



Molecular P_ψ pentaquarks from light-meson exchange saturation

Mao-Jun Yan

Institute of Theoretical Physics, Chinese Academy of Sciences, China

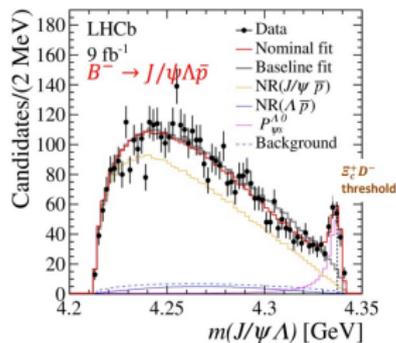
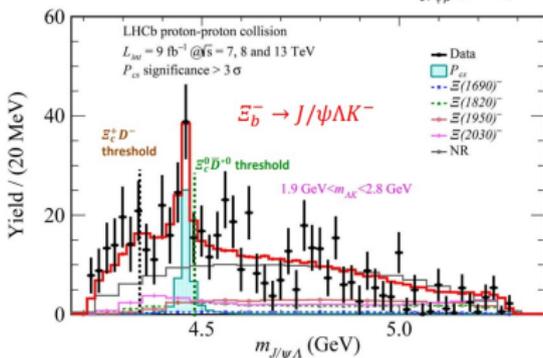
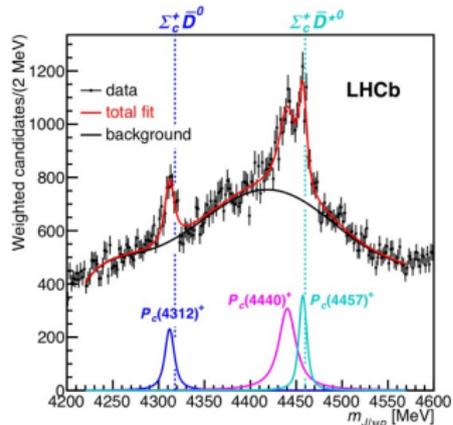
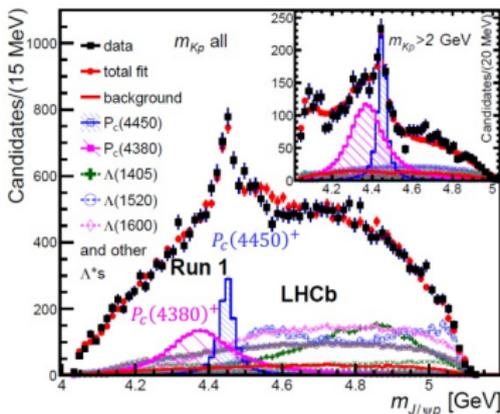
In collaboration with Z.Y. Yang, F.Z. Peng, M.S. Sanchez and M.P. Valderrama

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Outline

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Motivation: observations on $P_{c(s)}$



LHCb, PRL 115 (2015) 072001; PRL 122 (2019) 222001; Science Bulletin 66(2021)1278 ; PRL 131 (2023) 031901

Motivation: observations on $P_{c(s)}$

$P_{c(s)}$ masses and widths [MeV]

- $P_c(4312)$: $M = 4311.9 \pm 0.7_{-0.6}^{+6.8}$, $\Gamma = 9.8 \pm 2.7_{-4.5}^{+3.7}$
- $P_c(4440)$: $M = 4440.3 \pm 1.3_{-4.7}^{+4.1}$, $\Gamma = 20.6 \pm 4.9_{-10.1}^{+8.7}$
- $P_c(4457)$: $M = 4457.3 \pm 0.6_{-1.7}^{+4.1}$, $\Gamma = 6.4 \pm 2.0_{-1.9}^{+5.7}$
- $P_{cs}(4459)$: $M = 4458.8 \pm 2.9_{-1.1}^{+4.7}$, $\Gamma = 17.3 \pm 6.5_{-5.7}^{+8.0}$
- $P_{cs}(4338)$: $M = 4338.2 \pm 0.7$, $\Gamma = 7.0 \pm 1.2$.

Double-peak $P_{cs}(4459)$

- P_{cs1} : $M = 4454.9 \pm 2.7$, $\Gamma = 7.5 \pm 9.7$
- P_{cs2} : $M = 4467.8 \pm 3.7$, $\Gamma = 5.3 \pm 5.3$.

Thresholds

- $\Sigma_c \bar{D}$: 4320.7; $\Sigma_c \bar{D}^*$: 4462.0;
- $\Xi_c' \bar{D}$: 4446.0; $\Xi_c \bar{D}^*$: 4478.0; $\Xi_c^* \bar{D}$: 4513.2; $\Xi_c \bar{D}$: 4336.3

Motivation: explanations on $P_{c(s)}$

Non-molecular state

- Threshold cusps: F.K. Guo et al, PRD92(2015); M. Bayar et al, Phys.Rev.D 94 (2016); S.X. Nakamura et al, PRD 103 (2021) 111503; PRD104 (2021) 9; T.J. Burns et al, PLB838(2023) · · ·
- Hadrocharmonium: M.I. Eides et al, MPL A 35(2020)18; J. Ferretti et al, Sci.Bull. 67 (2022) 1209 · · ·
- Compact P_c : Fl. Stancu, PRD 101(2020)9; H. Garcilazo et al, PRD 105 (2022) 11 · · ·

Molecular state

- QCD sum rule: H.X. Chen et al, PRL115(2015)17 · · ·
- Cont. EFT: M.Z. Liu et al, PRL122(2019); M.L. Du et al, PRL124(2020)7, PRD104(2021)11 · · ·
- Pionful EFT: L. Meng et al, PRD100(2019)1; B. Wang et al, PRD101(2020)3 · · ·
- OBE: R. Chen et al, PRL115(2015)13; F.L. Wang et al, PLB 835 (2022) 137583 · · ·
- Vector-meson Exch.: J.J. Wu et al, PRL105(2010), PRC 84 (2011); C.W. Xiao et al, PRD88(2013), PRD100(2019); L. Roca et al, PRD92 (2015) 9; X.K. Dong et al, Progr.Phys. 41(2021) 65 · · ·

Model on EFT and light meson saturation

EFT inspired OBE

- The S-wave antimeson-baryon interaction in contact-range

$$\langle \vec{p}' | V_C | \vec{p} \rangle = C_0 + C_1 \hat{S}_{L1} \cdot \hat{S}_{L2},$$

with C_0 and C_1 their respective couplings, and with $\hat{S}_{Li} = \vec{S}_{Li} / |S_{Li}|$ the normalized light-quark spin operator of hadron $i = 1, 2$

Saturation: $\zeta = +1$ for P_c

- The C_0 and C_1 are saturated by scalar and vector meson exchange.

$$V_S = - \frac{g_{S1} g_{S2}}{m_S^2 + \vec{q}^2} |_{C_0}$$

$$V_V = (\zeta + T_{12}) \left[\frac{g_{V1} g_{V2}}{m_V^2 + \vec{q}^2} |_{C_0} + \frac{f_{V1} f_{V2}}{6M^2} \frac{m_V^2}{m_V^2 + \vec{q}^2} \hat{S}_{L1} \cdot \hat{S}_{L2} \right] + \dots,$$

Model on EFT and light meson saturation

Couplings

- $g_S = 3.4$ and 6.8 for the charmed mesons and baryons, respectively.
- $g_V = 2.9$ and 5.8 for the charmed mesons and $\Sigma_c^{(*)}$ baryons and $g_V = 2.9$ for the $\Xi_c^{('/*)}$ baryons.
- the magnetic-like couplings we define $f_V = \kappa_V g_V$ with $\kappa_V = \frac{3}{2} (\mu_u / \mu_N)$ for both charmed mesons and baryons ($\Sigma_c, \Sigma_c^*, \Xi_c', \Xi_c^*$), with μ_N the nuclear magneton and $\mu_u \simeq 1.9 \mu_N$.

Mass inputs

- $m_S = 475 \text{ MeV}$ (i.e. the average of its $400 - 550$) MeV .
- $m_V = (m_\rho + m_\omega) / 2 = 775 \text{ MeV}$.
- The extension to the other molecular configurations can be done by modifying the vector meson masses to the K^* and ϕ vector mesons.
- $M = m_N = 938.9 \text{ MeV}$ the nucleon mass.

Model on EFT and light meson saturation

Saturation

- σ -exchange: $C_0^S (\Lambda \sim m_S) \propto -\frac{g_{S1}g_{S2}}{m_S^2}$, $C_1^S (\Lambda \sim m_S) \propto 0$
- vector-meson Exch.:
 $C_0^V (\Lambda \sim m_V) \propto \frac{g_{V1}g_{V2}}{m_V^2} (\zeta + T_{12})$, $C_1^V (\Lambda \sim m_V) \propto \frac{f_{V1}f_{V2}}{6M^2} (\zeta + T_{12})$,

RG evolution

- Owing to $m_S \neq m_V$, the regularization scale for the saturation of the C_0 coupling by mesons is not identical. Considering RGE of V_C ,

$$\frac{d}{d\Lambda} \langle \Psi | V_C(\Lambda) | \Psi \rangle = 0,$$

- If $\Psi(r) \sim r^{\alpha/2}$ at distances $r \sim 1/\Lambda$, $\frac{d}{d\Lambda} \left[\frac{C(\Lambda)}{\Lambda^\alpha} \right] = 0$, with $\alpha = 1$.
- $C(\Lambda = m_V) = C_V(m_V) + \left(\frac{m_V}{m_S} \right)^\alpha C_S(m_S)$

Model on EFT and light meson saturation

Inputs

- $V_C = C_{\text{mol}}^{\text{sat}} f\left(\frac{p'}{\Lambda}\right) f\left(\frac{p}{\Lambda}\right)$ with $f(x) = e^{-x^2}$.
- LS Eq.: $1 + 2\mu_{\text{mol}} C_{\text{mol}}^{\text{sat}} \int \frac{d^3\vec{q}}{(2\pi)^3} \frac{f^2\left(\frac{q}{\Lambda}\right)}{M_{\text{th}} + \frac{q^2}{2\mu_{\text{mol}}} - M_{\text{mol}}} = 0$,
with a ref. $P_{\psi}^N(4312)$, $C_{\text{ref}}^{\text{sat}} = -0.80_{-0.01}^{+0.14} \text{ fm}^2$ for $\Lambda = 1 \text{ GeV}$.

Predictions

- Single channel: $1 + 2\mu_{\text{ref}} C_{\text{ref}}^{\text{sat}} R_{\text{mol}} \int \frac{d^3\vec{q}}{(2\pi)^3} \frac{f^2\left(\frac{q}{\Lambda}\right)}{M_{\text{th}} + \frac{q^2}{2\mu_{\text{mol}}} - M_{\text{mol}}} = 0$,
with the ratio between the relative strength of the coupling:

$$R_{\text{mol}} = \frac{\mu_{\text{mol}} C_{\text{mol}}^{\text{sat}}}{\mu_{\text{ref}} C_{\text{ref}}^{\text{sat}}}$$

- $R_{\text{mol}} > 0.72_{-0.01}^{+0.13} |B$, $(R_{\text{mol}} \leq 0.72_{-0.01}^{+0.13}) |V$

Model on EFT and light meson saturation

Errors

- The uncertainty of the mass of the $P_\psi^N(4312)$, $C_{\text{ref}}^{\text{sat}} = -0.80_{-0.01}^{+0.14} \text{fm}^2$
- The uncertainty of the mass of the scalar meson, $m_S = (475 \pm 75) \text{MeV}$, i.e., equivalent to the $m_S = (400 - 550) \text{MeV}$.
- Choosing the largest of these two sources:

$$\Delta M_{\text{mol}} = \text{minmax}\{M_{\text{mol}}(C_{\text{ref}}^{\text{sat}} \pm \Delta C_{\text{ref}}^{\text{sat}}) - M_{\text{mol}}, \\ M_{\text{mol}}(m_S \pm \Delta m_S) - M_{\text{mol}}\},$$

"minmax" refers to the maximum (for the positive error) and minimum (for the negative error) number within the list.

- Widths of charmed hadrons are not included. i.e. $\Gamma(\Sigma_c^*) \approx 15 \text{MeV}$.

Predictions on $P_{c(s)}$ spectrum

System	$I(J^P)$	R_{mol}	M_{mol}	Candidate	$M_{\text{candidate}}$
$\Lambda_c \bar{D}$	$\frac{1}{2} (1_2^-)$	0.69	$(4153.6_{-3.9}^{+0.1(B)})V$		
$\Lambda_c \bar{D}^*$	$\frac{1}{2} (1_2^-, 3_2^-)$	0.72	$(4295.0_{-2.7}^{+0.0(B)})V$		
$\Lambda_c \bar{D}_s$	$0 (1_2^-)$	0.86	$4252.5_{-2.0}^{+2.3}$		
$\Lambda_c \bar{D}_s^*$	$0 (1_2^-, 3_2^-)$	0.89	$4395.2_{-2.3}^{+3.2}$		
$\Xi_c \bar{D}$	$0 (1_2^-)$	1.00	$4327.4_{-0.9}^{+6.9}$	$P_{\psi s}^\Lambda(4338)$	4338.2 ± 0.7
$\Xi_c \bar{D}^*$	$0 (1_2^-, 3_2^-)$	1.04	$4466.7_{-1.0}^{+7.8}$	$P_{\psi s}^\Lambda(4459)$	$4458.9_{-3.1}^{+5.5}$
$\Xi_c \bar{D}$	$1 (1_2^-)$	0.72	$(4336.3_{-2.9}^{+0.0(B)})V$		
$\Xi_c \bar{D}^*$	$1 (1_2^-, 3_2^-)$	0.74	$4477.6_{-2.2}^{+0.0(V)}$		
$\Xi_c \bar{D}_s$	$\frac{1}{2} (1_2^-)$	0.82	$4436.3_{-2.5}^{+1.2}$		
$\Xi_c \bar{D}_s^*$	$\frac{1}{2} (1_2^-, 3_2^-)$	0.85	$4579.2_{-3.0}^{+2.0}$		

System	$I(J^P)$	R_{mol}	M_{mol}	Candidate	$M_{\text{candidate}}$
$\Sigma_c \bar{D}$	$\frac{1}{2} (1_2^-)$	1.00	Input	$P_\psi^N(4312)$	$4311.9_{-0.9}^{+6.8}$
$\Sigma_c^* \bar{D}$	$\frac{1}{2} (3_2^-)$	1.04	$4376.0_{-0.9}^{+7.1}$	$P_\psi^N(4380)$	-
$\Sigma_c \bar{D}^*$	$\frac{1}{2} (1_2^-)$	0.85	$4459.7_{-2.5}^{+2.3}$	$P_\psi^N(4457)$	$4457.3_{-1.8}^{+4.1}$
$\Sigma_c \bar{D}^*$	$\frac{1}{2} (3_2^-)$	1.13	$4445.2_{-2.9}^{+10.4}$	$P_\psi^N(4440)$	$4440.3_{-1.8}^{+4.3}$
$\Sigma_c^* \bar{D}^*$	$\frac{1}{2} (1_2^-)$	0.82	$4525.4_{-2.7}^{+1.3(V)}$	-	-
$\Sigma_c^* \bar{D}^*$	$\frac{1}{2} (3_2^-)$	0.96	$4520.3_{-1.7}^{+5.3}$	-	-
$\Sigma_c^* \bar{D}^*$	$\frac{1}{2} (5_2^-)$	1.19	$4505.8_{-4.7}^{+12.0}$	-	-

Predictions on $P_{c(s)}$ spectrum

Coupled channel scattering

- When there are two nearby meson-baryon thresholds with the same quantum numbers, we will use the previous equation in matrix form

$$\phi_A + 2\mu_A \sum_B \phi_B C_{\text{ref}}^{\text{sat}} R_{\text{mol}}^{AB} \int \frac{d^3\vec{q}}{(2\pi)^3} \frac{f^2\left(\frac{q}{\Lambda}\right)}{M_{\text{th}(B)} - \frac{q^2}{2\mu} - M_{\text{mol}}} = 0.$$

Channels	System	$I(J^P)$	S	M_{mol}	$\Delta M_{\text{mol}}^{\text{CC}}$	Candidate	$M_{\text{candidate}}$
(1) $\Lambda_c \bar{D}^* - \Sigma_c \bar{D}$	$\Lambda_c \bar{D}^*$	$0(\frac{1}{2}^-)$	0	$4291.0^{+3.9}_{-1.9}$	$-4.0^{(*)}$	-	-
	$\Sigma_c \bar{D}$	$0(\frac{1}{2}^-)$	0	$4315.8^{+4.8}_{-2.3} - (6.6^{+8.2}_{-4.2})i$	+3.9	$P_{\psi}^N(4312)$	$4311.9^{+6.8}_{-0.9} - (4.9^{+2.3}_{-2.6})i$
(2) $\Lambda_c \bar{D}_s - \Xi_c \bar{D}$	$\Lambda_c \bar{D}_s$	$0(\frac{1}{2}^-)$	-1	$4251.8^{+3.0}_{-1.6}$	-0.6	-	-
	$\Xi_c \bar{D}$	$0(\frac{1}{2}^-)$	-1	$4328.1^{+6.4}_{-0.8} - (1.1^{+1.3}_{-0.7})i$	+0.7	$P_{\psi}^{\Lambda}(4338)$	$4338.2 \pm 0.7 - (3.5 \pm 0.6)i$
(3) $\Lambda_c \bar{D}_s^* - \Xi_c' \bar{D} - \Xi_c \bar{D}^*$	$\Lambda_c \bar{D}_s^*$	$0(\frac{1}{2}^-)$	-1	$4393.5^{+4.7}_{-1.3}$	-1.7	-	-
	$\Xi_c' \bar{D}$	$0(\frac{1}{2}^-)$	-1	$4433.6^{+8.8}_{-2.6} - (0.5 \pm 0.2)i$	-4.3	-	-
	$\Xi_c \bar{D}^*$	$0(\frac{1}{2}^-)$	-1	$4468.8^{+6.4}_{-1.2} - (3.0^{+3.4}_{-1.9})i$	+1.1	$P_{\psi}^{\Lambda}(4459)$	$4458.9^{+5.5}_{-3.1} - (8.7^{+5.2}_{-4.3})i$
(4) $\Lambda_c \bar{D}_s^* - \Xi_c \bar{D}^* - \Xi_c' \bar{D}$	$\Lambda_c \bar{D}_s^*$	$0(\frac{3}{2}^-)$	-1	$4394.0^{+4.3}_{-1.6}$	-1.2	-	-
	$\Xi_c \bar{D}^*$	$0(\frac{3}{2}^-)$	-1	$4464.6^{+9.1}_{-2.3} - (0.6^{+0.4}_{-0.3})i$	-2.1	$P_{\psi}^{\Lambda}(4459)$	$4458.9^{+5.5}_{-3.1} - (8.7^{+5.2}_{-4.3})i$
	$\Xi_c' \bar{D}$	$0(\frac{3}{2}^-)$	-1	$4503.8^{+6.7}_{-0.9} - (2.0^{+2.3}_{-1.2})i$	+1.3	-	-

More predictions on $P_{c(s)}$ spectrum

System	$I(J^P)$	R_{mol}	M_{mol}	Candidate	$M_{\text{candidate}}$
$\Sigma_c \bar{D}$	$1(\frac{1}{2}^-)$	1.00	Input	$P_\psi^N(4312)$	$4311.9^{+6.8}_{-0.9}$
$\Sigma_c' \bar{D}$	$1(\frac{3}{2}^-)$	1.04	$4376.0^{+7.1}_{-0.9}$	-	-
$\Sigma_c \bar{D}^*$	$1(\frac{1}{2}^-)$	0.85	$4459.7^{+2.3}_{-2.5}$	$P_\psi^N(4457)$	$4457.3^{+4.1}_{-1.8}$
$\Sigma_c \bar{D}^*$	$1(\frac{3}{2}^-)$	1.13	$4445.2^{+10.4}_{-2.9}$	$P_\psi^N(4440)$	$4440.3^{+4.3}_{-1.8}$
$\Sigma_c' \bar{D}^*$	$1(\frac{1}{2}^-)$	0.82	$4525.4^{+1.3(V)}_{-2.7}$	-	-
$\Sigma_c' \bar{D}^*$	$1(\frac{3}{2}^-)$	0.96	$4520.3^{+3.3}_{-1.7}$	-	-
$\Sigma_c'' \bar{D}^*$	$1(\frac{3}{2}^-)$	1.19	$4505.8^{+2.0}_{-4.7}$	-	-
$\Xi_c' \bar{D}$	$0(\frac{1}{2}^-)$	1.02	$4435.9^{+7.3}_{-0.9}$	-	-
$\Xi_c'' \bar{D}$	$0(\frac{3}{2}^-)$	1.03	$4502.5^{+7.5}_{-1.0}$	-	-
$\Xi_c' \bar{D}^*$	$0(\frac{1}{2}^-)$	0.88	$4584.2^{+2.8}_{-2.7}$	-	-
$\Xi_c'' \bar{D}^*$	$0(\frac{1}{2}^-)$	1.16	$4568.8^{+10.3}_{-3.0}$	-	-
$\Xi_c' \bar{D}^*$	$0(\frac{3}{2}^-)$	0.84	$4652.5^{+1.7(V)}_{-3.0}$	-	-
$\Xi_c'' \bar{D}^*$	$0(\frac{3}{2}^-)$	0.98	$4646.9^{+5.8}_{-1.8}$	-	-
$\Xi_c''' \bar{D}^*$	$0(\frac{3}{2}^-)$	1.22	$4631.8^{+2.5}_{-4.9}$	-	-
$\Sigma_c \bar{D}_s - \Xi_c' \bar{D}$	$1(\frac{1}{2}^-)$	0.96	$4417.3^{+4.2}_{-1.6}$	-	-
$\Sigma_c' \bar{D}_s - \Xi_c' \bar{D}$	$1(\frac{3}{2}^-)$	0.98	$4481.6^{+4.4}_{-1.7}$	-	-
$\Sigma_c \bar{D}_s^* - \Xi_c' \bar{D}^*$	$1(\frac{1}{2}^-)$	0.84	$4581.9^{+4.4}_{-2.4} - (0.6^{+0.0}_{-0.5})i$	-	-
$\Sigma_c \bar{D}_s^* - \Xi_c' \bar{D}^*$	$1(\frac{3}{2}^-)$	1.08	$4556.4^{+7.4}_{-1.0}$	-	-
$\Sigma_c' \bar{D}_s^* - \Xi_c' \bar{D}^*$	$1(\frac{1}{2}^-)$	0.81	$4647.4^{+5.3}_{-0.0} - (1.8^{+1.7}_{-0.9})i$	-	-
$\Sigma_c' \bar{D}_s^* - \Xi_c' \bar{D}^*$	$1(\frac{3}{2}^-)$	0.94	$4625.5^{+4.3}_{-2.6}$	-	-
$\Sigma_c'' \bar{D}_s^* - \Xi_c' \bar{D}^*$	$1(\frac{3}{2}^-)$	1.13	$4618.6^{+8.7}_{-1.1}$	-	-
$\Xi_c' \bar{D}_s - \Omega_c \bar{D}$	$1(\frac{1}{2}^-)$	1.01	$4542.9^{+3.9}_{-1.8}$	-	-
$\Xi_c'' \bar{D}_s - \Omega_c \bar{D}$	$1(\frac{3}{2}^-)$	1.02	$4610.0^{+3.9}_{-2.0}$	-	-
$\Xi_c' \bar{D}_s^* - \Omega_c \bar{D}^*$	$1(\frac{1}{2}^-)$	0.87	$4699.6^{+3.8}_{-2.9} - (0.5^{+0.2}_{-0.3})i$	-	-
$\Xi_c'' \bar{D}_s^* - \Omega_c \bar{D}^*$	$1(\frac{3}{2}^-)$	1.13	$4681.0^{+8.3}_{-1.1}$	-	-
$\Xi_c' \bar{D}_s^* - \Omega_c \bar{D}^*$	$1(\frac{3}{2}^-)$	0.84	$4770.8^{+3.4}_{-3.4} - (1.0^{+0.1}_{-0.6})i$	-	-
$\Xi_c'' \bar{D}_s^* - \Omega_c \bar{D}^*$	$1(\frac{3}{2}^-)$	0.97	$4752.0^{+4.9}_{-2.7}$	-	-
$\Xi_c''' \bar{D}_s^* - \Omega_c \bar{D}^*$	$1(\frac{3}{2}^-)$	1.19	$4746.3^{+9.6}_{-2.3}$	-	-

TABLE II. Molecular octet pentaquark states predicted in this work.

System	$I(J^P)$	R_{mol}	M_{mol}
$\Sigma_c \bar{D}$	$1(\frac{1}{2}^-)$	0.57	$(4316.8^{+3.9(B)}_{-2.2(B)})^V$
$\Sigma_c' \bar{D}$	$1(\frac{3}{2}^-)$	0.58	$(4381.8^{+3.6(B)}_{-21.3})^V$
$\Sigma_c \bar{D}^*$	$1(\frac{1}{2}^-)$	0.96	$(4455.6^{+5.4}_{-1.5})^V$
$\Sigma_c \bar{D}^*$	$1(\frac{3}{2}^-)$	0.41	-
$\Sigma_c' \bar{D}^*$	$1(\frac{1}{2}^-)$	1.07	$(4514.6^{+8.3}_{-1.0})^V$
$\Sigma_c' \bar{D}^*$	$1(\frac{3}{2}^-)$	0.79	$(4526.1^{+0.5(V)}_{-2.6})^V$
$\Sigma_c'' \bar{D}^*$	$1(\frac{3}{2}^-)$	0.32	-
$\Sigma_c \bar{D}_s - \Xi_c' \bar{D}$	$1(\frac{1}{2}^-)$	0.65	$(4445.8^{+0.0}_{-2.2(R)} - (0.1^{+4.0}_{-0.1(R)})i)^V$
$\Sigma_c' \bar{D}_s - \Xi_c' \bar{D}$	$1(\frac{3}{2}^-)$	0.65	$(4513.0^{+0.0}_{-2.2(R)} - (0.0^{+4.0}_{-0.0(R)})i)^V$
$\Sigma_c \bar{D}_s^* - \Xi_c' \bar{D}^*$	$1(\frac{1}{2}^-)$	1.02	$(4560.2^{+4.8}_{-2.4(R)} - (4.6^{+17.0}_{-4.6(R)})i)^V$
$\Sigma_c \bar{D}_s^* - \Xi_c' \bar{D}^*$	$1(\frac{3}{2}^-)$	0.50	$(4647.5^{+5.5}_{-2.4(R)} - (4.6^{+17.0}_{-4.6(R)})i)^V$
$\Sigma_c' \bar{D}_s^* - \Xi_c' \bar{D}^*$	$1(\frac{1}{2}^-)$	1.12	$(4622.5^{+6.6}_{-1.1})^V$
$\Sigma_c' \bar{D}_s^* - \Xi_c' \bar{D}^*$	$1(\frac{3}{2}^-)$	0.85	$(4651.9^{+2.3}_{-3.0} - (0.0 \pm 0.0)i)^V$
$\Sigma_c'' \bar{D}_s^* - \Xi_c' \bar{D}^*$	$1(\frac{3}{2}^-)$	0.42	$(4656.2^{+11.2}_{-3.0(R)} - (9.8^{+23.5}_{-9.8(R)})i)^V$
$\Xi_c' \bar{D}_s - \Omega_c \bar{D}$	$1(\frac{1}{2}^-)$	0.72	$(4560.2^{+2.9}_{-2.8} - (1.5^{+0.6}_{-1.2})i)^V$
$\Xi_c'' \bar{D}_s - \Omega_c \bar{D}$	$1(\frac{3}{2}^-)$	0.73	$(4630.6^{+2.8}_{-2.8} - (1.5^{+0.6}_{-1.2})i)^V$
$\Xi_c' \bar{D}_s^* - \Omega_c \bar{D}^*$	$1(\frac{1}{2}^-)$	1.07	$(4680.3^{+7.6}_{-1.0})^V$
$\Xi_c'' \bar{D}_s^* - \Omega_c \bar{D}^*$	$1(\frac{3}{2}^-)$	0.59	$(4704.9^{+10.3}_{-5.7} - (4.7^{+3.4}_{-2.2})i)^V$
$\Xi_c' \bar{D}_s^* - \Omega_c \bar{D}^*$	$1(\frac{3}{2}^-)$	1.17	$(4742.8^{+9.9}_{-1.4})^V$
$\Xi_c'' \bar{D}_s^* - \Omega_c \bar{D}^*$	$1(\frac{3}{2}^-)$	0.92	$(4769.3^{+4.5}_{-2.7} - (0.1 \pm 0.0)i)^V$
$\Xi_c''' \bar{D}_s^* - \Omega_c \bar{D}^*$	$1(\frac{3}{2}^-)$	0.52	$(4778.2^{+15.4}_{-7.9} - (6.5^{+5.8}_{-3.1})i)^V$
$\Omega_c \bar{D}_s$	$0(\frac{1}{2}^-)$	0.78	$(4663.1^{+0.5(V)}_{-2.8})^V$
$\Omega_c' \bar{D}_s$	$0(\frac{3}{2}^-)$	0.79	$(4733.6^{+0.7(V)}_{-3.0})^V$
$\Omega_c \bar{D}_s^*$	$0(\frac{1}{2}^-)$	1.12	$(4792.9^{+9.2}_{-1.1})^V$
$\Omega_c' \bar{D}_s^*$	$0(\frac{3}{2}^-)$	0.67	$(4807.0^{+0.4(B)}_{-10.7})^V$
$\Omega_c \bar{D}_s^*$	$0(\frac{3}{2}^-)$	1.20	$(4857.9^{+11.5}_{-2.3})^V$
$\Omega_c' \bar{D}_s^*$	$0(\frac{3}{2}^-)$	0.98	$(4871.4^{+5.4}_{-2.8})^V$
$\Omega_c'' \bar{D}_s^*$	$0(\frac{3}{2}^-)$	0.60	$(4875.9^{+2.2(B)}_{-23.6})^V$

TABLE III. Molecular decuplet pentaquark states predicted in this work.

Predictions on charmed baryons

Couplings

- $a_0(\Lambda N, {}^1S_0) = -3.1 \text{ fm}$, $a_0(\Lambda N, {}^3S_1) = -2.9 \text{ fm}$, $a_0(\Lambda\Lambda) = -1.3 \text{ fm}$ obtained with $g_{\sigma\Lambda\Lambda} = 3/4 g_{\sigma NN}$.
- $g_S = 10.2$, $g_V = 2.9$, $\mu_u = 1.9$ and $\mu_s = -0.6$.

Hadron	g_σ	g_ρ	g_ω	g_ϕ	κ_ρ	κ_ω	κ_ϕ
D, D^*	$\frac{1}{3} g_S$	g_V	g_V	0	$\frac{3}{2} \mu_u$	$\frac{3}{2} \mu_u$	0
D_s, D_s^*	$\frac{1}{3} g_S$	0	0	$\sqrt{2} g_V$	0	0	$-3 \mu_s$
N	g_S	g_V	$3g_V$	0	$\frac{5}{2} \mu_u$	$\frac{1}{2} \mu_u$	0
Λ	$0.75 g_S$	0	$2g_V$	$\sqrt{2} g_V$	0	0	$-3 \mu_s$
Σ	g_S	$2g_V$	$2g_V$	$\sqrt{2} g_V$	μ_u	μ_u	$-\mu_s$
Ξ	g_S	g_V	g_V	$2\sqrt{2} g_V$	$-\frac{1}{2} \mu_u$	$-\frac{1}{2} \mu_u$	$2\mu_s$

Predictions on charmed baryons with $\Lambda = 0.75$ GeV

System	S	I	J^P	$R_{\text{mol}}(\Lambda_c^+)$	B_{mol}	M_{mol}	$R_{\text{mol}}(\Sigma_c^+)$	B_{mol}	M_{mol}	Candidate	M_{cand}
ND_s	+1	$\frac{1}{2}$	$\frac{1}{2}^+$	0.60	(2.2) ^V	(2905.0 ^{+2.2(B)} _{-6.9}) ^V	0.91	2.8	2904.5 ± 1.4	-	-
ND_s^*	+1	$\frac{1}{2}$	$\frac{3}{2}^+$	0.62	(1.7) ^V	(3049.4 ^{+1.7(B)} _{-1.5}) ^V	0.95	3.4	3047.7 ^{+1.6} _{-1.5}	-	-
ND_s^*	+1	$\frac{1}{2}$	$\frac{5}{2}^+$	0.62	(1.7) ^V	(3049.4 ^{+1.7(B)} _{-6.4}) ^V	0.95	3.4	3047.7 ^{+1.6} _{-1.5}	-	-
ND	0	0	$\frac{1}{2}^+$	0.79	0.6	2805.6 ^{+0.5} _{-1.4} ^V	1.20	17.7	2788.5 ^{+4.9} _{-6.8}	$\Lambda_c(2765)$	2766.6 ± 2.4 [31]
ND^*	0	0	$\frac{3}{2}^+$	0.44	(16) ^V	(2932 ⁺¹⁴ ₋₃₄) ^V	0.66	1.2	2946.3 ^{+1.0} _{-10.5}	$\Lambda_c(2940)$	2939.6 ^{+1.3} _{-1.5} [31]
ND^*	0	0	$\frac{5}{2}^+$	1 (Input)	7.9	2939.6	1.51	42	2906 ⁺¹⁵ ₋₂₂	$\Lambda_c(2940)$	2939.6 ^{+1.3} _{-1.5} [31]
ND	0	1	$\frac{1}{2}^+$	0.66	(0.6) ^V	(2805.6 ^{+0.6(B)} _{-3.1}) ^V	1 (Input)	6.2	2800.0	$\Sigma_c(2800)$	~ 2800 [31]
ND^*	0	1	$\frac{3}{2}^+$	0.72	(0.0) ^V	(2947.5 ^{+0.0(B)} _{-1.1}) ^V	1.09	10.4	2937.1 ^{+1.3} _{-1.7}	-	-
ND^*	0	1	$\frac{5}{2}^+$	0.66	(0.7) ^V	(2946.8 ^{+0.7(B)} _{-3.7}) ^V	0.99	5.7	2941.8 ^{+0.6} _{-0.5}	-	-
ΛD_s	0	0	$\frac{1}{2}^+$	0.54	(5.0) ^V	(3079.0 ^{+3.7} _{-1.9}) ^V	0.82	0.4	3083.6 ± 0.2	-	-
ΛD_s^*	0	0	$\frac{3}{2}^+$	0.51	(7.0) ^V	(3220.9 ^{+5.2} _{-2.1}) ^V	0.77	0.0	3227.9 ^{+0.0(V)} _{-0.3}	-	-
ΛD_s^*	0	0	$\frac{5}{2}^+$	0.59	(3.0) ^V	(3224.9 ^{+2.5} _{-2.4}) ^V	0.87	1.4	3226.5 ^{+0.0(V)} _{-0.3}	-	-
ΣD_s	0	1	$\frac{1}{2}^+$	0.74	0.0	3164.5 ^{+0.0(V)} _{-1.9}	1.12	10.6	3150.9 ^{+1.6} _{-1.3}	-	-
ΣD_s^*	0	1	$\frac{3}{2}^+$	0.74	0.0	3305.3 ^{+0.0(V)} _{-2.1}	1.13	10.7	3294.7 ^{+2.1} _{-1.8}	-	-
ΣD_s^*	0	1	$\frac{5}{2}^+$	0.77	0.2	3305.2 ^{+0.2(V)} _{-2.4}	1.16	12.5	3292.9 ^{+1.4} _{-1.2}	-	-
ΛD	-1	$\frac{1}{2}$	$\frac{1}{2}^+$	0.57	(3.4) ^V	(2979.6 ^{+2.5} _{-1.3}) ^V	0.87	1.3	2981.7 ^{+0.2} _{-0.3}	-	-
ΛD^*	-1	$\frac{1}{2}$	$\frac{3}{2}^+$	0.59	(2.6) ^V	(3121.6 ^{+2.1} _{-3.9}) ^V	0.89	1.8	3122.5 ^{+0.3} _{-0.4}	$\Xi_c(3123)$	3122.9 ± 1.3 [31]
ΛD^*	-1	$\frac{1}{2}$	$\frac{5}{2}^+$	0.59	(2.6) ^V	(3121.6 ^{+2.1} _{-3.9}) ^V	0.89	1.8	3122.5 ^{+0.3} _{-0.3}	$\Xi_c(3123)$	3122.9 ± 1.3 [31]
ΣD	-1	$\frac{1}{2}$	$\frac{1}{2}^+$	0.92	3.8	3056.6 ^{+1.9} _{-2.5} ^V	1.40	28.1	3023.3 ^{+6.1} _{-8.2}	$\Xi_c(3055)$	3055.9 ± 0.4 [31]
ΣD^*	-1	$\frac{1}{2}$	$\frac{3}{2}^+$	0.66	(0.6) ^V	(3201.1 ^{+0.6(B)} _{-5.8}) ^V	0.99	5.0	3196.8 ^{+3.0} _{-3.2}	-	-
ΣD^*	-1	$\frac{1}{2}$	$\frac{5}{2}^+$	1.10	11.5	3190.3 ^{+1.1} _{-1.2}	1.66	47	3155 ⁺¹⁴ ₋₁₈	-	-
ΣD	-1	$\frac{3}{2}$	$\frac{1}{2}^+$	0.69	(0.1) ^V	(3060.3 ^{+0.1(B)} _{-3.4}) ^V	1.05	7.3	3053.1 ^{+2.3} _{-2.1}	-	-
ΣD^*	-1	$\frac{3}{2}$	$\frac{3}{2}^+$	0.71	(0.0) ^V	(3201.7 ^{+0.0(B)} _{-2.7}) ^V	1.08	8.4	3193.3 ^{+2.5} _{-2.2}	-	-
ΣD^*	-1	$\frac{3}{2}$	$\frac{5}{2}^+$	0.71	(0.0) ^V	(3201.7 ^{+0.0(B)} _{-2.7}) ^V	1.08	8.4	3193.3 ^{+2.5} _{-2.2}	-	-
ΞD_s	-1	$\frac{1}{2}$	$\frac{1}{2}^+$	0.82	1.0	3286.7 ^{+0.9} _{-3.2}	1.25	16.6	3270.0 ^{+0.4} _{-0.3}	-	-
ΞD_s^*	-1	$\frac{1}{2}$	$\frac{3}{2}^+$	0.91	3.1	3427.4 ^{+2.3} _{-3.8}	1.38	24.2	3406.3 ^{+2.0} _{-2.5}	-	-
ΞD_s^*	-1	$\frac{1}{2}$	$\frac{5}{2}^+$	0.81	0.8	3429.8 ^{+0.6} _{-3.4}	1.23	15.3	3415.2 ^{+1.6} _{-1.4}	-	-
ΞD	-2	0	$\frac{1}{2}^+$	0.90	2.8	3182.7 ^{+2.1} _{-3.2}	1.36	23.9	3161.6 ^{+3.0} _{-3.9}	$\Omega_c(3185)$	3185.1 ^{+7.6} _{-1.9} [26]
ΞD^*	-2	0	$\frac{3}{2}^+$	1.03	7.6	3319.3 ^{+2.5} _{-3.0}	1.56	36.7	3290.2 ^{+7.0} _{-10.5}	-	-
ΞD^*	-2	0	$\frac{5}{2}^+$	0.87	2.0	3324.8 ^{+1.9} _{-3.6}	1.32	20.9	3306.0 ^{+1.1} _{-1.3}	$\Omega_c(3327)$	3327.1 ^{+1.2} _{-1.8} [26]
ΞD	-2	1	$\frac{1}{2}^+$	0.73	0.0	3185.5 ^{+0.0(V)} _{-2.1}	1.11	9.8	3175.8 ^{+2.6} _{-2.3}	-	-
ΞD^*	-2	1	$\frac{3}{2}^+$	0.76	0.1	3326.8 ^{+0.1(V)} _{-2.6}	1.15	11.1	3315.8 ^{+2.8} _{-2.5}	-	-
ΞD^*	-2	1	$\frac{5}{2}^+$	0.76	0.1	3326.8 ^{+0.1(V)} _{-2.6}	1.15	11.1	3315.8 ^{+2.8} _{-2.5}	-	-

Summary

- The contact range interaction is derived from the combination of the RG evolution and exchange of light-mesons (σ, ρ, ω).
 - Assuming $P_{\psi}^N(4312)$ is a $\bar{D}\Sigma_c$ bound state, we approximately reproduce the masses of the $P_{\psi}^N(4440)$ and $P_{\psi}^N(4457)$ as $J = \frac{3}{2}$ and $\frac{1}{2} \bar{D}^*\Sigma_c$ bound states, as well as the narrow $P_{\psi}^N(4380)$.
 - $\bar{D}\Xi_c$ and $\bar{D}^*\Xi_c$ states with masses of 4327MeV and 4467MeV, respectively, corresponding to the $P_{\psi_S}^{\Lambda}(4338)$ and the $P_{\psi_S}^{\Lambda}(4459)$.
-
- The ΣD and ΛD^* bound states corresponding to the $\Xi_c(3055/3123)$.
 - ΞD and ΞD^* bound states with masses matching those of the recently observed $\Omega_c(3185)$ and $\Omega_c(3327)$.

Summary

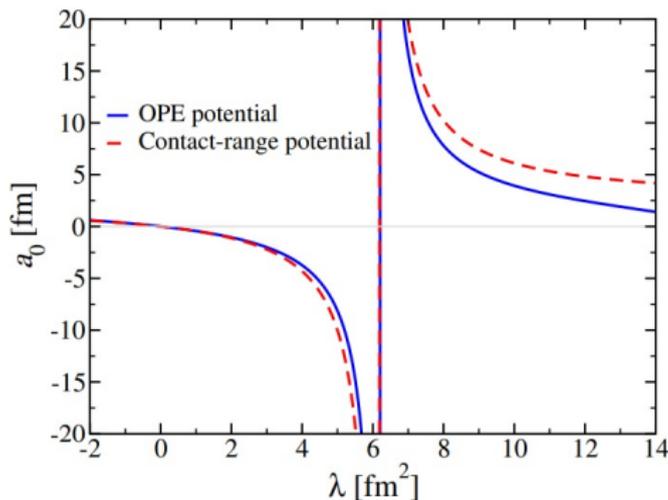
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 - Assuming $P_{\psi}^N(4312)$ is a $\bar{D}\Sigma_c$ bound state, we approximately reproduce the masses of the $P_{\psi}^N(4440)$ and $P_{\psi}^N(4457)$ as $J = \frac{3}{2}$ and $\frac{1}{2} \bar{D}^*\Sigma_c$ bound states, as well as the narrow $P_{\psi}^N(4380)$.
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- The ΣD and ΛD^* bound states corresponding to the $\Xi_c(3055/3123)$.
 - ΞD and ΞD^* bound states with masses matching those of the recently observed $\Omega_c(3185)$ and $\Omega_c(3327)$.

Thanks !

Backup: One-pion-exchange

OPE in $D^* \bar{D}$ scattering, F.Z. Peng et al, PRD 102 (2020)

- $V(\vec{q}) = -\frac{g^2}{6f^2} \tau \vec{\sigma}_1 \cdot \vec{\sigma}_2 \left[\frac{q^2}{m_\pi^2} - \frac{q^4}{m_\pi^4} + \dots \right]$
- $C_0^{\text{OPE}}(\mu \sim m_\pi) \sim 0, \quad C_1^{\text{OPE}}(\mu \sim m_\pi) \sim 0$
- $C_0(\mu) + C_1(\mu) \vec{\sigma}_1 \cdot \vec{\sigma}_2 = -\lambda$



Backup: Error from the input

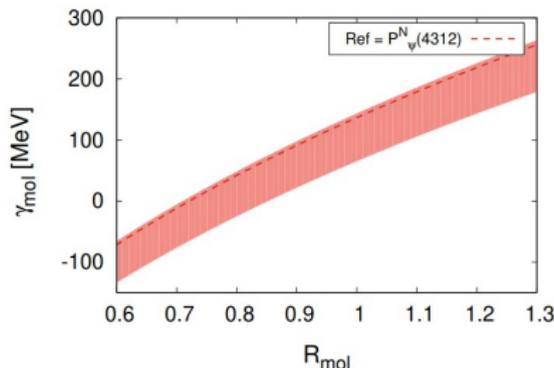


FIG. 1. Dependence of the wave number of a molecule (γ_{mol} , in MeV) on the molecular ratio R_{mol} (defined in Eq. (20)) when the reference state is the $P_{\psi}^N(4312)$. This dependence is calculated by solving Eq. (19) with $\Lambda = 1.0\text{ GeV}$, $C_{ref}^{sat} = -0.80^{+0.14}_{-0.01} \text{ fm}^2$ (given

Analytic loop integral

$$\bullet \int \frac{d^3\vec{q}}{(2\pi)^3} \frac{f^2\left(\frac{q}{\Lambda}\right)}{M_{th} + \frac{q^2}{2\mu_{mol}} - M_{mol}} = \frac{\mu_{mol}}{4\pi^2} \left[\sqrt{2\pi}\Lambda - 2e^{2\gamma^2/\Lambda^2} \pi\gamma \operatorname{erfc}\left(\frac{\sqrt{2}\gamma}{\Lambda}\right) \right],$$

with $\gamma = \sqrt{2\mu(M_{th} - M_{mol})}$. Solutions for which $\gamma > 0$ and $\gamma < 0$ correspond to bound and virtual states, respectively.

$SU(3)$ decomposition

$\bar{H}_c T_c$

$$|\bar{D}_s \Lambda_c\rangle = \frac{1}{\sqrt{3}}|\tilde{1}\rangle + \sqrt{\frac{2}{3}}|\tilde{8}\rangle, \quad |\bar{D}\Xi_c(0)\rangle = \sqrt{\frac{2}{3}}|\tilde{1}\rangle - \frac{1}{\sqrt{3}}|\tilde{8}\rangle,$$

$\bar{H}_c S_c$

$$|\bar{D}_s \Sigma_c\rangle = +\sqrt{\frac{2}{3}}|8\rangle + \frac{1}{\sqrt{3}}|10\rangle, \quad |\bar{D}\Xi'_c(1)\rangle = -\frac{1}{\sqrt{3}}|8\rangle + \sqrt{\frac{2}{3}}|10\rangle$$
$$|\bar{D}_s \Xi'_c\rangle = +\frac{1}{\sqrt{3}}|8\rangle + \sqrt{\frac{2}{3}}|10\rangle, \quad |\bar{D}\Omega_c\rangle = -\sqrt{\frac{2}{3}}|8\rangle + \frac{1}{\sqrt{3}}|10\rangle$$

The octet pieces of the $\bar{H}_c T_c$ and $\bar{H}_c S_c$ pentaquarks generate mixing. This requires the spin-spin / M_1 component of vector meson exchange, and it is expected to be weaker. Yet, the mass difference between the thresholds involved is not particularly large (e.g. the $\bar{D}^* \Xi_c$ channel with the $\bar{D}\Xi'_c$ or $\bar{D}\Xi_c^*$ channels), making this effect more important than naively expected.

Backup: coupled channel scattering

- The $J = \frac{1}{2} \bar{D}^* \Lambda_c - \bar{D} \Sigma_c$ potential

$$V(\bar{D}^* \Lambda_c - \bar{D} \Sigma_c) = \begin{pmatrix} \tilde{V}^O & -W^O \\ -W^O & V^O \end{pmatrix},$$

- The $J = \frac{1}{2} \bar{D}_s \Lambda_c - \bar{D} \Xi_c$ potential

$$V(\bar{D}_s \Lambda_c - \bar{D} \Xi_c) = \begin{pmatrix} \frac{1}{3} \tilde{V}^S + \frac{2}{3} \tilde{V}^O & \frac{\sqrt{2}}{3} (\tilde{V}^S - \tilde{V}^O) \\ \frac{\sqrt{2}}{3} (\tilde{V}^S - \tilde{V}^O) & \frac{2}{3} \tilde{V}^S + \frac{1}{3} \tilde{V}^O \end{pmatrix},$$

- The $J = \frac{1}{2} \bar{D}_s^* \Lambda_c - \bar{D} \Xi'_c - \bar{D}^* \Xi_c$ potential

$$V = \begin{pmatrix} \frac{1}{3} \tilde{V}^S + \frac{2}{3} \tilde{V}^O & -\sqrt{\frac{2}{3}} W^O & \frac{\sqrt{2}}{3} (\tilde{V}^S - \tilde{V}^O) \\ -\sqrt{\frac{2}{3}} W^O & V^O & \frac{1}{\sqrt{3}} W^O \\ \frac{\sqrt{2}}{3} (\tilde{V}^S - \tilde{V}^O) & \frac{1}{\sqrt{3}} W^O & \frac{2}{3} \tilde{V}^S + \frac{1}{3} \tilde{V}^O \end{pmatrix},$$