

Static energies from lattice QCD

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Plan of the talk

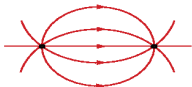
- ▶ The quark-mass dependence of the potential energy between static colour sources in the QCD vacuum with light and strange quarks [Bulava, Knechtli, Koch, Morningstar, Peardon, to appear soon]



The static potential - a probe of QCD at all scales

Potential energy levels (static energies) $V_n(r)$, $n = 0, 1, 2, \dots$ of a quark-anti-quark static pair at distance r

Pure gauge theory:



$r < 0.1 \text{ fm}$

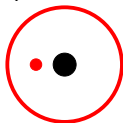
asymptotic freedom,
perturbation theory



$r \gg 1 \text{ fm}$

confinement, flux tube,
effective bosonic string theory

At large distances with dynamical (sea) quarks:



Formation of a pair of static-light mesons: string breaking



String breaking as a mixing phenomenon

String breaking

- ▶ Describes the flattening of the potential $V_0(r)$ at large r
- ▶ *“The ground state potential $V_0(r)$ can therefore be called a static quark potential or a static meson potential”*

[Sommer, Phys. Rept. 275(1) (1996)]

- ▶ Estimate $r_c \approx 1.5$ fm from $V_0(r_c) = 2E_{\bar{Q}Q}$ [Alexandrou et al., Nucl.Phys. B414 (1994)]

- ▶ It has been observed as a mixing of “string-like” and two-meson operators [Drummond, 9805012; Philipsen and Wittig, 9807020; Knechtli and Sommer, 9807022, 0005021; Bali et al., 0505012, Bulava et al., 1902.04006]



String breaking of the ground state potential

- ▶ $N_f=2+1$ QCD: degenerate up, down (light) quarks and a strange quark
- ▶ Two avoided level crossings at $2E_{\bar{Q}l}$ (static-light) and $2E_{\bar{Q}s}$ (static-strange)
- ▶ First calculation in [Bulava et al., 1902.04006]
- ▶ Here: quark-mass dependence and extrapolation to the physical point

$$C(\mathbf{r}, t) = \begin{pmatrix} \square & \sqrt{2} \times \square & \square \\ \sqrt{2} \times \square & 2 \times \square + \begin{matrix} \text{wavy} \\ \text{wavy} \end{matrix} & \sqrt{2} \times \square \\ \square & \sqrt{2} \times \square & \square + \begin{matrix} \text{wavy} \\ \text{wavy} \\ \text{wavy} \end{matrix} \end{pmatrix}$$

CLS ensembles

id	N_{conf}	N_{conf}^W	t_0/a^2	N_s	N_t	m_π [MeV]	m_K [MeV]	$m_\pi L$
N203	94	752	5.1433(74)	48	128	340	440	5.4
N200	104	1664	5.1590(76)	48	128	280	460	4.4
D200	209	1117	5.1802(78)	64	128	200	480	4.2

- ▶ $N_f = 2 + 1$ ensembles from CLS (Coordinated Lattice Simulations) [M. Bruno et al., 1411.3982; Bali et al. 1606.09039]
- ▶ Tree-level $O(a^2)$ improved gauge action [Lüscher and Weisz, *Comm. Math. Phys.* 97 (1985)] and $O(a)$ improved Wilson fermions with non-perturbative C_{SW} [Bulava and Schaefer, 1304.7093]
- ▶ Open b.c. in time and twisted-mass reweighting [Lüscher and Schaefer, 1206.2809]



CLS ensembles contd

- ▶ lattice spacing $a = 0.0633(4)(6)$ fm [Strassberger et al., 2112.06696]
- ▶ quark masses $m_u = m_d = m_l$ and m_s vary along $\sum_{f=u,d,s} m_{\text{bare},f} = \text{constant} \approx \text{physical value}$
- ▶ quark-mass parameter

$$\mu_l = \frac{3m_\pi^2}{m_\pi^2 + 2m_K^2} \approx \frac{3m_l}{2m_l + m_s} \propto m_l$$

$N_f = 3$ symmetric point: $\mu_l = 1$

physical point: $\mu_l = 0.1076$ ($m_\pi = 134.8$ MeV, $m_K = 494.2$ MeV isospin limit)

- ▶ **physical units:** we use the scale t_0 [Lüscher, 1006.4518] and $\sqrt{t_0} = 0.1443(7)(13)$ fm at the physical point [Strassberger et al., 2112.06696]



Techniques

Smearing

- ▶ Temporal gauge links: HYP2 static quark action [Hasenfratz and Knechtli, 0103029; Della Morte, Shindler, Sommer, 0506008; Donnellan et al, 1012.3037]
- ▶ Quarks: Distillation projection on a time-slice onto space spanned by lowest eigenmodes of the 3D gauge-covariant Laplace operator [Peardon et al, 0905.2160], stochastic estimation of quark propagators between distillation spaces [Morningstar et al, 1104.3870]



Techniques contd

Generalized Eigenvalue Problem (GEVP)

- ▶ For fixed inter-quark separation \mathbf{r} solve the GEVP [Lüscher and Wolff, Nucl.Phys.B 339(1) (1990); Blossier et al., 0902.1265]

$$C(t) v_n(t, t_0) = \lambda_n(t, t_0) C(t_0) v_n(t, t_0), n = 0, 1, 2, t \geq t_0$$

We use $t_0 = 5$, solve GEVP at $t = 10$ and project

$$\hat{C}_{ij} = (v_i(t_0, t_d), C(t)v_j(t_0, t_d))$$

We form the ratios

$$R_n(t) = \frac{\hat{C}_{nn}(t)}{C_{\bar{Q}l}^2(t)}, n = 0, 1, 2$$

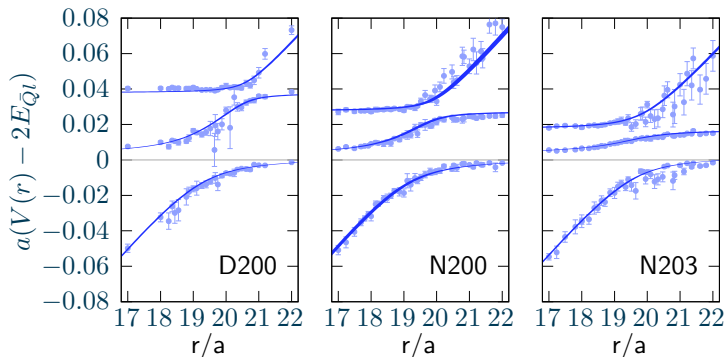
$C_{\bar{Q}l}$ is the correlator of a single static-light meson

- ▶ Correlated single-exponential fit to

$$R_n = \text{const.} \times \exp[-t(V_n - 2E_{\bar{Q}l})]$$



Potential levels



- ▶ Three lowest potential levels on each gauge ensembles
- ▶ Two avoided level crossings at twice the energy of static-light and static-strange mesons
- ▶ Off-axis distances to increase the resolution



Model

Model Hamiltonian

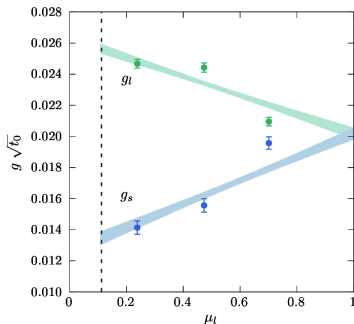
$$H(r) = \begin{pmatrix} \hat{V}(r) & \sqrt{2}g_l & g_s \\ \sqrt{2}g_l & \hat{E}_1 & 0 \\ g_s & 0 & \hat{E}_2 \end{pmatrix}, \hat{V}(r) = \hat{V}_0 + \sigma r + \gamma/r$$

- ▶ 7 fit parameters $\{\hat{E}_1, \hat{E}_2, g_l, g_s, \sigma, \hat{V}_0, \gamma\}$
- ▶ $\hat{V}(r)$, \hat{E}_1 , \hat{E}_2 are the asymptotic energy levels for $r \rightarrow \infty$ up to $O(r^{-1})$
- ▶ Flavour factor $\sqrt{N_l = 2}$ in front of g_l , cf. matrix $C(\mathbf{r}, t)$
- ▶ Model parameters by minimization of *correlated* χ^2
- ▶ Fit ranges start at $4a$ (ground state), when gap $> 2m_\pi$ for excited states, maximal distance is $L/2$ (or below)

In the following: Mass extrapolations of fit parameters in units of t_0 to physical point assuming simple linear behaviour in μ_l



Mass extrapolations of the mixing coefficients

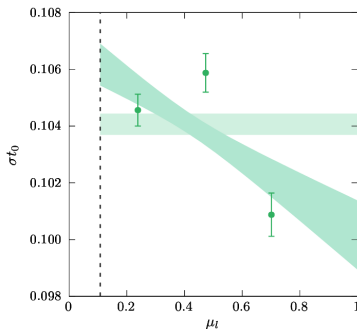


$$g_l \sqrt{t_0} = 0.0256(3), \quad g_l = 35.0 (5)_{\text{stat}} (3)_{\text{sys}} \text{ MeV}$$
$$g_s \sqrt{t_0} = 0.0134(4), \quad g_s = 18.3 (6)_{\text{stat}} (2)_{\text{sys}} \text{ MeV}$$

Mixing parameters g_l and g_s differ by \approx factor two
Some tension between data and model close to $\mu_l = 1$



Mass extrapolations of the string tension

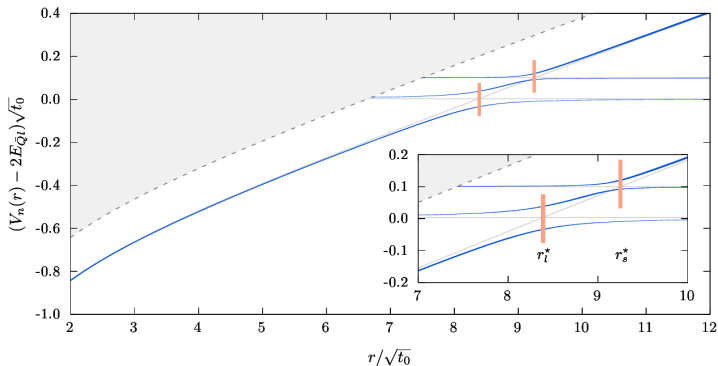


$$\sigma t_0 = 0.1061(7)(20), \quad \sqrt{\sigma} = 445(3)_{\text{stat}}(6)_{\text{sys}} \text{ MeV}$$

String tension has only a mild quark mass dependence
Systematic error includes uncertainty from the difference
between linear (central value) and constant fit



Potentials at the physical point



Gray region is ground state $+2m_\pi$

Crossings of $\hat{V}(r)$ with the asymptotic levels \hat{E}_1 and \hat{E}_2 resp.

$$r_l^* = 8.39(3)\sqrt{t_0} = 1.211(7)_{\text{stat}}(11)_{\text{sys}} \text{ fm}$$

$$r_s^* = 9.26(2)\sqrt{t_0} = 1.336(7)_{\text{stat}}(12)_{\text{sys}} \text{ fm}$$



Conclusions and Outlook

- ▶ Three lowest levels of the static potential up to $r \approx 1.7$ fm in QCD with up, down and strange quarks
- ▶ Simple model based on Cornell + mixing between string and static-meson states
- ▶ Parametrization of the physical potential
- ▶ Mixing parameter of light quarks $\approx 2\times$ of strange quark
- ▶ String tension at the physical point

$$\sigma t_0 = 0.1061(7)(20), \quad \sqrt{\sigma} = 445(3)_{\text{stat}}(6)_{\text{sys}} \text{ MeV}$$

For comparison: pure-gauge $\sigma t_0 \in [0.143, 0.159]$ from T_c/σ , cf. [Necco, 0309017]

TUMQCD [Brambilla et al., 2206.03156], $N_f = 2 + 1 + 1$:

$\sqrt{\sigma} = 467(7), 482(7)$ MeV from Cornell fits at $r < 1$ fm

- ▶ Outlook: hybrid potentials using Laplacian trial states

[Höllwieser et al., 2212.08485 and 2401.09453]

