

## **Photoproduction of X(3872) in the near-threshold region**

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Xiong-Hui Cao, Meng-Lin Du, FKG, <u>arXiv:2401.16112</u>



### $J/\psi$ -nucleon scattering and photoproduction



- OZI suppressed scattering
  - $\square$  Relatively suppressed by  $O(1/N_c)$
  - **\Box** General mechanisms (take  $J/\psi N$  as an example)
    - ➢ Gluon exchanges







≻ Coupled-channel:  $J/\psi N - \Lambda_c \overline{D}^{(*)} / \Sigma_c^{(*)} \overline{D}^{(*)} - J/\psi N$ 



Importance of coupled-channel mechanism in evading OZI suppression in mesonic sector: H. Lipkin, B.-S. Zou, PRD 53 (1996) 6693

#### $J/\psi$ -nucleon scattering and photoproduction

P

 $\bullet$  The  $J/\psi$  photoproduction probes the gluonic contribution to the nucleon mass

□ if the mechanism of gluon exchanges is dominant

 $\Box$  if the  $J/\psi$  photoproduction can be modeled by vector-meson dominance

D. Kharzeev, H. Satz, A. Syamtomov, G. Zinovjev, EPJC 9 (1999) 459



Scattering length from VMD and photoproduction: 3 - 25 am L. Pentchev, I. Strakovsky, EPJA 57 (2021) 56  $21.3 \pm 8.2$  am GlueX, PRC 108 (2023) 025201

> Nucleon gluonic gravitational form factor, mass radius measured with this assumption

Duran et al. [J/ψ-007], Nature 615 (2023) 813

• For the search for hidden-charm  $P_c$  pentaquarks

 $J/\psi$  in the VMD model would be highly off-shell, but the scattering length and cross section are defined for real  $J/\psi$ 

#### $J/\psi$ photoproduction





0.0

8500

9000

 $E_V$  [MeV]

9500

10000

Data: GlueX, PRL 123 (2019) 072001

#### $J/\psi$ photoproduction

• Coupled-channel mechanism

Hint of threshold cusps in the latest data, but still not conclusive



GlueX, PRC 108 (2023) 025201

See also JPAC analysis in D. Winney et al., PRD 108 (2023) 054018

#### X(3872) photoproduction

• Searched for at COMPASS, but not seen

 $\Box$  Evidence of  $\tilde{X}(3872)$  in  $\gamma^*N \to X^0\pi^{\pm}N'$  with 4.1 $\sigma$ 

 $\sqrt{s_{\gamma N}} \in [8, 18] \text{ GeV}$ 

COMPASS, PLB783(2018)334

 $M_{\tilde{x}} = (3860.4 \pm 10.0) \text{ MeV}$ 



**C**ross sections at  $\sqrt{s_{\gamma N}} = 13.7 \text{ GeV}$ :  $\sigma(\gamma N \to \tilde{X}\pi N') \times \mathcal{B}(\tilde{X} \to J/\psi \pi^+\pi^-) = (71 \pm 28 \pm 39) \text{ pb}$  $\sigma(\gamma N \to X(3872)N') \times \mathcal{B}(X(3872) \to J/\psi \pi^+\pi^-) < 2.9 \text{ pb} (\text{CL} = 90\%)$ 



#### X(3872) photoproduction with VMD

• Existing estimate for the X(3872) photoproduction

 $\Box$  Vector-meson-dominance (VMD) model:  $\gamma^* \rightarrow J/\psi$ 

**Low energy:**  $\rho$ ,  $\omega$  exchanges,  $\mathcal{O}(10 \text{ nb})$ 

□ High energy: Regge exchange

JPAC, PRD102(2020)114010





• Coupled-channel mechanism for the X(3872) photoproduction in the near-threshold region

D

X(3872)

D

 $\boldsymbol{v}$ 

 $\bar{D}^*$ 

 $\Sigma_{c}(2800)$ 

- $\square$  X(3872) couples strongly to  $D\overline{D}^*$
- □ Nearby open-charm thresholds:

 $\overline{D}^{0}\Lambda_{c}(2940)^{+}:4804^{+2}_{-1} \text{ MeV}$   $\overline{D}^{*0}\Sigma_{c}(2800)^{+}:4799^{+14}_{-5} \text{ MeV}$   $D^{*-}\Sigma_{c}(2800)^{++}:4811^{+4}_{-6} \text{ MeV}$   $\overline{D}^{*0}\Lambda_{c}(2860)^{+}:4863^{+2}_{-6} \text{ MeV}$   $\overline{D}^{*0}\Lambda_{c}(2940)^{+}:4946^{+1}_{-1} \text{ MeV}.$ 

 $Λ_c(2860) [3/2^+]$  $Λ_c(2940) [3/2^-]$ 

 $\Sigma_c(2800): \frac{1}{2}^-$  from the analysis of  $\Lambda_b \to \pi^- p D^0$ 

S. Sakai, FKG, B. Kubis, PLB 808 (2020) 135623

Measured widths

Coupling constants

- Heavy quark spin symmetry
- >  $XD\overline{D}^*$  coupling consistent with hadronic molecular picture

 $m_p + m_X = 4810 \text{ MeV}$ 

 $\bar{D}^{(*)0}$ 

 $\Lambda_c(2860)^+ \Lambda_c(2940)^+$ 

 $D^{(*)0}$ 

X(3872)

 $D^{(*)0}$ 





• Evaluate the box diagrams with a dispersive approach:

$$\mathcal{A}_{\ell S; \bar{\ell}\bar{S}(\gamma p \to \chi_{c1} p)}^{(j)}(s) = \sum_{\ell', S'} \frac{1}{\pi} \int ds'^{s_{\text{cut}}}_{s_{\text{th}}} \frac{\mathcal{A}_{\ell'S'; \bar{\ell}\bar{S}(\gamma p \to \bar{D}^* \Sigma_c / \Lambda_c)}^{(j)}(s') \rho(s') \mathcal{A}_{\ell'S'; \ell S(\chi_{c1} p \to \bar{D}^* \Sigma_c / \Lambda_c)}^{(j)}(s')}{s' - s}$$

consider S-wave for open-charm channels, S-, P- and D-waves for  $\gamma p$ , Xp

**D** Only limited channels, hard cutoff:  $s_{\text{cut}} = \sqrt{q_{\text{max}}^2 + m_{\Sigma_c/\Lambda_c}^2} + \sqrt{q_{\text{max}}^2 + m_{D^{(*)}}^2}$  with  $q_{\text{max}} = 1 \text{ GeV}$ 

 $\blacksquare$  Monopole form factor for exchanged particles with  $\Lambda=m_{\rm ex}+\eta\Lambda_{\rm QCD}$ 

$$F(t) = \frac{\Lambda^2 - m_{\rm ex}^2}{\Lambda^2 - t}$$





Here  $g_{+/-}$  is the ratio of the  $\Lambda(2860)D^{(*)}p$  and  $\Lambda(2940)D^{(*)}p$  coupling constants constrained by the measured widths



• Triangle singularity induced structure may be used to measure the X(3872) binding energy ( $\delta$ ) FKG, PRL 122 (2019) 202002



• However, here the sensitivity is smeared out by the  $\Lambda(2940)$  width (~ 20 MeV)



#### **Summary**



•  $\sigma(\gamma p \to X(3872)p) = \mathcal{O}(30 \text{ nb}) \text{ at } \sqrt{s_{\gamma p}} \sim 5 \text{ GeV} \Rightarrow \mathcal{O}(0.3 \text{ nb}) \text{ for electroproduction}$ 

**D** Plenty of X(3872) events will be collected with 50 fb<sup>-1</sup> for EicC, 300 fb<sup>-1</sup> for EIC

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