

# Current status of the GENEVA event generator and recent improvements

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on behalf of the GENEVA collaboration



# Overview of GENEVA

GENEVA consistently combines three theoretical ingredients of QCD calculations



## 1. fully differential fixed-order calculation

- ▶ up to **NNLO** through  $N$ -jettiness ( $\mathcal{T}_N$ ) subtraction

## 2. higher-order resummation

- ▶ up to **NNLL'** in  $\mathcal{T}_0$  using SCET
- ▶ not limited to SCET nor to  $\mathcal{T}_N$  (we have  $p_T$  at **N<sup>3</sup>LL** via RadISH)

## 3. parton shower, hadronization and MPI

- ▶ provided by a shower MC (currently PYTHIA8 or SHERPA)

⇒ produce a **NNLO+NNLL'+PS** event generator

## Advantages of higher-order resummation

- ▶ link between **NNLO** and **PS**
- ▶ consistently improve perturbative accuracy over full spectrum
- ▶ provide event-by-event systematic estimate of theory uncertainties

# Overview of GENEVA

1. slice the phase space into jet bins using a **resolution parameter** (like  $\mathcal{T}_0$  or  $p_T$ )
2. compute FO + resummed **matched cross section** up to NNLO+NNLL'
3. fill in additional radiation by **matching to a parton shower**, in a way that avoids spoiling the resummed accuracy
4. include **hadronization and MPI**

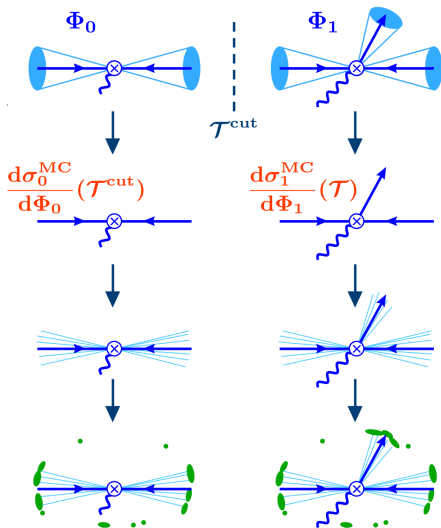


figure by F. Tackmann

## N-jettiness as resolution parameter

Separate phase space  $\Phi$  into jet bins  $\Phi_N$ , then let  $\mathcal{T}_N^{\text{cut}} \rightarrow 0$

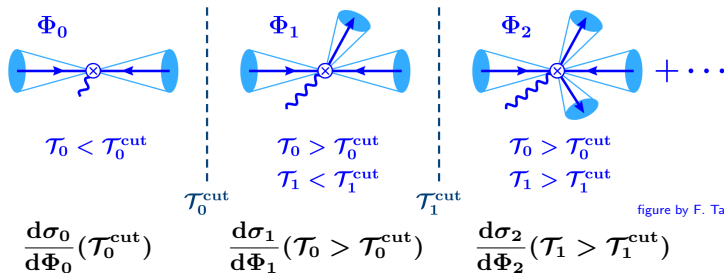


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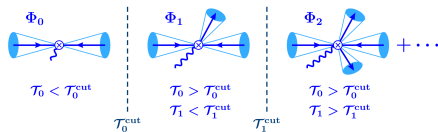
### Advantages of $\mathcal{T}_N$

Most GENEVA applications use  $N$ -jettiness  $\mathcal{T}_N$  [Stewart, Tackmann, Waalewijn '09, '10] as resolution parameter

- ▶ classify jet observables in an **IR-safe** way, can take  $\mathcal{T}_N \rightarrow 0$
- ▶ **resummation** of large  $\mathcal{T}_N$  logarithms **known** up to *at least NNLL'* in SCET for color-singlet final state

other possible res. variable  $p_T$  [Alioli, Bauer, Broggio, Gavardi, Kallweit, Lim, RN, Napoletano, Rottoli 2102.08390]  
 but also  $p_T^{\text{jet}}$ ,  $k_T$ -ness, ...

## Matching fixed-order and resummed calculations



At NNLO needed:

- ▶ 0/1-jet res. parameter
- ▶ 1/2-jet res. parameter

### Jet-bin cross sections (simplified)

$$\begin{aligned}
 \frac{d\sigma_0}{d\Phi_0}(\mathcal{T}_0^{\text{cut}}) &= \frac{d\sigma^{\text{NNLL}'}}{d\Phi_0} + \overbrace{\left. \frac{d\sigma_0^{\text{FO}}}{d\Phi_0} - \frac{d\sigma^{\text{NNLL}'}}{d\Phi_0} \right|_{\alpha_s^2}}^{\text{d}\sigma^{\text{nons}}} \\
 \frac{d\sigma_1}{d\Phi_1}(\mathcal{T}_0 > \mathcal{T}_0^{\text{cut}}) &= \left[ \frac{d\sigma^{\text{NNLL}'}}{d\Phi_0 d\mathcal{T}_0} - \left. \frac{d\sigma^{\text{NNLL}'}}{d\Phi_0 d\mathcal{T}_0} \right|_{\alpha_s^2} \right] \mathcal{P}(\Phi_1) + \frac{d\sigma_1^{\text{FO}}}{d\Phi_1} \\
 \frac{d\sigma_2}{d\Phi_2}(\mathcal{T}_1 > \mathcal{T}_1^{\text{cut}}) &= \left[ \frac{d\sigma^{\text{NNLL}'}}{d\Phi_0 d\mathcal{T}_0} - \left. \frac{d\sigma^{\text{NNLL}'}}{d\Phi_0 d\mathcal{T}_0} \right|_{\alpha_s^2} \right] \mathcal{P}(\tilde{\Phi}_1) \mathcal{P}(\Phi_2) + \frac{d\sigma_2^{\text{FO}}}{d\Phi_2}
 \end{aligned}$$

1. **splitting functions**  $\mathcal{P}$  make resummed piece differential in  $\Phi_N$   
 $\hookrightarrow$  new treatment of functions  $\mathcal{P}$
2. usually also resum 1/2-jet res. variable with  $\mathcal{T}_1$  **NLL Sudakov**

## $\mathcal{T}_0$ resummation in SCET

For color singlet final state, the  $\mathcal{T}_0$  spectrum is factorized in SCET

[Stewart, Tackmann, Waalewijn '09]

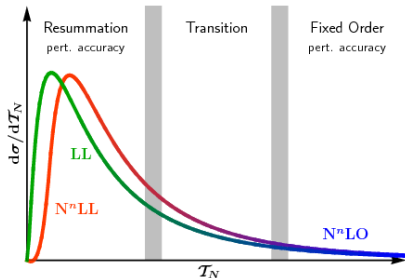
$$\frac{d\sigma}{d\mathcal{T}_0} = \sum_{ij} H_{ij}(\mu) \int dt_a dt_b B_i(t_a, \mu) B_j(t_b, \mu) S_{ij} \left( \mathcal{T}_0 - \frac{t_a + t_b}{Q}, \mu \right)$$

- ▶ **hard function**  $H_{ij}$  is process dependent, at NNLO contains the 2-loop virtual corrections
- ▶ **beam functions**  $B_i$  known up to N<sup>3</sup>LO
- ▶ **soft function**  $S_{ij}$  is process dependent, only perturbative part included

### Resummation

- ▶ compute each at its typical scale  $\mu_H \sim Q$ ,  $\mu_B \sim \sqrt{Q\mathcal{T}_0}$  and  $\mu_S \sim \mathcal{T}_0$
- ▶ resum by **RGE evolution**  $\mu_i \rightarrow \mu$
- ▶ **New:** we take the resummed results from `scetlib` [Ebert, Michel, Tackmann, et al.]
- ▶ similar factorization formulae for resummation in other variables (like  $p_T$ )

## $\mathcal{T}_0$ spectrum and profile scales



- ▶ **peak region**: nonsingulars are power-suppressed; pert. accuracy determined by **resummation**
- ▶ **transition region**: smooth connection; scale choice generates ambiguity
- ▶ **tail region**: power expansion breaks down, must turn off resummation; pert. accuracy determined by **fixed order**

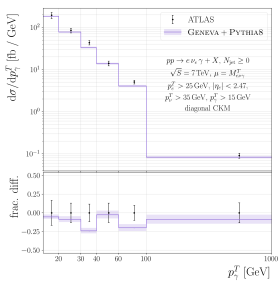
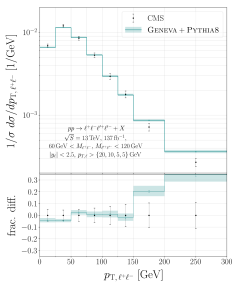
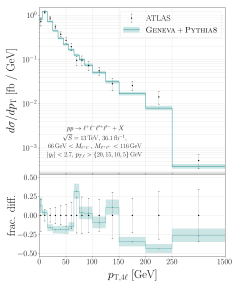
### Profile scales

This is achieved by appropriate choice of  $\mathcal{T}_N$ -dependent **profile scales**  $\mu_i(\mathcal{T}_N)$  so that while approaching the FO region, all scales  $\mu_i \rightarrow \mu_{FO}$

- ▶ profile scale variations **isolate different sources of pert. uncertainties**
- ▶ total theoretical uncertainty **combining FO and resummation** scale variations
- ▶ cross section evaluated at all sets of scales, so that pert. uncertainty are **provided event-by-event**

# List of available NNLO processes in GENEVA

- ▶ Drell-Yan ( $\mathcal{T}_0$ ) [Alioli, Bauer, Berggren, Tackmann, Walsch 1508.01475]
- ▶ Higgsstrahlung [Alioli, Broggio, Kallweit, Lim, Rottoli 1909.02026]
- ▶ Higgs decays ( $gg, b\bar{b}$ ) [Alioli, Broggio, Gavardi, Kallweit, Lim, RN, Napoletano, Rottoli 2009.13533]
- ▶  $\gamma\gamma$  [Alioli, Broggio, Gavardi, Kallweit, Lim, RN, Napoletano, Rottoli 2010.10498]
- ▶ Drell-Yan ( $p_T$ ) [Alioli, Bauer, Broggio, Gavardi, Kallweit, Lim, RN, Napoletano, Rottoli 2102.08390]
- ▶  $ZZ$  [Alioli, Broggio, Gavardi, Kallweit, Lim, RN, Napoletano, Rottoli 2103.01214]
- ▶  $W^\pm\gamma$  [Cridge, Lim, RN 2105.13214]
- ▶ **New:**  $gg \rightarrow HH$  [Alioli, Billis, Broggio, Gavardi, Kallweit, Lim, Marinelli, RN, Napoletano 2212.10489]
- ▶ **New:**  $gg \rightarrow H$  [Alioli, Billis, Broggio, Gavardi, Kallweit, Lim, Marinelli, RN, Napoletano 2301.11875]

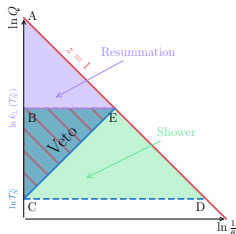




# Parton shower interface

## Two challenges

- ▶ avoid impacting pert. accuracy of  $\mathcal{T}_0$  spectrum
- ▶ avoid double counting between resummation in  $\mathcal{T}_0$  and PS (usually not  $\mathcal{T}$ -ordered)



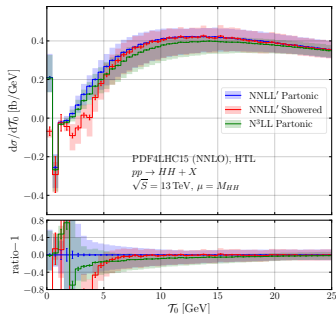
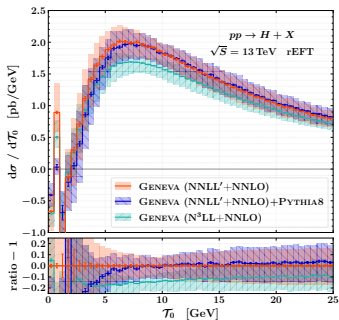
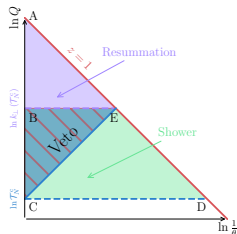
## Matching to a parton shower in GENEVA

- ▶ impose veto **discarding showered events with  $\mathcal{T}_N > \mathcal{T}_N^{\text{cut}}$**
- ▶ all  $\Phi_0$  events are showered with the only requirement  $\mathcal{T}_0 < \mathcal{T}_0^{\text{cut}}$
- ▶ **first shower emission  $\Phi_1 \rightarrow \Phi_2$**  performed by hand in GENEVA using  $\mathcal{T}_0$ -preserving mapping
- ▶ starting from  $\Phi_2$  events, PS affects  $\mathcal{T}_0$  spectrum only beyond NNLL'

# Parton shower interface

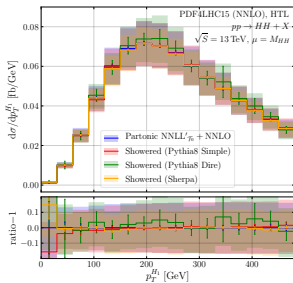
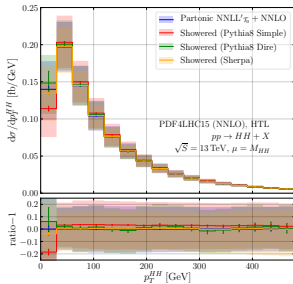
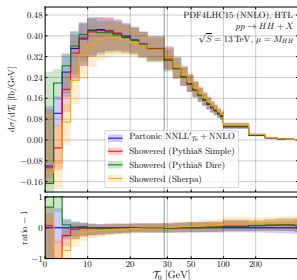
## Two challenges

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# New: interface with different parton shower MCs

- ▶ study effects of different **PS evolution variables**
- ▶ estimate shower matching uncertainty
- ▶ currently possible to interface
  - ▶ PYTHIA8
  - ▶ **New:** DIRE (in PYTHIA8)
  - ▶ **New:** SHERPA
- ▶ applied to  $gg \rightarrow HH$  process

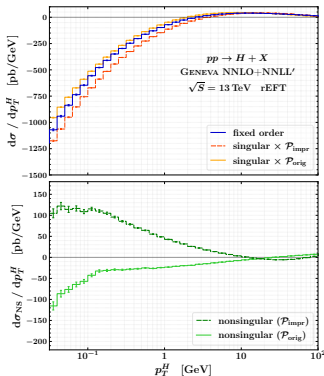
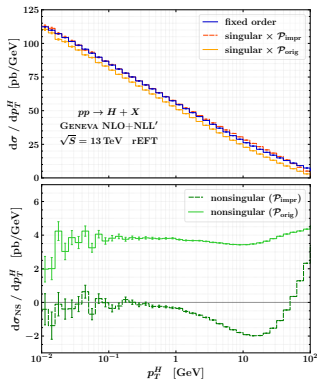


# Improved treatment of splitting functions $\mathcal{P}$

Reproduce higher-multiplicity phase space

$$\frac{d\sigma_1}{d\Phi_1}(\tau_0 > \tau_0^{\text{cut}}) = \left[ \frac{d\sigma^{\text{NNLL}'}}{d\Phi_0 d\tau_0} - \frac{d\sigma^{\text{NNLL}'}}{d\Phi_0 d\tau_0} \Big|_{\alpha_s^2} \right] \mathcal{P}(\Phi_1) + \frac{d\sigma_1^{\text{FO}}}{d\Phi_1}$$

- ▶ soft and collinear limit in  $\mathcal{P}_{0 \rightarrow 1} \Rightarrow$  now correctly reproduced
- ▶ soft and collinear limit in  $\mathcal{P}_{1 \rightarrow 2} \Rightarrow$  improved, now only miss single log
- ▶ **New:** tested on  $gg \rightarrow H$  and Drell-Yan

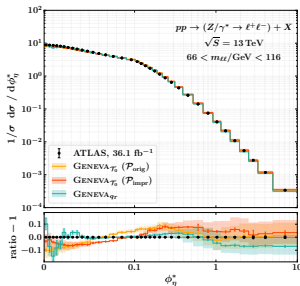
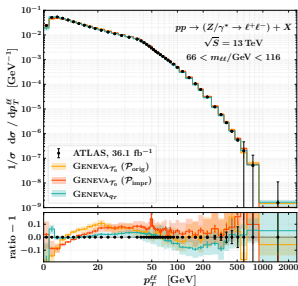


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# Complex scale prescription

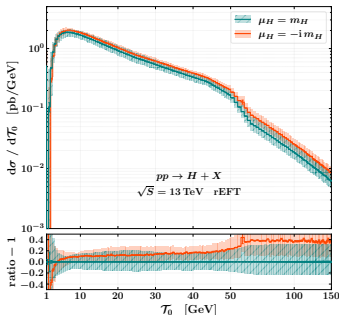
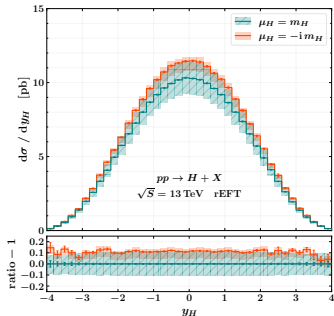
- ▶ allow choice of **complex hard scale**  $\mu_H$  in resummed cross section

$$\mu_H = Q \quad \rightarrow \quad \mu_H = -iQ$$

- ▶ include **additional source of theory uncertainty**  $\phi$

$$\mu_H = Q e^{i\phi}, \quad \phi \in [\pi/4, 3\pi/4]$$

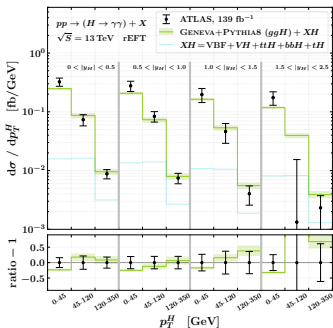
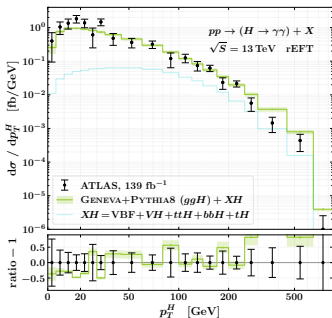
- ▶ “**resumming**” **large  $\pi^2$  terms** arising from  $\log^k(-q^2/\mu^2)$ , effectively reducing overall theory uncertainty
- ▶ **New:** implemented for  $gg \rightarrow H$  production



# Comparison to ATLAS and CMS Higgs data

Comparing to latest ATLAS [CERN-EP-2021-227] and CMS [CERN-EP-2022-142] data sets for Higgs differential cross sections in the **diphoton decay channel**

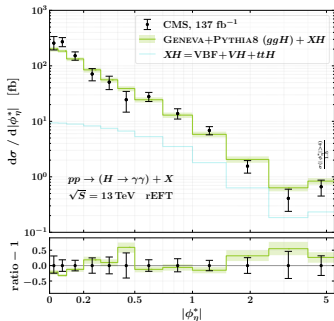
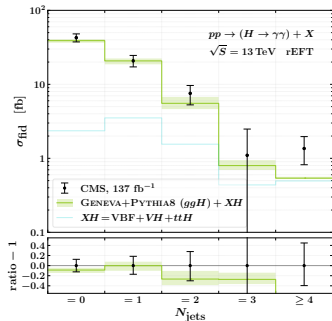
- ▶  $gg \rightarrow H$  computed at **NNLO+NNLL'<sub>0</sub>+NLL<sub>1</sub>+PS** in GENEVA, in the **rescaled EFT scheme**  
 $\hookrightarrow m_t \rightarrow \infty$ , rescaled by exact overall LO  $m_t$  dependence
- ▶  $H \rightarrow \gamma\gamma$  inserted by PYTHIA8 at LO in QCD
- ▶ include VBF,  $VH$ ,  $ttH$ , etc. taken from experimental simulations
- ▶ **New:** include **7 point scale variations**, nondiagonal in  $\{\mu_R, \mu_F\}$



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# WIP: $pp \rightarrow Z + j$ at NNLO in GENEVA

## 1-jet final state

- ▶  $Z + j$  needs **process-defining cuts** (e.g.  $\mathcal{T}_0$  or  $p_T$  cut)
- ▶ relevant  $\Phi_1/\Phi_2$  resolution variable is  $\mathcal{T}_1$
- ▶ at NNLO need a  $\mathcal{T}_1$ -preserving **phase-space mapping**  $\Rightarrow$  **still WIP!**

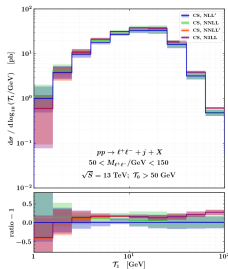
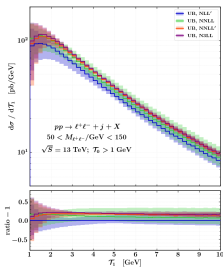
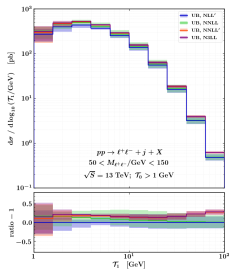
## $Z + j$ in GENEVA

- ▶ only  $\mathcal{T}_1$  spectra computed, with **resummation up to N<sup>3</sup>LL**
- ▶  $\mathcal{T}_1$  definition introduces **frame dependence**
  - ▶ underlying Born frame (**UB**)
  - ▶ colour singlet frame (**CS**)
  - ▶ laboratory frame (**LAB**)
- ▶ factorization formula in SCET (showing explicit frame dependence)

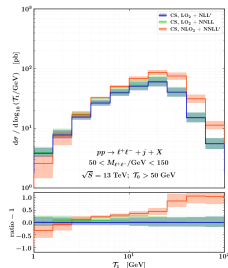
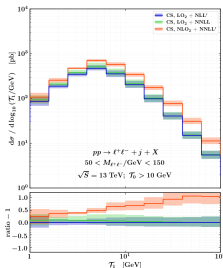
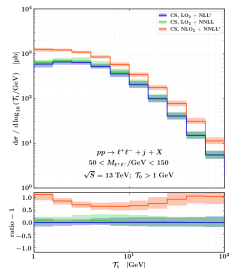
$$\frac{d\sigma}{d\mathcal{T}_1} = \sum_{\kappa} H_{\kappa} \int dt_a dt_b ds_J B_{\kappa_a}(t_a) B_{\kappa_b}(t_b) J_{\kappa_J}(s_J) \\ \times S_{\kappa} \left( n_{a,b} \cdot n_J, \mathcal{T}_1 - \frac{t_a}{Q_a} - \frac{t_b}{Q_b} - \frac{s_J}{Q_J} \right)$$

# WIP: Resummed and matched $\mathcal{T}_1$ spectrum

## Resummed spectra:



## Matched results:



## More processes

- ▶  $Z + j$  ( $\mathcal{T}_1$ -preserving  $1 \rightarrow 2$  mapping is WIP)
- ▶  $t\bar{t}$ ,  $t\bar{t}V$  ( $\mathcal{T}_0$  resummation for  $t\bar{t}$  already done [[Alioli, Broggio, Lim 2111.03632](#)])
- ▶  $Z\gamma$  (...  $W\gamma$  already there)
- ▶  $gg \rightarrow (H)H$  with exact  $m_t$  dependence

## More features

- ▶ implement different 0-jet resolution variable, e.g.  $p_T$ , for more processes
- ▶ include EW corrections to Drell-Yan
- ▶ revamp  $e^+e^- \rightarrow X$  generator

## Conclusions

- ▶ **GENEVA** performs matching of NNLO calculations with higher-order resummation and parton shower MCs
- ▶ versatile framework:
  - ▶ freedom to choose jet-resolution variables ( $\mathcal{T}_0$ ,  $p_T$  mappings tested)
  - ▶ freedom to choose resummation formalism (SCET, RadISH implemented)
  - ▶ interface to **different parton showers** (PYTHIA8, DIRE, SHERPA already in)
- ▶ implemented for several processes, **latest:  $gg \rightarrow H, HH$**
- ▶ covered a few technical improvements:
  - ▶ **better accuracy** of differential distributions **in the soft & collinear limits**
  - ▶ consolidated theory uncertainty by allowing **7-point scale variations**
  - ▶ included  $\pi^2$  **resummation** through new scetlib interface
- ▶ first steps towards NNLO+NNLL'+PS for **coloured final states**: stay tuned!

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Thank you!

## Backup: $N$ -jettiness

### General definition

For a phase space point with  $M$  final state particles ( $M > N$ )

$$\mathcal{T}_N(\Phi_M) = \sum_k \min\{\hat{q}_a \cdot p_k, \hat{q}_b \cdot p_k, \hat{q}_1 \cdot p_k, \dots, \hat{q}_N \cdot p_k\}$$

where  $\hat{q}_{a,b}$  are **beam directions** and  $\hat{q}_i$  are **jet direction** 4-vectors

[Stewart, Tackmann, Waalewijn '09, '10]

- ▶ physical meaning: as  $\mathcal{T}_N \rightarrow 0$  the final state looks more like a  $N$ -jet final state, i.e. **unresolved emissions** are **either soft or collinear** to one of the beam or jet directions
- ▶  $\mathcal{T}_0$  definition corresponds to **beam thrust**, happens to be boost invariant

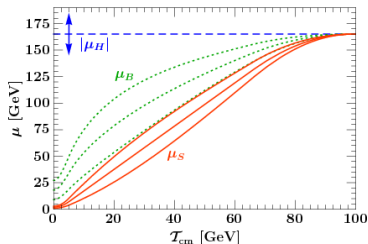
$$\mathcal{T}_0 = \sum_k |p_k^\perp| e^{-|\eta_k - Y|}$$

- ▶  $\mathcal{T}_1$  is **not boost invariant**, frame dependence through  $Q_a, Q_b, Q_J$

$$\mathcal{T}_1 = \sum_k \min \left\{ \frac{2q_a \cdot p_k}{Q_a}, \frac{2q_b \cdot p_k}{Q_b}, \frac{2q_J \cdot p_k}{Q_J} \right\}$$

## Backup: Profile scales

- ▶  $\mathcal{T}_0$  resummation is performed by RGE-evolving each piece  $H$ ,  $B_i$ ,  $S$  from their typical scales  $\mu_H$ ,  $\mu_B$ ,  $\mu_S$  to the common  $\mu$
- ▶ when  $\sigma^{\text{nons}} \sim \sigma^{\text{sing}}$ , resummation must be switched off



- ▶ achieve this via smooth  $\mathcal{T}_0$ -dependent profile scales

$$\mu_H(\mathcal{T}_0) = \mu_{\text{FO}}$$

$$\mu_B(\mathcal{T}_0) = \mu_{\text{FO}} \sqrt{f(\mathcal{T}_0/Q)}$$

$$\mu_S(\mathcal{T}_0) = \mu_{\text{FO}} f(\mathcal{T}_0/Q)$$

- ▶ for  $\mathcal{T}_0$  in the resummation region  $f(\mathcal{T}_0/Q) = \mathcal{T}_0/Q$
- ▶ for  $\mathcal{T}_0$  in the fixed-order region  $f(\mathcal{T}_0/Q) = 1$

$$\Rightarrow \mu_H = \mu_B = \mu_S = \mu_{\text{FO}}$$

- ▶ in nonperturbative region  $f(\mathcal{T}_0/Q)$  freezes all scales to a minimum value
- ▶ scale uncertainty estimated by independently varying  $\mu_{\text{FO}}$ ,  $\mu_S$ ,  $\mu_B$ , and varying the functions  $f(\mathcal{T}_0/Q)$  themselves

## Backup: Ingredients of $N^3LL \mathcal{T}_1$ resummation

### Hard, beam, jet, and soft functions

- ▶ **hard function** known analytically up to two loops [Gehrmann, Tancredi, et al. '11, '22]
  - ▶ at NNLL' include one-loop squared  $gg \rightarrow Zg$
  - ▶ **beam and jet functions** known at  $N^3LO$  [Ebert, Mistlberger, Vita '20]  
[Bruser, Liu, Stahlhofen '18]
  - ▶ only NNLO beam [Gaunt, Stahlhofen, Tackmann '14] and jet [Becher, Neubert '06; Becher, Bell '11] are needed for our purpose
  - ▶ **soft function** at NLO implemented as on-the-fly integrals using [Jouttenus, Stewart, Tackmann, Waalewijn '13]
  - ▶ NNLO boundary terms provided by interpolation over SoftSERVE grids [Bell, Rahn, Talbert '18, '20]
- 
- ▶ NLO results validated in all frames
  - ▶ NNLO results validated in underlying Born frame against MCFM [Campbell, Ellis, Mondini, Williams '17]; for other frames ours are first results