



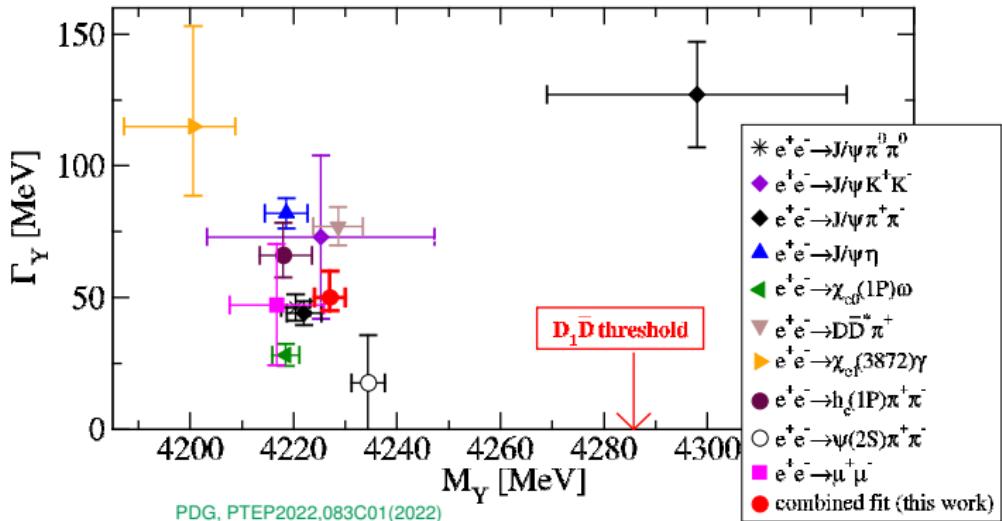
HOW MANY VECTOR STATES SIT IN THE ENERGY RANGE FROM 4.2 TO 4.35 GEV?

February 28, 2024 | Leon von Detten | IAS-4/IKP-3

QWG 2024

based on arXiv:2402.03057 in collaboration with Christoph Hanhart, Vadim Baru, Qian Wang, Daniel Winney and Qiang Zhao

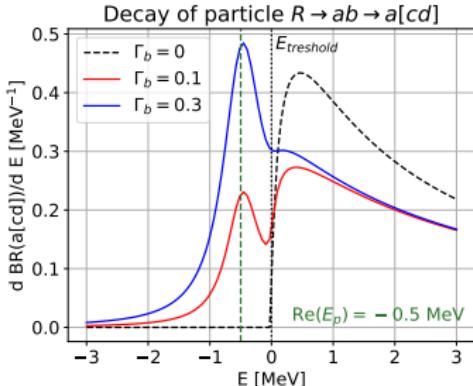
Motivation



- Why do mass and decay width scatter so much for different reactions?
- How big is the influence of the $D_1\bar{D}$ threshold?
- Why $Y(4320)$ only seen in $e^+e^- \rightarrow J/\psi\pi\pi$?

Simultaneous fit of almost all final states.

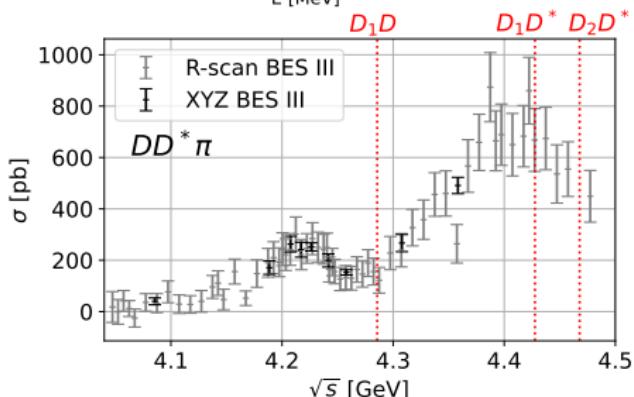
Hadronic Molecules



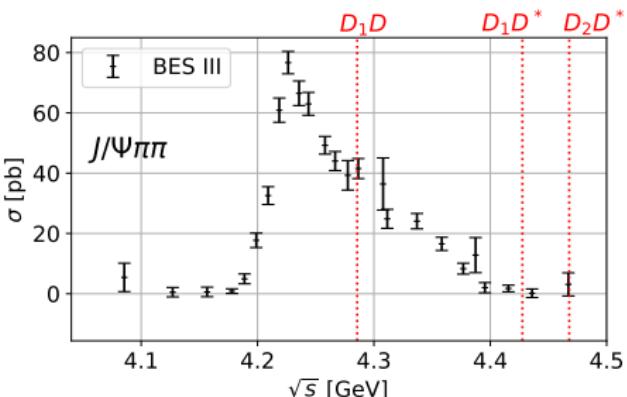
- Coupling to 2 hadron continuum maximal
→ Loop contributions at leading order
→ Possible imprint of threshold on lineshape

Braaten and Lu PRD76(2007)094028, Hanhart et al., PRD81(2010)094028

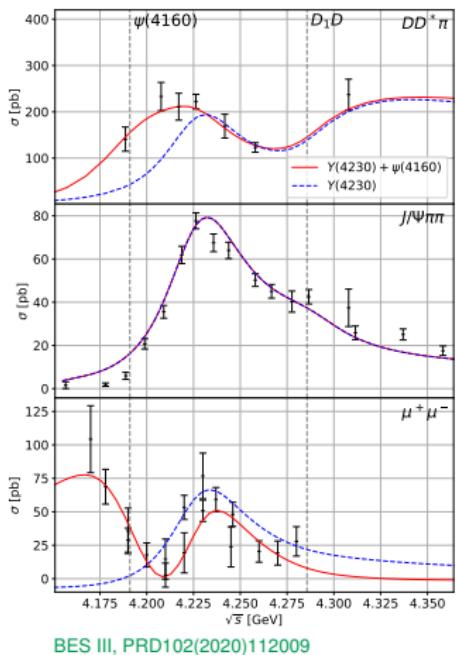
- $Y(4230) \rightarrow D_1 \bar{D} \rightarrow [D^* \pi] \bar{D}$
main decay channel
- Also $D_1 \bar{D}^*$ and $D_2 \bar{D}^*$ 1^{--} bound states in molecular scenario



BES III, PRD106(2022)7,072001 BES III, PRL130(2023)12,1219011

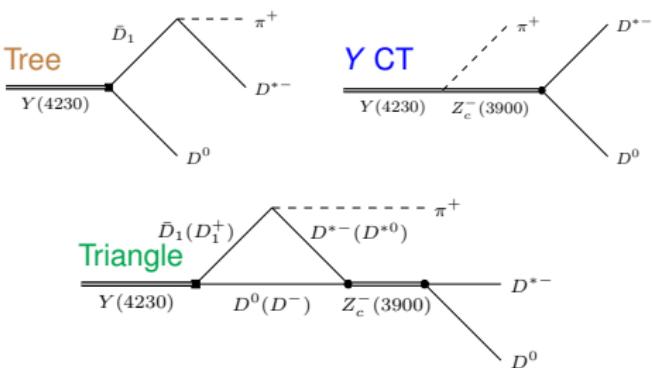
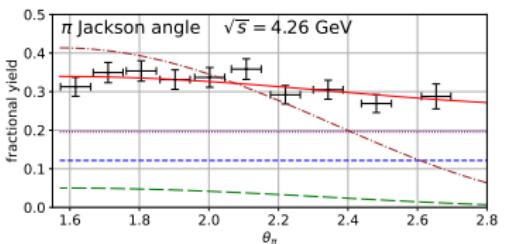
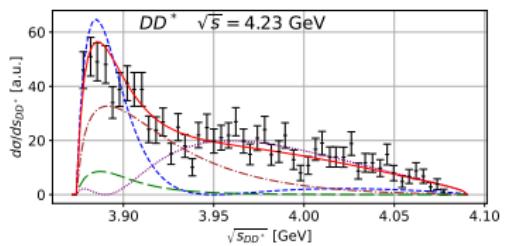
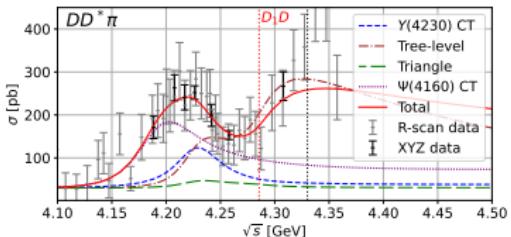


$\psi(4160)$



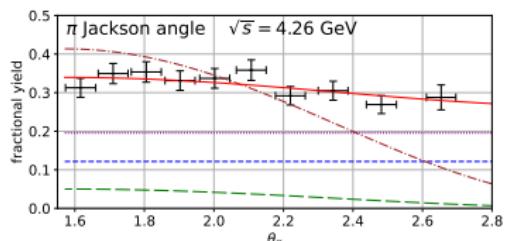
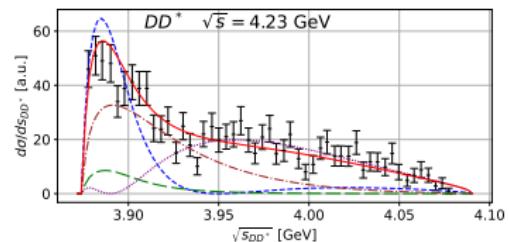
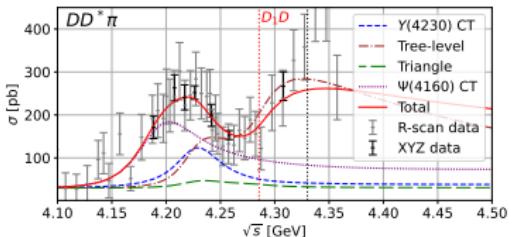
- Narrow structure in $J/\psi\pi\pi$ incompatible with broad structure in $D\bar{D}^*\pi$
- $\mu^+\mu^-$ channel shows destructive interference
- Values from RPP
 $m_{\psi(4160)} = (4191 \pm 5) \text{ MeV}$
 $\Gamma_{\psi(4160)} = (70 \pm 10) \text{ MeV}$
- experimental $Y(4230)$ extraction:
 $D^0 D^{*-} \pi^+ : \Gamma_Y = (77 \pm 6.3 \pm 6.8) \text{ MeV}$
 $J/\psi \pi^+ \pi^- : \Gamma_Y = (41.8 \pm 2.9 \pm 2.7) \text{ MeV}$
 $\mu^+ \mu^- : \Gamma_Y = (47.2 \pm 22.8 \pm 10.5) \text{ MeV}$

$D^0 D^{*-} \pi^+$



- Strong enhancement at $D_1 \bar{D}$ threshold, mainly driven by D_1 D -wave decay
- $D_1 \rightarrow D^* \pi$ decay in $S-$ and $D-$ wave
- Phase of contact terms fixed by $D^* \bar{D}$, $J/\psi \pi$ rescattering into Z_c

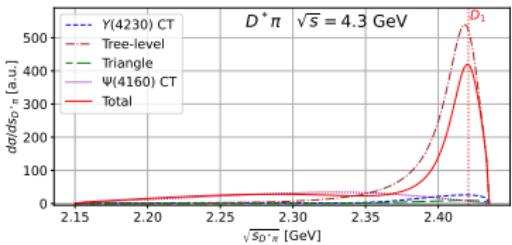
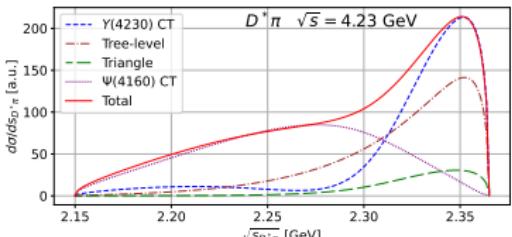
$D^0 D^{*-} \pi^+$



BES III, PRD92(2015)9,092006; PRL112(2014)2,022001

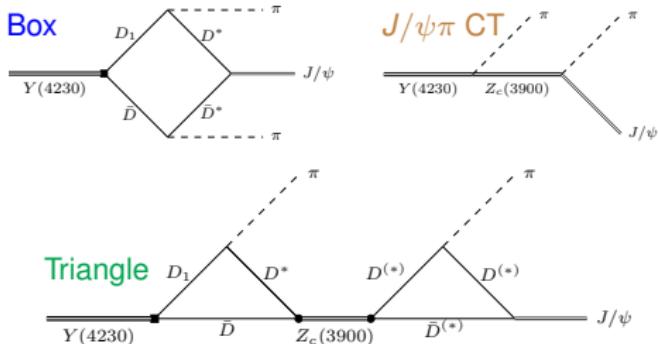
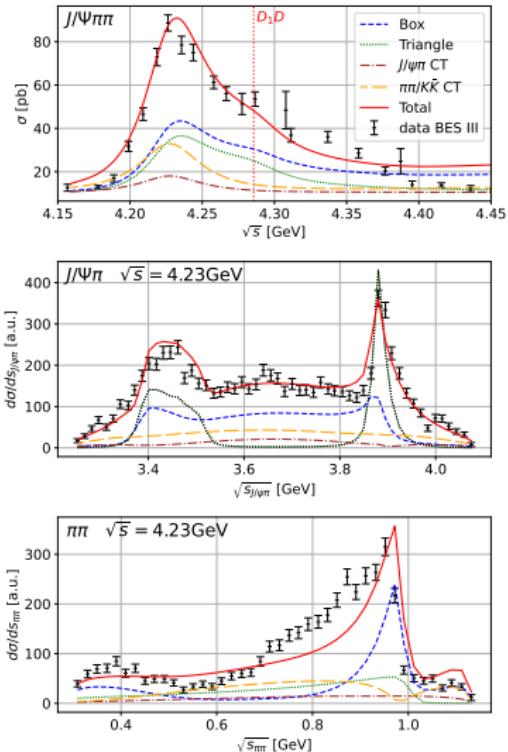
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- Molecular scenario predicts strong D_1 signal in $D^* \pi$ invariant mass distribution

$J/\psi\pi^+\pi^-$

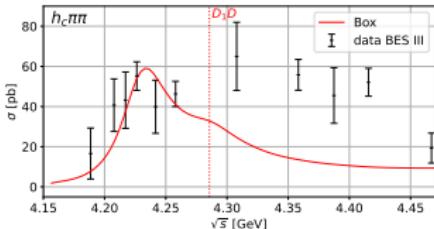


+ different box and triangle topologies

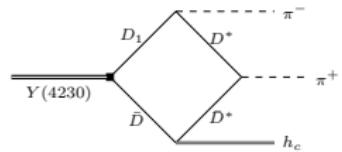
- Asymmetry in total crosssection generated by $D_1\bar{D}$ intermediate state
- $\pi\pi/KK$ s-wave final state interaction approximate via coupled channel Omn  s

Full FSI: [Chen et al., PRD99\(2019\)7,074016](#)
[Danilkin et al., PRD 102\(2020\)1,016019](#)

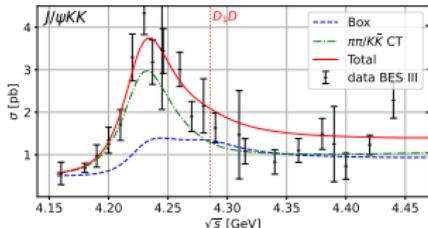
$h_c\pi^+\pi^-$ and $J/\psi K^+K^-$



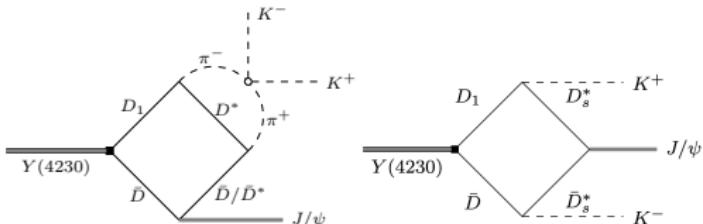
BES III, PRD106(2022)7,072001



- No contact term due to HQSS
- Triangle requires $D_1 \bar{D}^*$ coupling

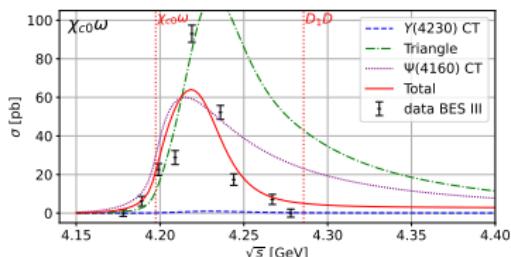


BES III, CPC46(2022)111002

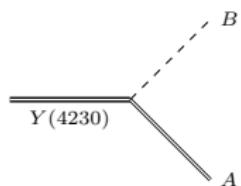
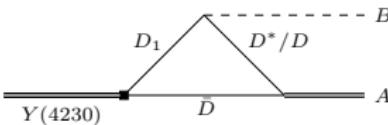


Full coupled channel analysis will include complete $\{D, D^*\} \otimes \{D_1, D_2\}$ multiplets with SU(3) flavor

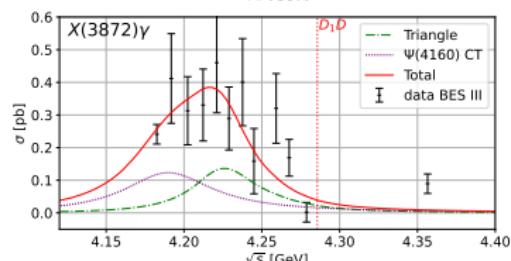
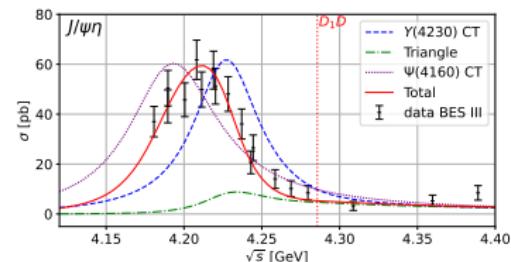
$\chi_{c0}\omega$, $J/\psi\eta$ and $X(3872)\gamma$



$Y(4230) \rightarrow AB$ two body decay

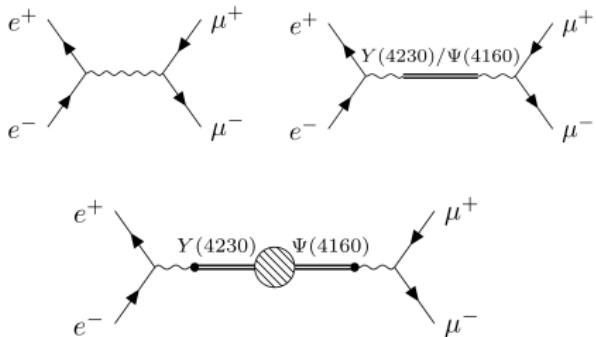
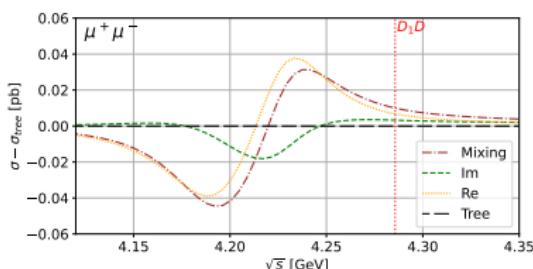
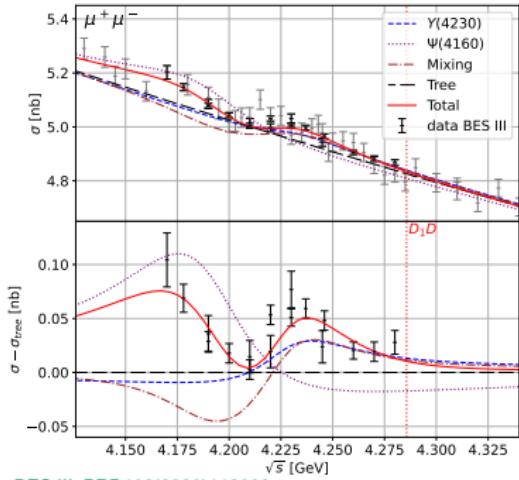


- Main contribution of $Y(4230)$ from triangle loop
- Destructive interference needed to create narrow structure in $\chi_{c0}\omega$
- $Y \rightarrow X(3872)\gamma$ contact term subleading



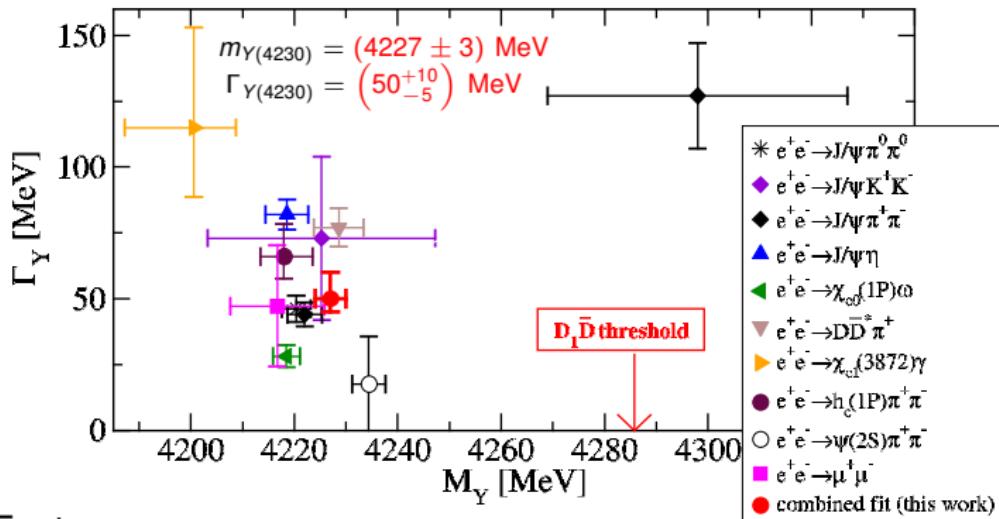
exp. data
 BES III, PRD99(2019)9,091103
 BES III, PRD102(2020)3,031101
 BES III, PRL122(2019)23,232002

$\mu^+ \mu^-$



- Sensitive to one free real parameter and one phase
- Imaginary part of Mixing amplitude fixed by optical theorem
- Reproduce destructive interference between $Y(4230)$ and $\Psi(4160)$ observed in data

Conclusion



Key Features:

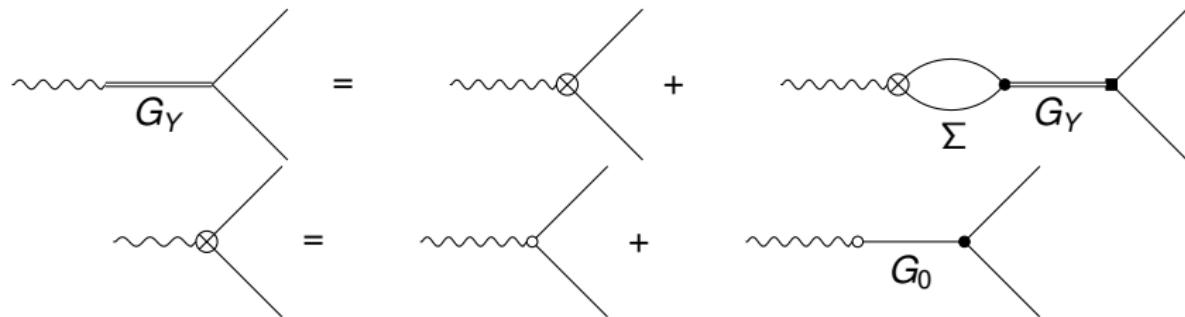
$Y(4230)$ as $D_1 \bar{D}$ molecule and inclusion of $\psi(4160)$

Outlook:

Coupled channel analysis with complete $\{D, D^*\} \otimes \{D_1, D_2\}$ multiplets

Appendix

Y4230



$$G_Y = G_0 + G_0 g_{Y0} (2\omega_Y \Sigma_{D_1 D}) g_{Y0} G_Y .$$

$J^{PC} = 1^{--}$ isosinglet wavefunction

$$|Y(4230)(C = -1, I = 0)\rangle = -\frac{1}{2} (|D_1^+ D^- \rangle + |D^+ D_1^- \rangle + |D_1^0 \bar{D}^0 \rangle + |D^0 \bar{D}_1^0 \rangle)$$

Lagrangian constructed by imposing invariance under heavy-quark spin and chiral symmetry

Lagrangian

Define superfields for light-quark spin doublets:

$$H_a^{(Q)} = \frac{1 + v}{2} [D_a^{*\mu} \gamma_\mu - D_a \gamma_5]$$
$$T_a^{(Q)\mu} = \frac{1 + v}{2} \left[D_{2a}^{\mu\nu} \gamma_\nu - \sqrt{\frac{3}{2}} D_{1a\nu} \gamma_5 (g^{\mu\nu} - \frac{1}{3} \gamma^\nu (\gamma^\mu - v^\mu)) \right]$$

Relevant terms for the interaction

$$\begin{aligned} \mathcal{L}_{\text{int}} = & g \langle H_b^{(Q)} \mathcal{A}_{ba} \gamma_5 \bar{H}_a^{(Q)} \rangle + k \langle T_b^{(Q)\mu} \mathcal{A}_{ba} \bar{T}_b^{(Q)} \rangle \\ & + \frac{h_1}{\Lambda_\chi} \langle T_b^{(Q)\mu} (D_\mu \mathcal{A})_{ba} \gamma_5 \bar{H}_a^{(Q)} \rangle + \frac{h_2}{\Lambda_\chi} \langle T_b^{(Q)\mu} (D \mathcal{A}_\mu)_{ba} \gamma_5 \bar{H}_a^{(Q)} \rangle \\ & + g \langle \bar{H}_a^{(\bar{Q})} \mathcal{A}_{ab} \gamma_5 H_b^{(\bar{Q})} \rangle + k \langle \bar{T}_a^{(\bar{Q})\mu} \mathcal{A}_{ab} T_b^{(\bar{Q})} \rangle \\ & + \frac{h_1}{\Lambda_\chi} \langle \bar{T}_a^{(\bar{Q})\mu} (\mathcal{A} \overleftarrow{D}_\mu)_{ab} \gamma_5 H_b^{(\bar{Q})} \rangle + \frac{h_2}{\Lambda_\chi} \langle \bar{T}_a^{(\bar{Q})\mu} (\mathcal{A}_\mu \overleftarrow{D})_{ab} \gamma_5 H_b^{(\bar{Q})} \rangle + \text{h.c.} \end{aligned}$$

Parameters

	Name	Value
Υ	m_Υ	$(4227 \pm 0.4) \text{ MeV}$
	$g_{\gamma 0}$	$(10.4 \pm 0.2) \text{ GeV}$
	$\Gamma_{\text{in}}^\Upsilon$	$(54 \pm 1) \text{ MeV}$
	$1/f_\Upsilon$	(0.012 ± 0.001)
	$\delta_{\Upsilon\gamma}$	$(17.1 \pm 0.1)^\circ$
ψ	$1/f_\psi$	(0.023 ± 0.003)
	$\delta_{\psi\gamma}$	$(67 \pm 2)^\circ$
Z_c	m_Z	$(3884 \pm 1) \text{ MeV}$
	g_{Z0}	$(4.15 \pm 0.06) \text{ GeV}$
	Γ_{in}^Z	$(48 \pm 1) \text{ MeV}$
$D\bar{D}^*\pi$	$\alpha_1^{(1)}$	(128 ± 12)
	$\alpha_2^{(1)}$	$(3.95 \pm 0.01) \text{ GeV}$
	$\beta_1^{(1)}$	(202 ± 18)
	$\beta_2^{(1)}$	$(3.89 \pm 0.1) \text{ GeV}$

* $\psi(4160)$ mass $m_\psi = 4191 \text{ MeV}$ and width $\Gamma_\psi = 70 \text{ MeV}$ fixed from RPP

	Name	Value
$J/\psi\pi^+\pi^-$	$\alpha_1^{(2)}$	$-(133.9 \pm 4)$
	g_1	$-(14.9 \pm 0.9) 10^{-3}$
	g_8	$(24 \pm 1) 10^{-3}$
	h_1	$-(16.8 \pm 2.4) 10^{-3}$
	h_8	$(15 \pm 0.7) 10^{-3}$
	$\beta_1^{(2)}$	(0 ± 0.1)
	c_{CT}^Δ	$(0.381 \pm 0.1) \text{ GeV}^2$
	$f_{J/\psi}$	456 MeV
$\chi_{c0}\omega$	$c_{\chi_{c0}\omega}^\Delta$	$(1.469 \pm 0.015) \text{ GeV}^2$
	$c_{\chi_{c0}\omega}^Y$	$(0.36 \pm 0.07) 10^{-3}$
	$c_{\chi_{c0}\omega}^\psi$	$-(16 \pm 0.5) 10^{-3}$
$J/\psi\eta$	$c_{J/\psi\eta}^Y$	$(67.3 \pm 3.4) 10^{-3} \text{ GeV}^{-1}$
	$c_{J/\psi\eta}^\psi$	$(298 \pm 11) 10^{-3} \text{ GeV}^{-1}$
$X\gamma$	$c_{X\gamma}^Y$	$(0.71 \pm 0.15) \text{ GeV}^2$
	$c_{X\gamma}^\psi$	$(0.017 \pm 0.003) \text{ GeV}$
$\mu^+\mu^-$	c_{mix}	(0.6 ± 0.01)

$\pi\pi - K\bar{K}$ final state interaction

$$\text{disc}\mathcal{M}_j^I(s) = 2i \sum_k T_{jk}^*(s) \sigma_k(s) \mathcal{M}_k^I(s)$$

Solution given by Muskhelishvili Omn'es function

Full amplitude expressed by:

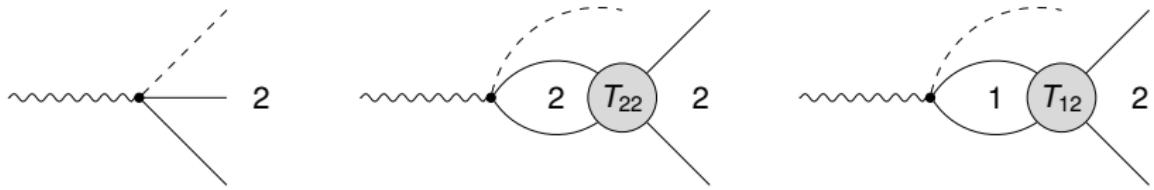
$$\begin{aligned}\mathcal{M}_j^{\text{full}}(s) &= \mathcal{M}_j + \Gamma_j = \mathcal{M}_j + \sum_k \Omega_{jk} \left[(\mathcal{P}_{n-1})_k + \sum_{lm} \frac{s^n}{\pi} \int \frac{z}{z^n} \frac{\Omega_{kl}^{-1}(z) T_{lm}(z) \sigma_m(s) \mathcal{M}_m^0(z)}{z-s} \right] \\ &= \left[\mathcal{M}_j^{>0} + \mathcal{M}_j^0 \right] + \sum_k \left[\Omega_{jk} \left((\mathcal{P}_{n-1})_k + \frac{s^n}{\pi} \text{ P.V.} \int [...] \right) + iT_{jk} \sigma_k \mathcal{M}_k^0 \right]\end{aligned}$$

Modified S-wave approximately given by

$$(\mathcal{M}_j^0)_{\text{mod}} = \mathcal{M}_j^0 + \sum_k \Omega_{jk} (\mathcal{P}_{n-1})_k + \left(iT_{jk} \sigma_k + \frac{1}{\pi} \ln \left(\frac{1}{s/s_{\text{th}} - 1} \right) T_{jk} \sigma_k \right) \mathcal{M}_k^0$$

$$(\mathcal{M}_j)_{\text{mod}} = \mathcal{M}_j^{>0} + (\mathcal{M}_j^0)_{\text{mod}}$$

Production with coupled channel FSI



$$\vec{F} = G_Z \begin{pmatrix} g_1 \left[(E - m_0) \hat{M}_1 + g_2^2 \Sigma_2 (\hat{M}_1 - \hat{M}_2) \right] \\ g_2 \left[(E - m_0) \hat{M}_2 + g_1^2 \Sigma_1 (\hat{M}_2 - \hat{M}_1) \right] \end{pmatrix}$$

$$\rightarrow \vec{F} = G_Z \begin{pmatrix} g_1 \alpha_1^{(1)} (\alpha_2^{(1)} + E) \\ g_2 \alpha_1^{(2)} (\alpha_2^{(2)} + E) \end{pmatrix}$$