

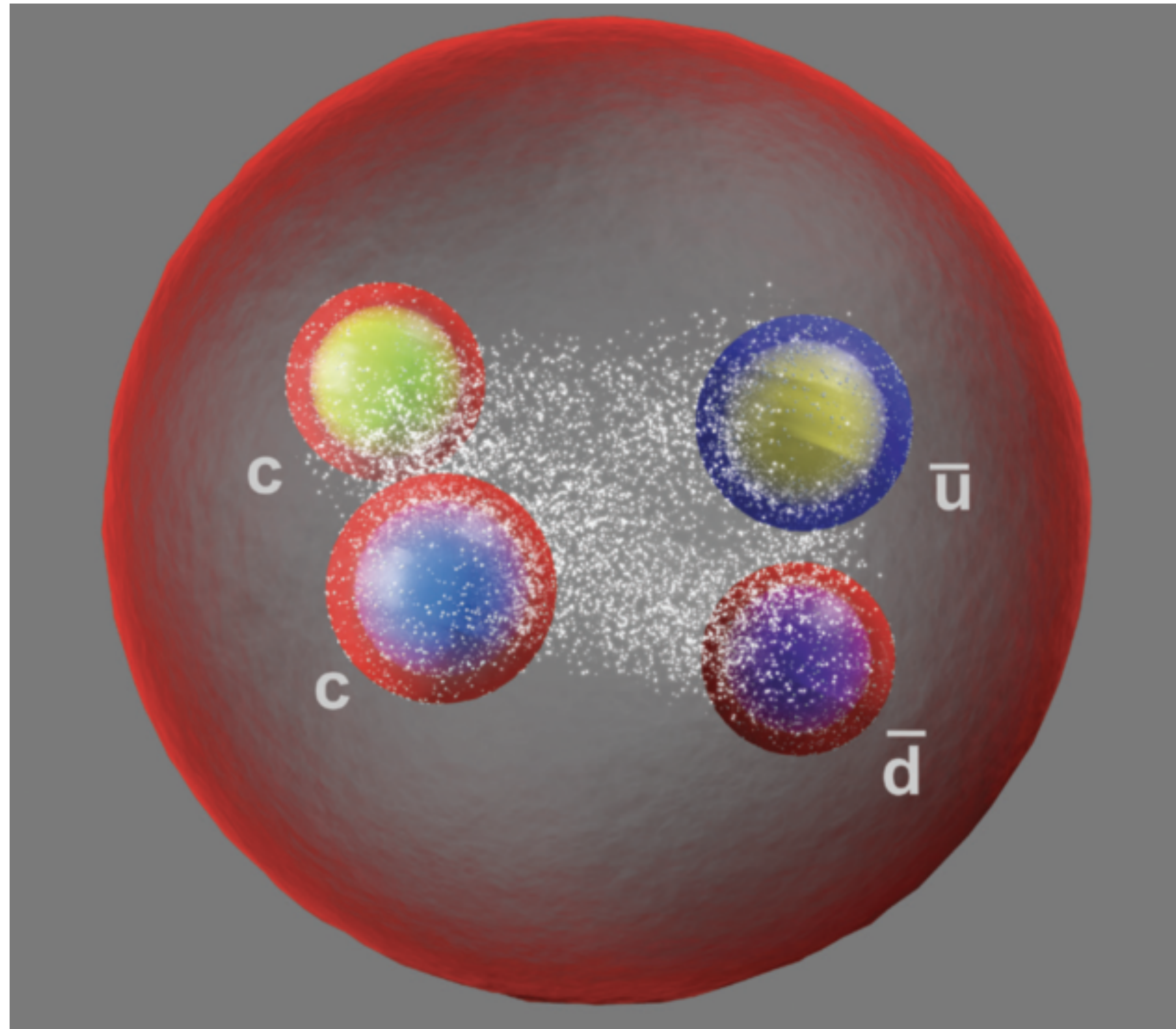
# Strong Decays of $T_{cc}^+$ at NLO in an Effective Field Theory

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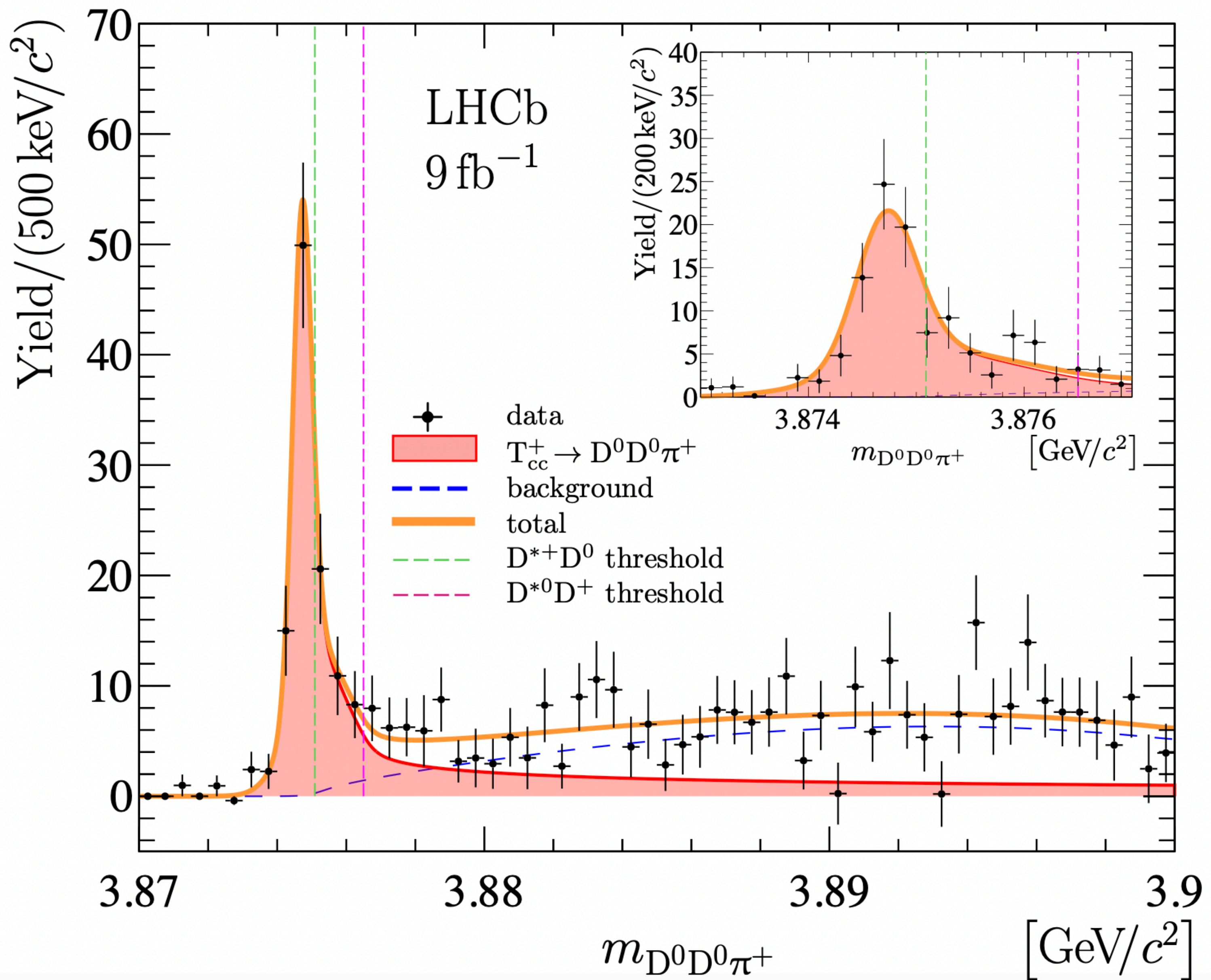
**IISER Mohali, India, 3/1/2024**

# Discovered 7/2021: $T_{cc}^+$ doubly charm tetraquark



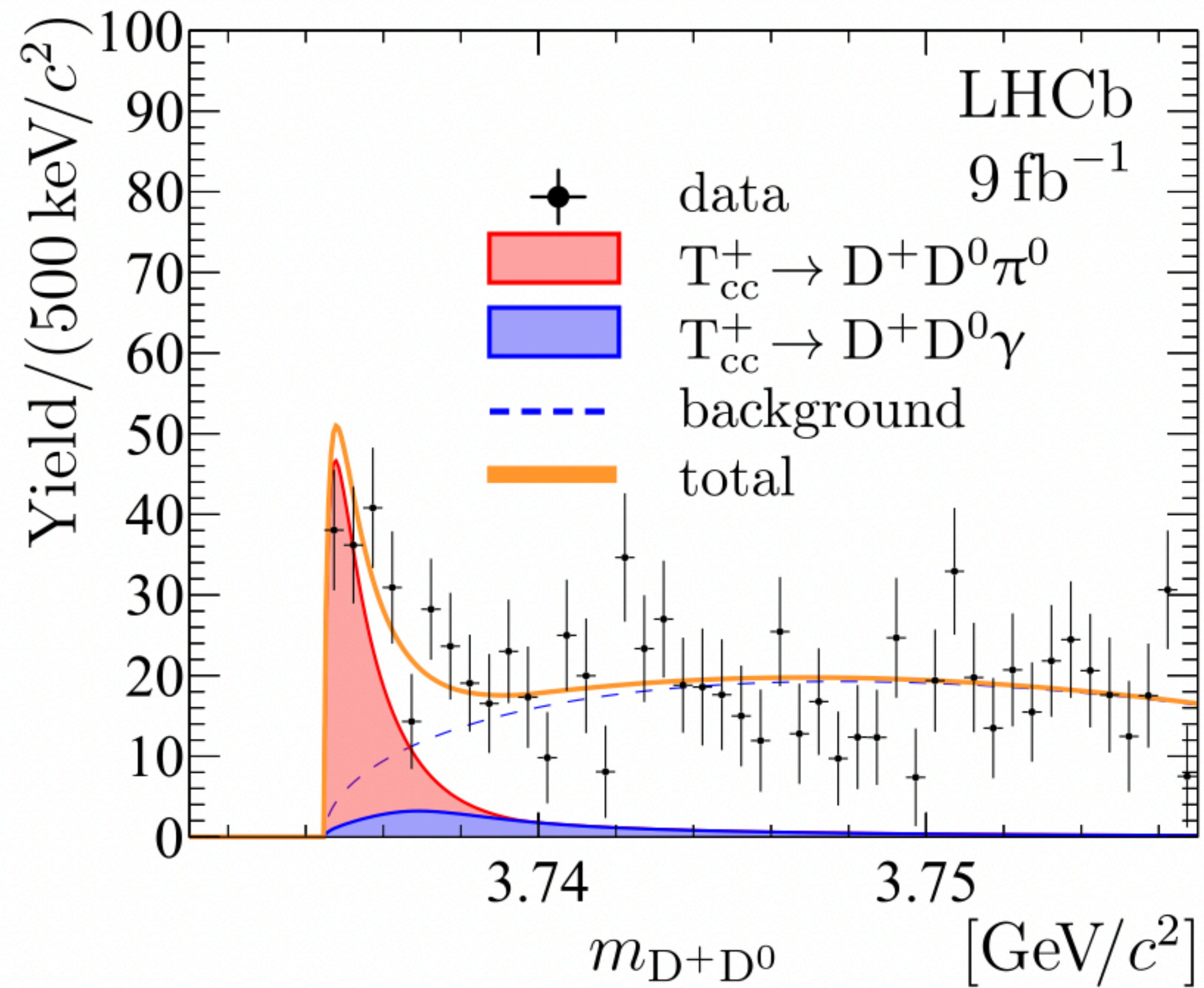
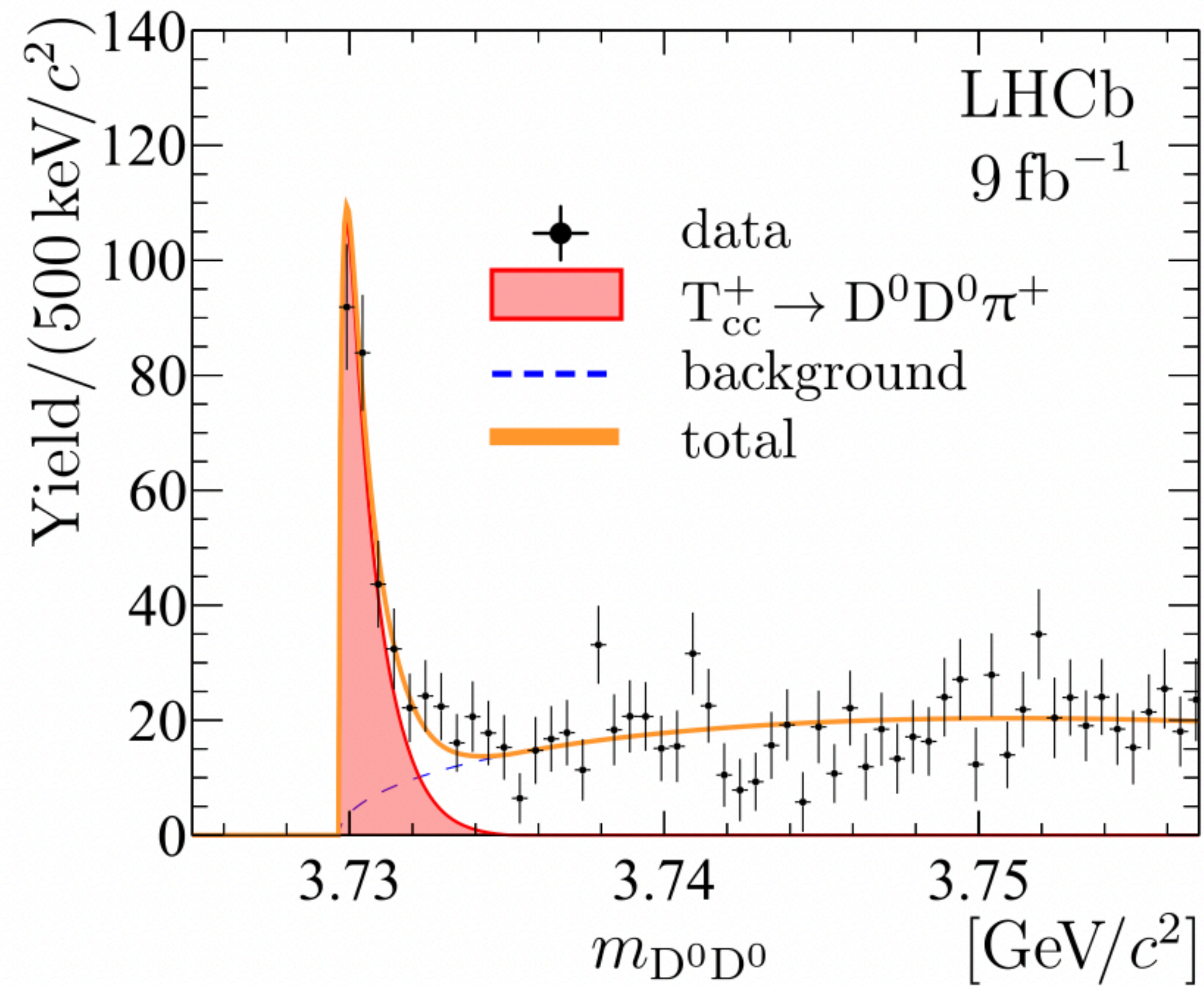
R. Aaij *et. al.* (LHCb), Nature Phys. 18 (2022) 7, 751-754, arXiv 2109.01038 [hep-ex]. (1)

R. Aaij *et. al.* (LHCb), Nature Commun. 13 (2022) 1, 3351, arXiv 2109.01056 [hep-ex]. (2)



$$\delta m_{pole} = -360 \pm 40_{-0}^{+4} \text{ keV},$$

$$\Gamma_{pole} = 48 \pm 2_{-14}^{+0} \text{ keV}. \quad (2)$$



Heavy Hadron Chiral Perturbation Theory  
plus contact terms for S-wave  $D^*D$  scattering

$$\begin{aligned} \mathcal{L} = & H^{*i\dagger} \left( i\partial^0 + \frac{\nabla^2}{2m_{H^*}} - \delta^* \right) H^{*i} + H^\dagger \left( i\partial^0 + \frac{\nabla^2}{2m_H} - \delta \right) H \\ & + \frac{g}{f_\pi} H^\dagger \partial^i \pi H^{*i} + \text{h.c.} + \frac{1}{2} H^\dagger \mu_D \vec{B}^i H^{*i} + \text{h.c.} \\ & - C_0 (H^{*T} \tau_2 H)^\dagger (H^{*T} \tau_2 H) - C_1 (H^{*T} \tau_2 \tau_a H)^\dagger (H^{*T} \tau_2 \tau_a H). \end{aligned}$$

$$H = \begin{pmatrix} D^0 \\ D^+ \end{pmatrix}, \quad H^{*i} = \begin{pmatrix} D^{*0i} \\ D^{*+i} \end{pmatrix}$$

S. Fleming, R. Hodges, TM, Phys. Rev. D. 104 (2021) 11, 116010, arXiv:2109.02188

Similar to XEFT for  $X(3872)$ , new feature is coupled channels

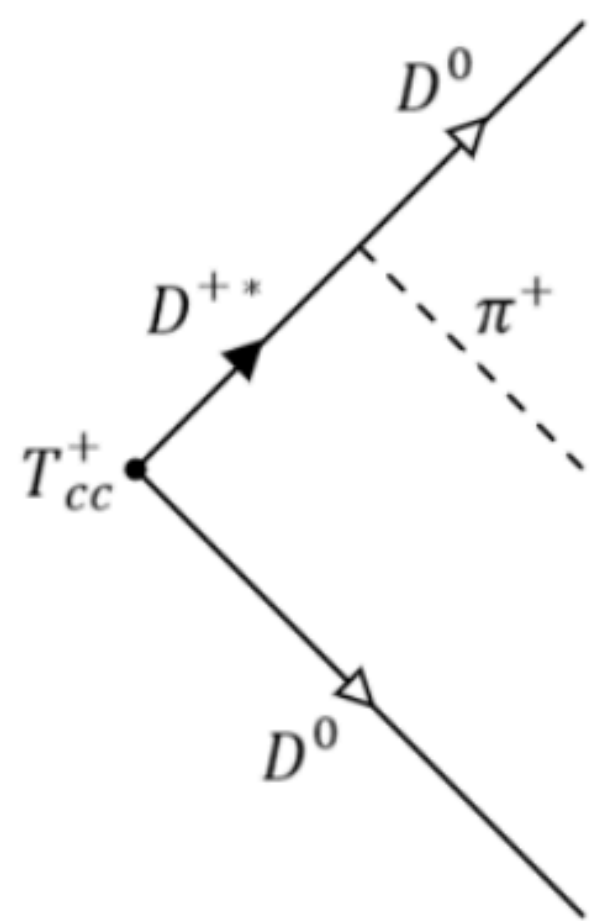
S. Fleming, M. Kusunoki, TM, U. van Kolck, Phys. Rev. D. 76 (2007) 034006, hep-ph/0703168

## T-Matrix for D\*D scattering

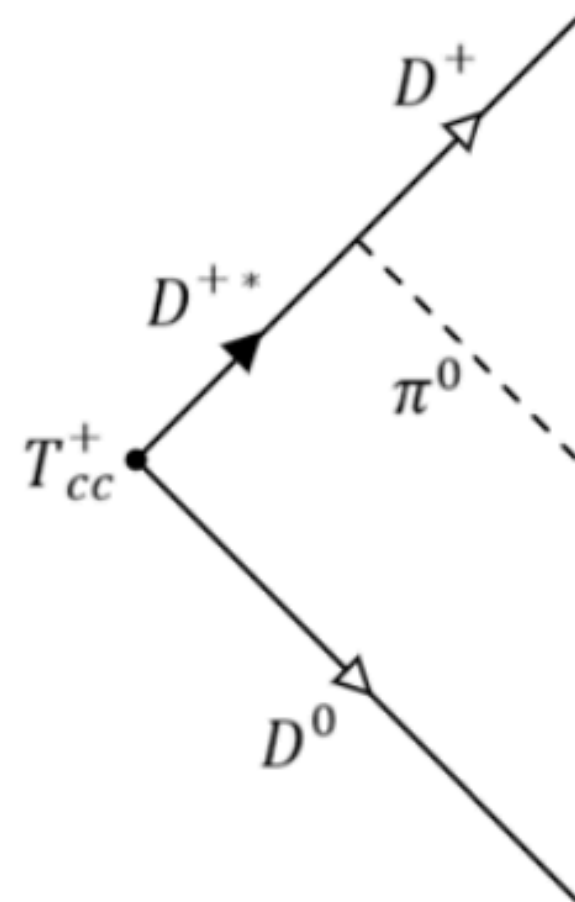
$$T = \frac{1}{E + E_T} \begin{pmatrix} g_0^2 & g_0 g_+ \\ g_0 g_+ & g_+^2 \end{pmatrix} \quad g_0^2 \Sigma'_0(-E_T) + g_+^2 \Sigma'_+(-E_T) = 1;$$

Tune interactions to produce pole at  $T_{cc}^+$

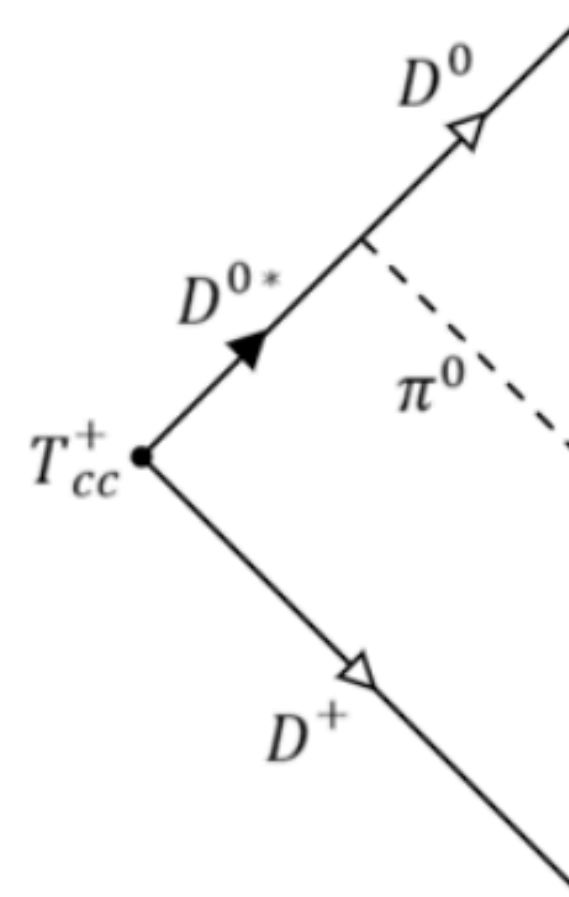
# Decay Diagrams



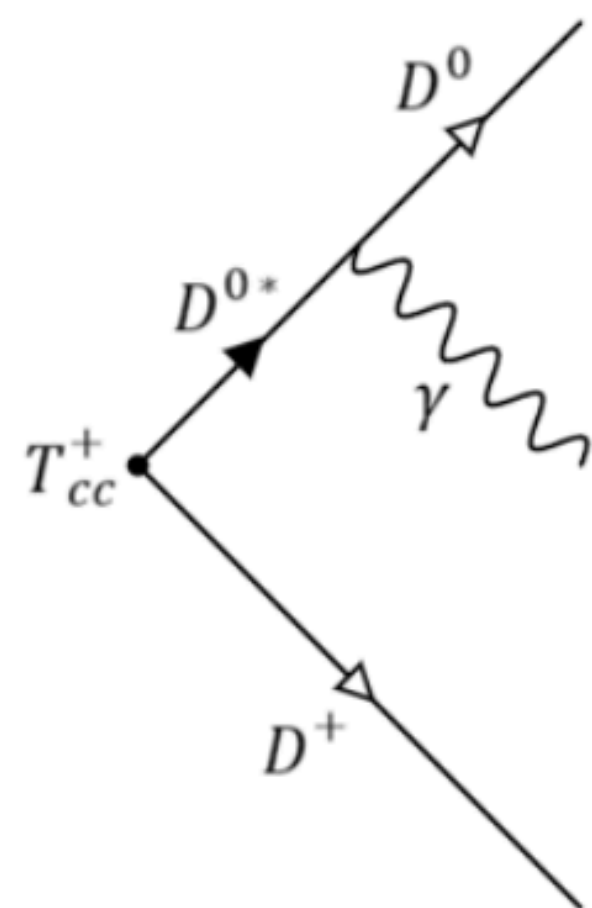
(a)



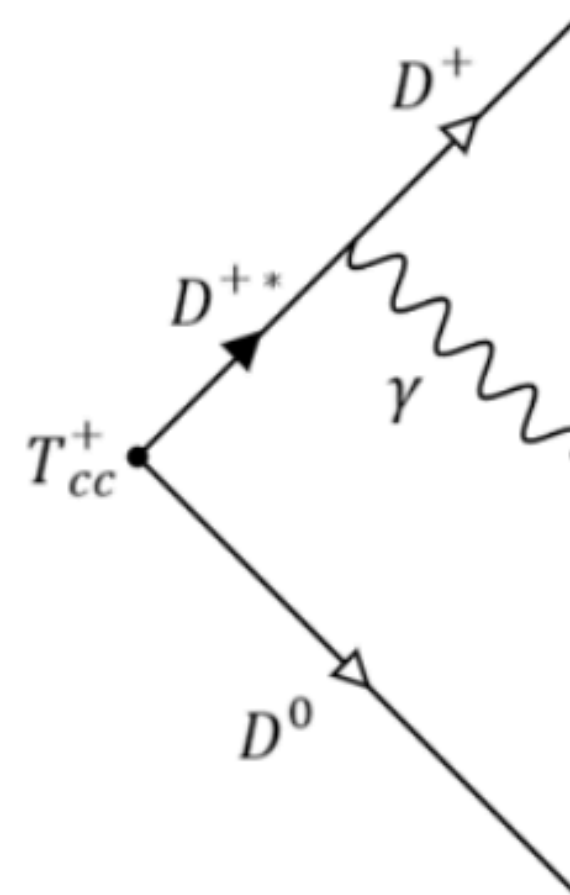
(b)



(c)



(d)



(e)

## Decay rate formulae

$$\frac{d\Gamma[T_{cc}^+ \rightarrow D^0 D^0 \pi^+]}{dp_{D_1^0}^2 dp_{D_2^0}^2} = c_\theta^2 \frac{g^2}{(4\pi f_\pi)^2} \frac{2\gamma_0 p_\pi^2}{3} \left[ \frac{1}{p_{D_1^0}^2 + \gamma_0^2} + \frac{1}{p_{D_2^0}^2 + \gamma_0^2} \right]^2 ,$$

$$\frac{d\Gamma[T_{cc}^+ \rightarrow D^+ D^0 \pi^0]}{dp_{D^+}^2 dp_{D^0}^2} = \frac{g^2}{(4\pi f_\pi)^2} \frac{2p_\pi^2}{3} \left[ \frac{\sqrt{\gamma_0} c_\theta}{p_{D^+}^2 + \gamma_0^2} - \frac{\sqrt{\gamma_+} s_\theta}{p_{D^0}^2 + \gamma_+^2} \right]^2 ,$$

$$\frac{d\Gamma[T_{cc}^+ \rightarrow D^+ D^0 \gamma]}{dp_{D^+}^2 dp_{D^0}^2} = \frac{E_\gamma^2}{6\pi^2} \left[ \frac{\sqrt{\gamma_0} c_\theta \mu_{D^0}}{p_{D^+}^2 + \gamma_0^2} - \frac{\sqrt{\gamma_+} s_\theta \mu_{D^+}}{p_{D^0}^2 + \gamma_+^2} \right]^2 .$$

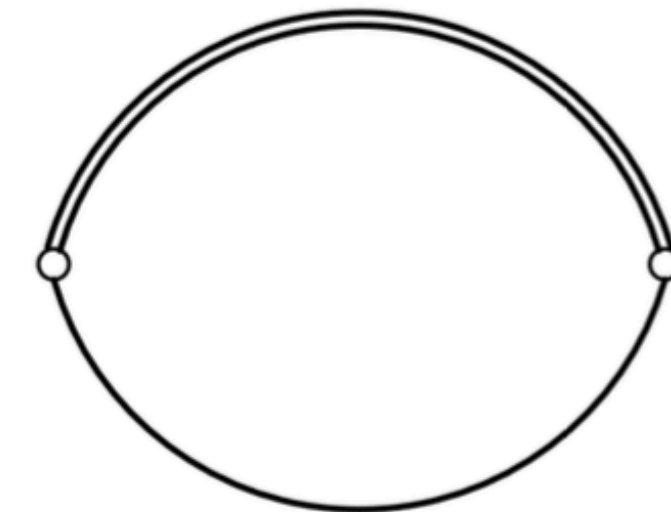
$$\gamma_0^2 = 2\mu_0(m_{D^0} + m_{D^{*+}} - m_T)$$

$$\gamma_+^2 = 2\mu_+(m_{D^+} + m_{D^{*0}} - m_T)$$

$$g_0^2 = \frac{\cos^2 \theta}{\Sigma'_0(-E_T)}$$

$$g_+^2 = \frac{\sin^2 \theta}{\Sigma'_+(-E_T)} ;$$

$$i\Sigma_i(E) =$$





# LO Predictions for Decay Rate

	I=0	I=1	$\Gamma_{\max}$
$\theta$	$-32.4^\circ$	$32.4^\circ$	$-8.34^\circ$
$\Gamma[T_{cc}^+ \rightarrow D^0 D^0 \pi^+]$	32	32	44
$\Gamma[T_{cc}^+ \rightarrow D^+ D^0 \pi^0]$	15	3.8	13
$\Gamma[T_{cc}^+ \rightarrow D^+ D^0 \gamma]$	6.1	2.8	1.9
$\Gamma[T_{cc}^+]$	52	38	58

$$I = 0 : g_0 = -g_+$$

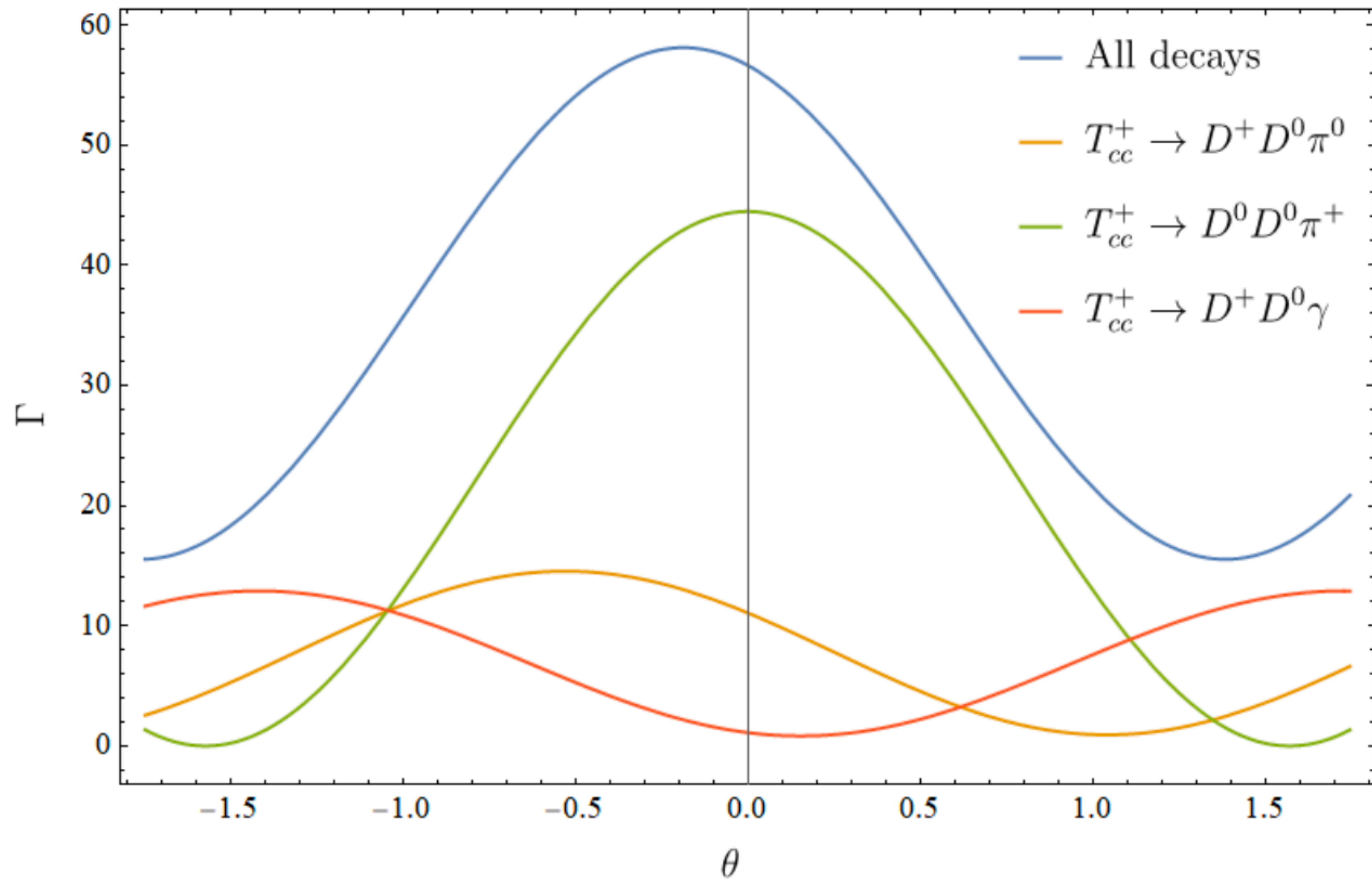
$$I = 1 : g_0 = g_+$$

## Other Predictions

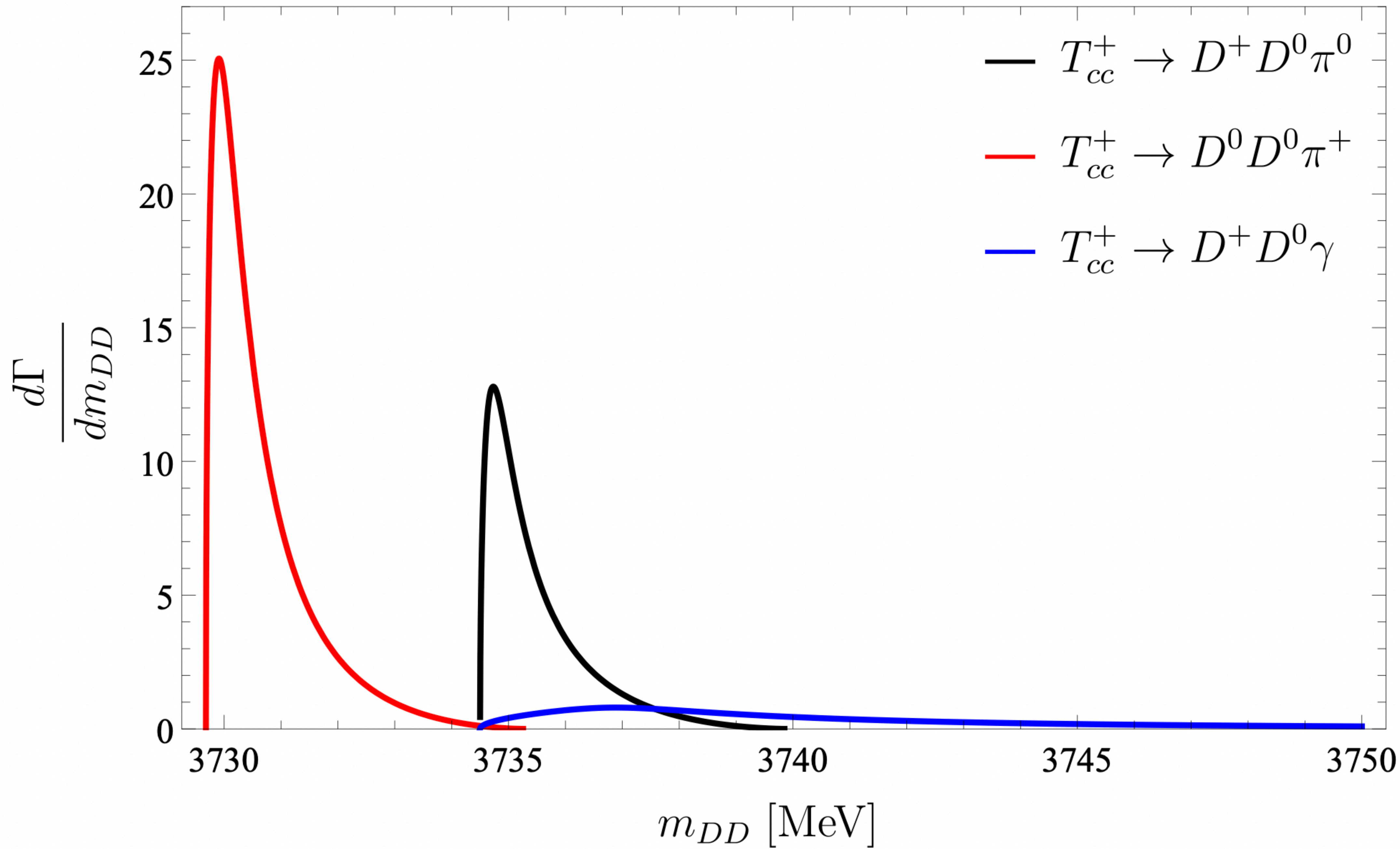
	$\Gamma(T_{cc}^+) \text{ (keV)}$
<b>Fleming et al.</b>	<b>52</b>
Meng et al.	$46.7_{-2.9}^{+2.7}$
Ling et al.	53
Feijoo et al.	43
Yan & Valderrama	$49 \pm 16$
Albaladejo	77

$$\delta m_{BW} = -273 \pm 61 \pm 5_{-14}^{+11} \text{ keV},$$

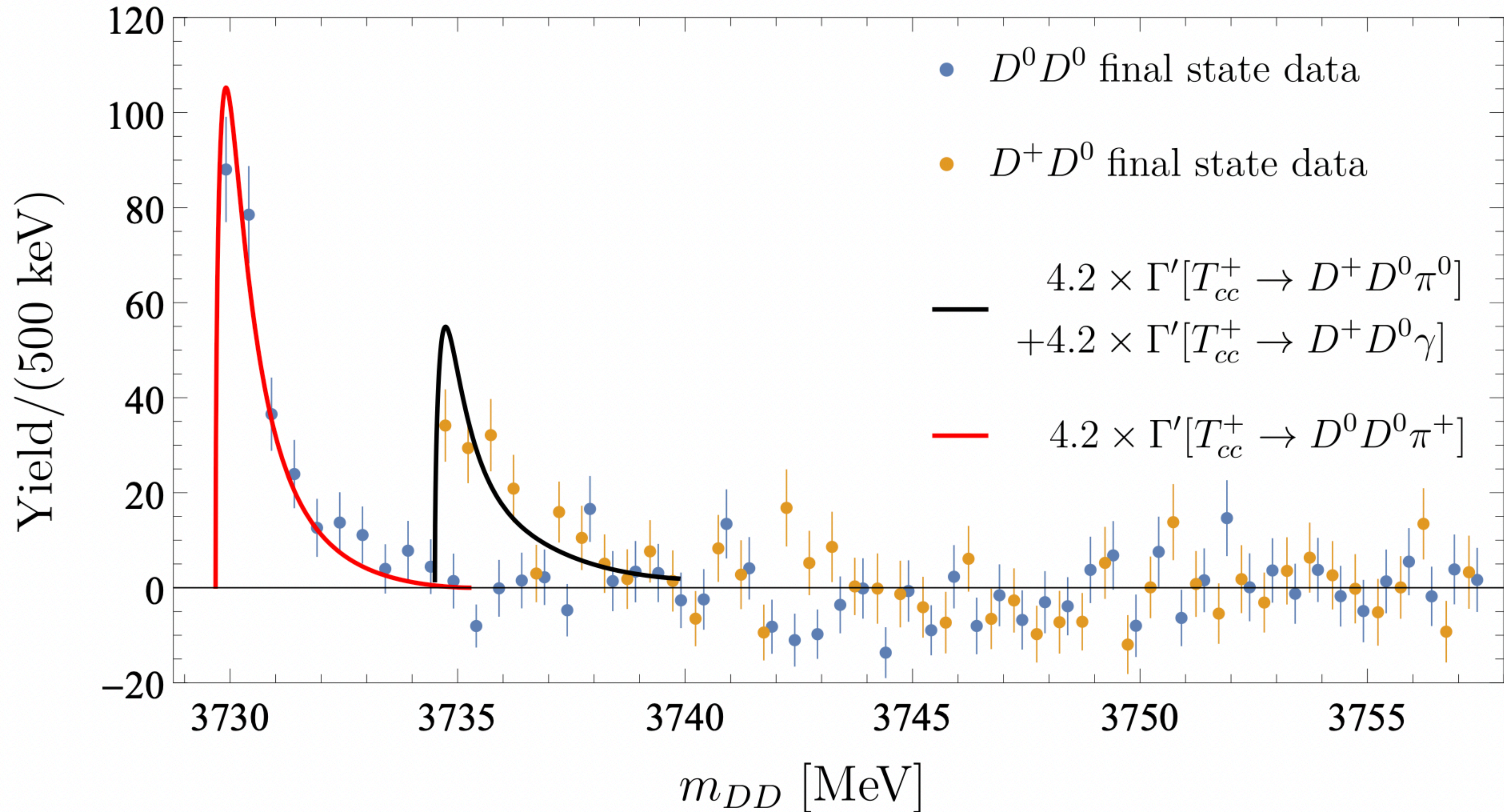
$$\Gamma_{BW} = 410 \pm 165 \pm 43_{-38}^{+18} \text{ keV}. \quad (1)$$



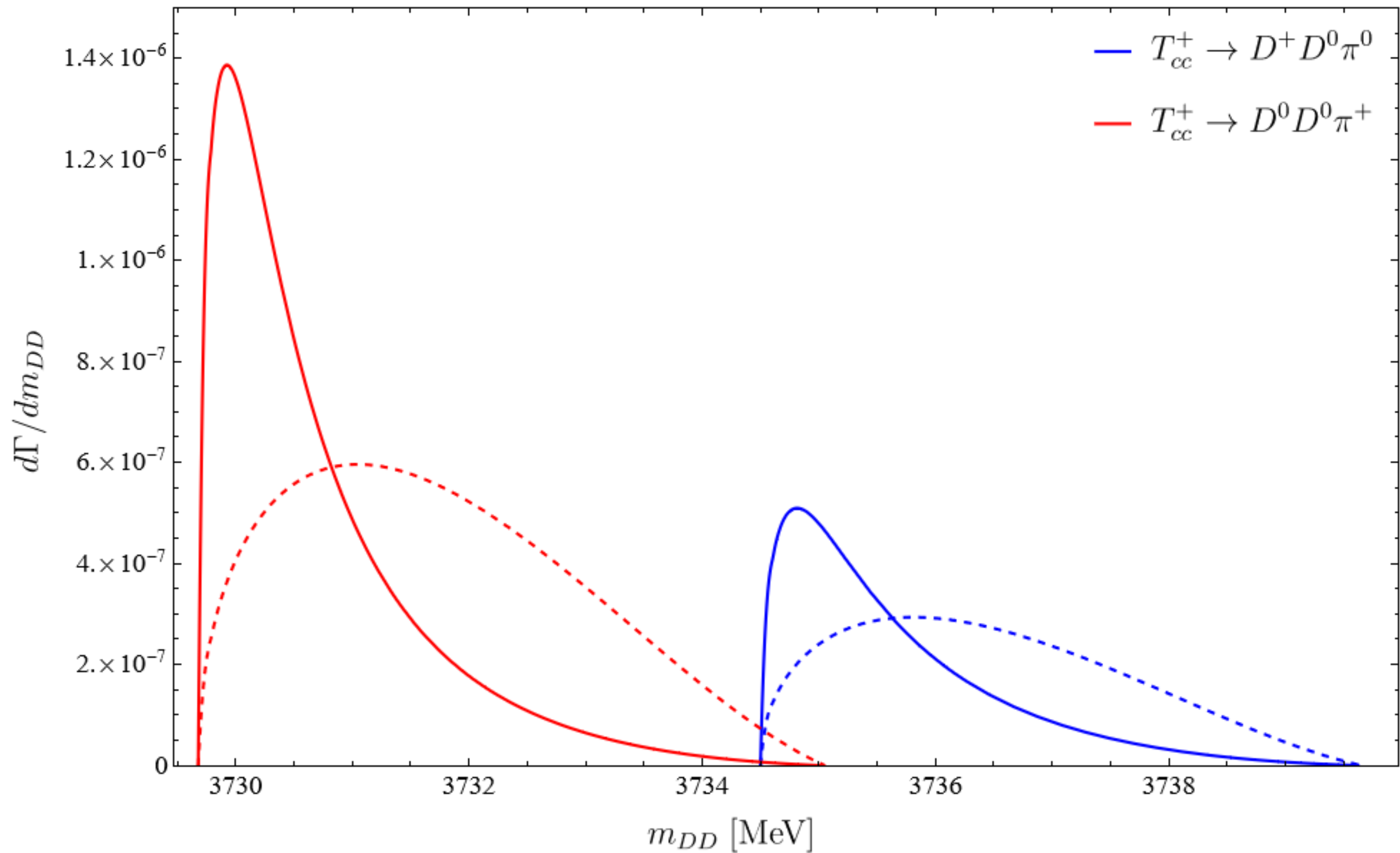
this agrees well with the same plot by Meng et al. (2107.14784)



# $d\Gamma/dm_{DD}$ vs. data

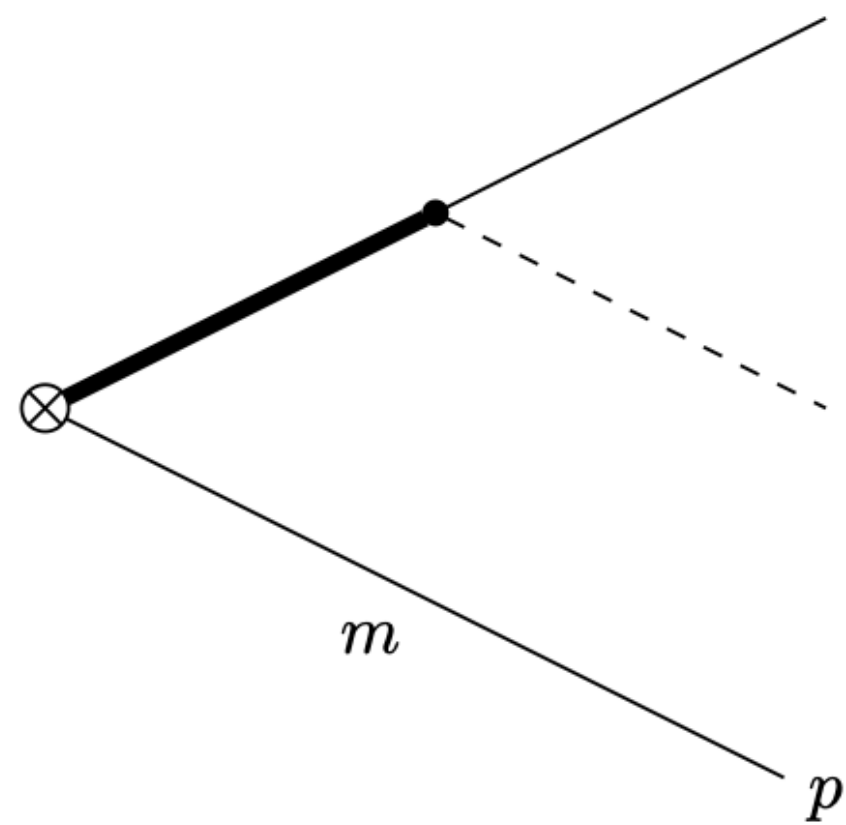


see also M.-L. Du, et.al., Phys. Rev. D, 105 (2022) 1, 014024

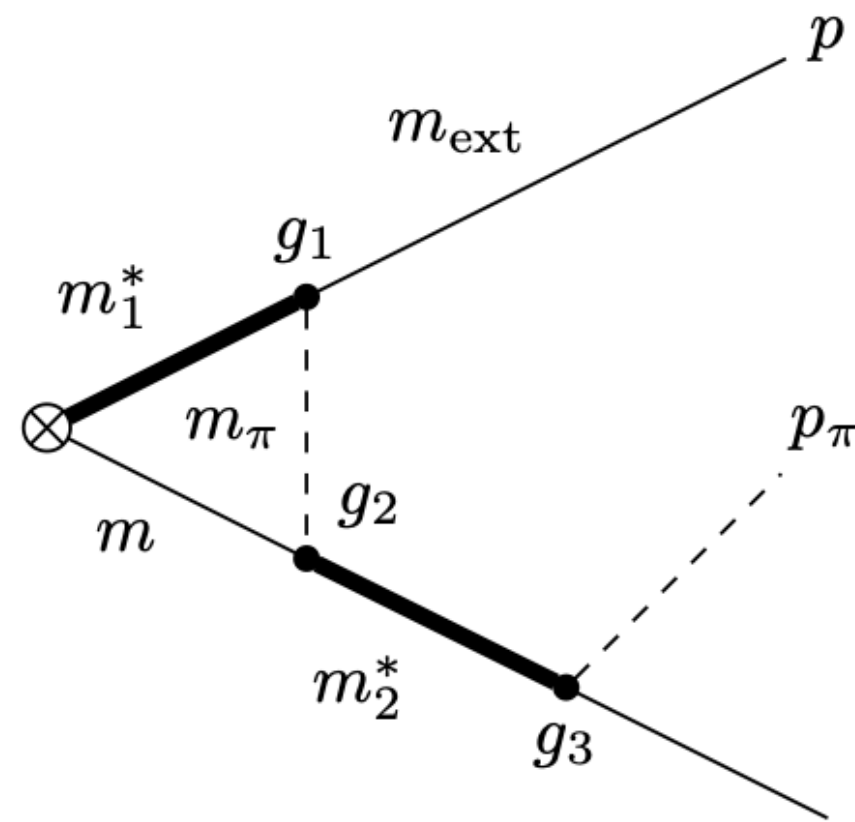


comparison with  $p_\pi^2 \times$  phase space (dashed)

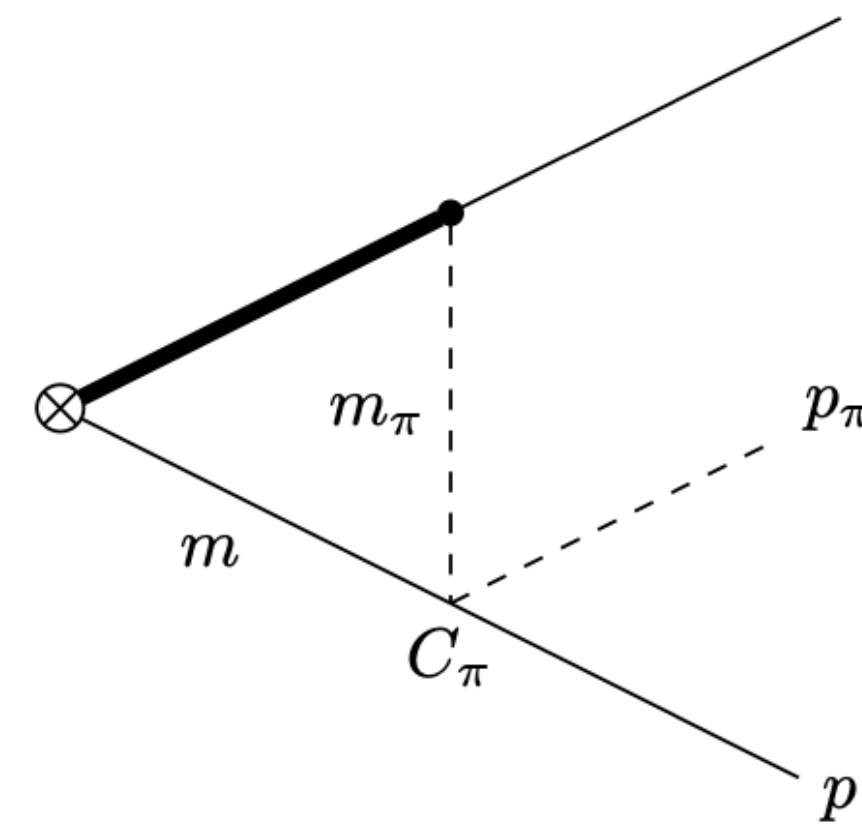
# NLO Corrections to Decay Rate - assume $l = 0$ state



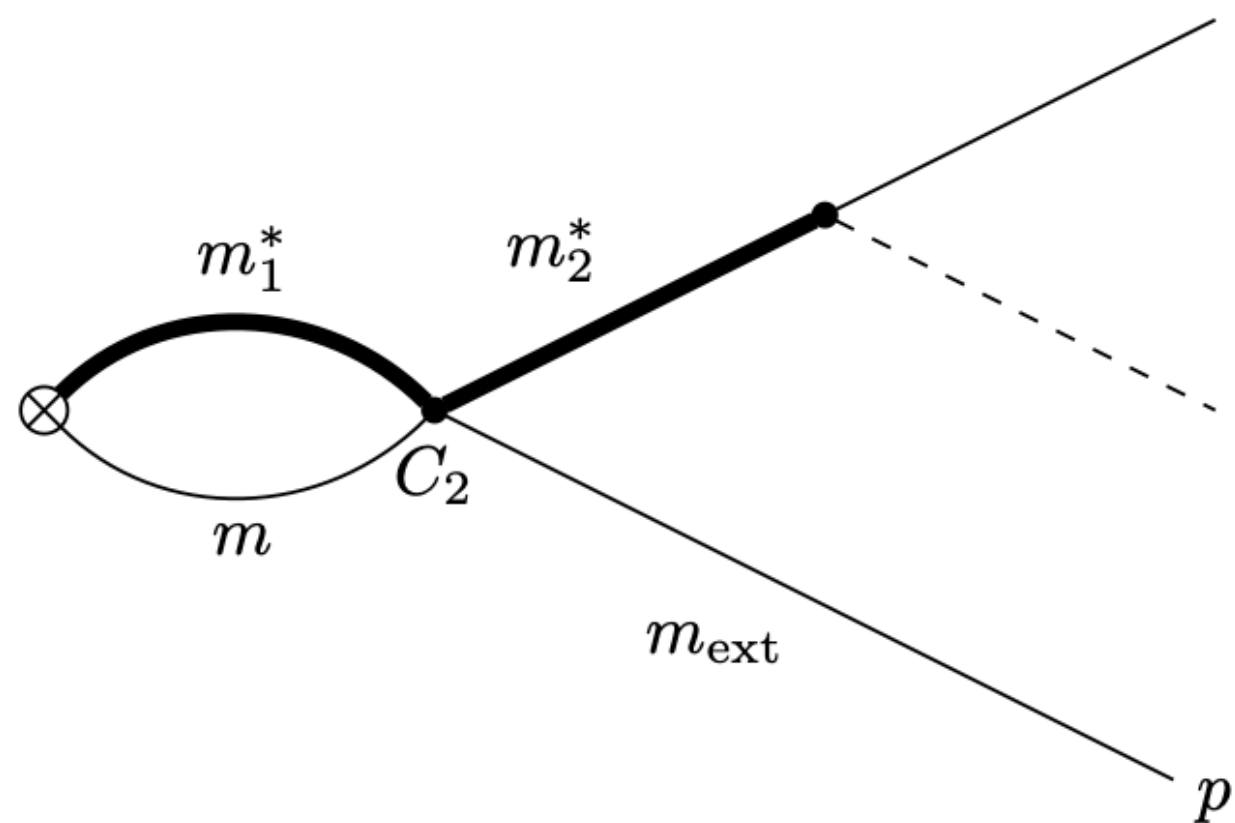
(a)



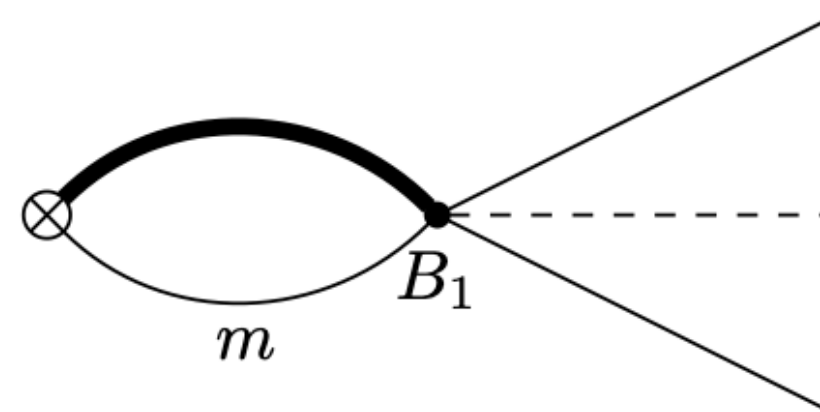
(b)



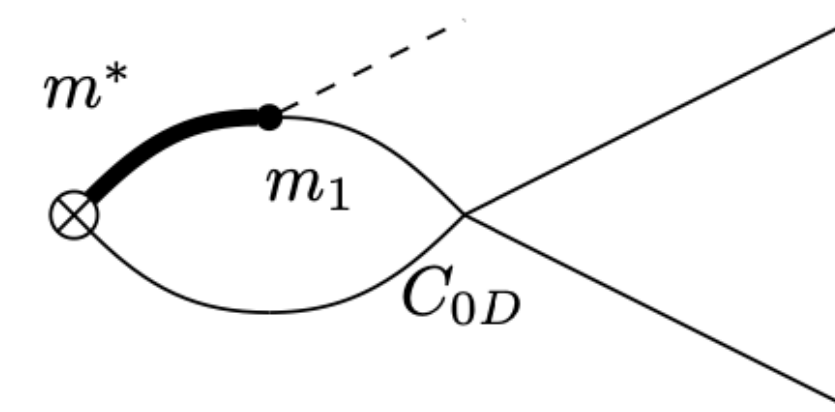
(c)



(d)

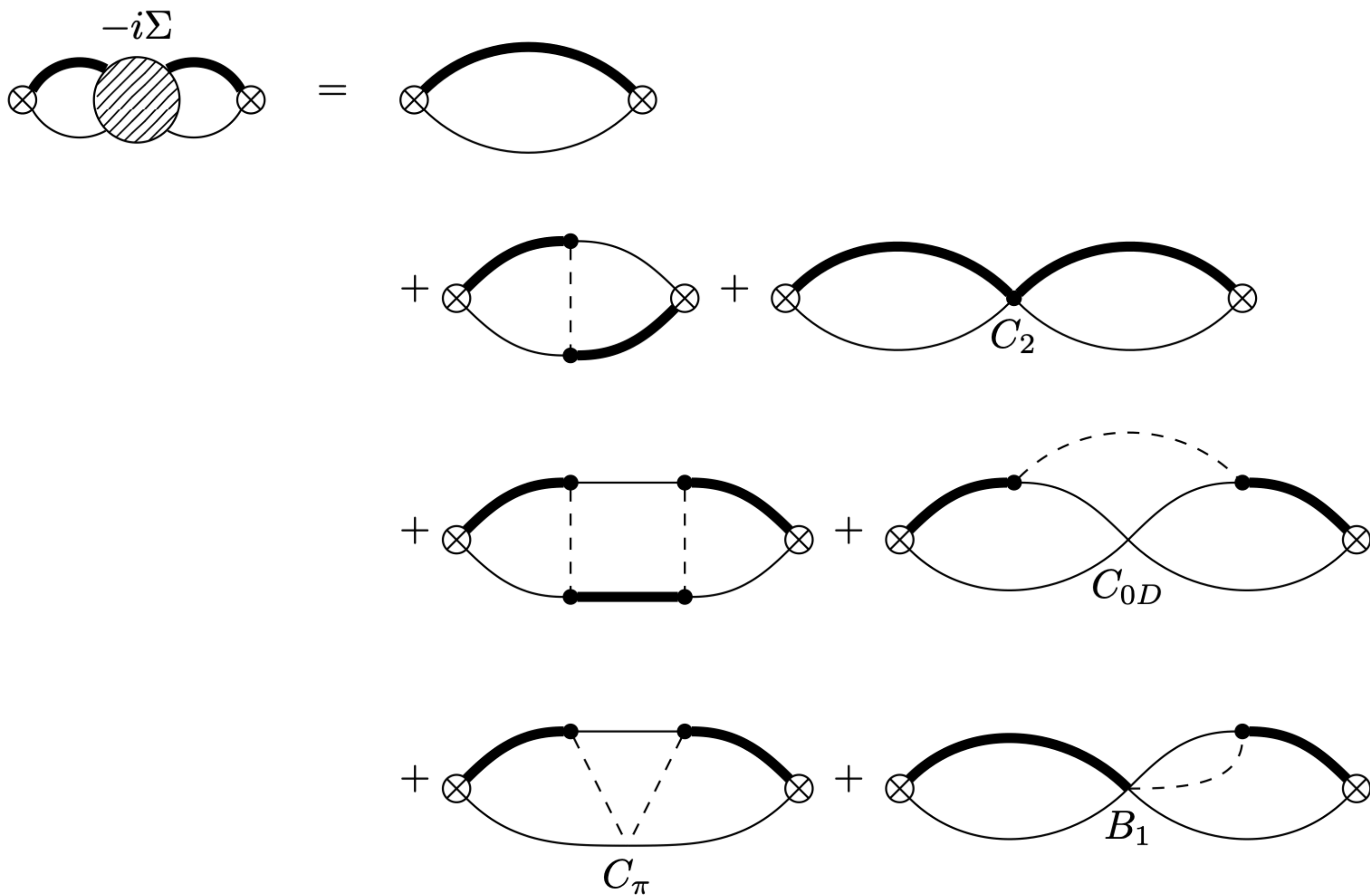


(e)



(f)

$$\Gamma_0 \approx \Gamma^{LO} \left( 1 - \frac{\text{Re } \Sigma_0'^{NLO}(-E_T)}{\text{Re tr } \Sigma'^{LO}(-E_T)} \right) + \frac{2 \text{Im } \Sigma_0^{NLO}(-E_T)}{\text{Re tr } \Sigma'^{LO}(-E_T)}$$



# NLO Corrections to Decay Rate

TABLE I: Partial and total widths in units of keV at LO and NLO.

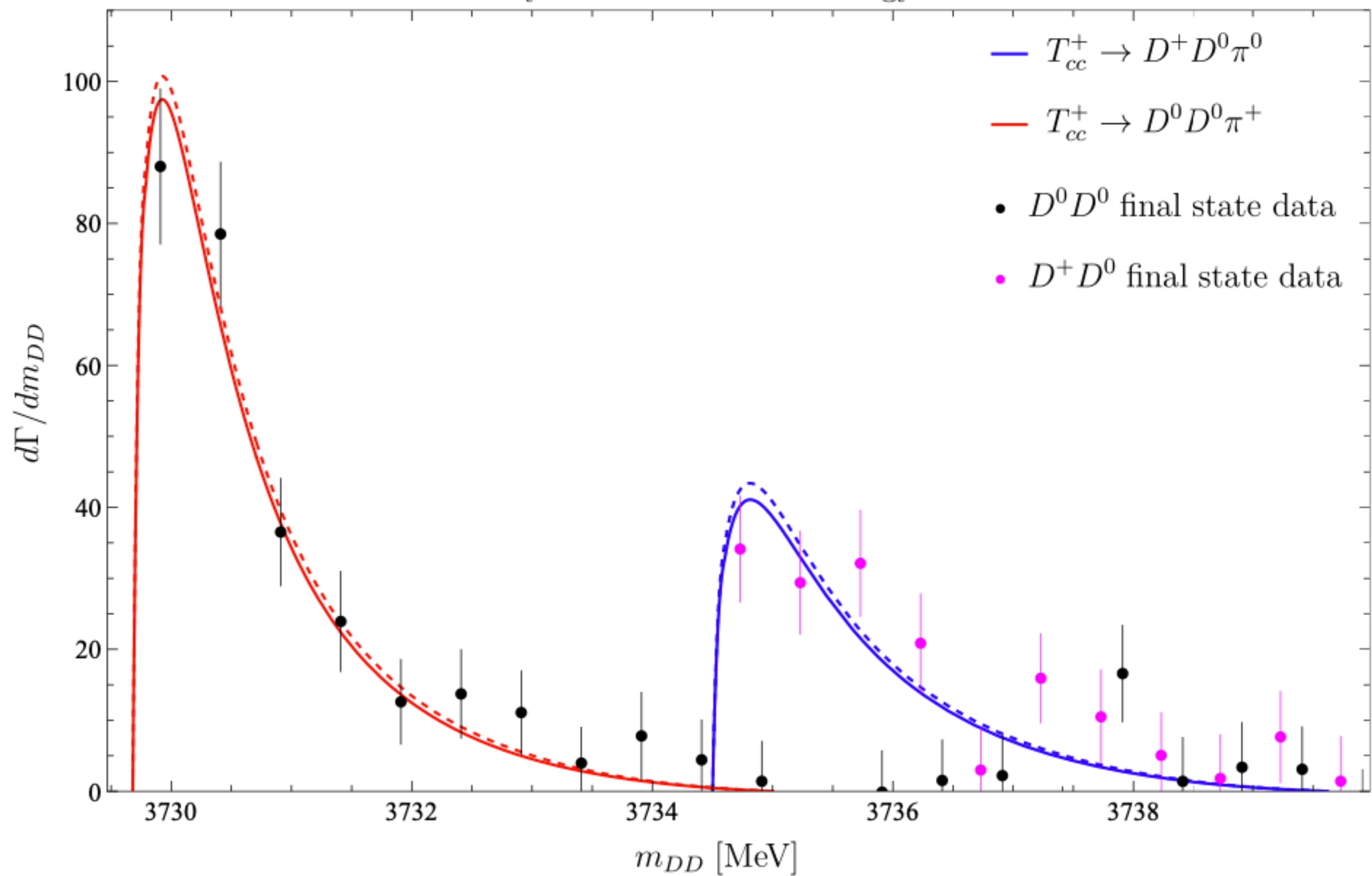
	LO result	NLO lower bound	NLO upper bound
$\Gamma[T_{cc}^+ \rightarrow D^0 D^0 \pi^+]$	28	21	44
$\Gamma[T_{cc}^+ \rightarrow D^+ D^0 \pi^0]$	13	7.8	21
$\Gamma_{\text{strong}}[T_{cc}^+]$	41	29	66
$\Gamma_{\text{strong}}[T_{cc}^+] + \Gamma_{\text{EM}}^{LO}[T_{cc}^+]$	47	35	72

$$\delta m_{pole} = -360 \pm 40_{-0}^{+4} \text{ keV} ,$$

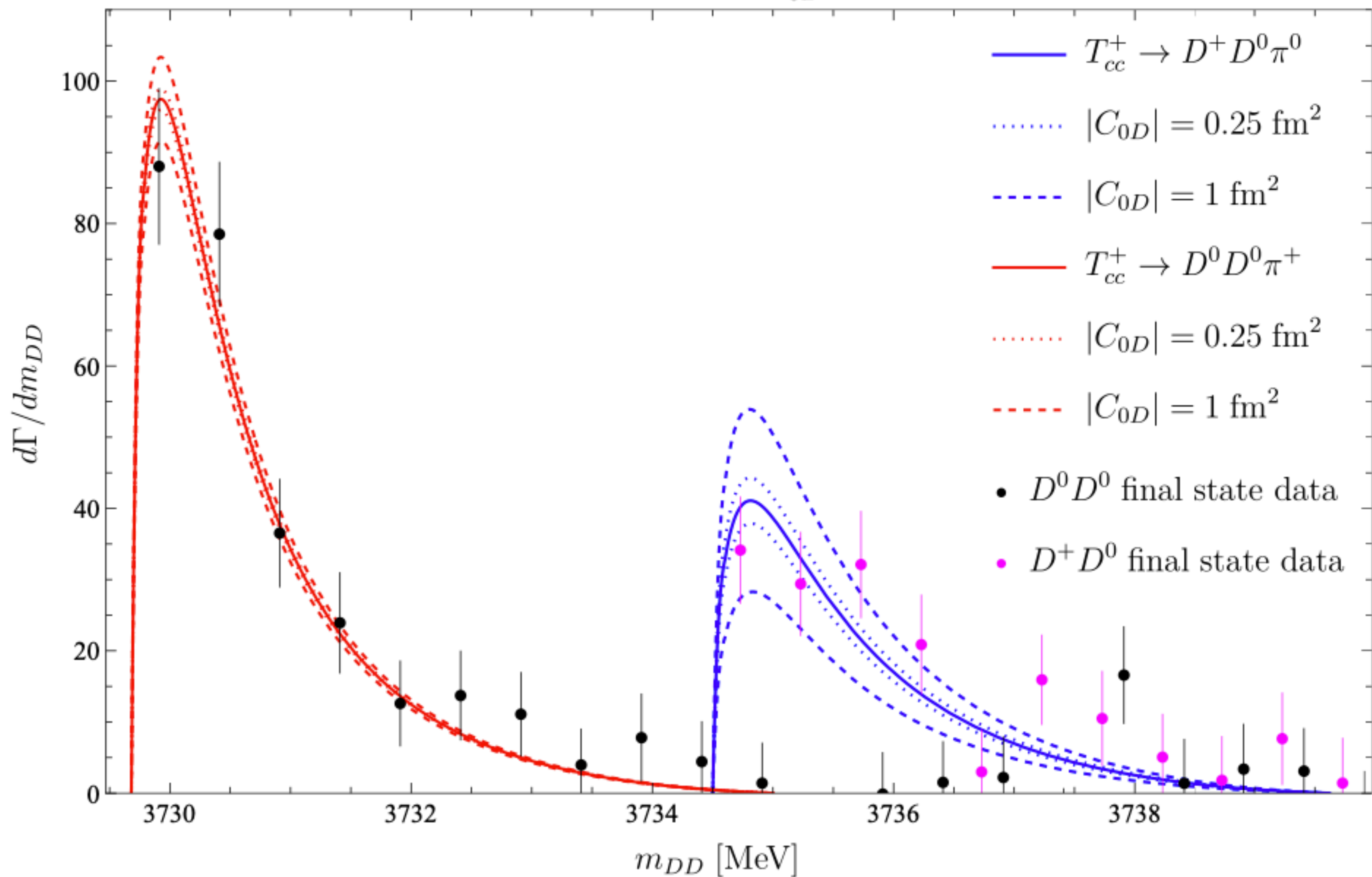
$$\Gamma_{pole} = 48 \pm 2_{-14}^{+0} \text{ keV} . \quad (2)$$

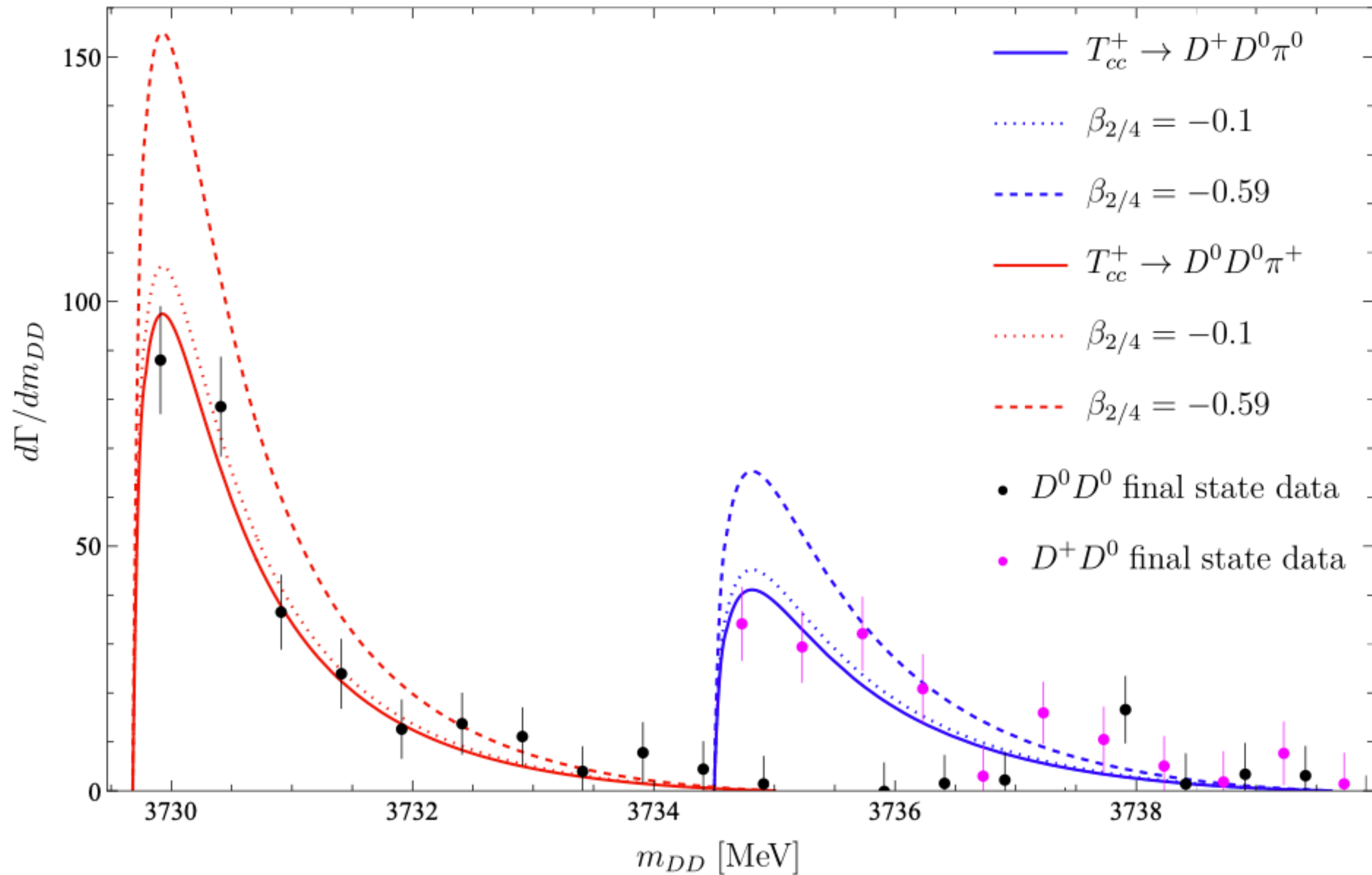


Non – analytic and NLO self – energy contributions

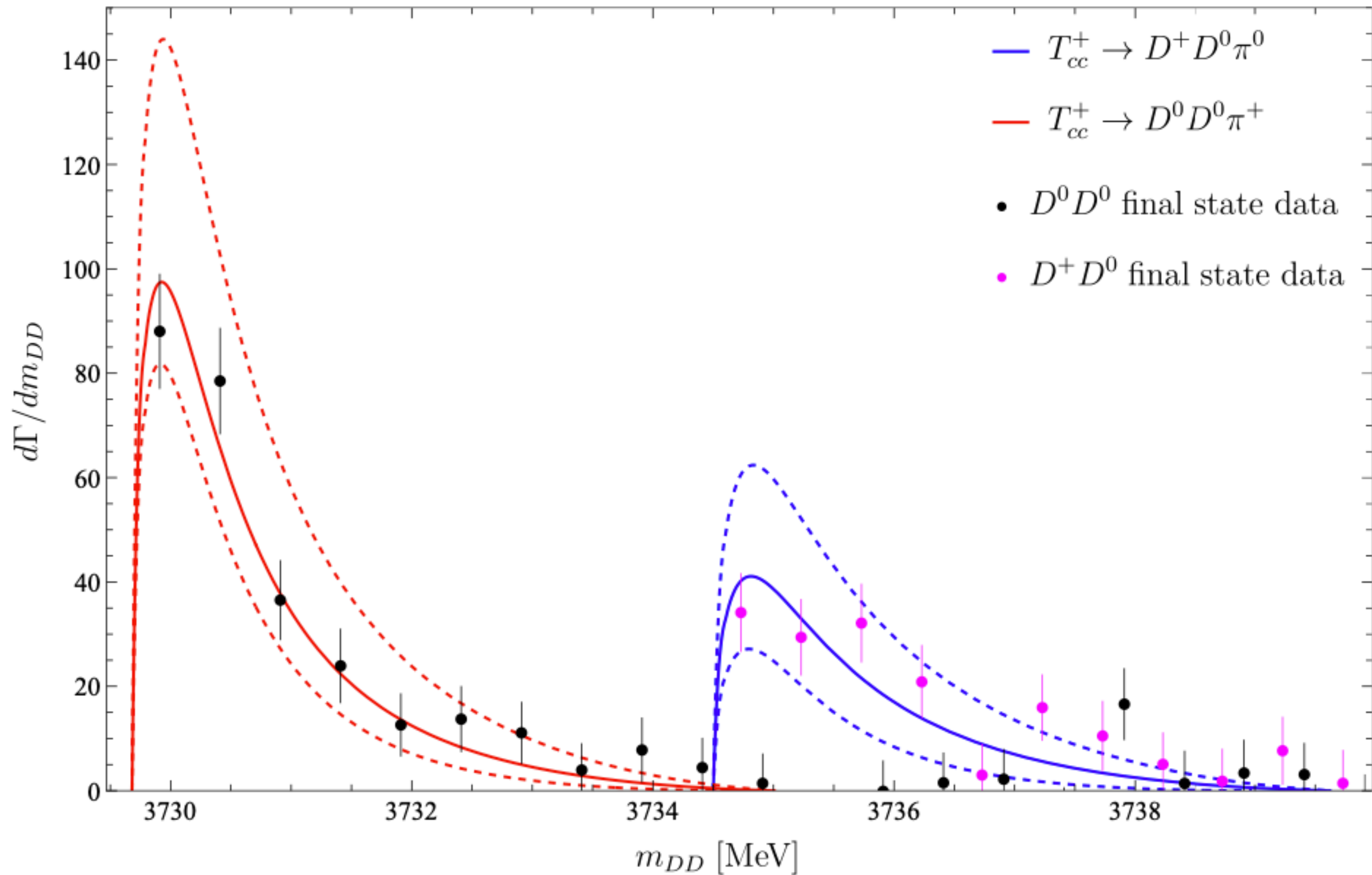


### Contributions from $C_{0D}$ interactions



Contributions from  $\beta_2$  and  $\beta_4$ 

## All NLO contributions



# Summary

$T_{cc}^+$  decays described by XEFT-like theory w/ coupled channels

$\Gamma[T_{cc}^+], d\Gamma[T_{cc}^+]/dm_{DD}$  excellent agreement w/ LO EFT calculations

NLO corrections: pion loops

negligible

$\pi D$  rescattering

negligible

range corrections

dominant

$DD$  rescattering

significant

NLO calculations agree with data, could constrain D interactions