## Masses of D–mesons and narrow Υ states

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# <span id="page-2-0"></span>VEPP–4M collider



VEPP-4(M) started to operate in 1981



Beam energy determination with resonant depolarization technique:

 $\bullet$  Touschek polarimeter (intrabeam scattering in dedicated runs),  $E < 2$  GeV Instant measurement accuracy  $\simeq 1 \times 10^{-6}$ Energy interpolation accuracy  $(5 \div 15) \times 10^{-6}$   $(10 \div 60 \text{ keV})$ 

Laser polarimeter (polarized light scattering asymmetry). At 4.73 GeV statistical accuracy  $\simeq 3 \times 10^{-6}$  / 15 minutes correctable systematic uncertainty 3  $\times$  10<sup>-6</sup> (30 keV)

# KEDR detector



Vacuum chamber Vertex detector Drift chamber Aerogel threshold counters 5 ToF–counters Liquid krypton calorimeter Superconducting coil (0.6 T) Magnet yoke Muon tubes CsI-calorimeter Compensation solenoid <sup>2</sup> VEPP-4M quadrupole

- **Q** Luminosity monitoring by single Bremsstrahlung in  $e^+$ and e<sup>−</sup> directions
- **O** Scattering electron tagging system for two-photon studies

#### <span id="page-4-0"></span>Current status D masses



Main goal of new KEDR experiment was to improve accuracy of  $D^+$  mass.

#### <span id="page-5-0"></span>Data set and analysis method

- $\phi \psi(3770) \rightarrow D\overline{D}$  with reconstruction of one D meson as in MARK expriments at SPEAR and the previous KEDR experiment
- In addition to 0.9  $pb^{-1}$  in 2004 (KEDR-2010), 4 pb $^{-1}$  in 2016–2017 at  ${\sim}1$  MeV vicinity of  $\psi(3770)$  peak
- Resonant depolarization method for beam energy determination
- Decays  $D^0 \to K^-\pi^+$ ,  $D^+ \to K^-\pi^+\pi^-$  and c.c. were reconstructed
- Variables

$$
M_{bc} = \sqrt{\left(\frac{W}{2}\right)^2 - \left(\sum_i \vec{p}_i\right)^2}, \quad \Delta E = \sum_i \sqrt{(m_i^2 + p_i^2)} - E_{beam}
$$

$$
\sigma^2(M_{bc}) = \frac{\sigma_W^2}{4} + \left(\frac{p_D}{M_D}\right)^2 \sigma_{p_D}^2 = \frac{\sigma_W^2}{4} + 0.02 \sigma_{p_D}^2, \quad \Delta E \approx 0
$$

• Difference of  $W = 2E_{beam}$  and actual  $D\overline{D}$  mass is accounted in M.C. simulation (ISR, machine energy spread)

# Analysis, data 2016-2017

• Unbinned maximum likelihood fit of  $M_{bc}$  and  $\Delta E$  in selected events with signal and background (*uds*,  $D\overline{D}$ ) PDFs obtained with M.C.

 $D^{0}$ , 169 signal events,  $M_{D^{0}}$ =1864.910±0.288 MeV:



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## Analysis, data 2004

2004 data were reanalyzed to reduce systematic uncertainties

 $D^{0}$ , 84 signal events,  $M_{D^{0}}$ =1865.341 $\pm$ 0.309 MeV:



# Systematic uncertainties (MeV)



# <span id="page-9-0"></span> $D^0\!\!-$  and  $D^\pm\!\!-\!\!$ mass results (preliminary!)

Data set 2004, published in 2010:

 $\mathcal{M}_{D^0} = 1865.300 \pm 0.330 \pm 0.230$   $M_{D^+} = 1869.530 \pm 0.490 \pm 0.200$  MeV Data set 2004<sup>.</sup>

 $\mathcal{M}_{D^0} = 1865.341 \pm 0.309 \pm 0.061$   $M_{D^+} = 1869.487 \pm 0.490 \pm 0.110$  MeV Data set 2016:

 $\mathcal{M}_{D^0} = 1864.910 \pm 0.288 \pm 0.061$   $M_{D^+} = 1869.587 \pm 0.357 \pm 0.078$  MeV Combined result (should supersede KEDR-2010):



 $M_{D^0} = 1865.110 \pm 0.210 \pm 0.058$   $M_{D^+} = 1869.550 \pm 0.288 \pm 0.066$  MeV



 $D<sup>+</sup>$  mass measurements

# <span id="page-10-0"></span>Revision of results on  $\Upsilon(1S) - \Upsilon(3S)$  masses

#### $T(1S)$  MASS



<sup>1</sup> Reanalysis of BARU 92B and ARTAMONOV 84 using new electron mass (COHEN 87).

#### <span id="page-11-0"></span>Introduction: reasons of reanalysis

Problems:

- Incorrect radiative correction accounting in
	- W. W. MacKay et al., PRD 29(1984),2483 (CUSB@CESR),
	- D. P. Barber et al., PLB 135(1984),498 (ARGUS+CB@DORIS)
- Use of obsolete electron mass value in these two works
- Ignoring of the interference effect in all three measurements
- Discrepancy of MD-1 and CUSB results on  $\Upsilon(1S)$  mass due to difference in calculation of special functions

Goal:

To urge PDG update values of masses as was done for quarkonia electronic widths after J. P. Alexander et al., Nucl.Phys.B 320(1989)45

Why now?

- Preparation to new experiment of KEDR@VEPP-4M with expected accuracy of about 50 keV
	- A.G.Shamov and O.L.Rezanova, Phys.Lett.B 839 (2023) 137766 Experimental aspects of the works were not considered!

# <span id="page-12-0"></span>Analysis of in Mackay et al. [CUSB<sup>∗</sup>@CESR]



Our fit of published CUSB data performed in 1986 with identical radiative corrections accounting gave the mass value 0.375 MeV higher than the published one. Attempts to clarify situation with authors had failed.

?Did a misprint in the data occur?

The data form the journal figure were restored as good as possible, points coincide, unlike the curves.

The mass difference is due to calculation of the resonance curve. We have tried a few independent implementations, the result was stable.

One misprint has been found and fixed in the table of assignment of 22 runs of the experiment to 13 points of the fit. The influence on the mass was negligible.

#### <span id="page-13-0"></span>Correction to radiative corrections

The first published paper on r.c. to the production of narrow resonances: Ya.I.Azimov, A.I.Vainshtein, L.N.Lipatov, V A.Khoze, JETP Lett. 21(1975)172, in a few months a good alternative appeared: M.Greco, G.Pancheri-Srivastava, Y.Srivastava, Nucl.Phys. B101(1975)234 However, the most analysis of  $\psi$  and  $\Upsilon$  before 1985 were performed according to J.D.Jackson and D.L.Scharre, NIM 128(I975)13



The 'radiative gaussian'  $G_R$  was derived with convolution of the gaussian energy spread G and the probability of energy radiation in the approximation of zero resonance width.

The radiation of additional soft photons we accounted in the case (a), but not in (b)

Fig. 1. e<sup>+</sup>e<sup>-</sup> annihilation via one-photon exchange. (a) Lowest order diagram; (b) Higher order diagrams, the top two involving real (soft) photon emission and the next four each involving one additional virtual photon.

 $\sigma(W) \propto G_R(W-M) + \delta_V \cdot G(W-M)$  instead of  $(1+\delta_V) \cdot G_R(W-M)$ Mass shift depends on the energy spread,  $\sim$  100 keV for  $\Upsilon$ 

<span id="page-14-0"></span>For beam energy determination in experiments the resonant depolarization method was employed. It gives the mean Lorentz factor of beam electrons thus the beam energy is proportional to electron's mass  $m_e$ . In 1983 the accuracy of  $m_e$  was about 2.8 ppm, that corresponds to 26 keV uncertainty is  $\Upsilon(15)$  mass.

«The 1986 adjustment of the fundamental physical constants», E.R.Cohen and B.N.Taylor, Rev.Mod.Phys. 59(1987)1121: the value of  $m_e$  was shifted to -8.5 ppm with reduction of uncertainty to 0.3 ppm due to refining of  $e/h$  value

The results from VEPP-4 were recalculated in Phys.Lett. B474(2000)427 The shifts of masses were -80, -85 и -88 keV for  $\Upsilon(1S)$ ,  $\Upsilon(2S)$  и  $\Upsilon(3S)$ , respectively

The results from CESR and DORIS stayed unchanged

#### <span id="page-15-0"></span>Resonance-continuum interference

Ya.l. Azimov et al., JETP Lett. 21(1975)172

contribution of a resonance to a final state  $f$  in soft photon approximation (needs small modifications nowadays):

$$
\sigma^{\Upsilon \to f}(W) = \frac{12\pi}{M^2} \left( 1 + \frac{3}{4}\beta \right) \left[ \frac{\Gamma_{ee}\Gamma_f}{\Gamma M} \ln f(W) - \frac{2\alpha \sqrt{R \Gamma_{ee}\Gamma_f}}{3W} \lambda \text{ Re } \frac{f(W)}{1 - \beta/6} \right]
$$
  
where  $f(W) = \left( \frac{M/2}{M - W - i\Gamma/2} \right)^{1 - \beta}, \quad \beta = \frac{4\alpha}{\pi} \left( \ln \frac{W}{m_e} - \frac{1}{2} \right)$ 

The parameter  $\lambda$  determines the strengths of interference effects,  $\lambda = 1$  for  $f=\mu^+\mu^-$ . For the sum of hadronic modes  $(b_m$  and  $\mathcal{B}_m^{(\mathsf{s})}$  are relative mode probabilities in electromagnetic and strong decays, respectively,  $\phi$  is the interference phase of electromagnetic and strong amplitudes)

<span id="page-15-1"></span>
$$
\lambda = \sqrt{\frac{R\mathcal{B}_{ee}}{\mathcal{B}_h}} + \sqrt{\frac{1}{\mathcal{B}_h}} \sum_m \sqrt{b_m \mathcal{B}_m^{(s)}} \langle \cos \phi_m \rangle . \tag{1}
$$

At the parton model level the strong 3g decays and the electromagnetic  $q\bar{q}$ decays do not interfere thus the sum in the left part of [\(1\)](#page-15-1) must be zero.

T mass shifts grows with the energy spread,  $\sim 100$  keV

<span id="page-16-0"></span> $\Upsilon(15)$  mass (MeV)



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#### <span id="page-17-0"></span>Summary on masses

- o It was demonstrated that the analysis of CUSB/CESR data on mass of  $\Upsilon(1S)$  was not fully correct, the mass was shifted by -0.375 MeV
- When necessary, published mass values were corrected to:
	- Improper radiative correction accounting
	- Use of obsolete electron mass value in these two works
	- Resonance-continuum interference



 $\bullet$  The discrepancy in MD-1 and CUSB results on  $\Upsilon(1S)$  mass reduced from 3.25  $\sigma$  to 1.83  $\sigma$ 

The average  $\Upsilon(1S)$  mass value calculated according PDG rules is  $9460.29 + 0.15$  MeV



#### $9460.40 \pm 0.10$  MeV



<sup>1</sup> Reanalysis of MD1 data using the electron mass from COHEN 1987, the radiative corrections from KURAEV 1985 and interference effects.

<sup>2</sup> Obtained by reanalysing CUSB data (MACKAY 1984), but not authored by the CUSB collaboration.

<sup>3</sup> Reanalysis of BARU 1992B and ARTAMONOV 1984 using new electron mass (COHEN 1987).

- <sup>4</sup> Superseded by SHAMOV 2023.
- <sup>5</sup> Supersedes BARU 1986.
- <sup>6</sup> Superseded by ARTAMONOV 2000.
- <sup>7</sup> Value includes data of ARTAMONOV 1982.
- 8 Reanalysed by SHAMOV 2023.

IMHO: Using of the revised CUSB result into averaging could increase reliability of the mass value (question of accelerator-related uncertainties)

# Thanks for attention!

# <span id="page-20-0"></span> $D^0$ – and  $D^\pm$ –mass backup



Experience of KEDR@VEPP-4M:

"Final analysis of KEDR data on  $J/\psi$  and  $\psi(2S)$  masses" PLB 749(2015)50

- 6  $J/\psi$  and 7  $\psi$ (25) high precision scans in 2002–2008
- systematic uncertainty in one scan  $7 \div 10$  keV (2.5 ppm)
- more than 15 sources of the uncertainty were considered

Difference between  $\psi$  and  $\Upsilon$  conditions:

 $\psi$ : Injection from VEPP-3 to VEPP-4M at the energy of scan point

Υ: Acceleration at VEPP-4M from 1.9 to 4.73 GeV

 $\Upsilon$ : Some systematic uncertainties  $\propto E_{beam}^2$ 

Goals for  $\Upsilon(15)$ :

- $\circ$  Systematic uncertainties  $<$  30 keV (6.3 ppm)
- $\circ$  Statistical uncertainty on mass  $M < 40$  keV
- **O** Statistical uncertainty on electronic width  $\Gamma_{ee} < 1\%$

Luminosity  $\simeq 10$  pb $^{-1}$ ,  $\simeq 200$  runs, optimistic time estimate  $\simeq 2$  months

 $\Upsilon(2S)$ ,  $\Upsilon(3S)$ : much more difficult,  $\Delta M \simeq 100$  keV?

# Forthcoming experiment KEDR@VEPP-4M (2)

Status of preparation:

- Polarization was obtained around  $E_{beam} = 4.73$  GeV
- Short test scan of  $\Upsilon(1S)$  with energy calibrations was done
- Works to improve energy stability of VEPP-4M
- Laser polarimeter is in good operation:



#### Expected systematic uncertainties



Systematic uncertainties in  $I/y$  scans  $(ka)$ .

<sup>∗</sup> — correction uncertainty

#### Resonant depolarization method

The electron beam in the accelerator spontaneously becomes polarized, the spin precesses around guiding field with the frequency

 $\Omega_{\mathsf{spin}} = \omega_{\mathsf{rev}} \left( 1+ \mu'/\mu \, \gamma \right) \,$  depending on the beam energy  $\mathcal{E} = \gamma \, m_e$ External electromagnetic field of variable frequency  $f_d$  depolarizes the beam at the resonance condition

 $\Omega_{spin} = m \cdot \omega_{rev} + n \cdot f_d$  (Υ(1S) at VEPP-4:  $m = 11, n = 1$ )

Measurement  $f_d$  and  $\omega_{\text{rev}}$  at the moment of depolarization allows for the beam energy determination with accuracy  $\sim$ 10<sup>-6</sup>

The moment of depolarization was detected using asymmetry in scattering of longitudinally polarized photons on the transversely polarized electron beam:



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#### VEPP-4M polarimeter

#### Layout of the laser polarimeter:



V. E. Blinov et al., JINST 15 (2020) C08024

 $\Upsilon(2\mathcal{S})$  mass (MeV)



 $\Upsilon(3S)$  mass (MeV)

