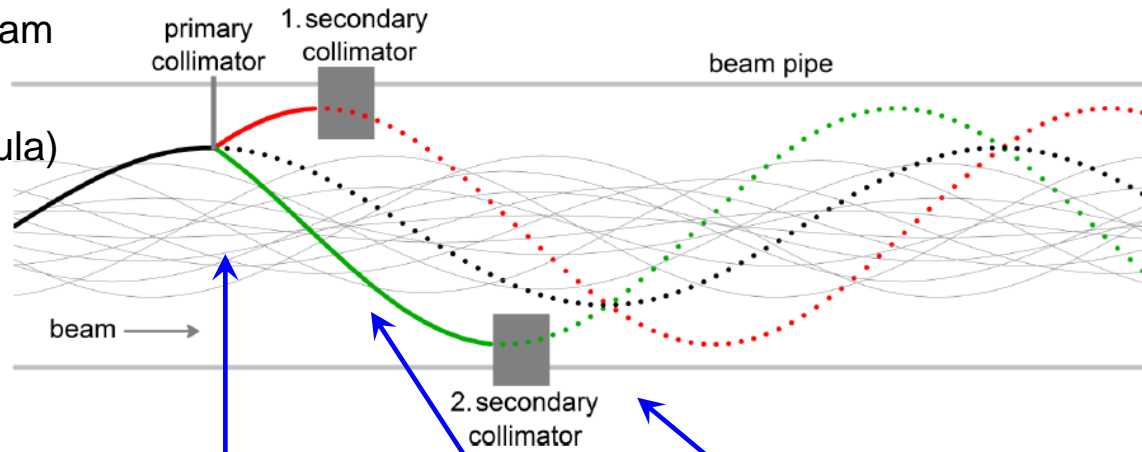
$J \sim 1500-2000 \text{ A/mm}^2$
 $I \sim 400 \text{ A}, B = 8-9 \text{ T}$

See lecture 'Superconducting Magnets' by Gijs van Rijk

Two Stage Betatron Collimation System = active Collimation

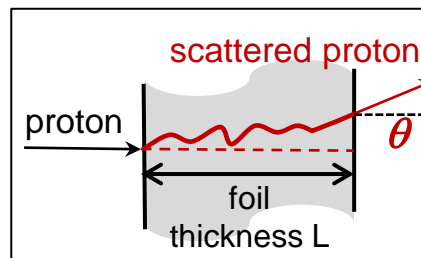
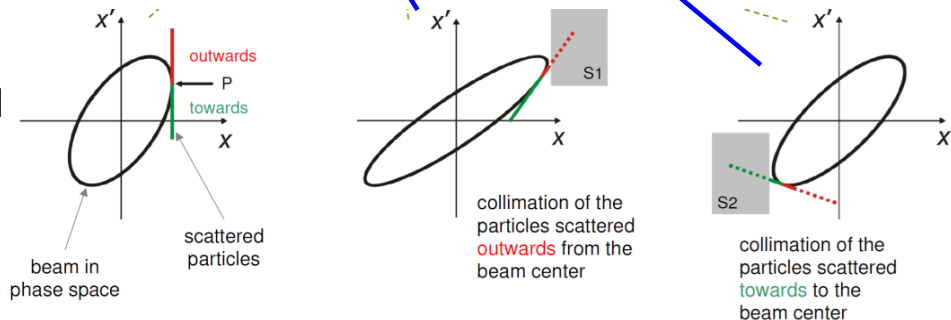
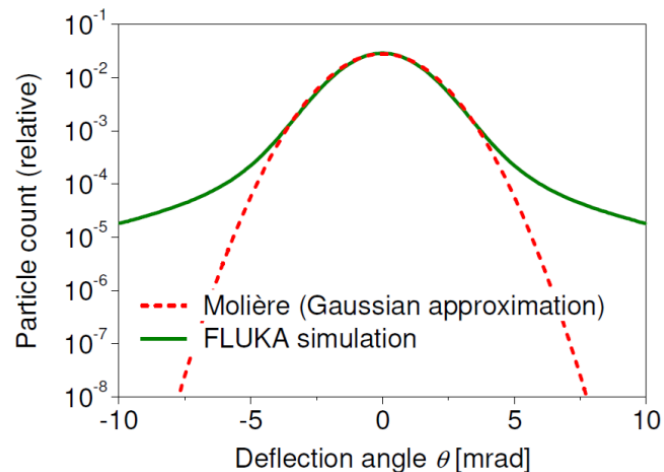
General functionality of cleaning:

- Primary stage as **thin foil close** to beam
 ⇒ scattering of halo particles
 (Coulomb scattering by Moliere formula)
- Betatron amplitude increases
- Max. extension after $\mu \approx 90^\circ$ or 270° betatron phase
- Secondary collimator as absorber more distant to beam



Example:

4.7 GeV scattering in L=1 mm Tungsten foil



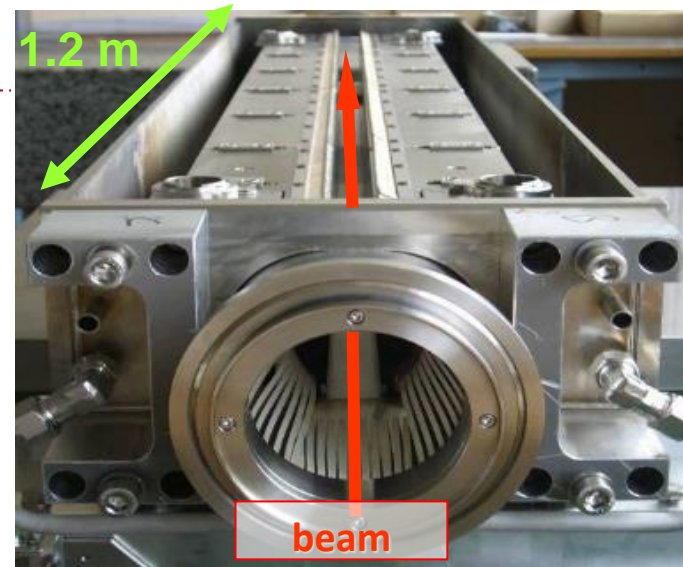
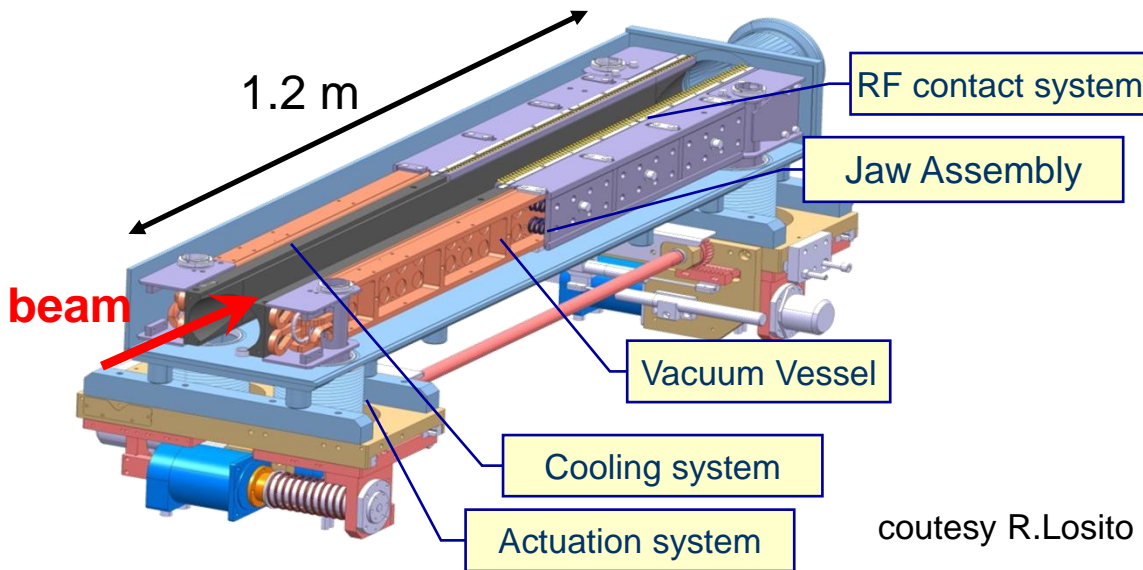
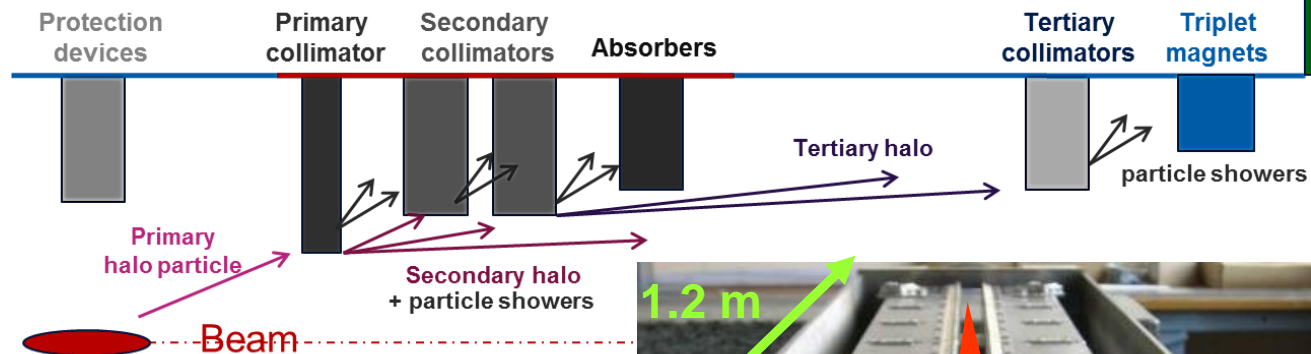
Courtesy I. Strasik CAS 2016

LHC Collimator Hardware

LHC Collimator system:

- Primary stage
- Secondary & tertiary stage
- Absorbers

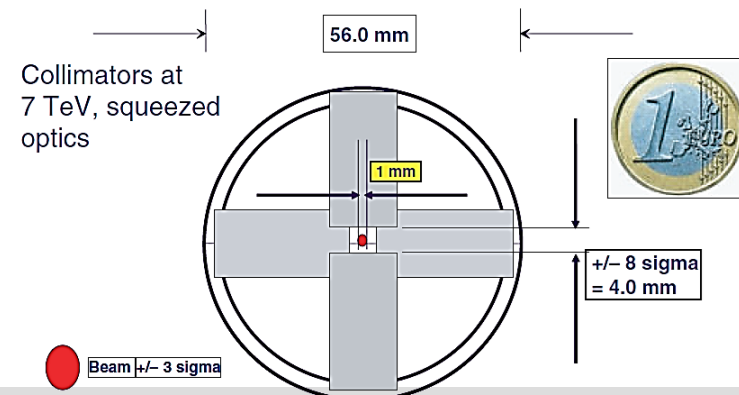
in total 110 movable devices



courtesy R.Losito

LHC maximal losses for 6.5 TeV protons:

- Total stored power 300 MJ
- Max. energy deposition in sc magnet: 0.1 J/cm²
- Corresponding to 6x10⁷ protons
- Or 2x10⁻⁷ of the stored beam of 3x10¹⁴ protons



LHC Collimator System

LHC Collimator system:

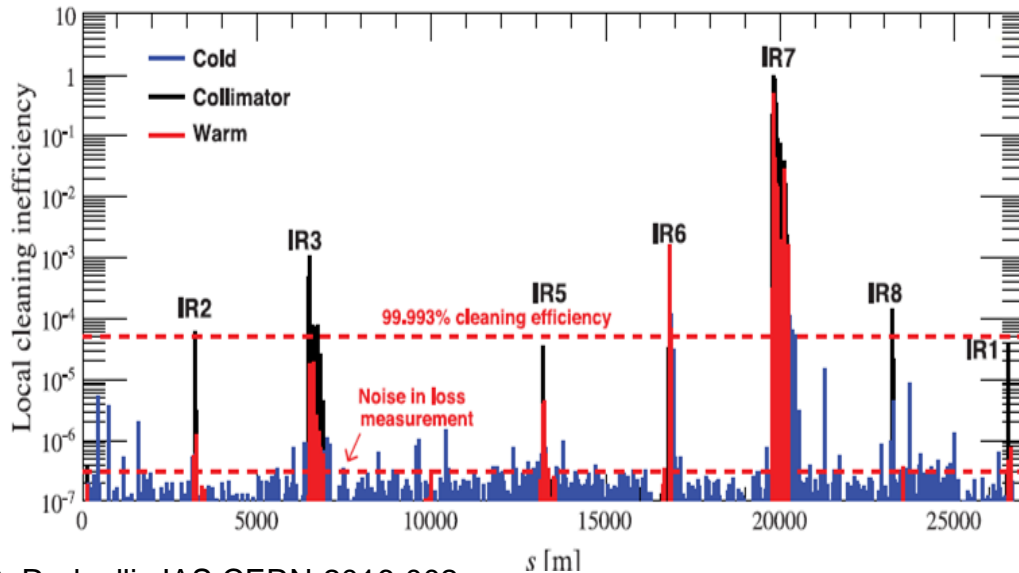
- Primary stage as close as $\approx 5\sigma_{\text{beam}} \approx 1 \text{ mm}$
 - Secondary & tertiary stage made of carbon
 - Absorbers made of tungsten alloy
- in total 110 movable devices moving e.g. from injection $r = 5 \text{ mm} \rightarrow 1 \text{ mm}$

Test of functionality:

- Loss concentrated at collimators

Experimental verification: Single bunch excitation

Result: Main losses concentrated at collimators



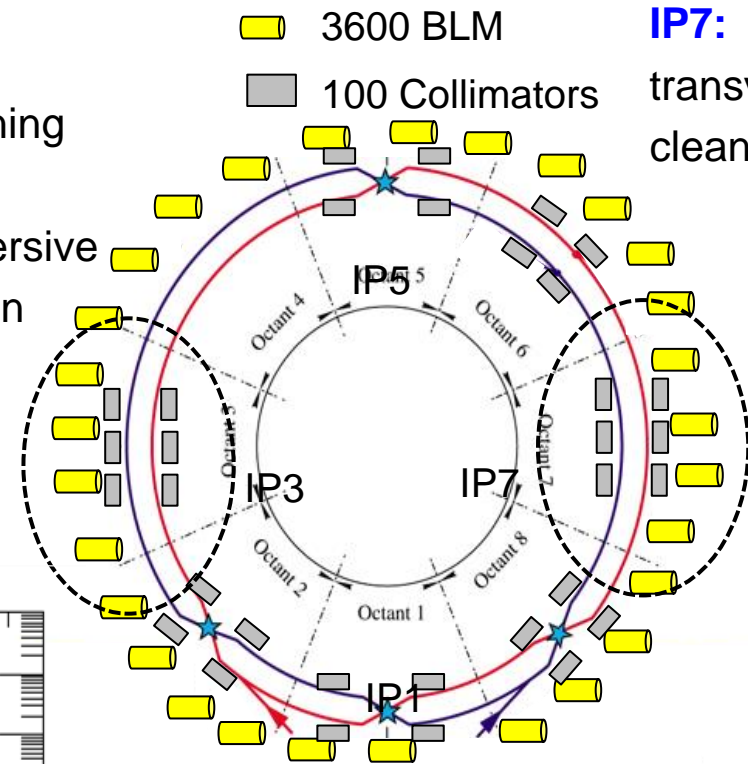
IP3:

long. cleaning at dispersive region

- 3600 BLM
- 100 Collimators

IP7:

transverse cleaning



Cleaning efficiency:

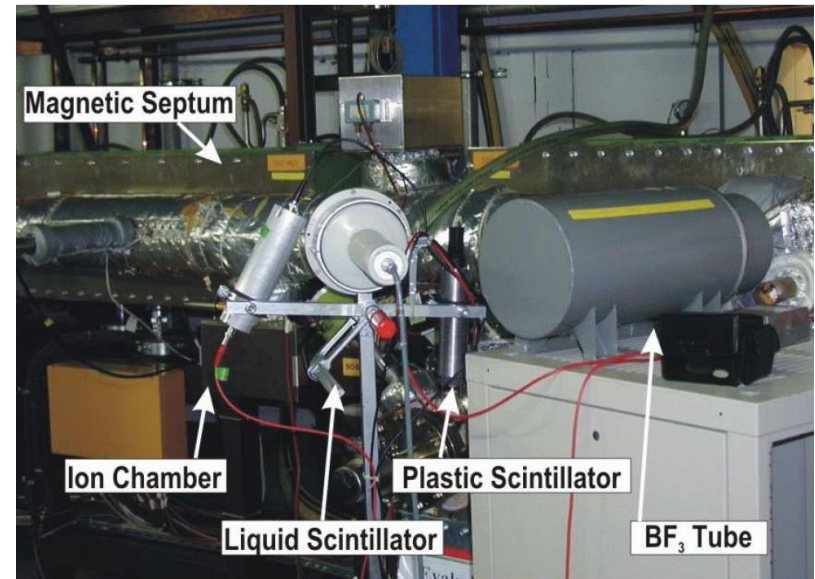
$$\eta = (\text{protons lost at collimator}) / (\text{total beam loss})$$

Result: $\eta = 99.8 \%$ reached

Courtesy M. Zerlauth, CAS 2018

Outline of this talk:

1. Introduction to risk & destruction potential
2. Important atomic and nuclear physics
3. Collimation as passive protection
4. **Measurements by Beam Loss Monitors**
5. Design of Machine Protection System
6. Overview of personal safety



Basic Idea of Beam Loss Monitors

Basic idea for Beam Loss Monitors BLM:

A loss beam particle must collide with the vacuum chamber or other insertions

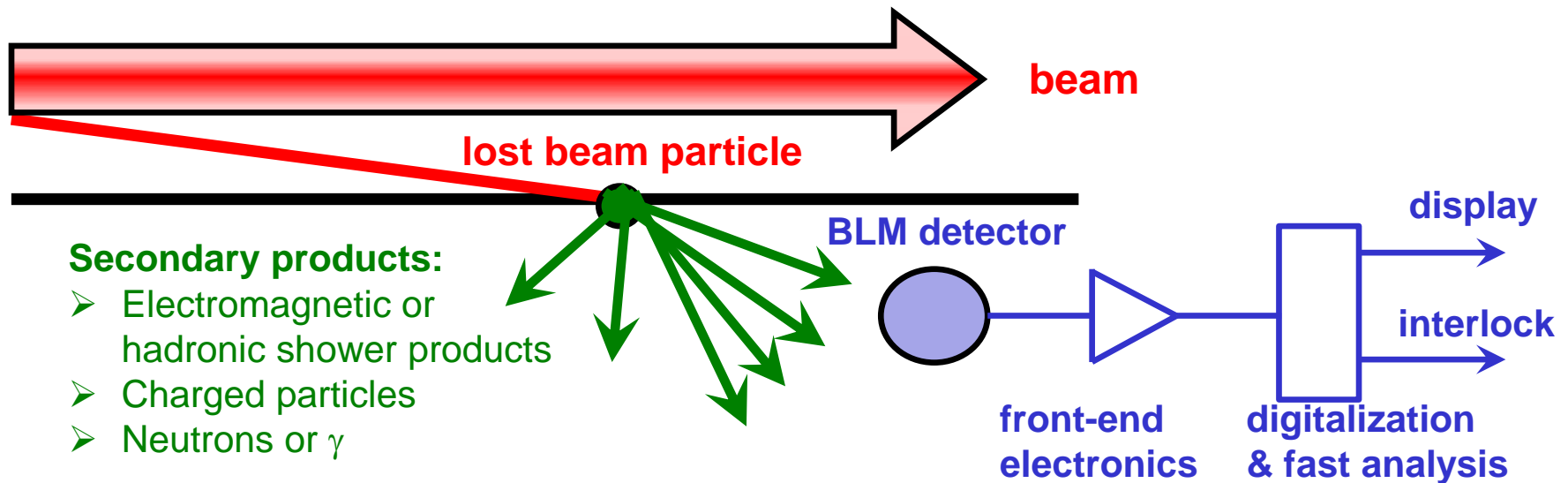
⇒ Interaction leads to some shower particle:

e^- , γ , protons, neutrons, excited nuclei, fragmented nuclei

→ Detection of these secondaries by an appropriate detector outside of beam pipe

→ Relative cheap detector installed at many locations

Remark: Due to grazing angle a thin vacuum chamber might be a ‘thick target’
vacuum pipe



Scintillators as Beam Loss Monitors

Plastics or liquids are used:

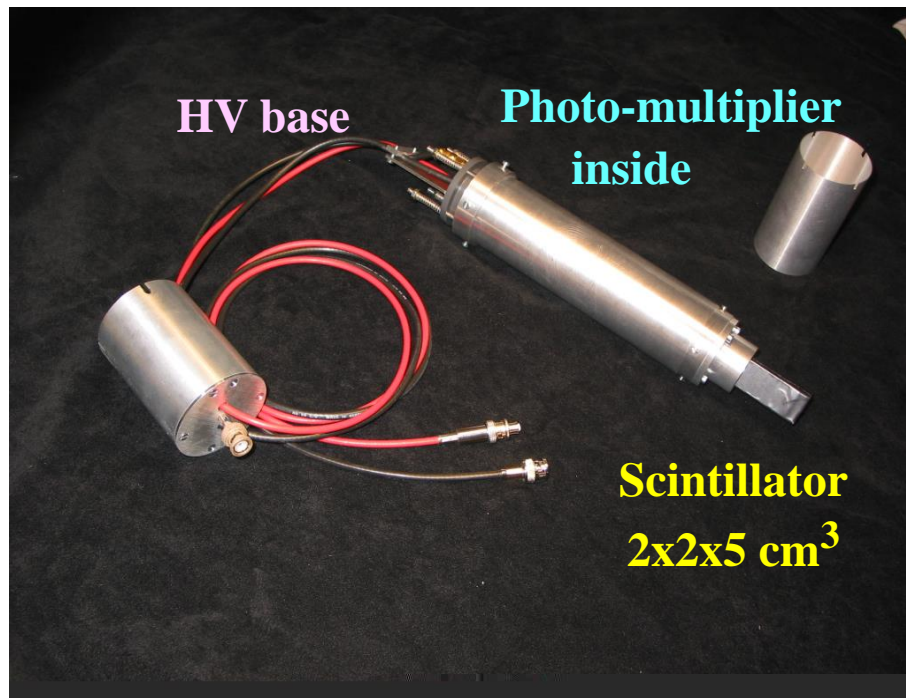
- Detection of **charged particles** by electronic stopping
- Detection of **neutrons** by elastic collisions n on p in plastics and fast p electronic stopping.

Scintillator + photo-multiplier:

counting (large PMT amplification)
or analog voltage ADC (low PMT amplification)

Radiation hardness:

plastics 1 Mrad = 10^4 Gy
liquid 10 Mrad = 10^5 Gy

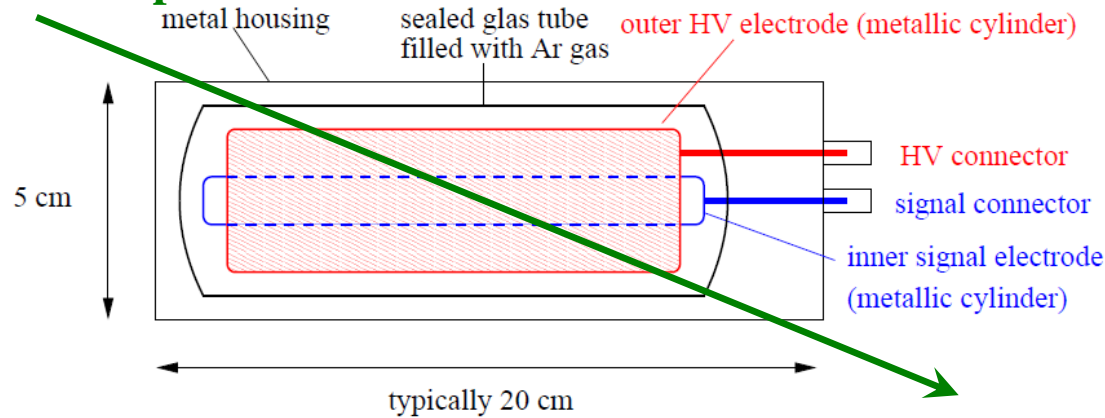


Ionization Chamber as Beam Loss Monitors

Energy loss of charged particles in gases → electron-ion pairs → current meas.

$$I_{\text{sec}} \propto \frac{1}{W} \cdot \frac{dE}{dx} \Delta x$$

shower particle



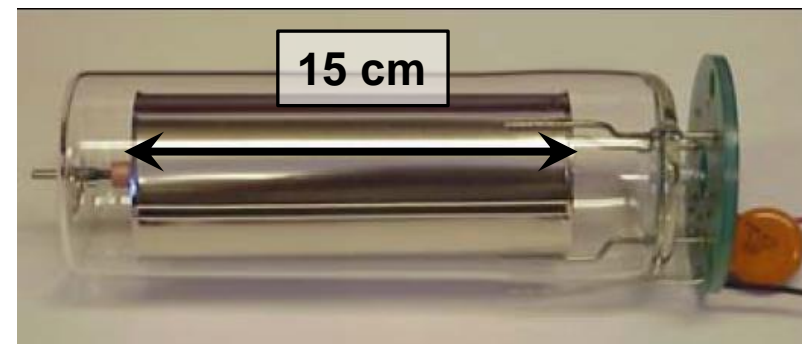
W is average energy for creation for one e^- -ion pair:

| Gas | Ionization Pot. [eV] | W-Value [eV] |
|----------------|----------------------|--------------|
| Ar | 15.7 | 26.4 |
| N ₂ | 15.5 | 34.8 |
| O ₂ | 12.5 | 30.8 |
| Air | | 33.8 |

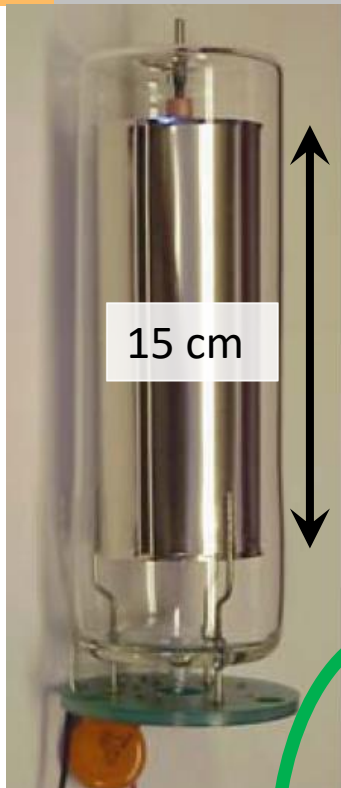
Sealed tube Filled with Ar or N₂ gas:

- Creation of Ar⁺-e⁻ pairs, average energy $W = 32$ eV/pair
- measurement of this current
- Slow time response due to $\approx 10 \mu\text{s}$ drift time of Ar⁺.

Per definition: Direct measurement of dose !



Ionization Chamber as BLM: TEVATRON and CERN Type

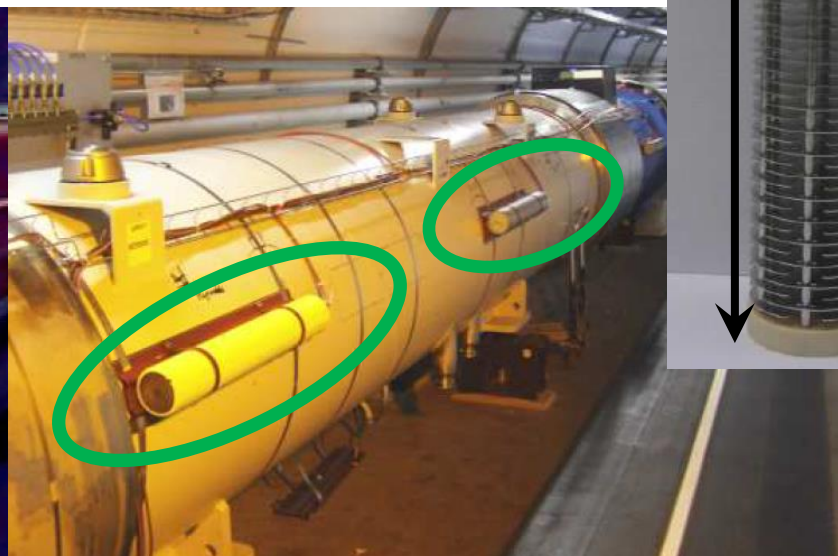
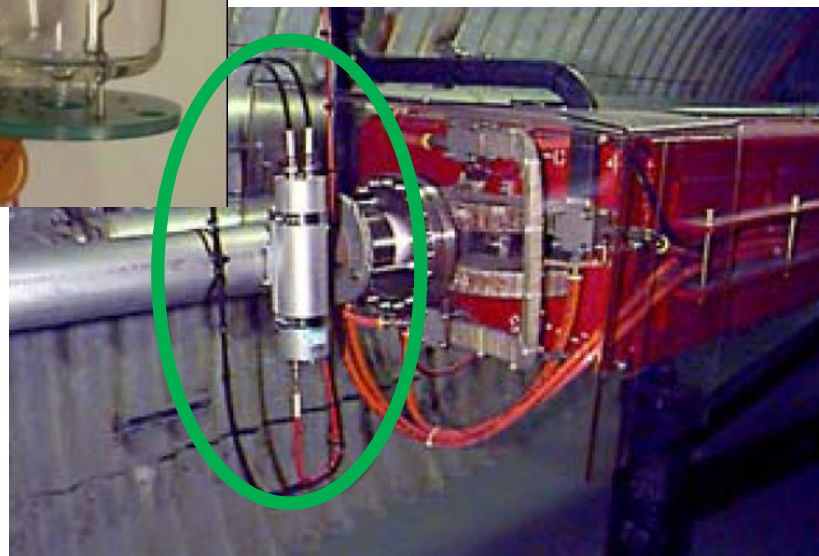
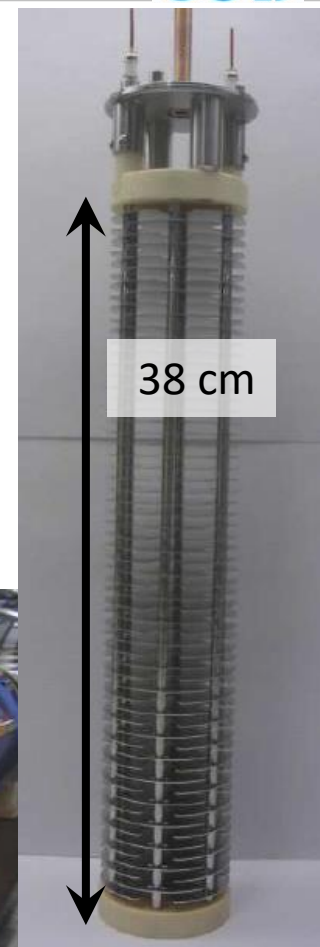


| Parameter | TEVATRON, RHIC | CERN type |
|-----------------|--------------------------|---------------------------|
| Size | 15cm, \varnothing 6 cm | 50 cm, \varnothing 9 cm |
| Gas | Ar at 1.1 bar | N ₂ at 1.1 bar |
| # of electrodes | 3 | 61 |
| Voltage | 1000 V | 1500 V |
| Reaction time | 3 μ s | 0.3 μ s |

4000 BLMs at LHC \leftrightarrow each \approx 6m

TEVATRON, RHIC type

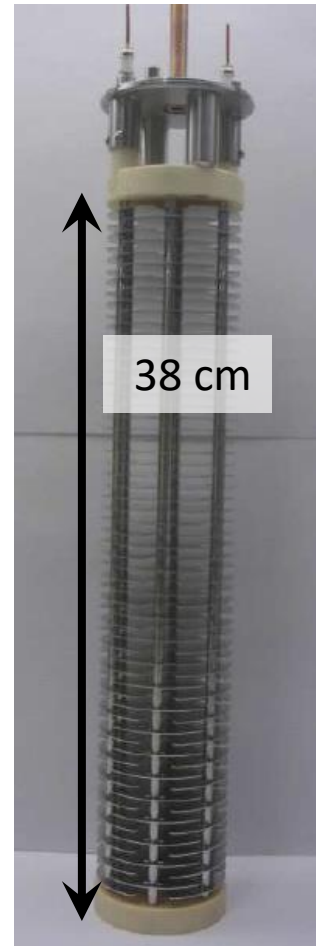
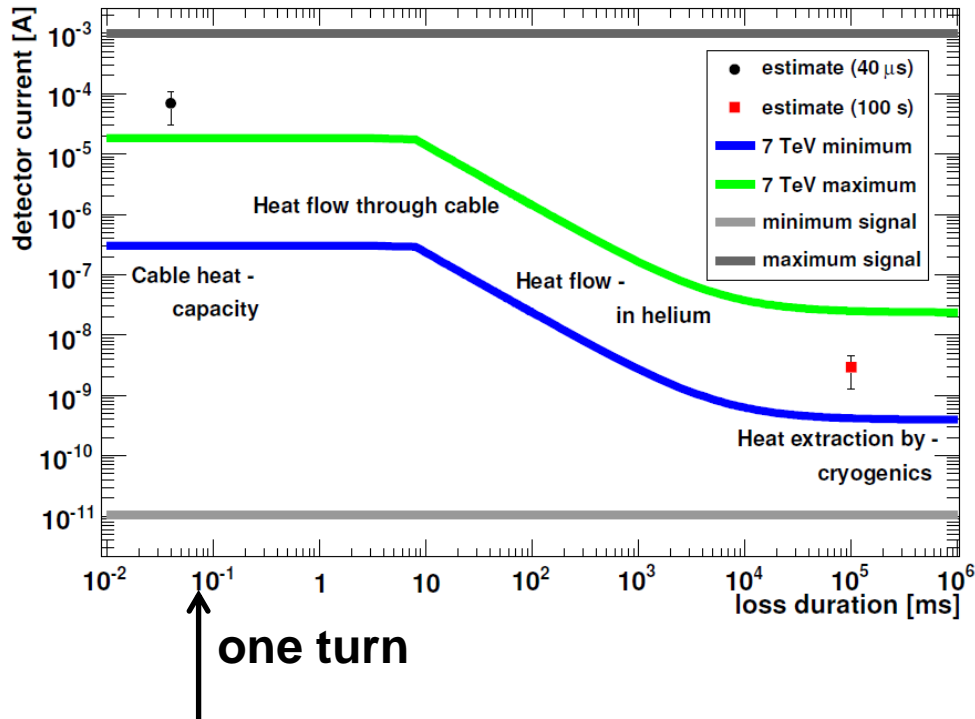
CERN type



Ionization Chamber as BLM: CERN Type

Simulation of det. efficiency by Geant4:

- Most sensitive to protons, electrons & high energy γ
- Low sensitive to neutrons
- ⇒ Calculation of lost protons by integrating of shower composition
- ⇒ **Quench limit estimation**



M. Stockner: PhD-Thesis, Tech. Uni. Vienna
 A. North et al., HB 2010

Comparison of different Types of BLMs

Different detectors are sensitive to various physical processes very different count rate, but basically proportional to each other

Typical choice of the detector type:

➤ Ionization Chamber:

Advantage:

- Measurement of absolute dose

Disadvantage:

- Low signal (low γ , eff, no neutron detection),
- Sometimes slow, ion drift time 10 ... 100 μ s

⇒ Often used at **proton** accelerators

➤ Scintillator, Cherenkov detector:

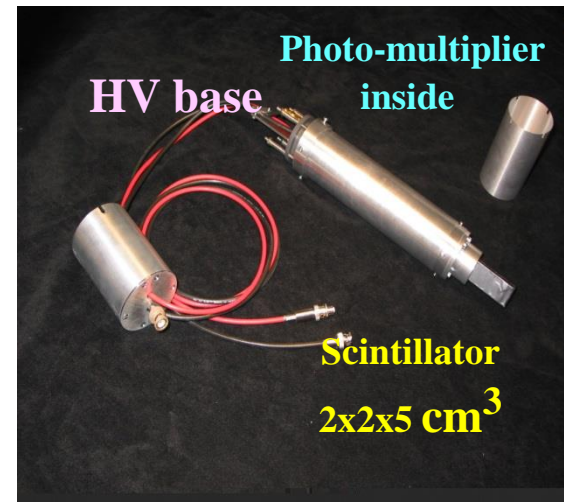
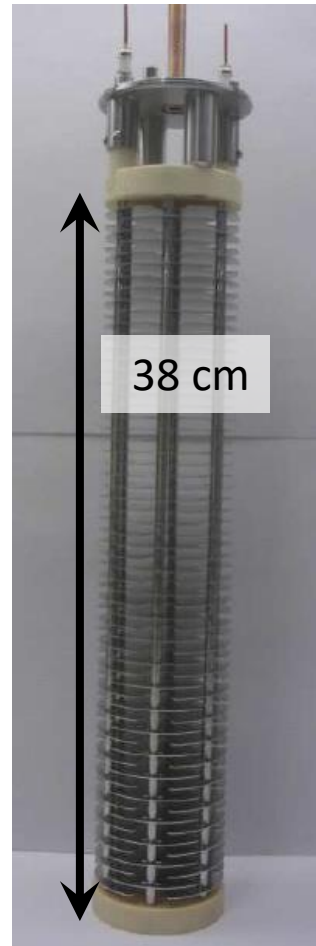
Advantage:

- Fast current reading or particle counting
- Can be fabricated in any shape, cheap

Disadvantage:

- Need calibration in many cases
- Might suffer from radiation

⇒ Often used at **electron** accelerators



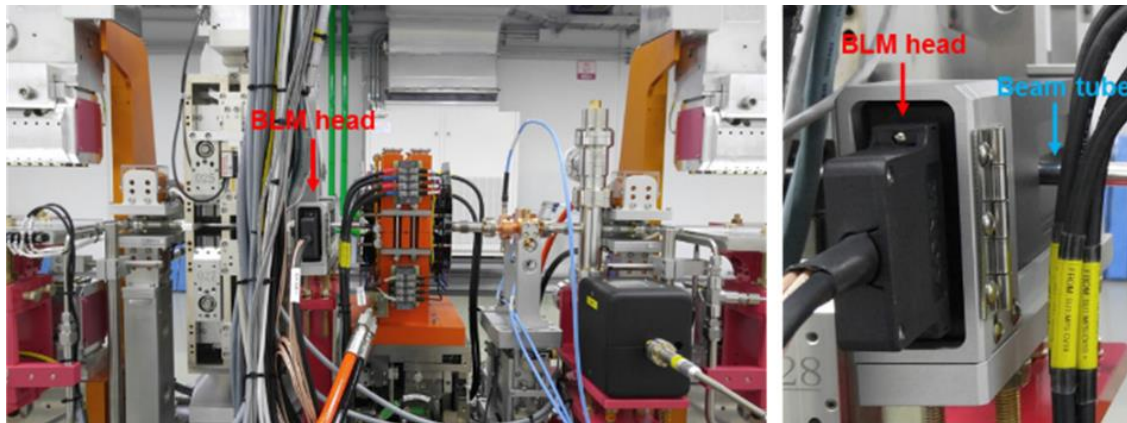
Cherenkov Light Detectors as Beam Loss Monitors

Cherenkov detectors:

Passage of a charged particle v faster than propagation of light $v > c_{medium} = c/n$

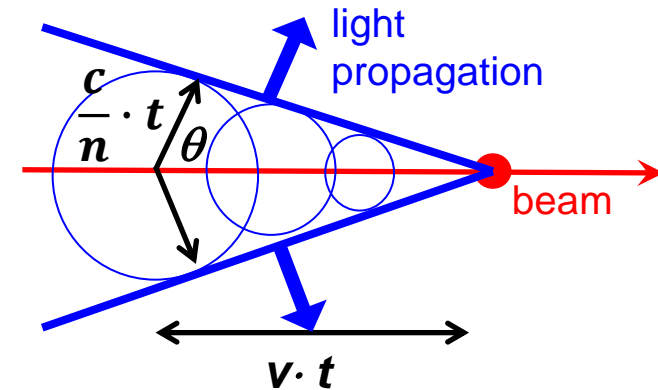
Technical: Quartz rod $n=1.5$ & photomultiplier

Example: Korean XFEL behind undulator



Cherenkov light emission:

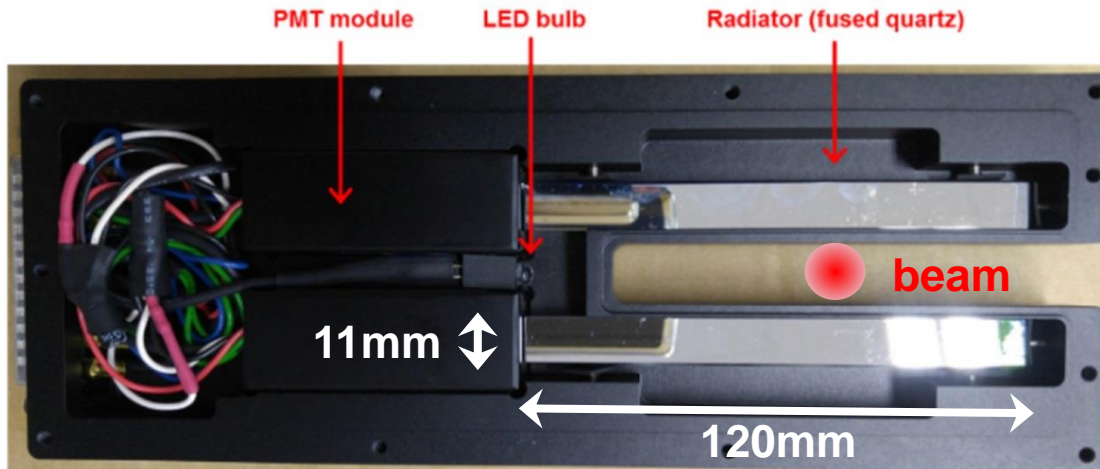
For $v > c_{medium} = c/n$
light wave-front like a wake
broadband light emission



Advantage:

- **Detection of fast electrons only**
not sensitive to γ & synch. photons
- No saturation effects
- Prompt light emission

Usage: Mainly at FELs for short and intense pulses



H. Yang, D.C. Shin, FEL Conf. 2017

Outline of this talk:

1. Introduction to risk & destruction potential
2. Relevant atomic and nuclear physics
3. Collimation as passive protection
4. Measurements by Beam Loss Monitors
- 5. Design of Machine Protection System**
6. Overview of personal safety

Relevant Losses for Machine Protection

Type of losses:

1. **Regular losses** or slow losses \Leftrightarrow **unavoidable losses**

- Caused by lifetime inside synchrotron (residual gas, Touschek ...),
- Caused by aperture limitation, beam manipulations
- Usually a few % of the beam intensity

2. **Irregular losses** or fast losses by malfunction \Leftrightarrow **avoidable losses**

- Occurs only seldom
- A significant fraction of the beam is lost
- Usually within a short period of the operational cycle (e.g. injection, acceleration, extraction, ...)
 - ⇒ Requirement for detector system: large dynamic range
- Usually caused by
 - Hardware failures, inaccurate settings or control errors (magnets, cavities ...)
 - Beam instabilities (wake-fields, resonances, ...)
 - Manually initialized improper beam alignment

⇒ **Beam abortion required to prevent for destruction via interlock generation**

Remark:

Machine protection: Appropriate BLMs, device specific loss threshold → might be more complex

Personal safety system: Simple devices, reliable technology → based on dose threshold [Gy/s]

General Layout of a Machine Protection System: Design

Design criteria for a Machine Protection System:

1. Beam based: *Choice of BLM detector type*

- Main type of radiation (protons, neutrons, electrons, muons.....)
- Expected radiation level at foreseen location
- Required time response (fast particle counts or short beam delivery ↔ medium fast IC ↔ slow IC)
- Required dynamic range to detect **irregular** losses e.g. 6 orders of magnitude!
- **Proton accelerators:** Most often IC are used for interlock-generation & particle counters for relative measurements (after calibration suited for interlock generation)

Electron accelerators: Scintillators and Cherenkov counters (partly due to short pulse operation)

2. Equipment based: **Functionality of any relevant device must be guaranteed**

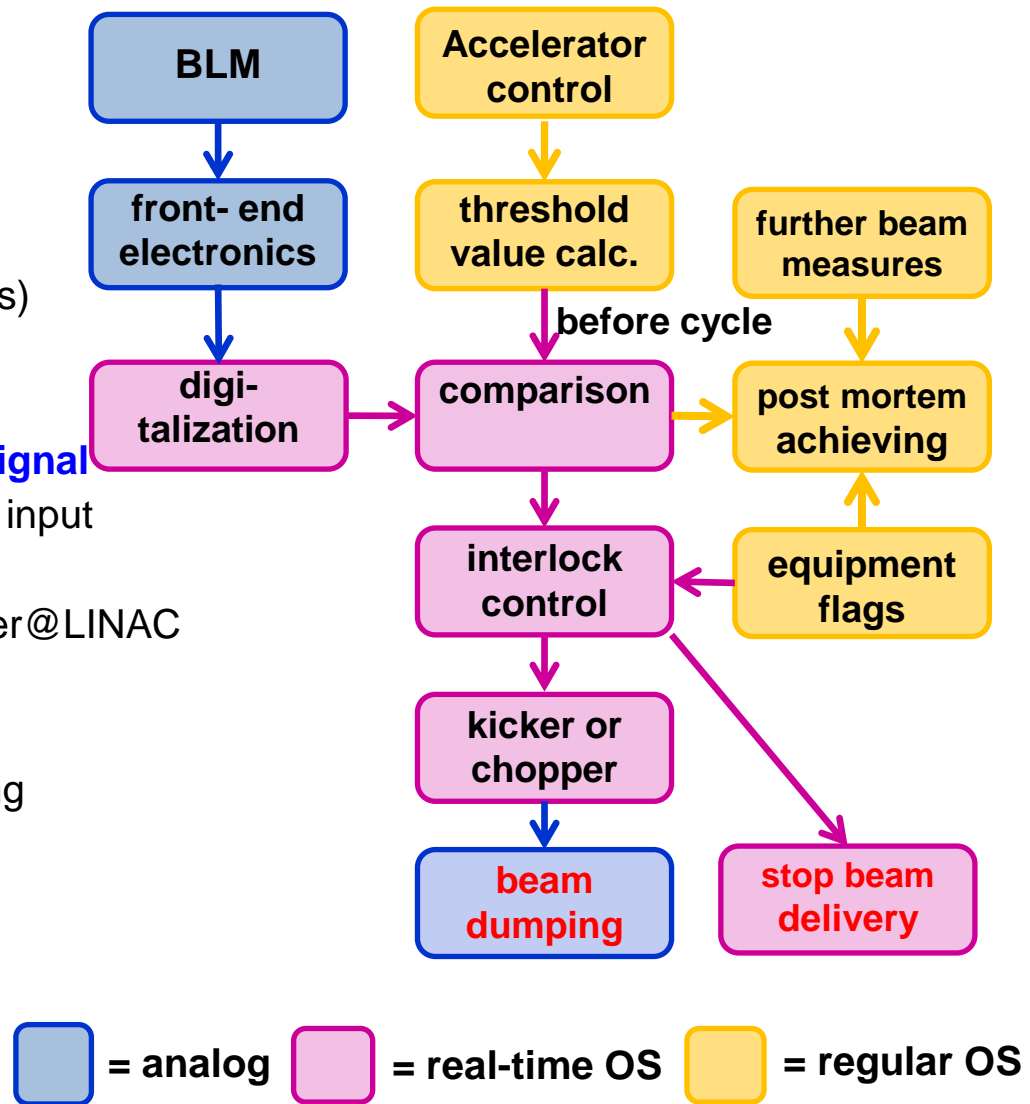
- Magnet power supplier, rf-generators, cavity properties
 - Super-conducting state of magnet or cavity
 - Vacuum conditions
 - Relevant diagnostics instruments
- ⇒ **Control system watchdog for all important devices**

Remark: In exceptional cases an interlock-source can be masked to allow for acc. operation

General Layout of a Machine Protection System: Hardware

Design of a protection system:

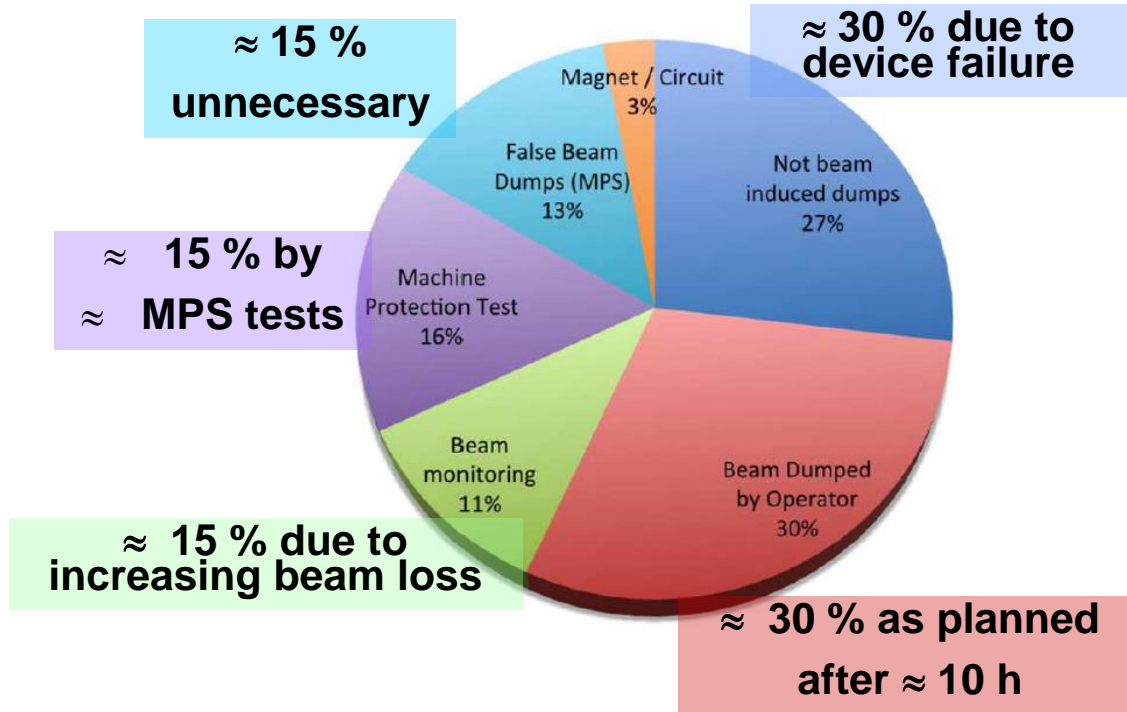
- **BLM detector & analog front-end**
low input signal under regular losses
large dynamic range for irregular losses
e.g. current-frequency converter
- **Digitalization**
high time resolution (e.g. LHC 1 turn = 89 μs)
- **Comparison to threshold values**
fast, real-time calculation (FPGA, DSP)
- **Generation & broadcasting of interlock signal**
real-time operation required, equipment ok input
- **Beam permit:** if not ok:
→ beam abortion kicker@synchr. or chopper@LINAC
→ disable next beam production
- **Data logging**
→ detailed 'post mortem' storage & archiving
→ error display
- **Generally**
robust & fail-safe system required!
challenge: large dynamic range



Statistics for Interlock Generation

Beam dump statistics at LHC in year 2015 and 2012 (above injection):

Beam dump LHC year 2015



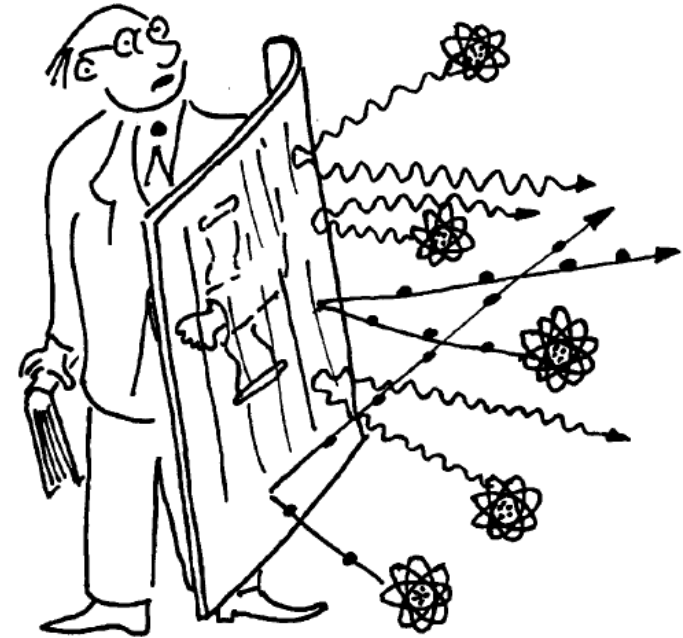
Sum: 442 dumps

B. Todd et al., CERNACC- 2014-0041

D. Wollmann et al., IPAC 2016, Busan, p. 4203 (2016)

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"Radiation Protection"

© by Claus Grupen

Cartoons from C. Grupen
Introduction to Radiation Protection,
Springer Verlag 2010

Radiological Quantities and Units for Humans

Basic quantities & units for personal safety:

➤ **Absorbed dose:** $D_{R,T} = \frac{1}{m} \int_{V_T} \frac{dE_R}{dV} \cdot dV$
 (physical quantity) $= \left[\frac{\text{J}}{\text{kg}} \right] = [\text{Gy}] = [100\text{rad}]$

for each radiation type **R** and each tissue **T**

➤ **Equivalent Dose:** $H_T = \sum_R w_R D_{R,T} = [\text{Sv}] = [100\text{rem}]$

with weight factor w_R for the radiation type **R**

➤ **Effective Dose:** $E = \sum_T w_T H_T = [\text{Sv}] = [100\text{rem}]$

with weight factor w_T for the absorption of each tissue **T**

whole-body irradiation $\Leftrightarrow \sum_T w_T = 1$

| Rad. type R | w_R |
|--|-------|
| γ all energies | 1 |
| e^- , e^+ , μ^\pm all energies | 1 |
| Protons $E > 2$ MeV | 5 |
| α , heavier nuclei | 20 |
| Neutrons: $E < 10$ keV | 5 |
| $10 \text{ keV} < E < 100 \text{ keV}$ | 10 |
| $100 \text{ keV} < E < 2 \text{ MeV}$ | 20 |
| $2 \text{ MeV} < E < 20 \text{ MeV}$ | 10 |
| $E > 20 \text{ MeV}$ | 5 |

Neutrons: Since 2007 smooth function

| Example: Organ or tissue | Sensi. | w_T |
|--|--------------|--------------|
| Gonads | High | 0.20 |
| Lung, stomach, colon, lens, hematopoietic & lymphatic system | Intermediate | 0,12 |
| Liver, esophagus, chest, skin, muscle, hart, bone surface | Low | 0.05 to 0.01 |

Shielding of Accelerators

Shielding of accelerator by rough rule of thumb:

Estimation of shielding by 10th-value λ_{10}

with $H(l) = H_0 10^{-l/\lambda_{10}}$

(disregarding any secondary particle transport)

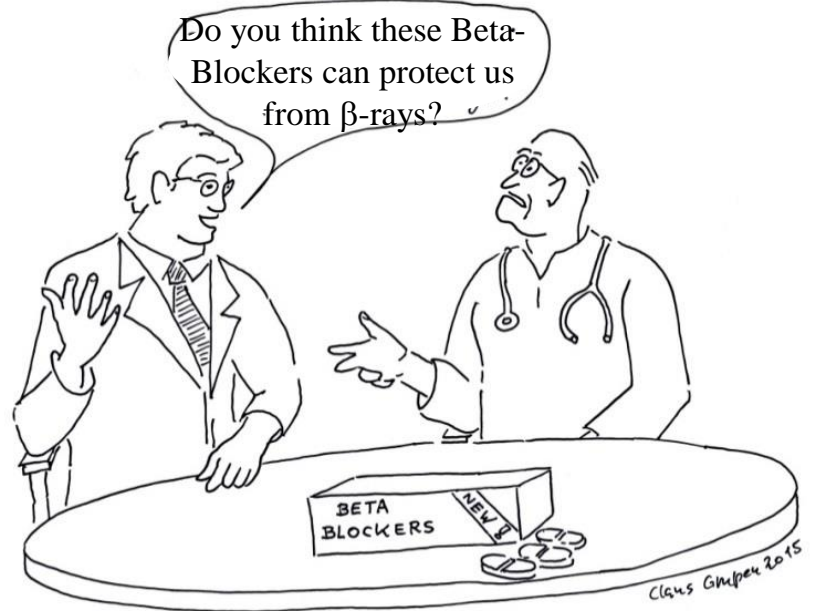
| Material | ρ [$\frac{g}{cm^3}$] | λ_{10} [cm] |
|----------------|-----------------------------|---------------------|
| Earth | 1.8 | 128 |
| Concrete | 2.4 | 100 |
| Heavy concrete | 3.2 | 80 |
| Iron | 7.4 | 41 |
| Lead | 11.3 | 39 |

Further rough rule of thumb:

- Protons, electrons & γ are att. by heavy materials
- Neutrons are scattered by hydrogen due to same mass
Concrete contains $\approx 10\%_{\text{weight}} \text{H}_2\text{O}$
- Nuclear reactions produces further particles

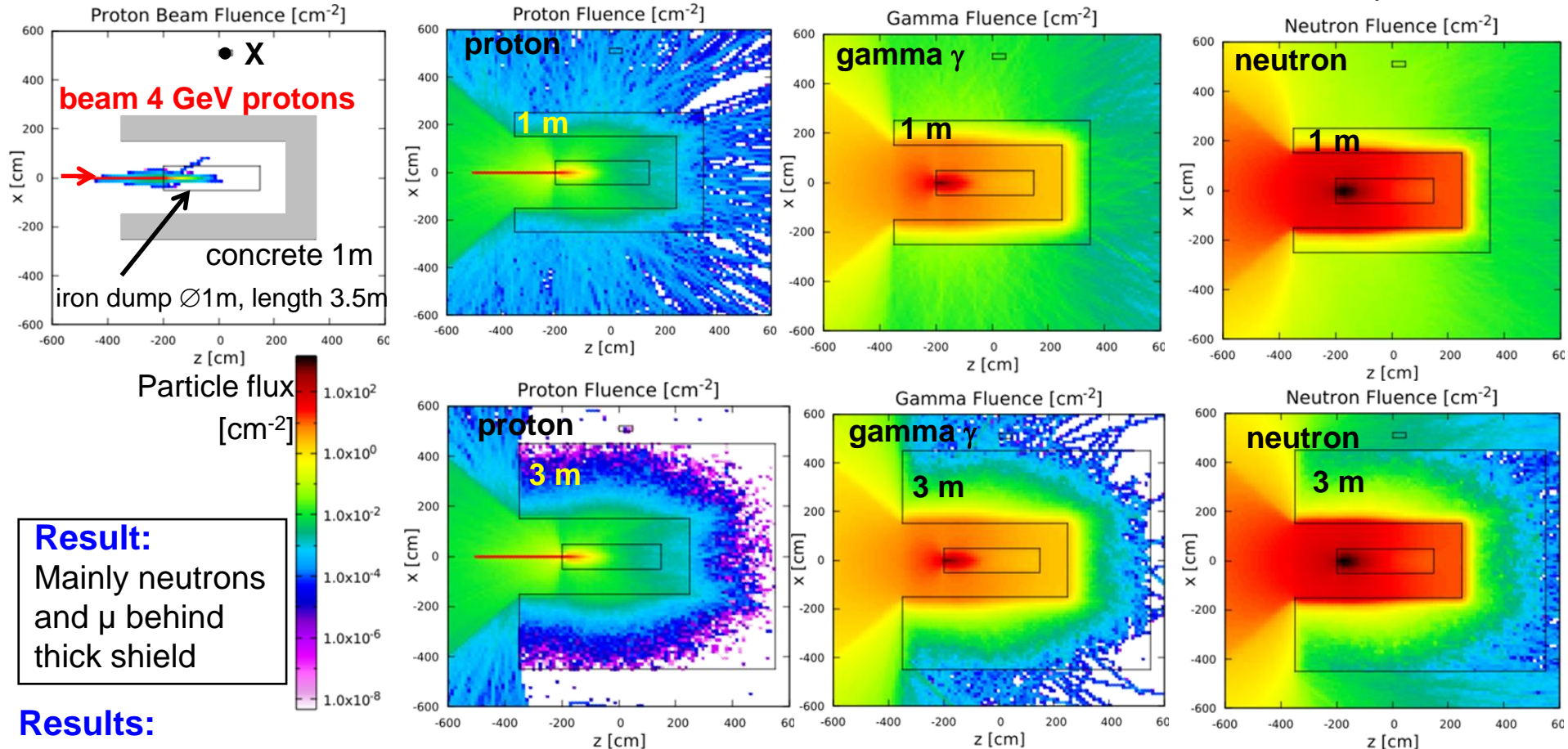


© by Clat



Simplified Model Shielding of Accelerators

Simplified FLUKA calculation: 4GeV protons, iron beam dump $\varnothing 1\text{m}$, $l=3.5\text{m}$, concrete 1 & 3 m, $5 \cdot 10^5$ particles
 Courtesy S. Udrea



Result:
 Mainly neutrons and μ behind thick shield

Results:

- Primary protons are stopped in dump
- **Protons** produced from neutrons, but partly stopped in the wall
- γ are from beam & neutrons in the wall $\approx 10^{-3}$ attenuation at X
- **Neutrons** produced, scattered at wall $\approx 10^{-3}$ atten. at X by distance & concrete
- 'Leakage' through opening
- **Protons** stopped in wall
- γ well shielded
- **Neutrons** at X $\approx 0.3\%$ of 1m.

Realistic Example for Shielding of Accelerators

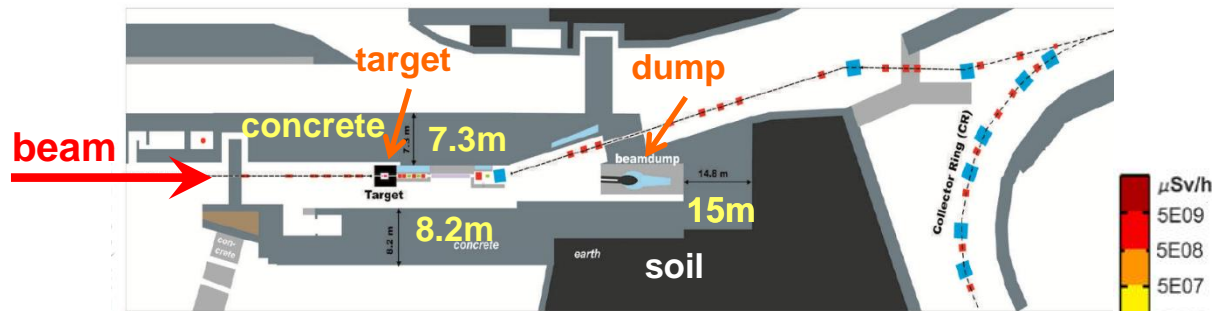
Example shielding of accelerator: Proton beam of 29 GeV for anti-proton production

Assumption $2.5 \cdot 10^{13}$ protons on 11cm long copper target

Shield: Iron (1.6 m downstream and 1 m transverse)

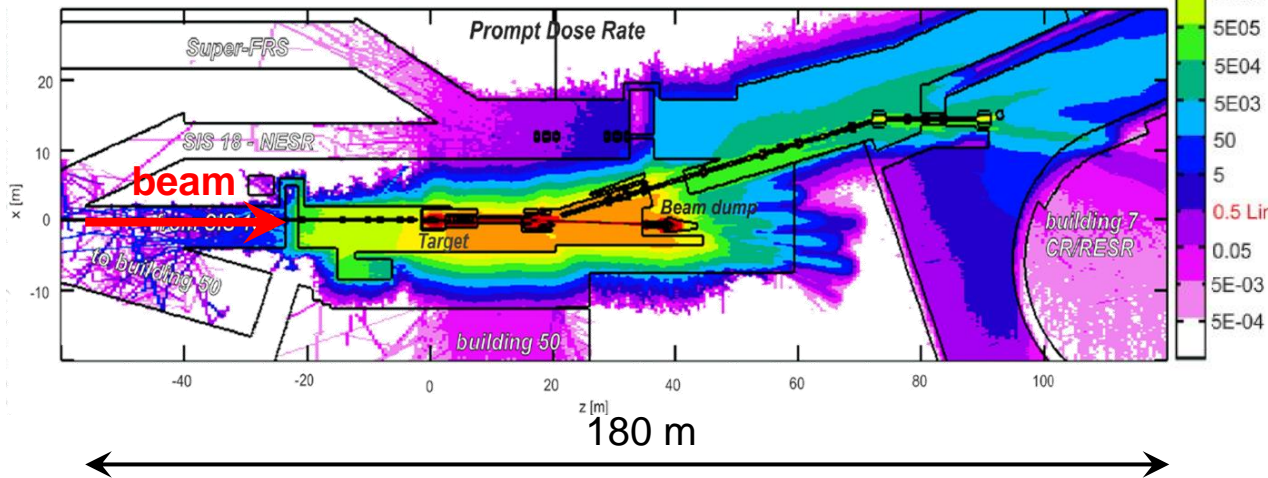
Concrete ≈ 8 m around beam pipe

Goal: Free access region outside i.e. equivalent dose rate $H/t < 0.5 \mu\text{Sv/h}$



Shielding calculations:

Required for safety procedure
 Numerical calculation required
 atomic, nuclear & particle
 physics models
 e.g. FLUKA, MARS, PHITS
 see lecture by Dan Faircloth



free access $H/t < 0.5 \mu\text{Sv/h}$

see lecture 'Secondary Beams and Targets' by K. Knie

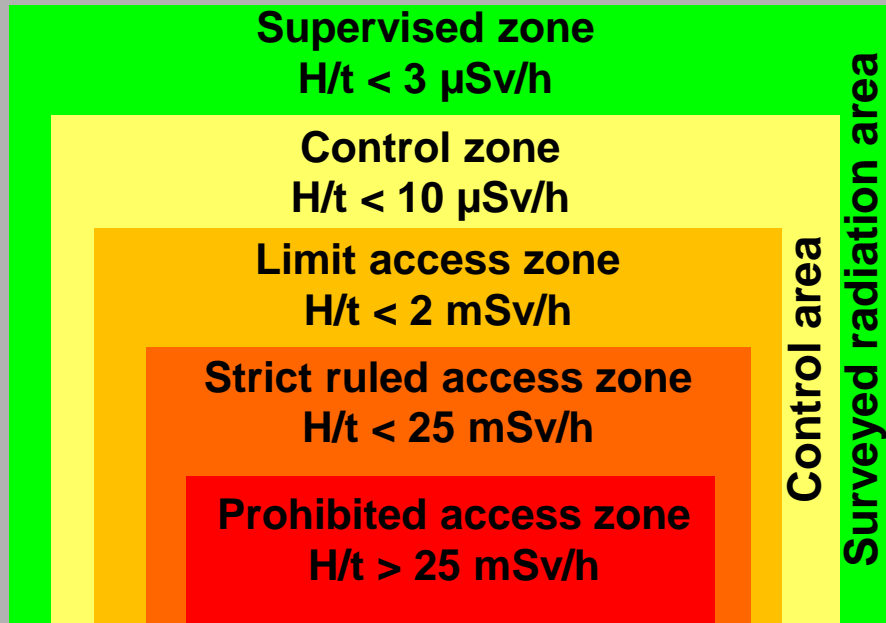
Categories of Locations & maximal Doses

Simplified categories of radiation areas:

For workers: Assumption 2000 h/a of access

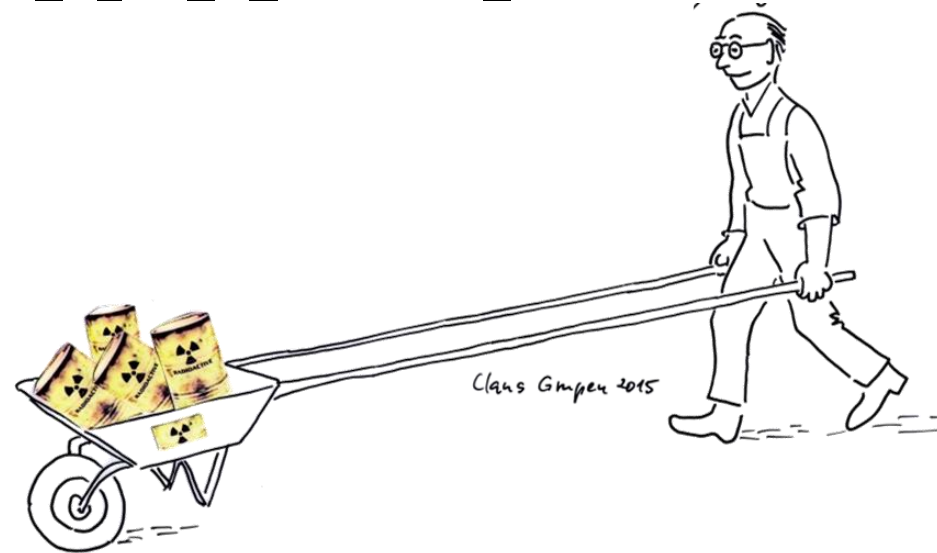
Non-designated, free access

$H/t < 1 \text{ mSv/a (full year)} = 0.5 \mu\text{Sv/h (for 2000 h)}$



ALARA principle:

As Low As Reasonable Achievable



Maximal dose for an radiation exposed worker:

Maximum dose for one year: 20 mSv/a

Maximum total life dose: 400 mSv

(Lethal dose for short term exposure: $\approx 4000 \text{ mSv}$)

Remark: Actual limits are given by national laws.

Categories of Locations & maximal Doses

Simplified categories of radiation areas:

For workers: Assumption 2000 h/a of access

Moderated prop. tube for n
 $1 \text{ eV} < E_n < 20 \text{ MeV}$

Proportional tube for γ :
 $30 \text{ keV} < E_{ph} < 1.3 \text{ MeV}$

Non-designated, free access
 $H/t < 1 \text{ mSv/a (full year)} = 0.5 \mu\text{Sv/h (for 2000 h)}$

Supervised zone
 $H/t < 3 \mu\text{Sv/h}$

Control zone
 $H/t < 10 \mu\text{Sv/h}$

Limit access zone
 $H/t < 2 \text{ mSv/h}$

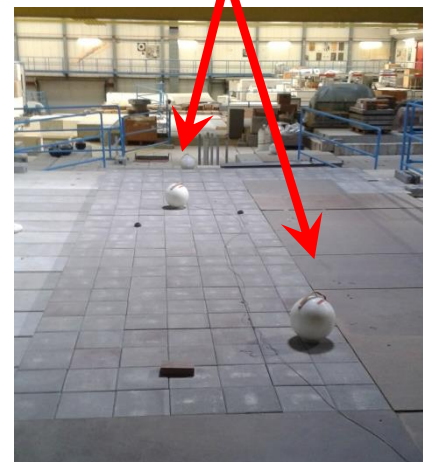
Strict ruled access zone
 $H/t < 25 \text{ mSv/h}$

Prohibited access zone
 $H/t > 25 \text{ mSv/h}$

Control area
 Surveyed radiation area



Moderated thermo-luminescence detector for passive n-detection



Maximal dose for an radiation exposed worker:

Maximum dose for one year: 20 mSv/a

Maximum total life dose: 400 mSv

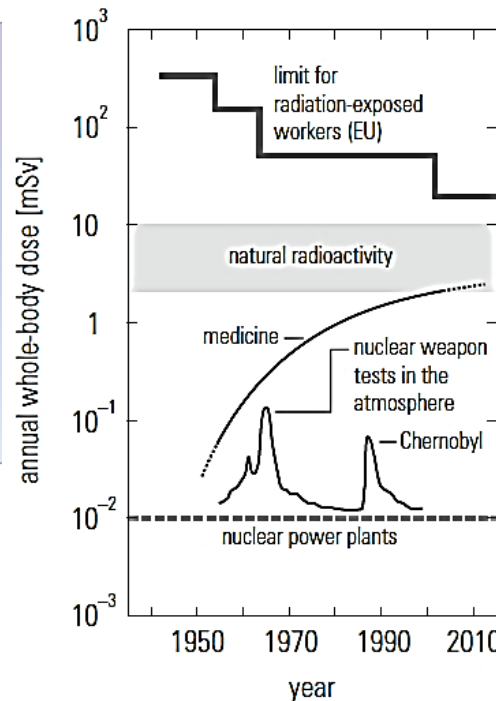
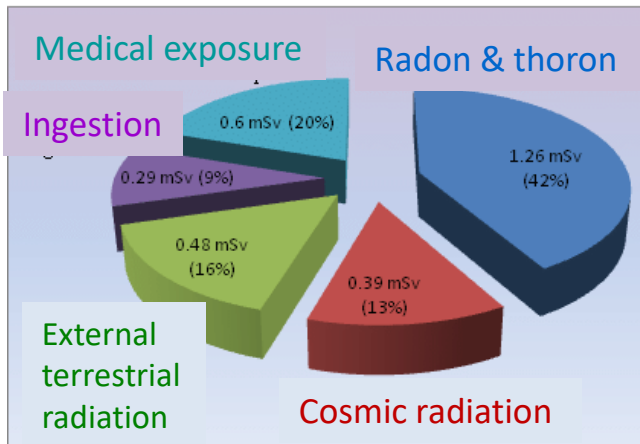
(Lethal dose for short term exposure: $\approx 4000 \text{ mSv}$)

Remark: Actual limits are given by national laws.

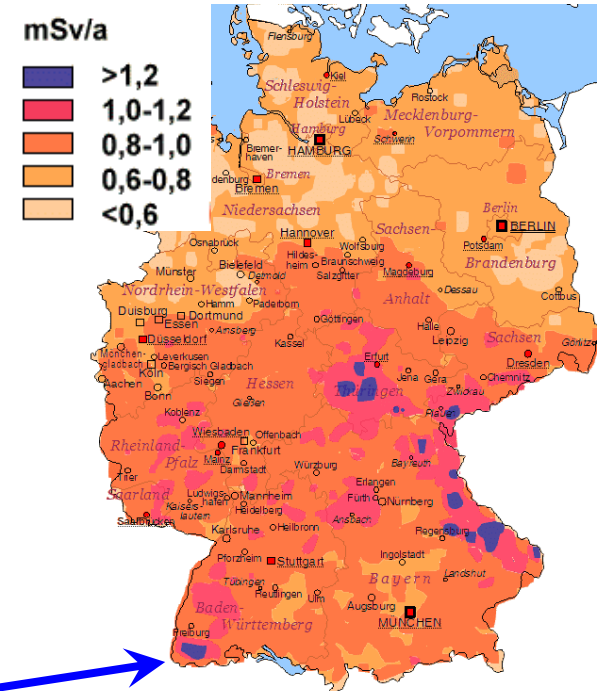
Natural Radiation Exposure

Example of radiation level:

- **Natural geological dose up to 10 mSv/a:**
Mainly due to noble gas radon from decay chain
without significant increase of diseases
- **Typical dose natural and artificial composition:**



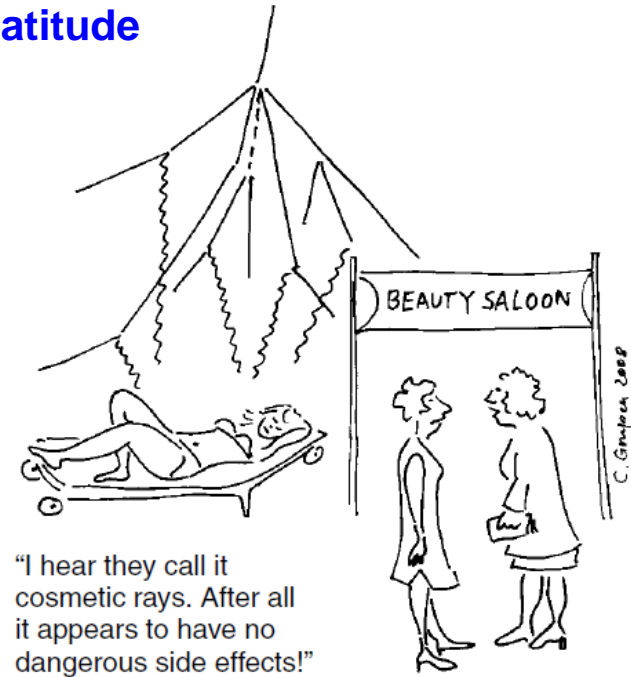
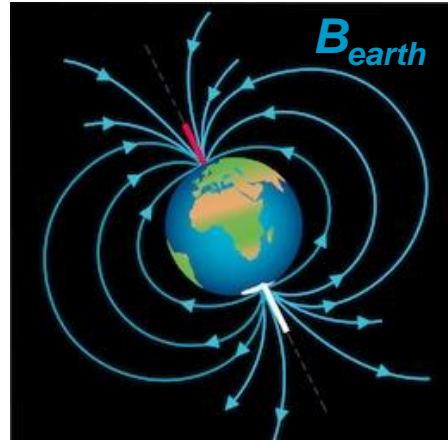
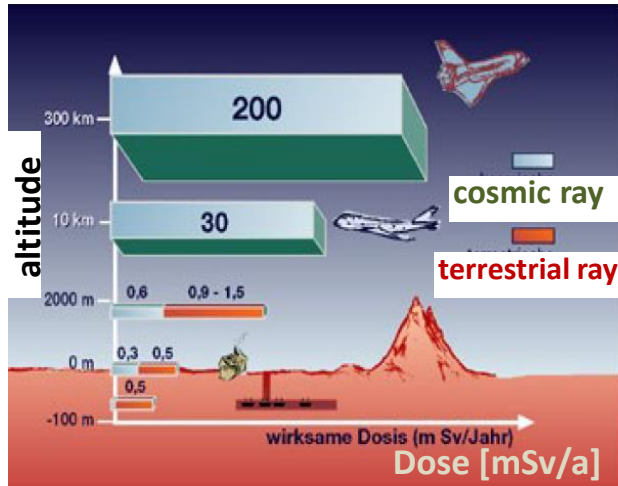
Natural dose in Germany:



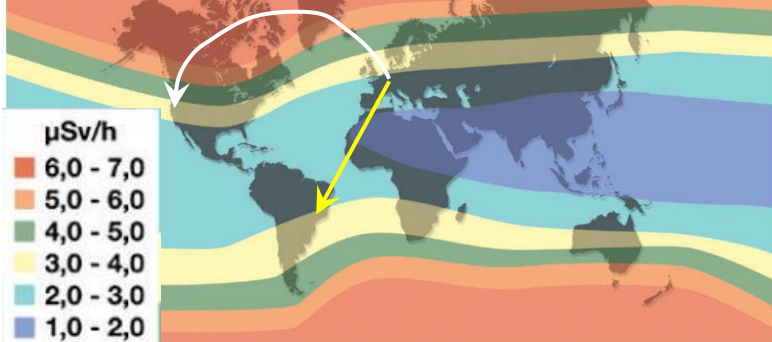
Source: German Bundesamt für Strahlenschutz
C. Grupen, Introduction to Radiation Protection

Avoidable, but wildly accepted Radiation Exposure

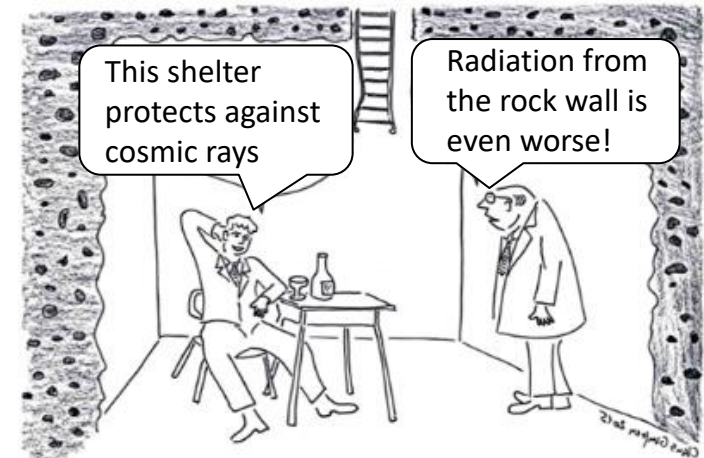
Cosmic ray based radiation effects depend on altitude and latitude



Radiation at 11 km altitude, year 2013



Zonen mit unterschiedlicher Höhenstrahlung [11 km Höhe, Ende 2013, Mikrosievert pro Stunde]



| Departure | Arrival | Duration | Dose |
|-----------|----------------|----------|--------------|
| Frankfurt | San Francisco | 11.5 h | 45 - 110 µSv |
| Frankfurt | Rio de Janeiro | 11.5 h | 17 - 28 µSv |

Source: German Bundesamt für Strahlenschutz

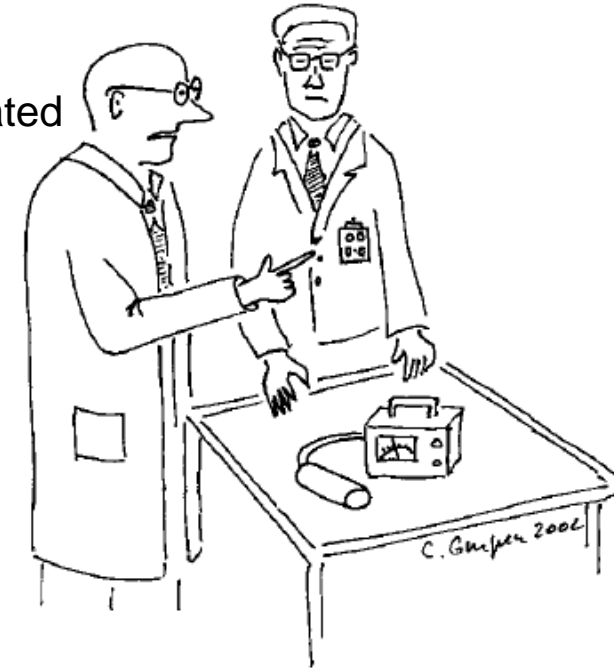
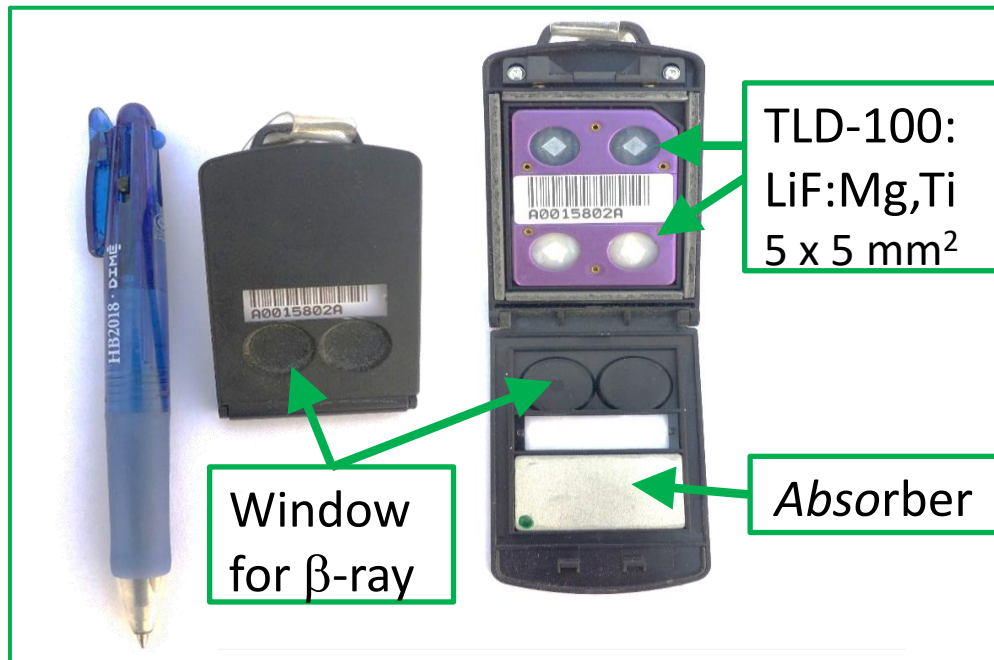
Passive Film Badge Dosimeter and TLD

For personal safety a dosimeter should be worn!

Thermo-luminescence dosimeter TLD:

- Crystal e.g. LiF is excited by radiation and emit light when heated
- neutron sensitive via ${}^6\text{Li}(n,\alpha)\text{T}$

Sensitivity for β & γ : 0.1 mSv to 10 Sv



“And these badges are supposed to protect us effectively from radiation?”

© by Claus Grupen

Advantage: Can be archived

Disadvantage: Limited sensitivity, **no** online display

Film badge: X-ray sensitive films photons (typ. 5keV... 9MeV) & β^\pm (typ. > 0.3MeV)
Sensitivity for β & γ : 0.1 mSv to 5 Sv

Active personal Dosimeter

Active dosimeters for online display

Dose measurement with alarm function,
has to be worn when entering a protected area

Ionization chambers or proportional chambers:

Alternative: PIN-diode solid state detector

Photons: typ. 10 keV... 10 MeV

β^\pm : 0.25 1.5 MeV

Sensitivity for β & γ : 0.05 $\mu\text{Sv/h}$ to 1 Sv/h

(TLD sensitivity: 100 μSv to 5 Sv, flight above pole: 45...110 μSv)

'Pocket meter' for γ -rays:

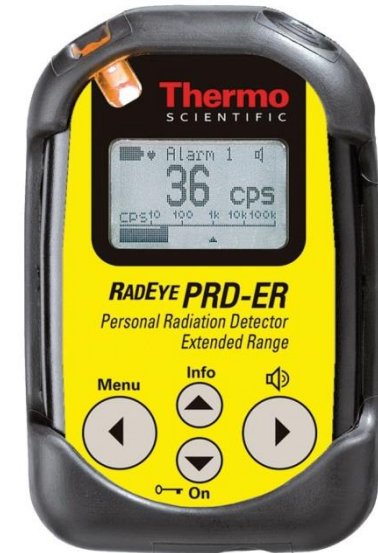
Scintillator NaI(Tl) + photo-multiplier for γ detection
photons (typ. 60 keV... 1.5 MeV)

Sensitivity for γ : 0.01 $\mu\text{Sv/h}$ to 100 mSv/h

Older versions: Proportional tube

Advantage: Alarm functionality, sensitive
can be archived with some efforts

Disadvantage: Expensive



Summary

- **Many accelerators are built to produce radiation, some risk remains**
- **Accelerator components must be protected from overheating ('atomic physics')**
 - e.g. super-conducting magnet & cavities
 - Particles' energy loss must be limited and/or steered to dedicated locations
 - Passive protection by collimators for protection or localizing
 - Active Machine Protection System based on Beam Loss Monitors
- **Accelerator components must be protected from activation ('nuclear physics')**
 - Losses must be limited to certain locations e.g. collimators & beam dump
 - '1 W/m criterion' to limit activation for hand-on maintenance
- **Shield of the accelerator required**
 - p, ion & γ best shield by high density material, but care for nuclear reactions
 - e⁻ shield for light material (lower Bremsstrahlung)
 - n light material preferred
- **ALARA principle: Unnecessary radiation exposure to people should be avoided**

Thank you for your attention!

- R. Schmidt (Ed.), *Beam Loss and Accelerator Protection*, Proc. Joint International Accelerator School CERN-2016-002
- US Particle Accelerator School – Beam Loss & Machine Protection, January 2017
<http://uspas.fnal.gov/materials/17UCDavis/davis-machineprotection.shtml>
- D. Kiselev , *Activation and radiation damage in the environment of hadron accelerators & D. Forkel-Wirth et al., Radiation protection at CERN in R. Bailey (Ed.) Proc. CAS CERN-2013-001*
- A. Zhukov, *BLMs: Physics, Simulation and Application in Accelerator*, Proc. BIW 2010, www.jacow.org
- C. Grupen, *Introduction to Radiation Protection*, Springer Verlag 2010
- Proceedings of several CERN Acc. Schools (introduction & advanced level, special topics).
- Contributions to conferences, in particular to IPAC & IBIC.

Backup slides

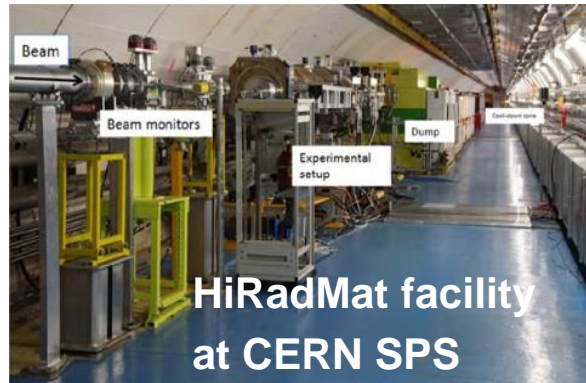
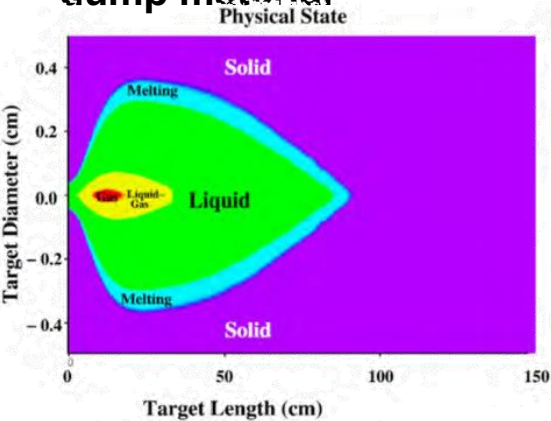
Energy Loss and Heating: Experiment

Verification of material interaction by 440 GeV protons:

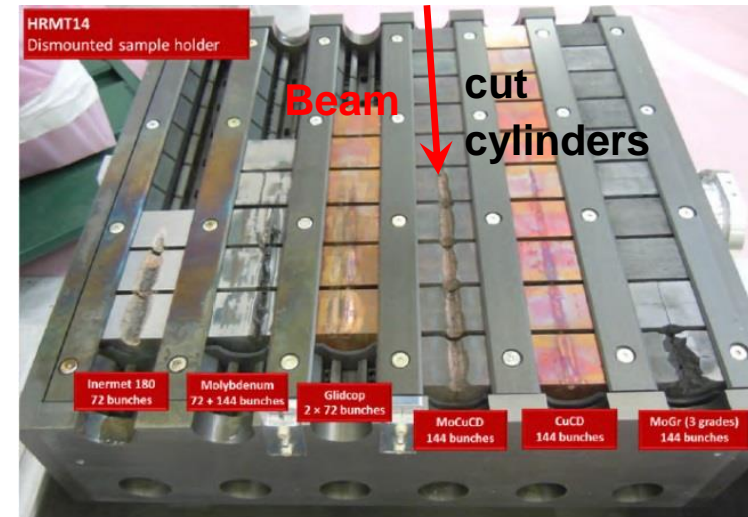
Destruction of material due to temperature rise

- melting, sublimation plasma formation
 - mechanical stress
- ⇒ verification of simulation

⇒ finding proper dump material



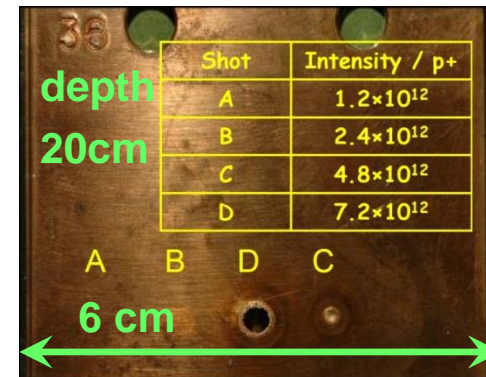
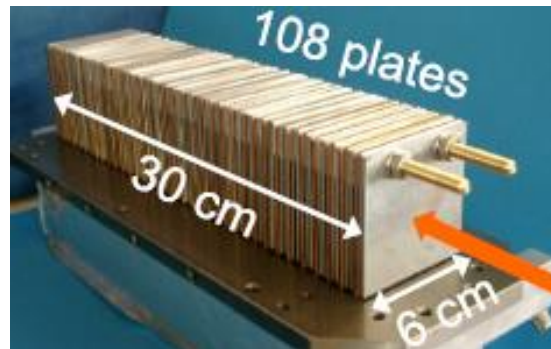
Beam: 440 GeV $\approx 10^{13}$ protons,
 $\sigma_x = \sigma_y \approx 2$ mm within $t = 50 \mu s$
 $\Rightarrow E_{tot} \approx 1$ MJ



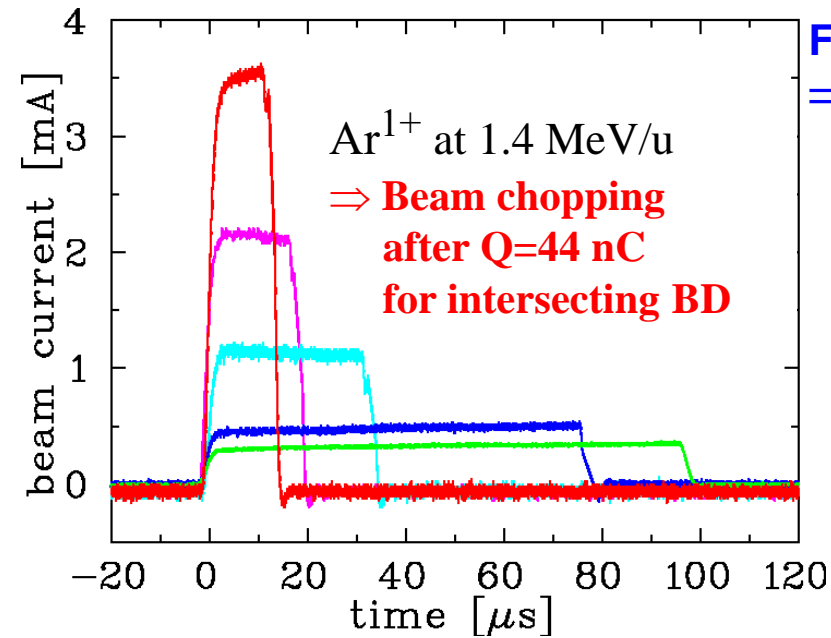
A. Bertarelli, JAS CERN-2016-002.

Experiment with 450 GeV protons:

V. Kain et al., PAC'05, 1607 (2005)



Dynamic Machine Protection by Transmission Measurement

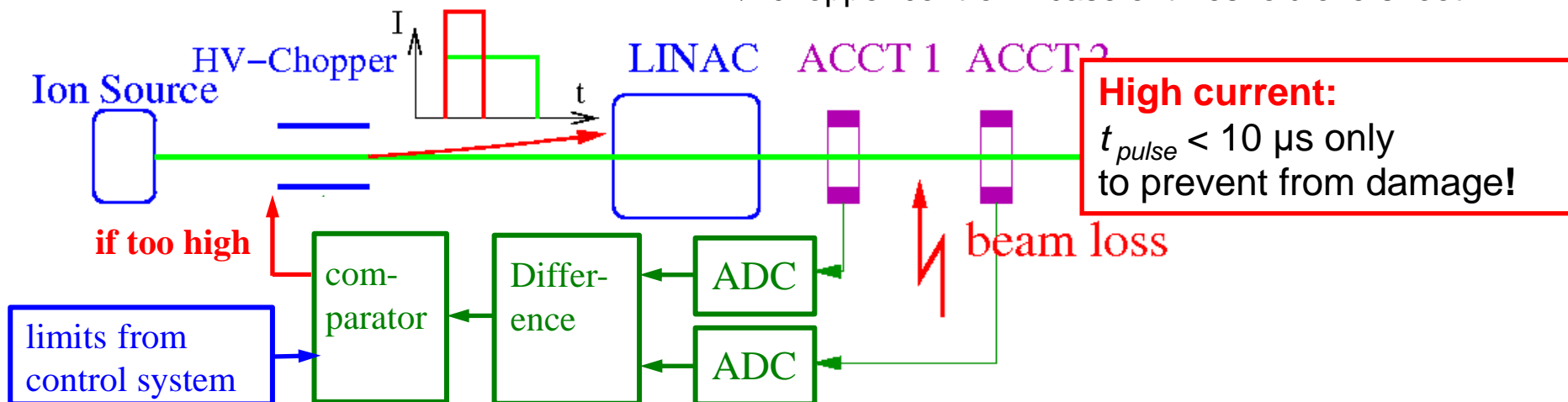


For $E > 50$ MeV protons: nuclear σ_{nucl} quite low
 \Rightarrow machine protection by **active transmission control**

Determination of maximal loss between consecutive transformers by 'differential current measurement'
 \rightarrow **dynamic** beam interruption in case of software-given threshold overshoot.

FPGA-electronics:

- \rightarrow ADC digitalization
- \rightarrow calculation of difference
- \rightarrow digital comparator
- \rightarrow chopper control in case of threshold overshoot



H. Reeg (GSI) et al., Proc. EPAC'06

Neutron Energy Spectrum

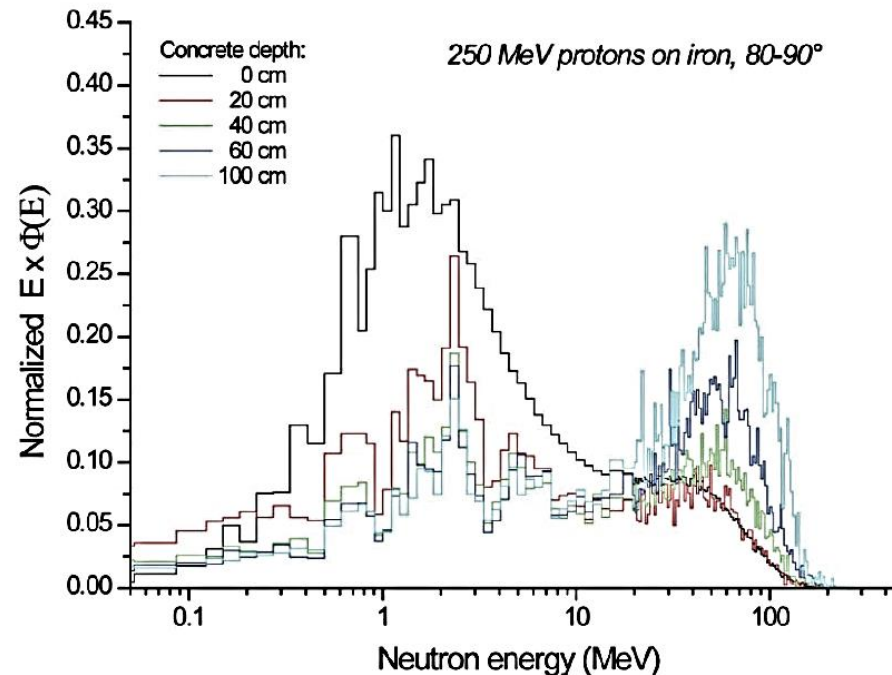
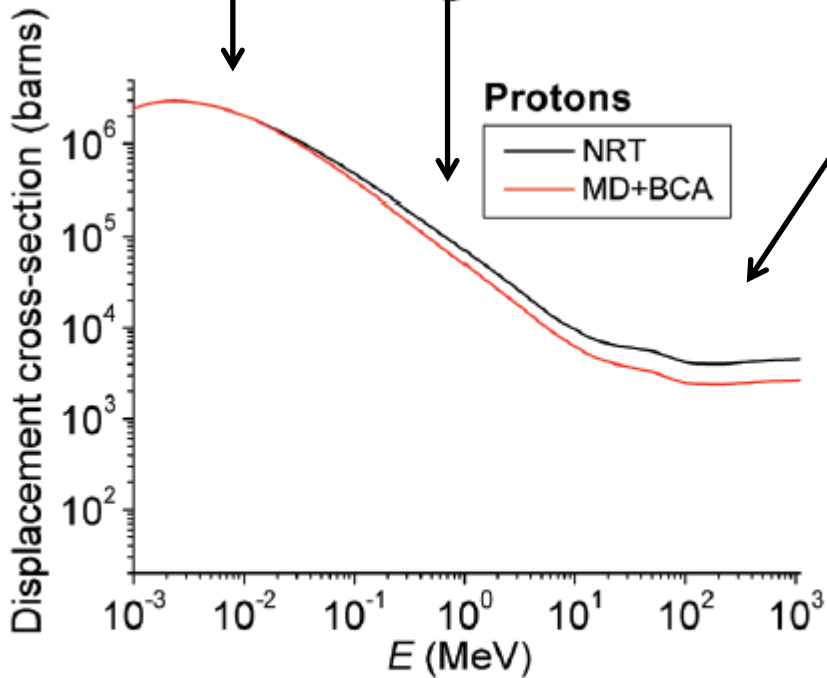


Fig. 6: Neutron energy distributions $E\Phi(E)$ in the transverse direction generated by 250 MeV protons impinging on an iron target thicker than the proton range. The distributions are for source neutrons and behind concrete shields of thicknesses ranging from 20 cm to 1 m. The distributions have been normalized to unit area in order to show better the change in the shape of the spectrum with increasing shield thickness.

Radiation Damage Displacements of Atoms

Low energy protons: Nuclear stopping (collision of protons with target nucleus results in recoil energy above binding energy to stopping)

Electronic stopping range



For $E_{kin} > 100$ MeV nearly equal cross section

Large capture cross section results in recoil energy

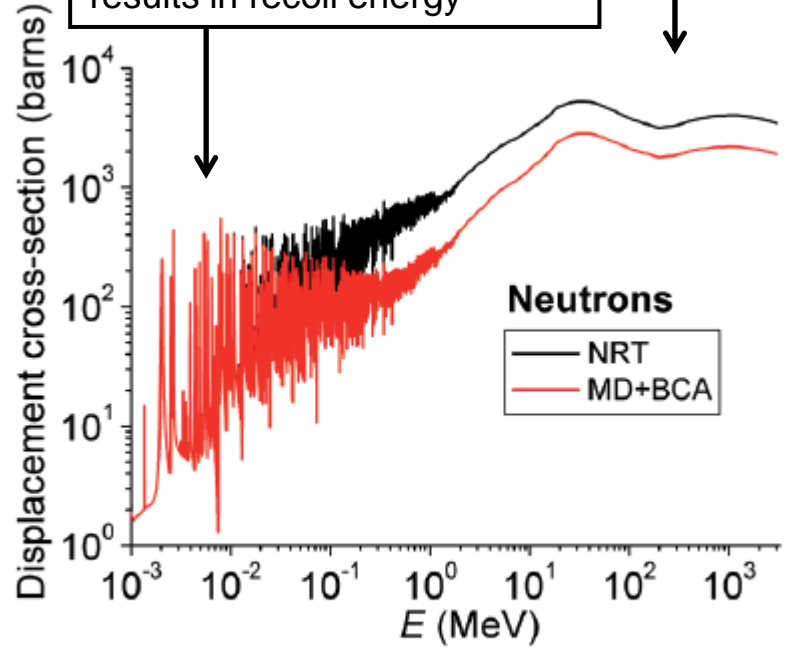


Fig. 12: Displacement cross-sections of protons (left) and neutrons (right) in copper obtained by two different approaches (see legend).

D. Kiselev, CAS 2011, CERN-2013-001

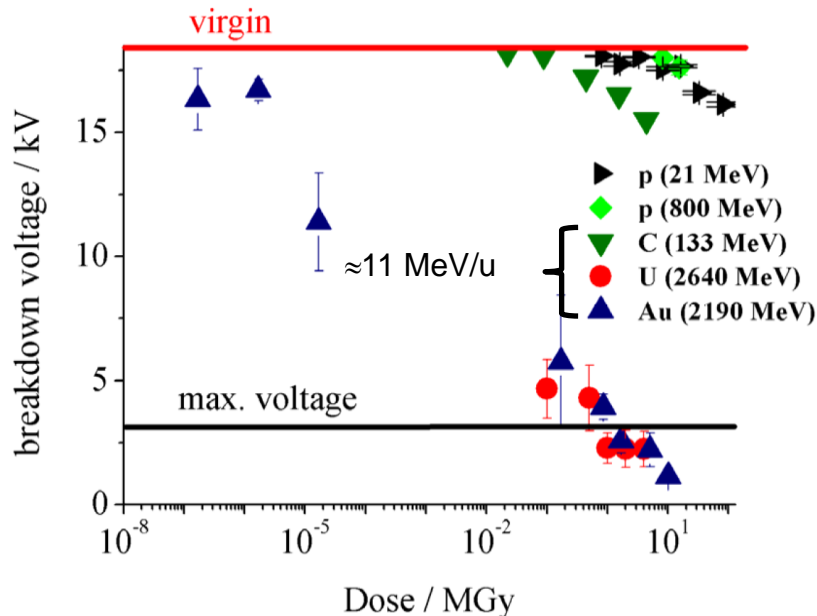
Radiation Damage of organic Materials

Radiation damage in plastic by ionizing radiation:

- Break of chemical bonds and displacement of atoms
- Microscopic defects in the chemical bonds
- Displacement of atoms in the structural material

Example: Kapton foil of 125 μm thickness

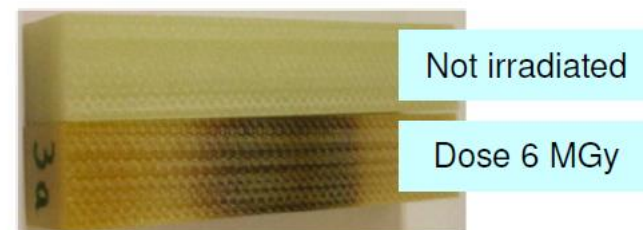
Direct irradiation by ion beam's energy loss dE/dx increases for heavy ions



T. Seidl et al, HB 2010

Rough estimation of maximal dose

| Material | Dose [Gy] |
|--------------------|----------------|
| Teflon (PTEE) | 10^3 |
| Mylar | $5 \cdot 10^4$ |
| Cable insulation | $5 \cdot 10^4$ |
| Magnet coil insul. | 10^6 |
| Kapton (Polyamide) | 10^7 |



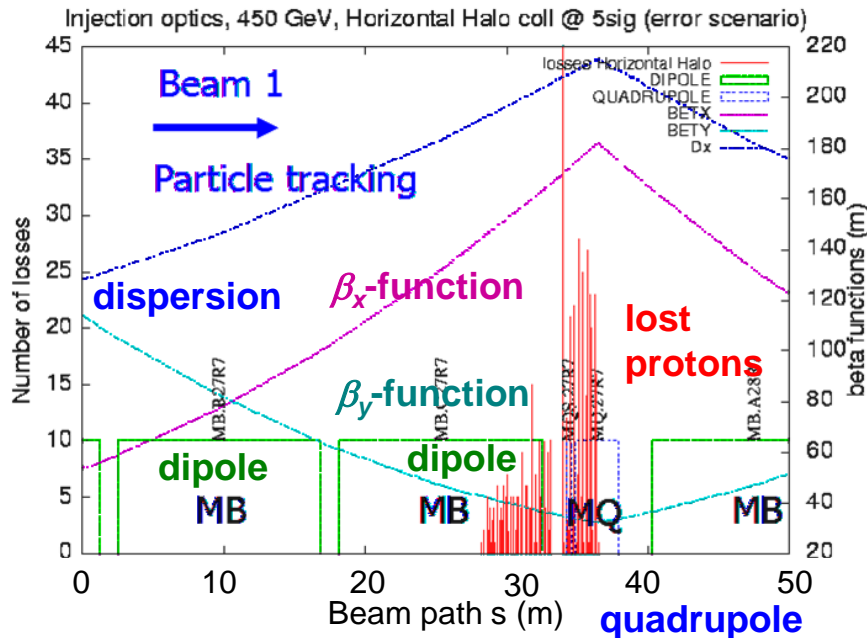
Placement of Beam Loss Monitors

Secondary particles and shower produces are emitted within a forward cone (in rest-frame isotropically but due to Lorentz-transformation forward in lab-frame).

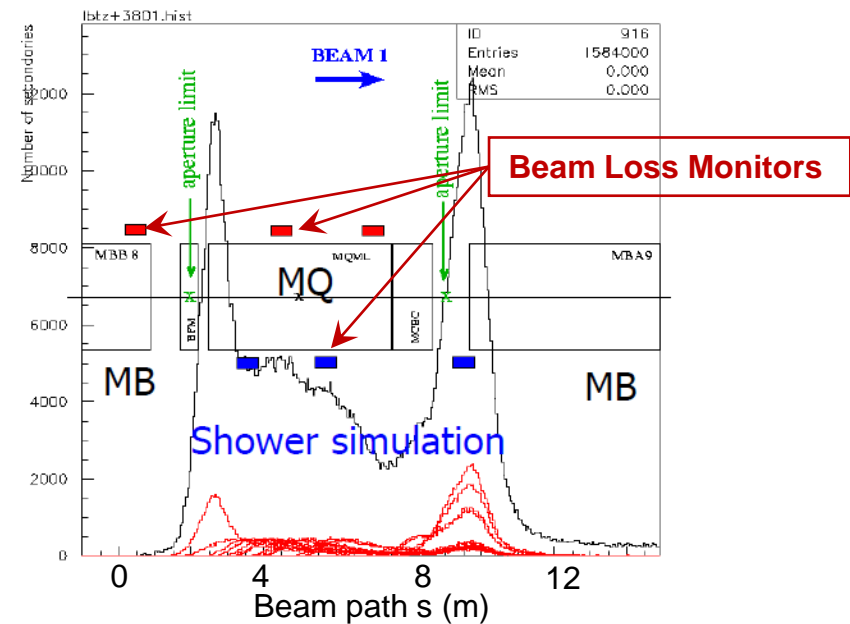
Position of detector at quadruples due to maximal beam size.

High energy particles leads to a shower in forward direction → Monte-Carlo simulation.

Example: Simulation of lost protons at LHC at 450 GeV of lost protons:
→ at focusing quad. D & β_x maximum



Example: Simulation of number of shower particles



Secondary Particle Production for Electron Beams

Processes for interaction of electrons

For $E_{kin} < 10$ MeV:

Mainly electronic stopping \Rightarrow X-rays, slow e^-

For $E_{kin} > 10$ MeV:

Bremsstrahlungs- γ , forward peaked $E_\gamma = 5-50$ MeV

$\Rightarrow \gamma \rightarrow e^+ + e^-$ or $\mu^\pm \dots \rightarrow$ electro-mag. showers

\Rightarrow Excitation of giant resonances $E_{res} \approx 10-30$ MeV

via (γ, n) , (γ, p) or (γ, np) with $\sigma_{giant} \approx \frac{1}{10} \sigma_{geo}$

\rightarrow Fast neutrons emitted (by nuclear interaction)

\rightarrow Neutrons: Long ranges in matter

Photo-Pion reaction: $d(\gamma, \pi^0)pn$ or $d(\gamma, \pi^-)pp$

\Rightarrow **activation at electron accelerators**

R.H. Thomas, in Handbook on Acc. Phys. & Eng.

Processes for interaction of electrons

For $E_{kin} < 10$ MeV:

Mainly elec

For $E_{kin} > 10$ MeV:

Bremsstrah

$\Rightarrow \gamma \rightarrow e^+ +$

\Rightarrow Excitatio

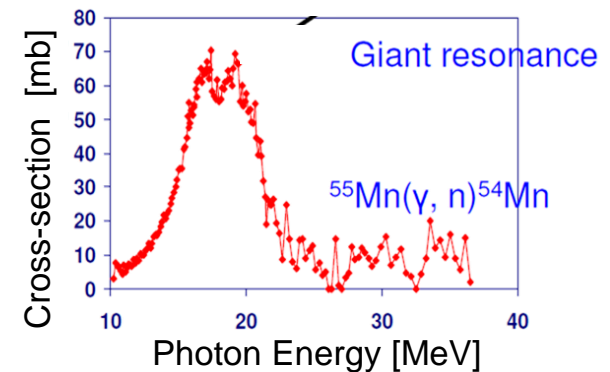
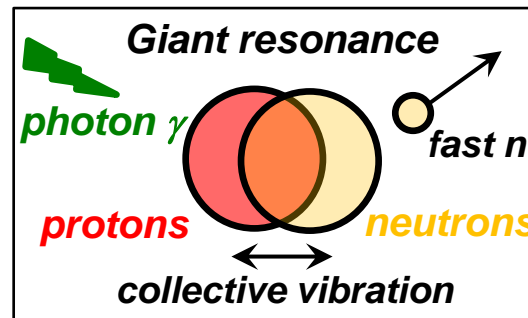
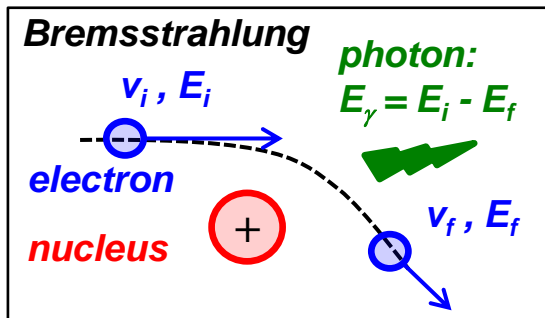
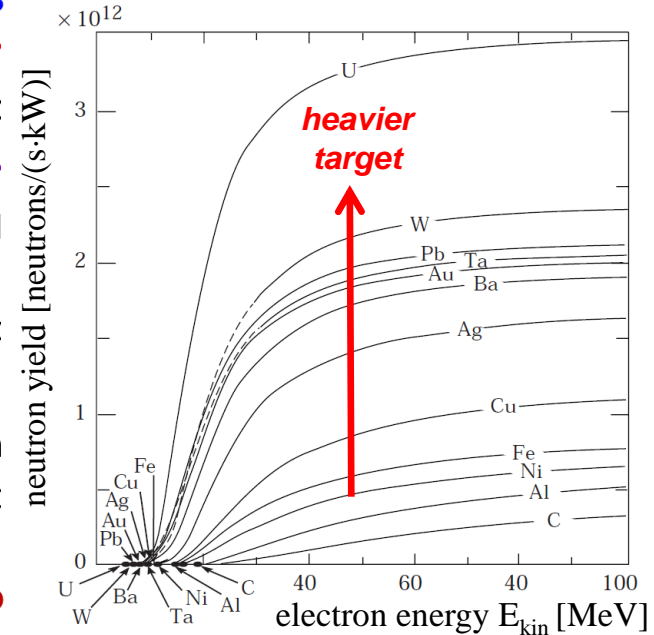
via (γ, n) ,

\rightarrow Fast n

\rightarrow Neutrc

Photo-Pion

\Rightarrow **activatio**



Collimation at LINACs

Halo development at LINACs caused by:

- Higher order magnet fields (e.g. aberration)
- Transverse mis-match
- Off-momentum particles due to wrong acceleration
- Space charge forces

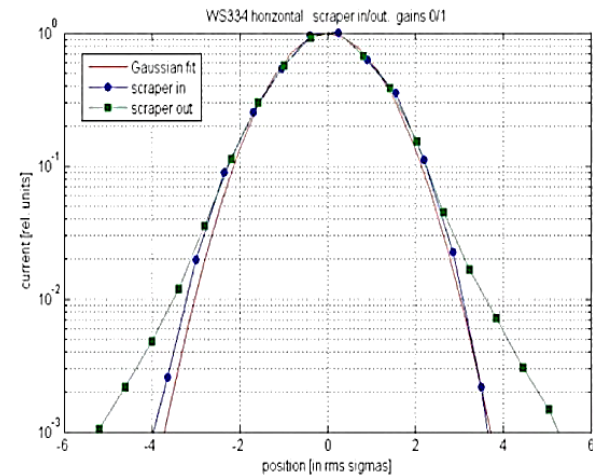
Goal: Halo cutting at low energy to prevent for activation

Collimators:

Cut the beam tail in space

$\mu = 90^\circ$ or $\mu = 45^\circ$ betatron phase to cut angle

⇒ at least two locations required



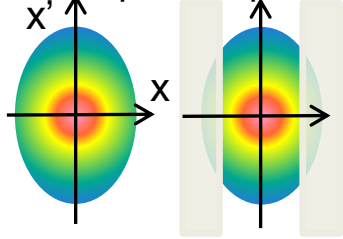
Example: SNS LINAC

Scraping at 3 MeV

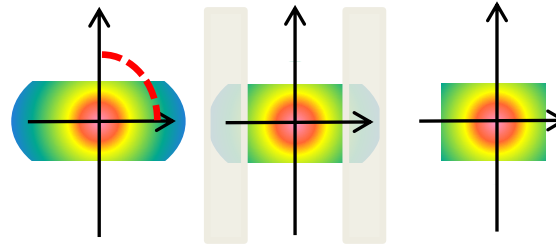
profile measurement at 40 MeV

M. Plum, CERN-2016-002

horizontal phase space



Betatron
phase
 $\mu = 90^\circ$



beam path s

i.e. phase space distribution is not completely cut

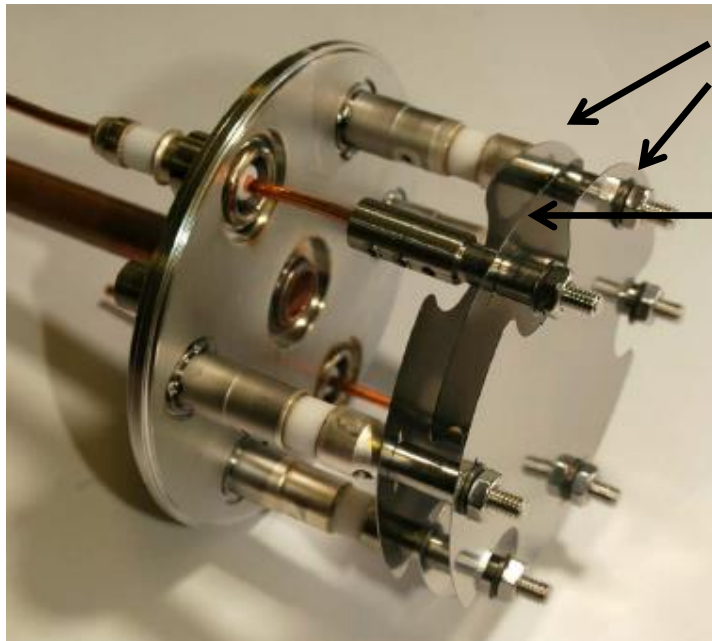
Secondary Electron Monitor as BLM

Ionizing radiation liberates secondary electrons from a surface.

Working principle:

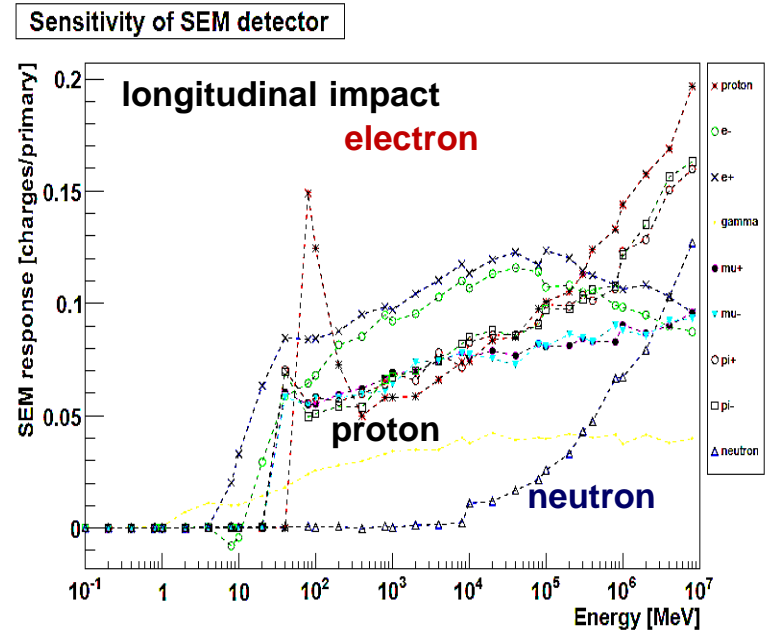
- Three plates mounted in a vacuum vessel (passively NEG pumped)
- Outer electrodes: biased by $U \approx +1$ kV
- Inner electrode: connected for current measurement (here current-frequency converter)

→ **small and cheap detector, very insensitive.**



HV electrodes

Electrode for measured current



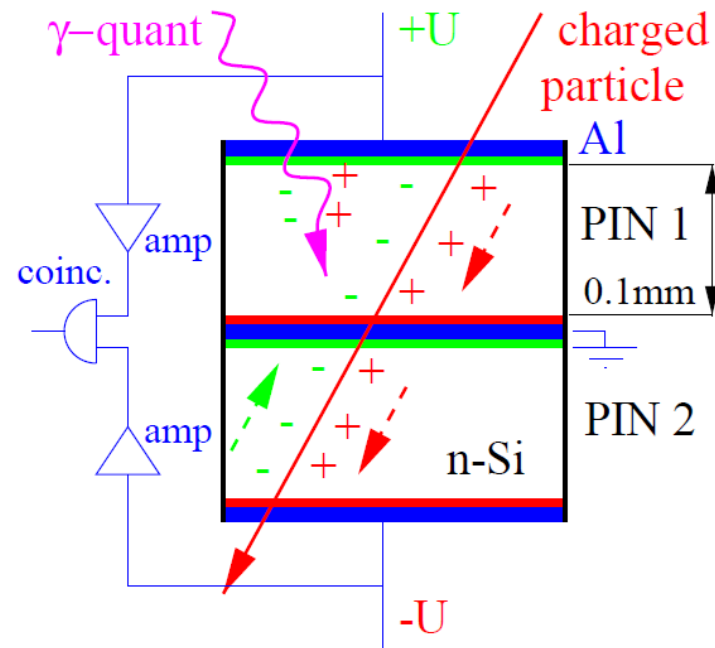
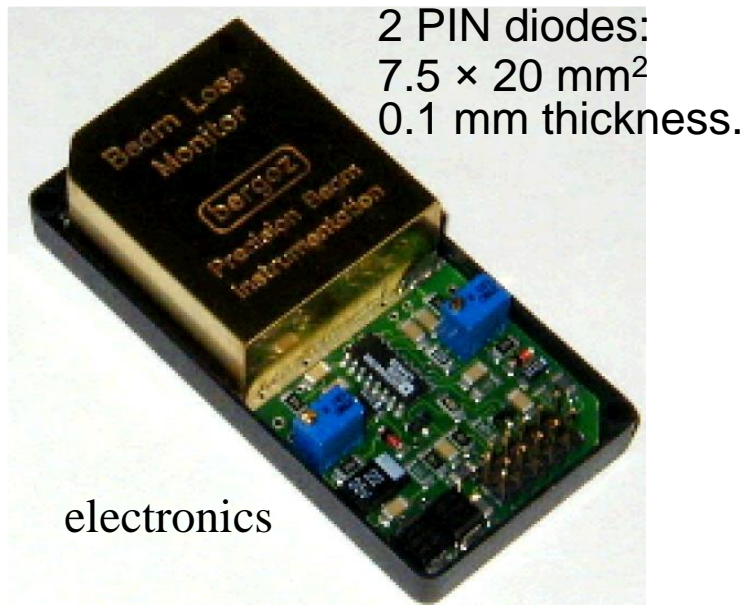
B. Dehning et al., PAC 2007

PIN-Diode (Solid State Detector) as BLM

Solid-state detector: Detection of charged particles.

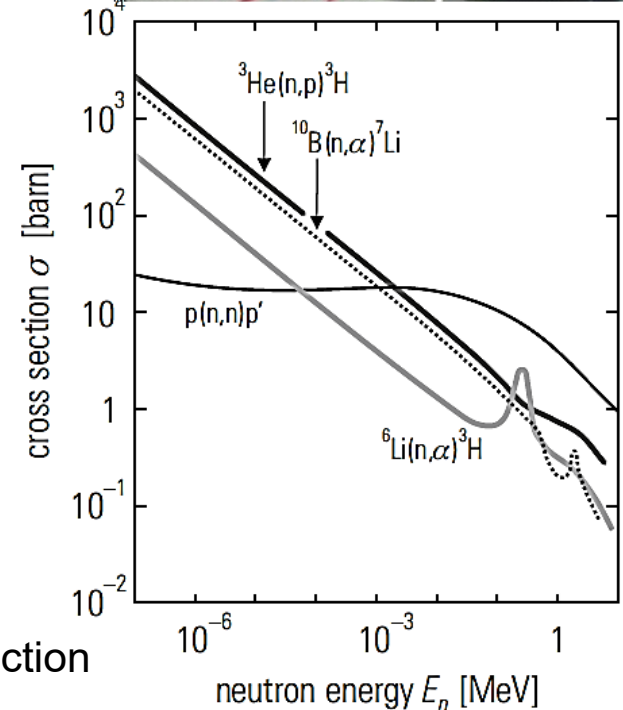
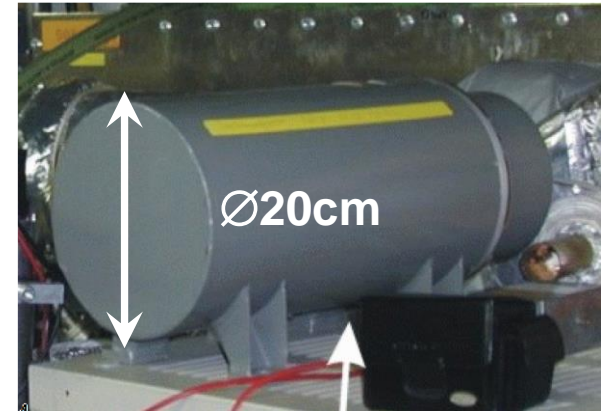
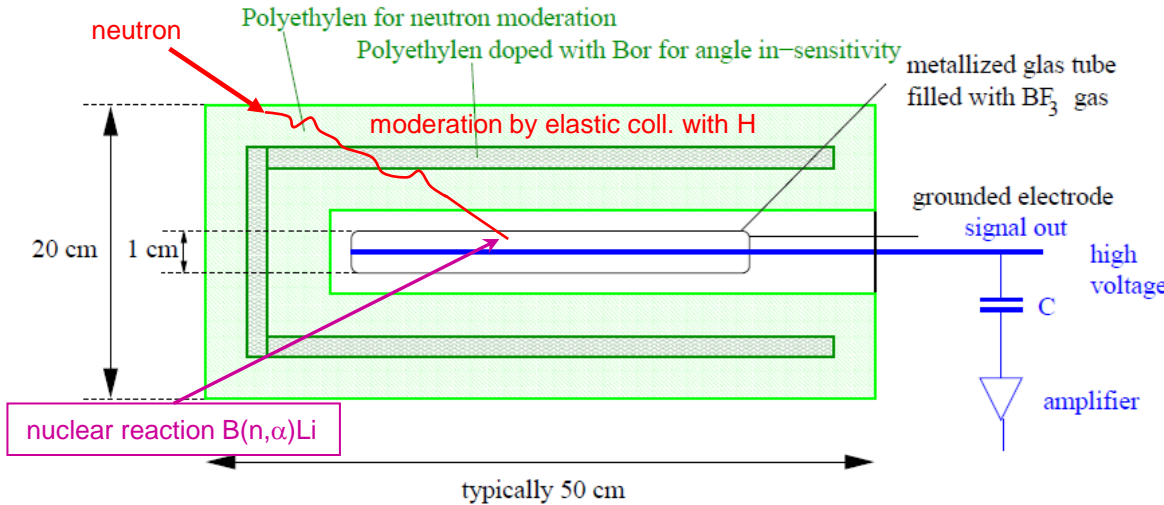
Working principle

- About 10^4 e⁻-hole pairs are created by a Minimum Ionizing Particle (MIP).
 - A coincidence of the two PIN reduces the background due to low energy photons.
 - A counting module is used with threshold value comparator for alarming.
- **small and cheap detector.**



BF₃ Proportional Tubes as BLM and for personal Protection

Detection of neutrons **only** with a 'REM-counter':



Physical processes of signal generation:

1. Slow down of fast neutrons by elastic collisions with p
2. Nuclear reaction inside BF₃ gas in tube:



3. Electronic stopping of ⁷Li and α leads to signal.

Remark: 'REM-counters' are frequently used for neutron detection outside of the concrete shield & in nuclear power plants