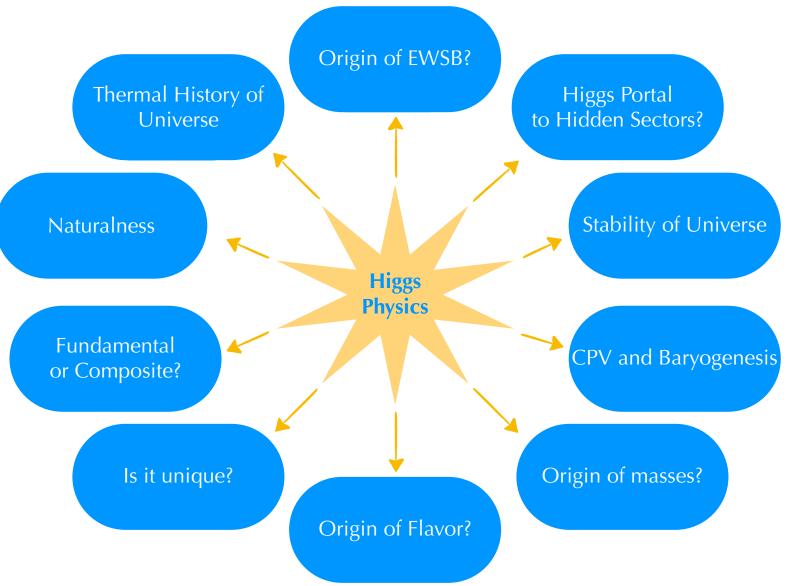
# Di-Higgs at the LHC

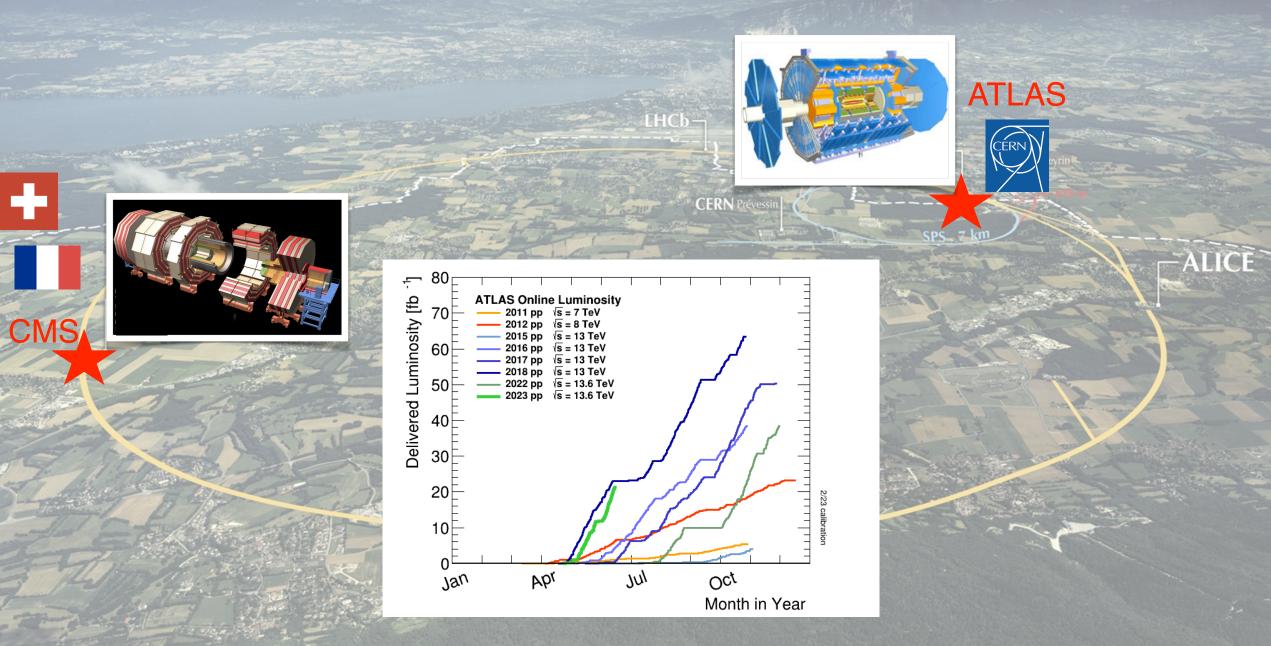
Caterina Vernieri

**SLAC** June, 26 2023



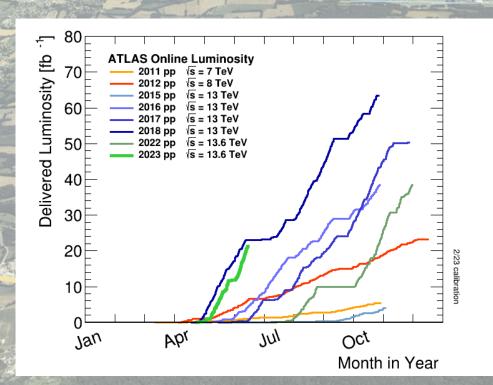


# The Large Hadron Collider (LHC)



# The Large Hadron Collider (LHC)

CMS

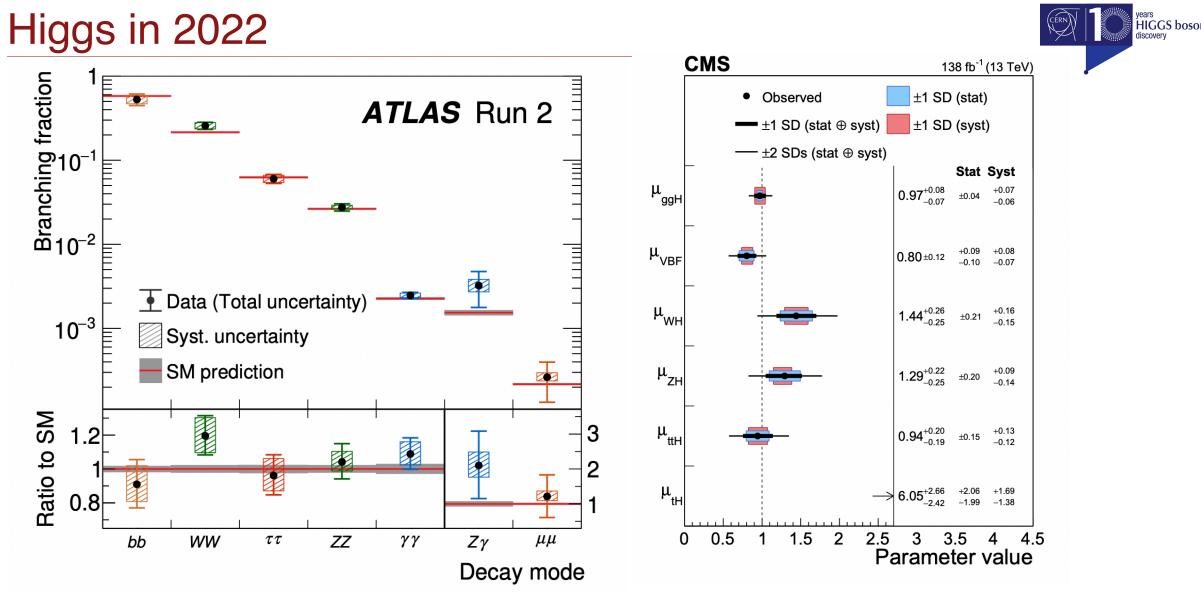


**CERN** Prévessi

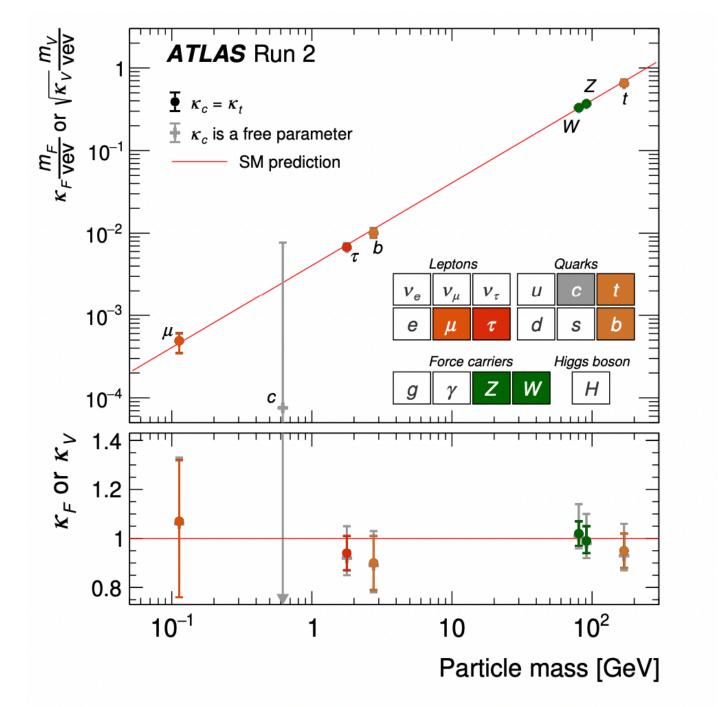
A factor 20x more data than at the time of the Higgs discovery

**ATLAS** 

CMS-Nature 607, 60–68 (2022 ATLAS-Nature 607, 52–59 (2022







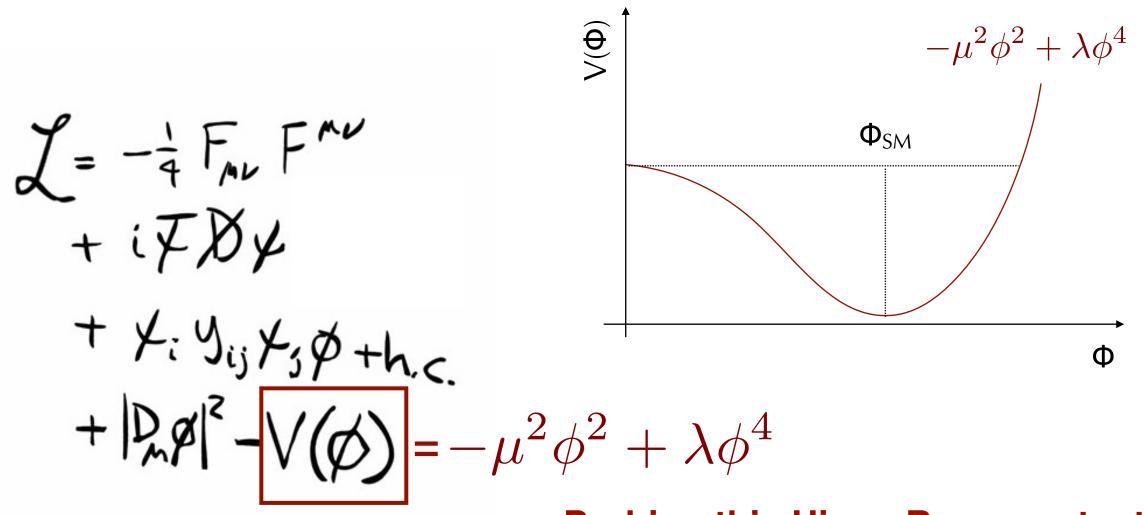
CMS-Nature 607, 60–68 (2022) ATLAS-Nature 607, 52–59 (2022)

BR(inv.) < 0.17 (0.11) CMS-PAS-HIG-20-003

lĸcl<3.4

CMS-PAS-HIG-21-008

Caterina Vernieri



### **Probing this Higgs Boson potential**

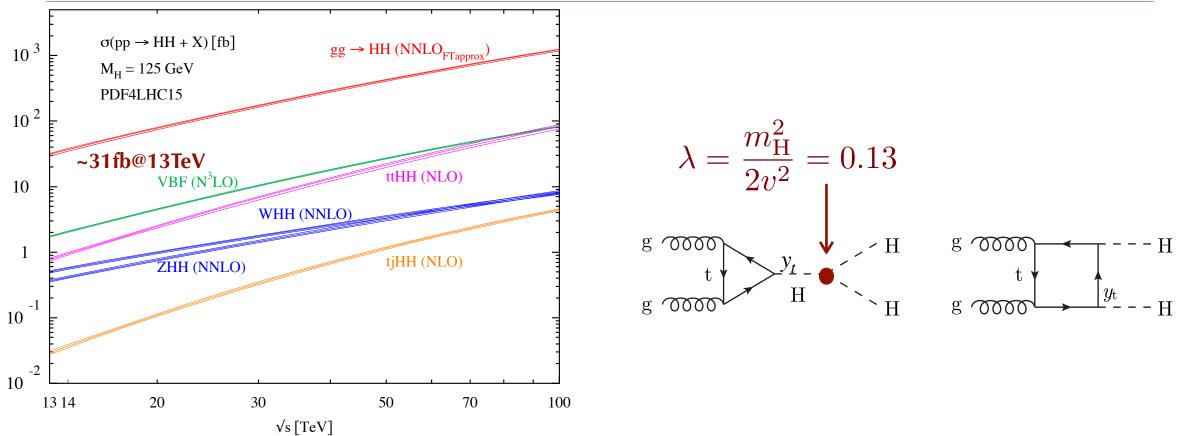
Caterina Vernieri

# Testing the shape <u><(</u>Ф) $V(\phi) = -\mu^2 \phi^2 + \lambda \phi^4$ $\Phi_{SM}$ $V(v+h) = V_0 + \frac{1}{2}m_h^2h^2 + \frac{m_h^2}{2v^2}vh^3 + \frac{1}{4}\frac{m_h^2}{2v^2}h^4$ H H Φ $\lambda = \frac{m_h^2}{2n^2} = 0.13$

### Testing the shape <u><(</u>Ф) $V(\phi) = -\mu^2 \phi^2 + \lambda \phi^4$ $\Phi_{SM}$ $V(v+h) = V_0 + \frac{1}{2}m_h^2h^2 + \frac{m_h^2}{2v^2}vh^3 + \frac{1}{4}\frac{m_h^2}{2v^2}h^4$ H H Φ OR ? Higgs potential Our $\lambda = \frac{m_h^2}{2n^2} = 0.13$ vacuum **Higgs field**

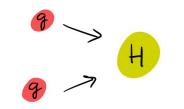
arXiv:1910.00012

### Higgs boson self-coupling

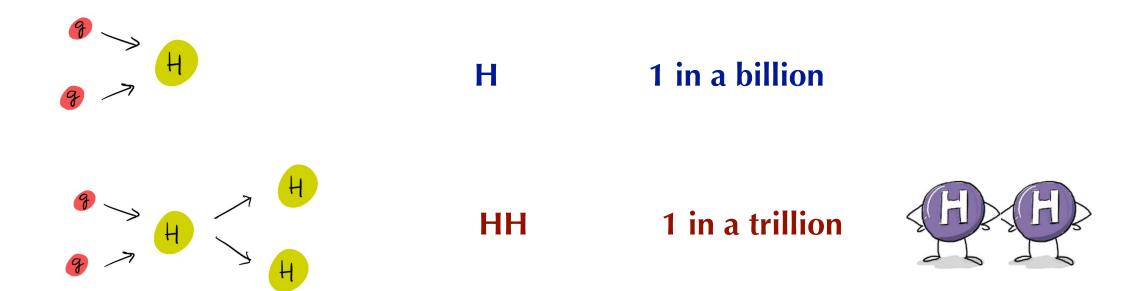


HH production allows to probe the self-coupling:  $\Delta\sigma/\sigma \sim \Delta\lambda/\lambda$  if  $\lambda \sim \lambda_{SM}$ 

Extremely challenging measurement at the LHC, but it can be sensitive to large deviations from BSM:  $\kappa_{\lambda} = \lambda / \lambda_{SM}$ 



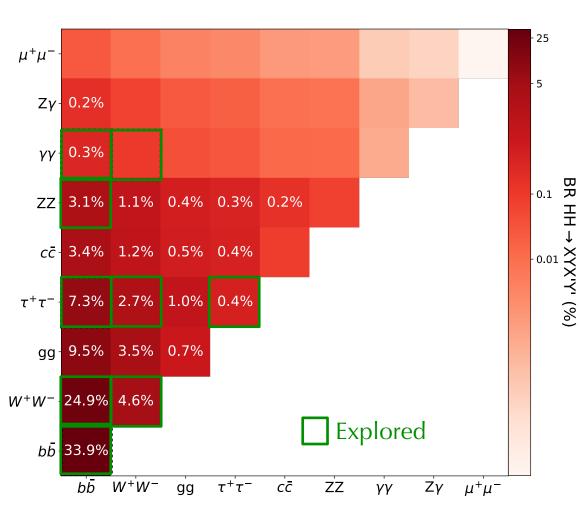
### H 1 in a billion





- At the present time, we have no experimental evidence that the Higgs boson results from the scalar potential of the SM
  - Observing double Higgs Boson production is crucial.
- Precision measurements of double Higgs boson production can then be combined with single Higgs measurements for a better understanding of the structure of the Higgs potential.
- Models of new physics that contain multiple Higgs bosons can lead to the production of Higgs bosons with different masses, leading to new experimental signatures.

### HH, a variety of final states



**H(bb)** is a key element in the exploration of HH at the LHC

#### highest BR

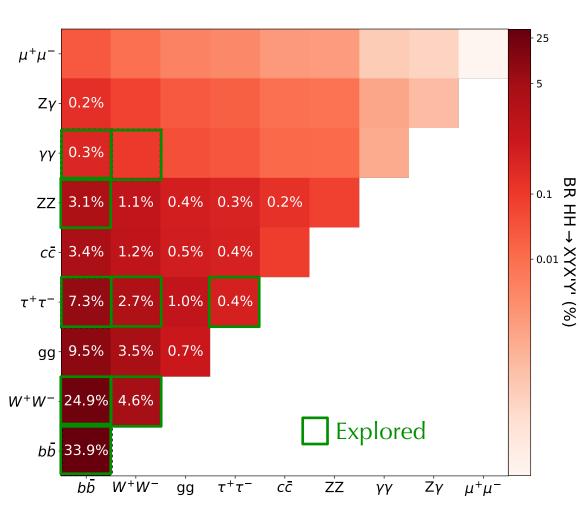
good b-jets identification performance: 70% efficiency at 0.3-1% q/g mistag probability

```
H(vv)
```

ΒR

clean final state excellent mass resolution, ~1%

### HH, a variety of final states



**H(bb)** is a key element in the exploration of HH at the LHC

#### highest BR

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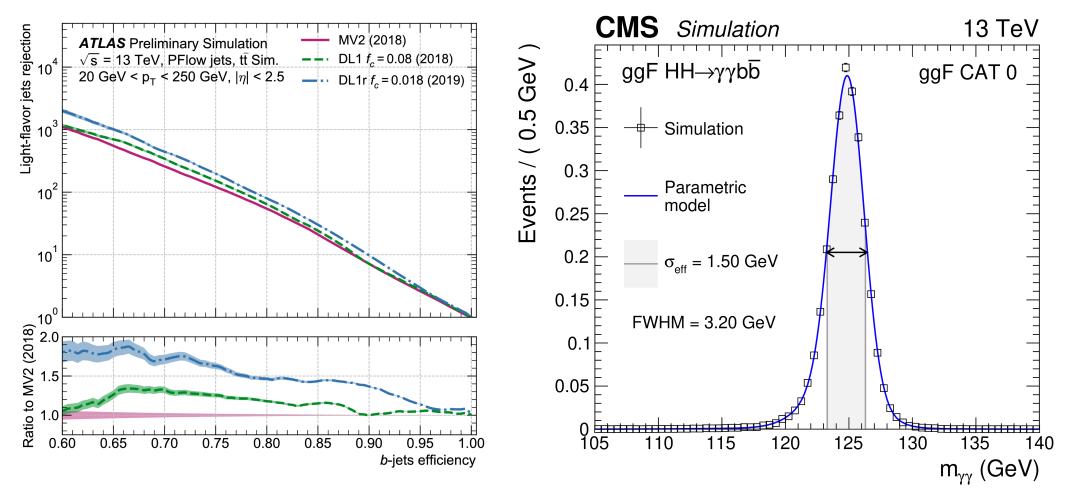
clean final state excellent mass resolution, ~1%

Several channels have been investigated:

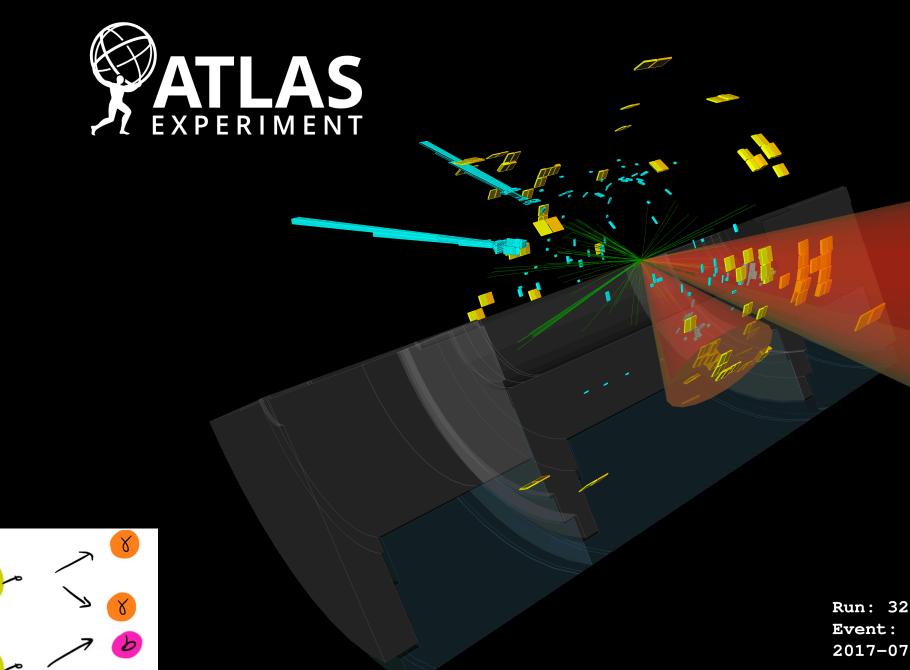
- Similar sensitivity to SM HH
- Different coverage on m<sub>HH</sub>

CMS Twiki ATLAS Twiki

# $H(b\bar{b})$ and $H(\gamma\gamma)$



- good b-jets identification performance: 70% efficiency at 0.7-1% q/g (10-20% c) mistag probability
- excellent diphoton ( $\gamma\gamma$ ) mass resolution, ~1%

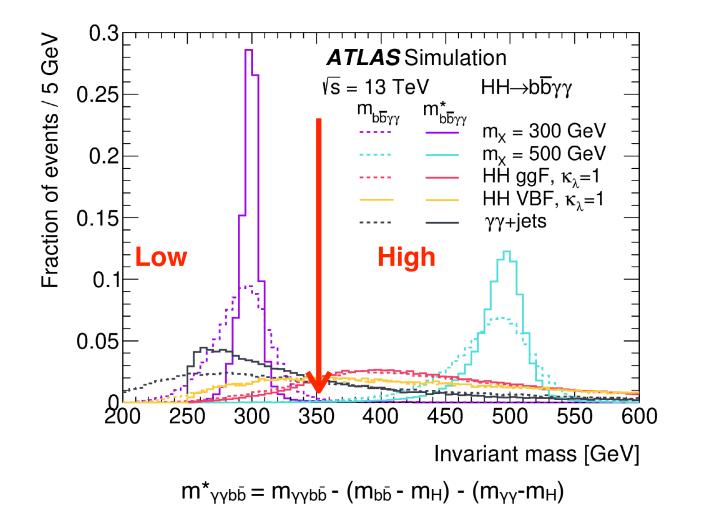


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6

Run: 329964 Event: 796155578 2017-07-17 23:58:15 CEST

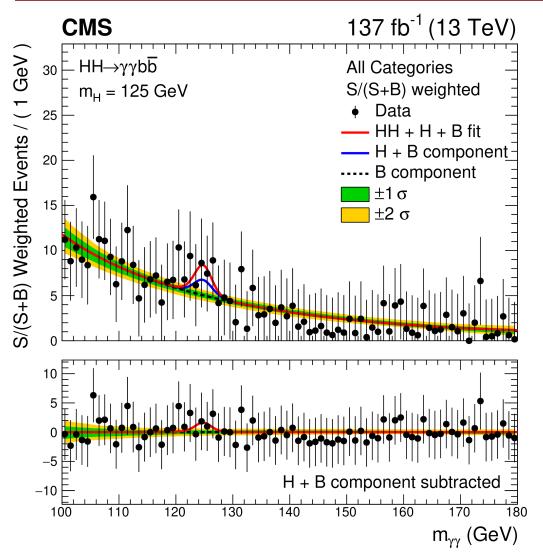
# HH to bbyy



- Photon selection similar to H(γγ) measurements
  - m<sub>γγ</sub> and m(bb) compatible with the Higgs boson mass
- Mx and BDT (includes angular correlations) used to categorize events
- Main **backgrounds** are:
  - ·  $\gamma\gamma$ +jets (prompt or jets misidentified as  $\gamma$ )
  - $\cdot$  SM single Higgs
  - Sensitivity to Low/High  $m_{\gamma\gamma b\bar{b}}$

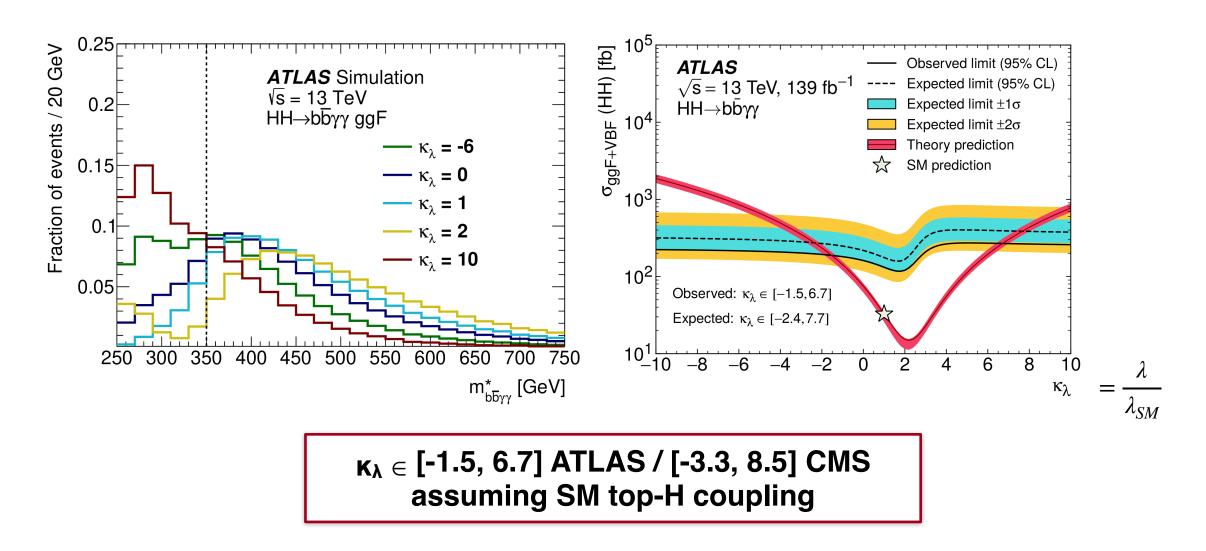
#### ATLAS-Phys. Rev. D 106, 052001 CMS-JHEP 03 (2021) 257

# HH to bbyy

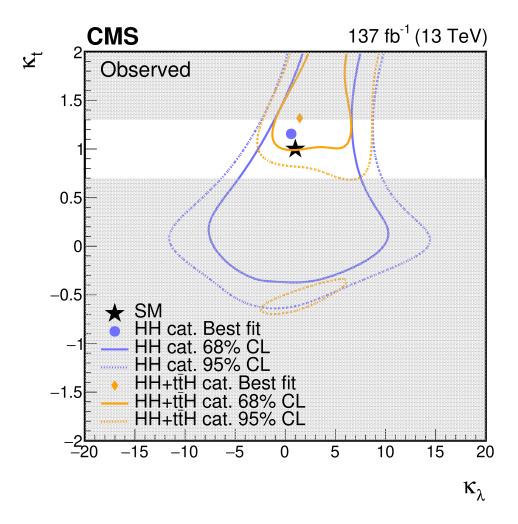


- CMS performs a likelihood fits simultaneous to  $m(b\bar{b})$  and  $m(\gamma\gamma)$  while ATLAS only uses the  $m(\gamma\gamma)$  distribution
- ATLAS 95% CL limits on HH production observed (expected) at 4.1 (5.5) times the SM prediction
- 5× improvement over previous result and current best available limits
- CMS observes 7.7 times SM predictions (expected 5.2x)

# Constraints on $\kappa_{\lambda}$ from $b\bar{b}\gamma\gamma$



### It is impossible to extract $\kappa_{\lambda}$ constraints from HH production without assumptions on $\kappa_t$



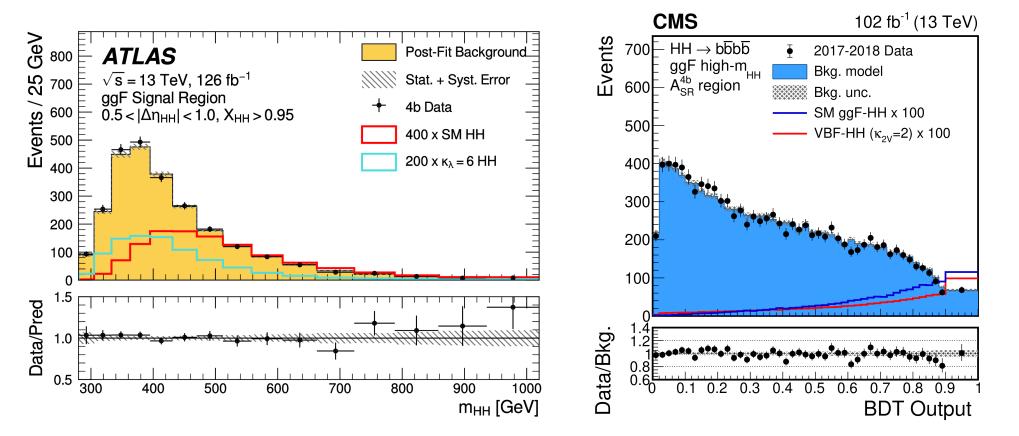
- A simultaneous likelihood scan for  $\kappa_{\lambda}$  and  $\kappa_t$  is performed to account for the modifications to
  - production HH and single Higgs cross sections
  - $B(H \rightarrow b\bar{b})$  and  $B(H \rightarrow \gamma\gamma)$
- Combination with ttH (lep. & had.) categories improves the constraints on  $\kappa_t$

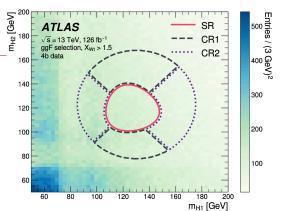
A combination of single H and HH measurements would provide a more model independent determination of  $\kappa_{\lambda}$ 

# HH to $b\bar{b}b\bar{b}$

### Large background from QCD multi-jet processes

Online b-tagging requirements to suppress multi-jet background Dedicated b-jet momentum corrections, to improve mass resolution (10–13%) Data driven background modeling with multi-variate kinematic reweighing from control regions

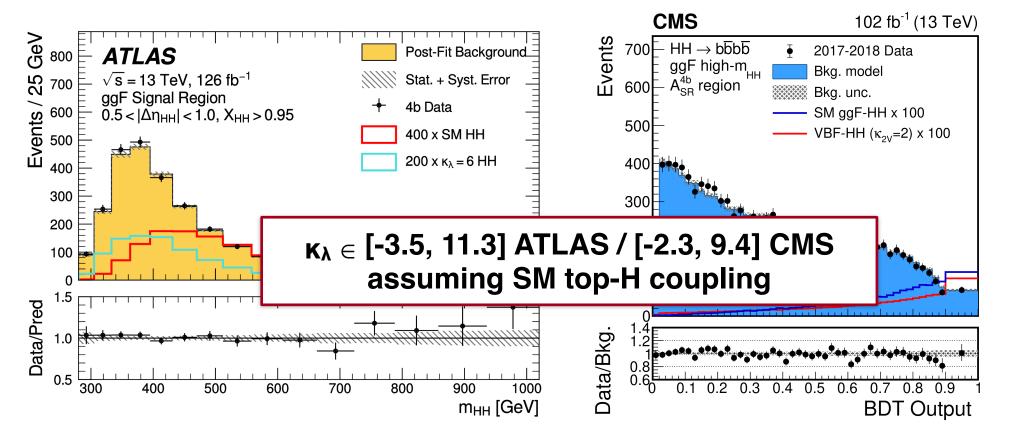




# HH to $b\bar{b}b\bar{b}$

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> CMS-Phys. Rev. Lett. 129 arXiv:2301.03212 29 (2022) 081802

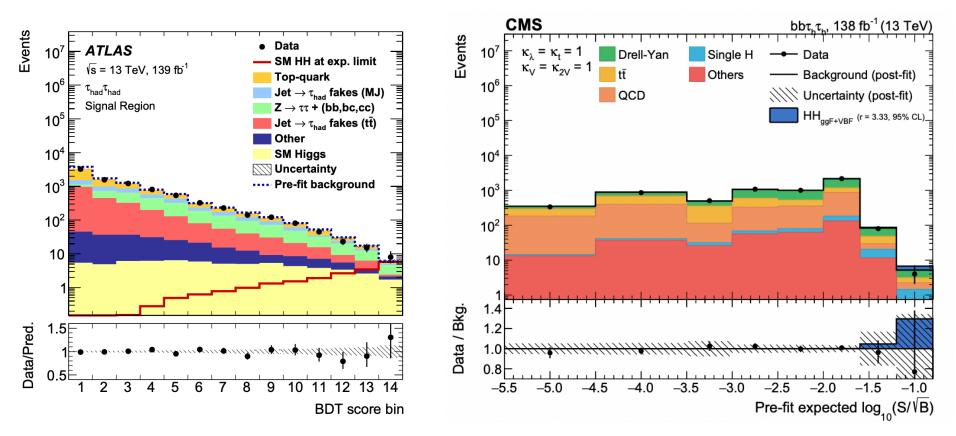
m<sub>H1</sub> [GeV]

#### ATLAS-arXiv:2209.10910 CMS-arXiv:2206.09401

# HH to $b\bar{b}\tau\tau$

### $\tau_h \tau_\mu + \tau_h \tau_e + \tau_h \tau_h$ , 88% of H( $\tau \tau$ )

- 2 jets (resolved) or 1 large-R jet (boosted) [CMS]
- Events are then categorized by number of b-tags (1 b-tag category adds 10% to sensitivity in CMS)
- NNs to separate signal from background based on angular separation of leptons and reconstructed invariant mass

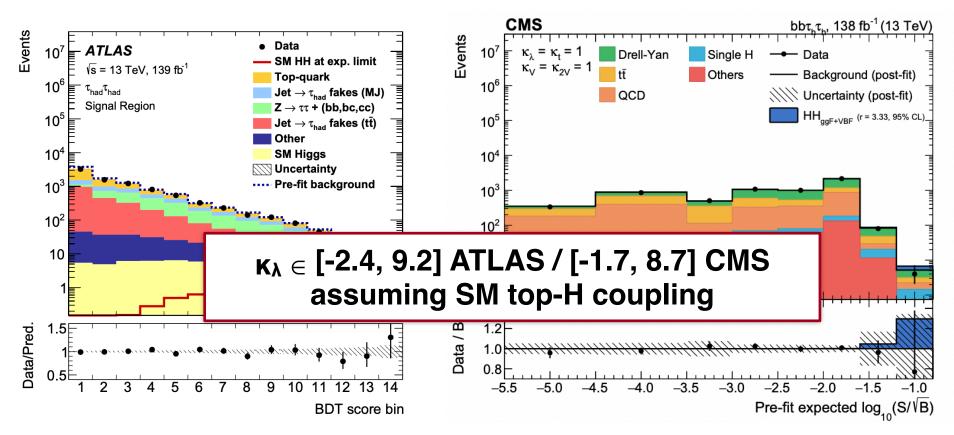


#### ATLAS-arXiv:2209.10910 CMS-arXiv:2206.09401

# HH to $b\bar{b}\tau\tau$

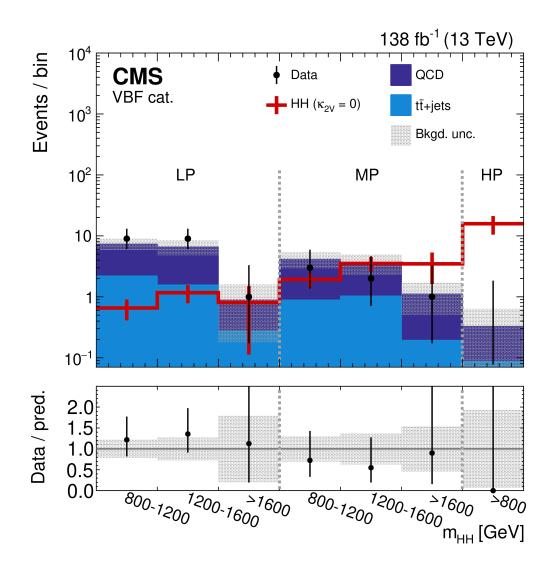
### $τ_h τ_μ + τ_h τ_e + τ_h τ_h$ , 88% of H(ττ)

- 2 jets (resolved) or 1 large-R jet (boosted) [CMS]
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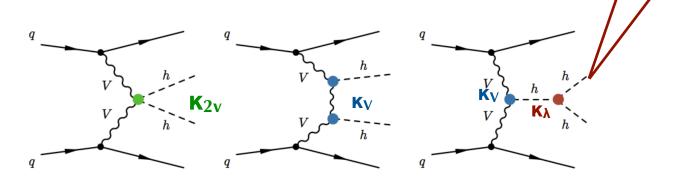
#### ATLAS-JHEP 05 (2021) 207 CMS-ArXiv:2205.06667

# VBF production



VBF production has a small cross-cross-section (1.73 fb @ N<sup>3</sup>LO)

- two high  $p_T$  forward jets provide a very specific topology b
- Dependencies on  $\kappa_{2v}$  (VVHH),  $\kappa_v$  and  $\kappa_\lambda$

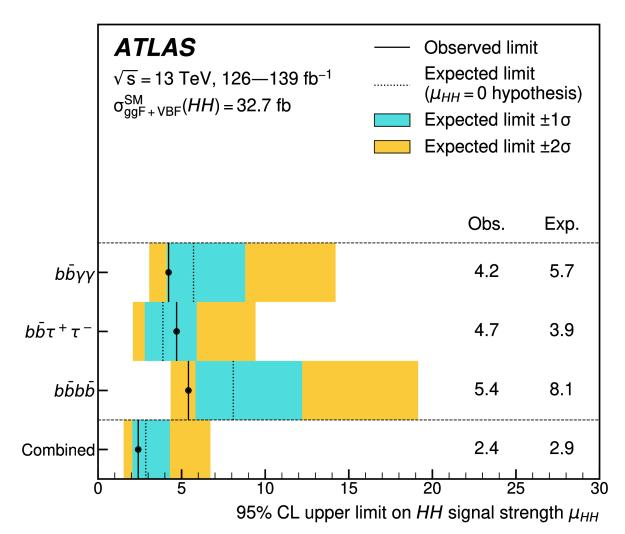


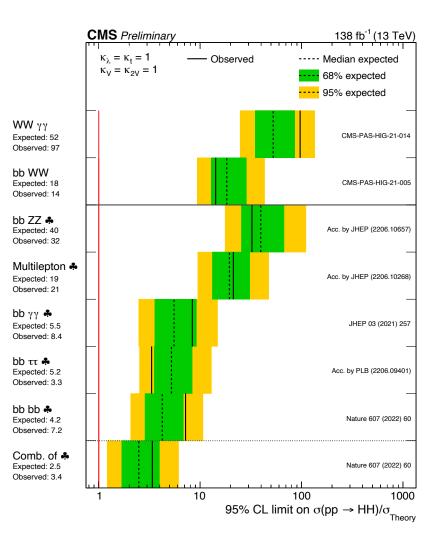
- New results from CMS for  $b\bar{b}b\bar{b}$  which targets Higgs bosons with high Lorentz boost
  - It probes large deviations of  $\kappa_V$  and  $\kappa_{2V}$  from their SM predictions, which result in a harder m<sub>HH</sub> spectrum and higher momentum for the b-jet from the Higgs boson decay.
- Each  $H \rightarrow b\bar{b}$  decay is reconstructed as a large-radius jet with characteristic substructure
- The observed (expected) limits constraint  $\kappa_{2V}$  within 0.6 <  $\kappa_{2V}$  < 1.4 at 95% CL *strongest limit on*  $\kappa_{2V}$ 
  - x30 improvement wrt early Run 2 results

b

# **Double Higgs Results**

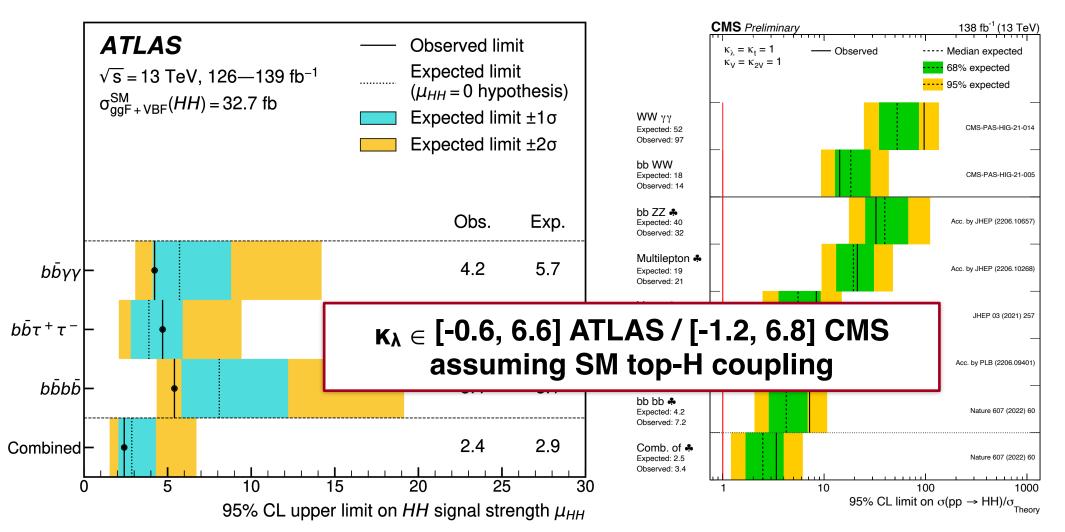
### Similar sensitivity from $b\bar{b}b\bar{b}$ , $b\bar{b}\tau\tau$ and $b\bar{b}\gamma\gamma$ to SM HH production



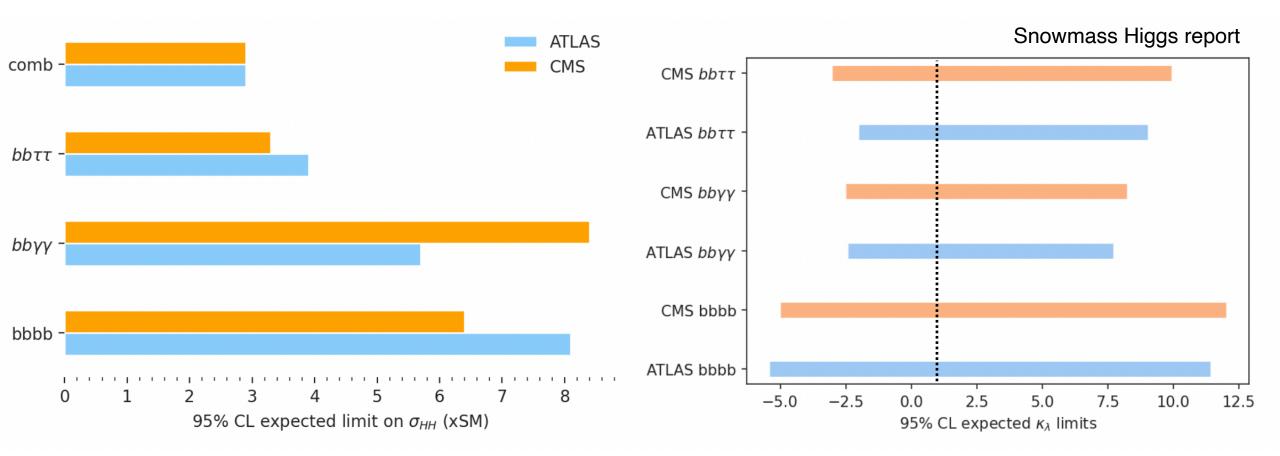


# **Double Higgs Results**

### Similar sensitivity from $b\bar{b}b\bar{b}$ , $b\bar{b}\tau\tau$ and $b\bar{b}\gamma\gamma$ to SM HH production



### Self-coupling : HH searches at the LHC

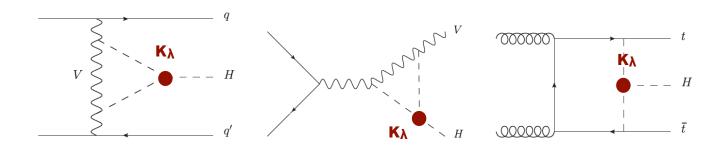


Main channels only based on full Run 2 (126-139/fb) analyses

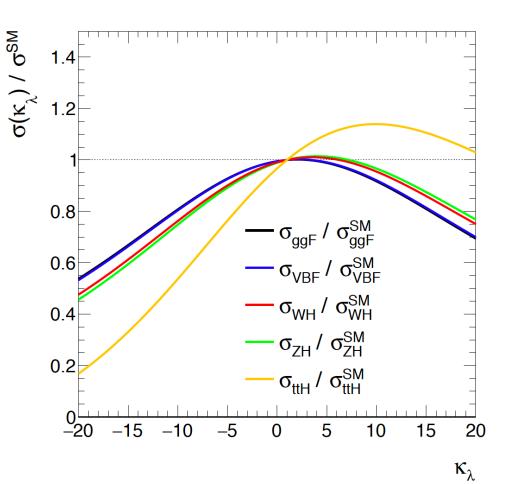
# Single Higgs to probe $\kappa_\lambda$

Single Higgs processes are sensitive to  $\lambda$  via NLO EW  $_{K_\lambda}\text{-}$  dependent corrections to

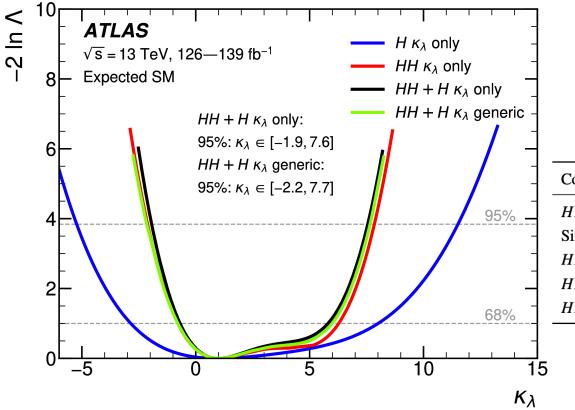
- cross-sections (ttH, ggF, VH, VBF )
- Higgs boson decay rates
- kinematics properties of the event (differential distributions)



JHEP 12 (2016) 080, Eur. Phys. J. C77 (2017) 887



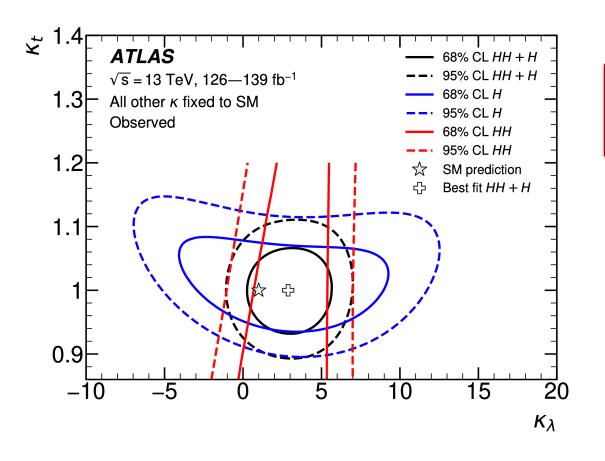
### First double and single Higgs combination



**First constraints** derived from 13 TeV single H + double H measurements with 126-139/fb

Combination assumption	Obs. 95% CL	Exp. 95% CL	Obs. value $^{+1\sigma}_{-1\sigma}$
HH combination	$-0.6 < \kappa_\lambda < 6.6$	$-2.1 < \kappa_\lambda < 7.8$	$\kappa_{\lambda} = 3.1^{+1.9}_{-2.0}$
Single- <i>H</i> combination	$-4.0 < \kappa_\lambda < 10.3$	$-5.2 < \kappa_\lambda < 11.5$	$\kappa_{\lambda} = 2.5^{+4.6}_{-3.9}$
<i>HH</i> + <i>H</i> combination	$-0.4 < \kappa_\lambda < 6.3$	$-1.9 < \kappa_\lambda < 7.6$	$\kappa_{\lambda} = 3.0^{+1.8}_{-1.9}$
<i>HH</i> + <i>H</i> combination, $\kappa_t$ floating	$-0.4 < \kappa_\lambda < 6.3$	$-1.9 < \kappa_\lambda < 7.6$	$\kappa_{\lambda} = 3.0^{+1.8}_{-1.9}$
<i>HH</i> + <i>H</i> combination, $\kappa_t$ , $\kappa_V$ , $\kappa_b$ , $\kappa_\tau$ floating	$-1.4 < \kappa_\lambda < 6.1$	$-2.2 < \kappa_\lambda < 7.7$	$\kappa_{\lambda} = 2.3^{+2.1}_{-2.0}$

### First double and single Higgs combination



First constraints derived from 13 TeV single H + double H measurements with 126-139/fb

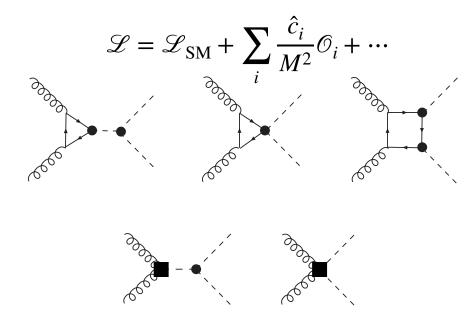
They are compatible to those derived from HH direct searches but an EFT global analysis is needed to proper account of BSM contributions

# Towards an (SM)EFT analysis for the self-coupling

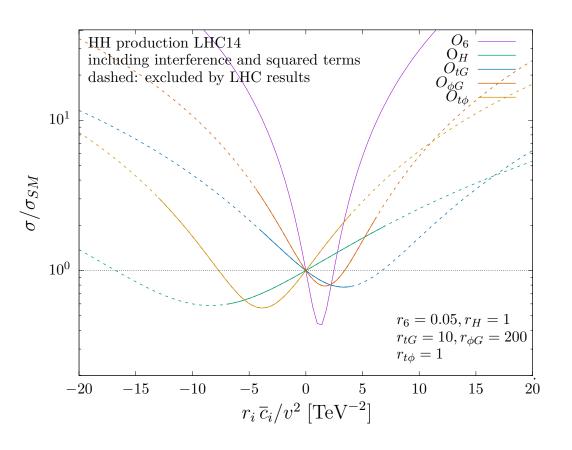
The EFT formalism summarizes deviations that might appear in a very wide class of models beyond the SM

#### The most general set of BSM perturbations will modify other interactions and not just the self-coupling

Within the EFT formalism, data from several different processes can be combined to constrain the new physics parameters BSM contributions to HH production cross sections are at most of size  $E^2/M^2$ , if E~m<sub>H</sub> and M~ O(TeV):

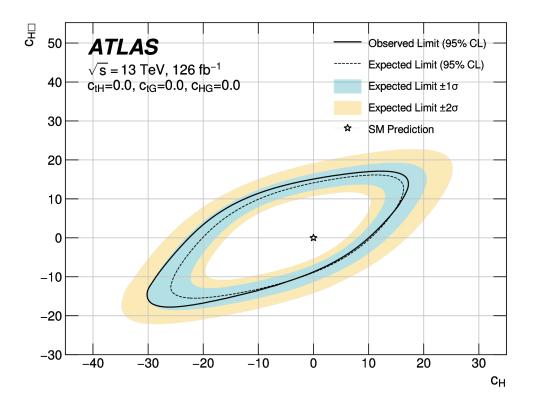


To determine the self-coupling at O(10%) we could consider only a subset of operators contributing at LO to  $gg \rightarrow HH$ • 5 operators contribute to the process - if CP conserved



# An example, HH to bbbb

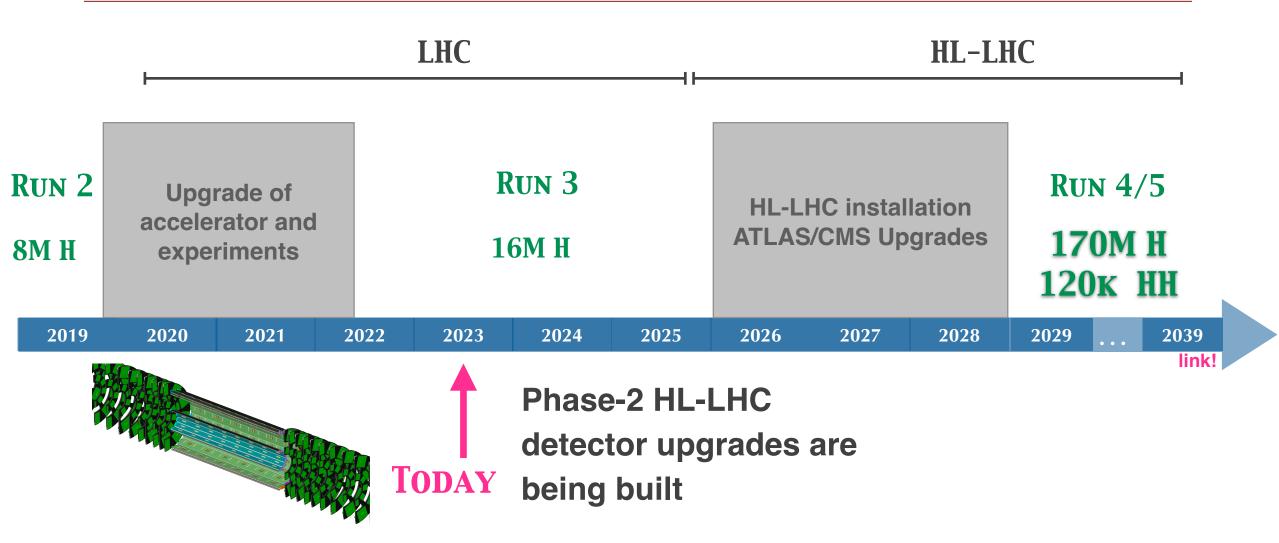
### 1D & 2D constraints on Wilson Coefficients



Wilson Coefficient	Operator
$c_H$	$(H^{\dagger}H)^{3}$
$c_{H\Box}$	$(H^{\dagger}H)\Box(H^{\dagger}H)$
$c_{tH}$	$(H^{\dagger}H)(\bar{Q}\tilde{H}t)$
$c_{HG}$	$ \begin{array}{c} H^{\dagger}HG^{A}_{\mu\nu}G^{\mu\nu}_{A} \\ (\bar{Q}\sigma^{\mu\nu}T^{A}t)\tilde{H}G^{A}_{\mu\nu} \end{array} $
$c_{tG}$	$(\bar{Q}\sigma^{\mu\nu}T^At)\tilde{H}G^A_{\mu\nu}$

Parameter	Expected Constraint		<b>Observed Constraint</b>	
	Lower	Upper	Lower	Upper
$c_H$	-20	11	-22	11
$c_{HG}$	-0.056	0.049	-0.067	0.060
$c_{H\Box}$	-9.3	13.9	-8.9	14.5
$c_{tH}$	-10.0	6.4	-10.7	6.2
$c_{tG}$	-0.97	0.94	-1.12	1.15

# LHC → High Luminosity LHC





Pileup increases with luminosity

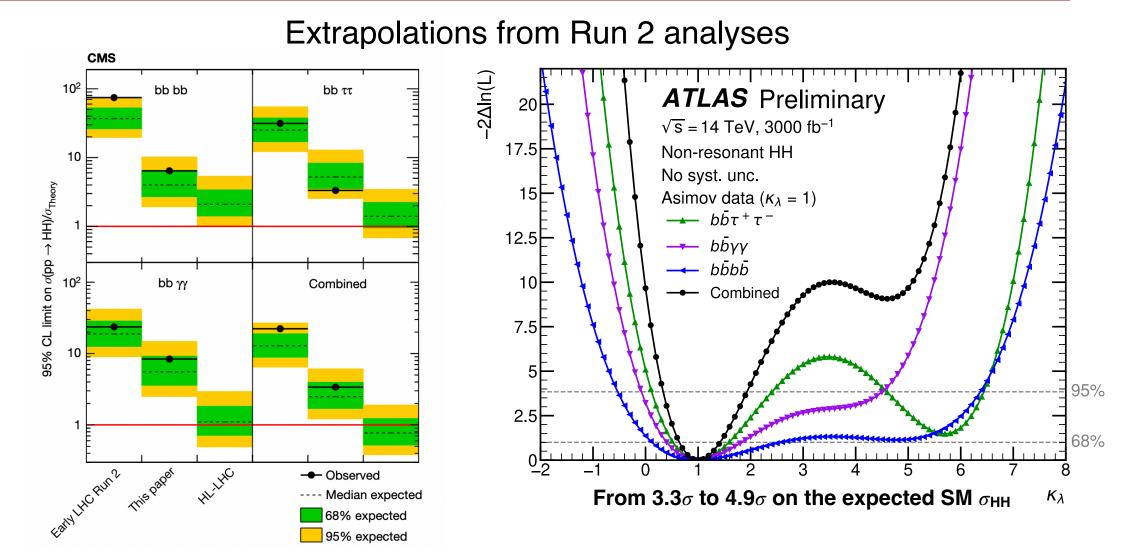
ref

z position is critical to identify primary vertex

High granularity tracking detector is critical

LHC (2023): 50-60 HL-LHC: 140-200

## Higgs physics at the HL-LHC



#### SLAC Caterina Vernieri · LOOPFEST · June 26, 2023

## **Conclusions and perspectives**

Experiments at LHC are approaching SM sensitivity to non-resonant HH production

- Several final states investigated, more to exploit with more data
  - new improved analysis techniques
  - · ML approaches being used in most final states
- More to be learned on the self-coupling from differential mHH, VBF topology and single H measurements
- First H+HH and EFT combined analyses for a more general interpretation
- At HL-LHC more data will be available to test the self-coupling and probe rare (new?) processes
  - · Challenging experimental environment for online selections and b-tagging
  - Improved tracker detectors should allow us to maintain or even yield to increased acceptance

### **Conclusions and perspectives**

#### Experiments at LHC are approaching SM sensitivity to non-resonant HH production

- · Several final states investigated, more to exploit with more data
  - new improved analysis techniques
  - ML approaches being used in most final states
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September 2018 - Science Magazine	)
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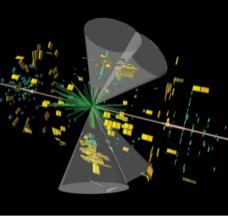
Careers -

The LHC experiments may need years to see a signal. Later this year, the LHC will idle for 2 years for upgrades. In 2026 it will undergo another 2-year hiatus to boost its collision rate. The so-called High-Luminosity LHC would then run until 2034. On paper, only the full run will yield enough data to validate the standard model prediction. However, some physicists think they can beat that timetable as their Higgs-spotting algorithms continue to improve. "Even before the High-Luminosity LHC, I think we could get close to the standard model prediction," says Caterina Vernieri, a CMS member at Fermilab.

Contents -

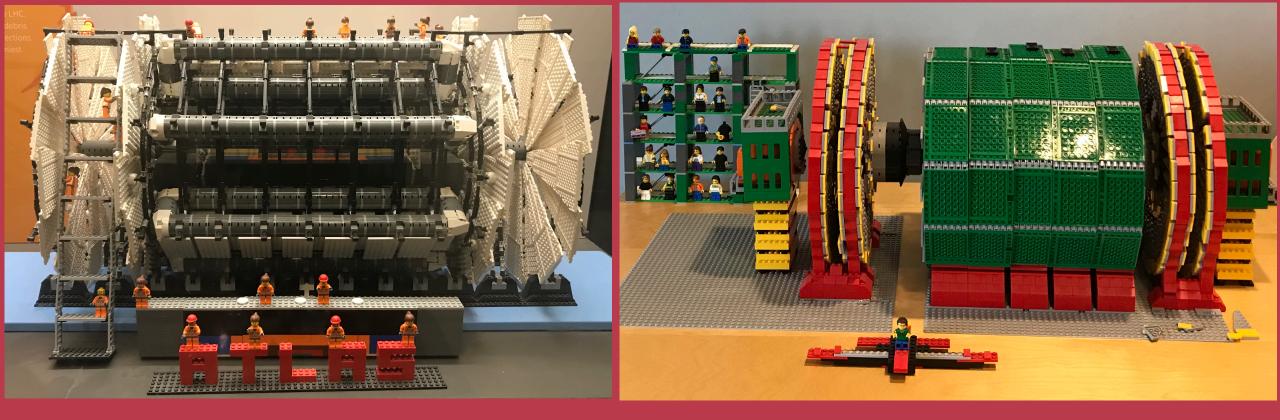
News

Of course, all LHC experimenters hope the rate for double-Higgs events will exceed the standard model prediction. It cannot be sky



Journals -

Two Higgs bosons may have decayed into bottom quarks in this 2016 collision in the ATLAS detector. ATLAS EXPERIMENT © 2018 CERN



# thank you!

#### arXiv:1910.00012

## Which precision on the self-coupling is needed?



#### Sensitivity to:

models where we expect new particles of few hundred GeV mass mixing of the Higgs boson with a heavy scalar with a mass of order 1 TeV loop diagram effects created by any new particle with strong coupling to the H typical quantum corrections to the Higgs self-coupling generated by loop diagrams

#### arXiv:1910.00012

## Which precision on the self-coupling is needed?

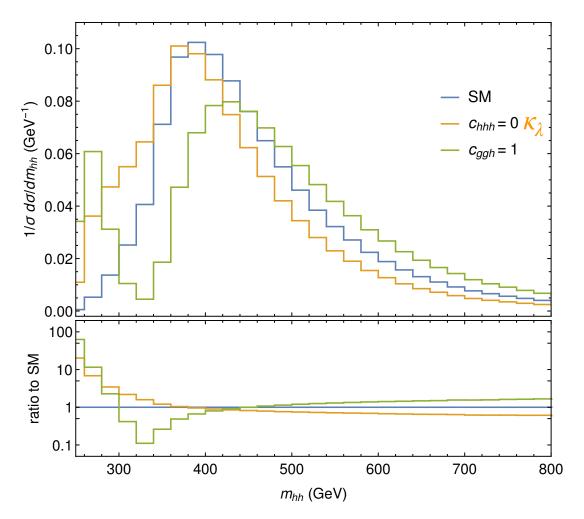


#### Sensitivity to:

models where we expect new particles of few hundred GeV mass

Interplay between precisions inference and direct searches for new particles.

#### Impact on m<sub>HH</sub>



- A global SMEFT interpretation could constrain the  $\mathcal{O}_6$  coefficient at levels of order 1 with HL-LHC dataset
- All these effective operators induce different distortions in the  $m_{\text{HH}}$  distribution
  - A shape analysis can thus help in disentangling the various operators in a **global fit**

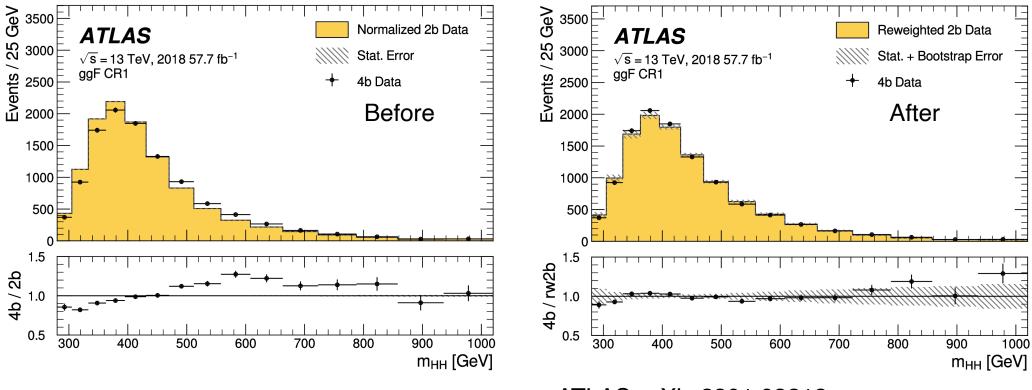
## HH to bbbb

#### Large background from QCD multi-jet processes

Online b-tagging requirements to suppress multi-jet background

Dedicated b-jet momentum corrections, to improve mass resolution (10–13%)

Data driven background modeling with multi-variate kinematic reweighing from control regions

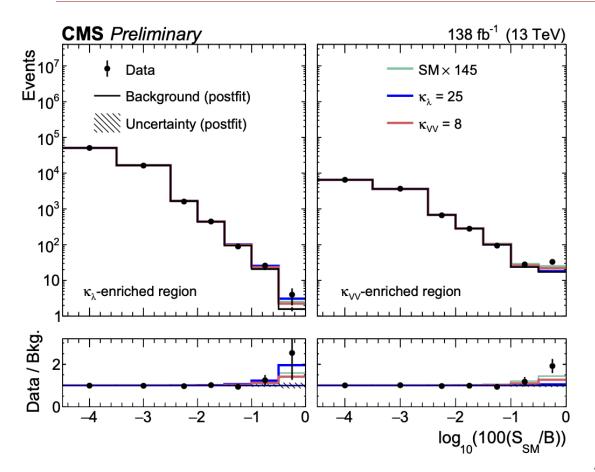


ATLAS-arXiv:2301.03212 CMS-Phys. Rev. Lett. 129 (2022) 081802

#### intries / (3 GeV)<sup>2</sup> 50 40 ATLAS $\sqrt{s} = 13 \text{ TeV}, 126 \text{ fb}^{-1}$ CR1 ggF selection, X<sub>wt</sub> > 1.5 CR2 4b data 300 120 200 100 100 100 120 180 m<sub>H1</sub> [GeV]

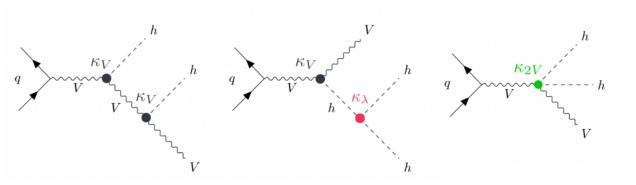
#### ATLAS-ArXiv:2210.054157 CMS-HIG-22-006

## VHH production



VHH production has a small cross-cross-section (0.86 fb @ NNLO)

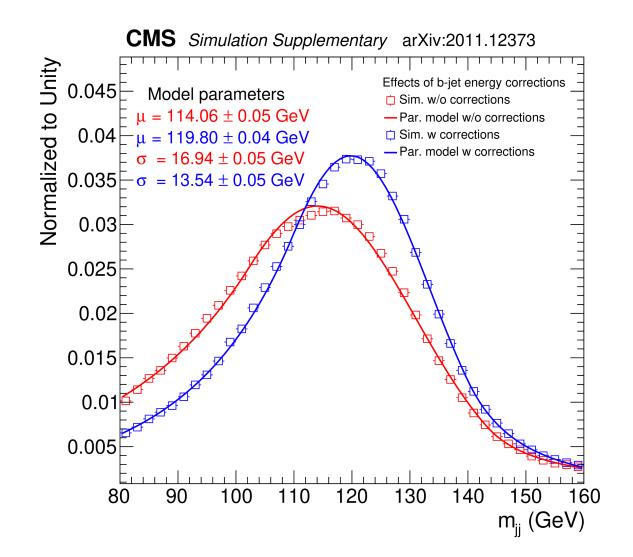
• Dependencies on  $\kappa_{2\nu}$  (VVHH),  $\kappa_{\nu}$  and  $\kappa_{\lambda}$ 



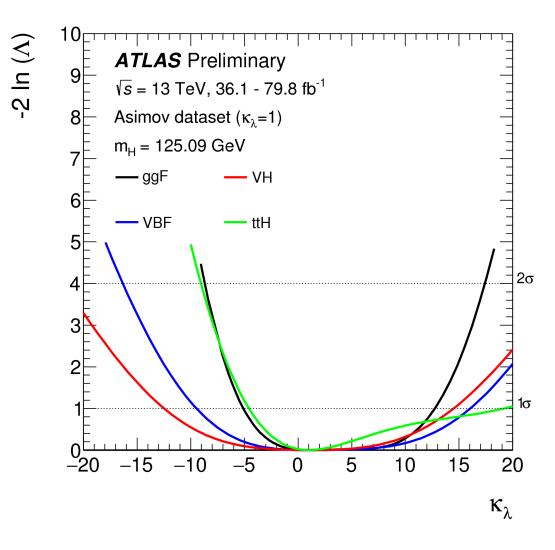
VHH production allows separation of WHH and ZHH

- Three leptonic channels (0L, 1L, 2L), one for each leptonic decay mode of the W and Z bosons
- HH in the bbbb final state
- The observed (expected) limits constraint  $\kappa_{2V}$  within -8.6 <  $\kappa_{2V}$  < 10 at 95% CL (ATLAS) and -12.2 <  $\kappa_{2V}$  < 13.5 at 95% CL (CMS)

#### Impact of the regression for b-jets



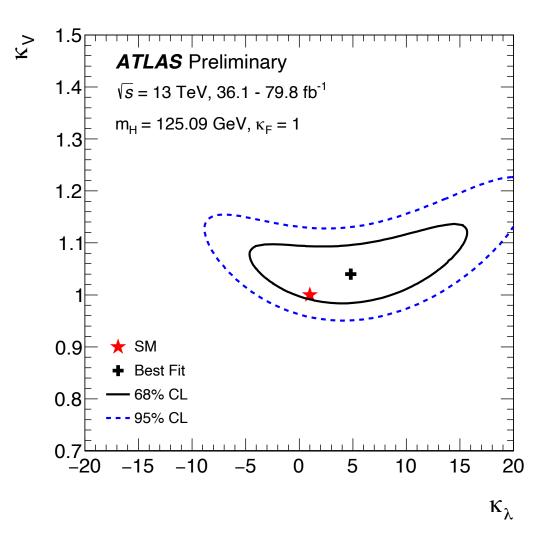
## Single Higgs constraints



First constraintsderived from 13 TeV single Hmeasurements with 36-80/fb $\kappa_{\lambda} = 4.0^{+4.3}$ -4.1

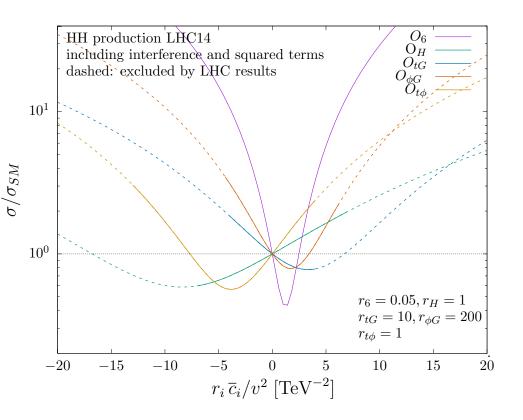
37

## Single Higgs constraints



First constraintsderived from 13 TeV single Hmeasurements with 36-80/fb $\kappa_{\lambda} = 4.0^{+4.3}$ -4.1

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The LO  $gg \rightarrow HH$  cross section depends rather strongly on **five coefficients** 

- Top quark measurements would constrain the dipole operator  $\mathcal{O}_{tG}$
- The top Yukawa operator  $\mathcal{O}_u$  can be constrained by **ttH**
- $\mathcal{O}_{\phi G}$  can be determined by **ggF**
- +  $\mathcal{O}_H$  can be extracted as a uniform rescaling of all Higgs couplings
- $\mathcal{O}_6$  can lead to deviations of order 10 in the HH cross section
- A global SMEFT interpretation could constrain the  $\mathcal{O}_6\,$  coefficient at levels of order 1 with HL-LHC dataset

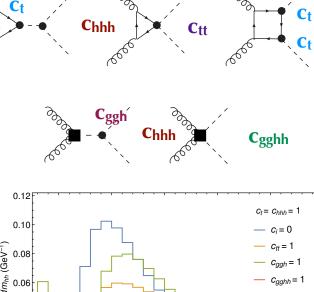
$$\kappa_{\lambda} = 1 + c_6 - \frac{3}{2}c_H + \cdots$$

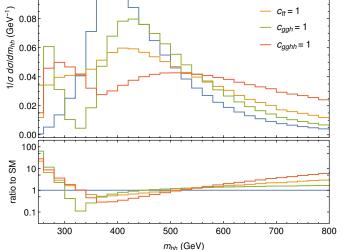
## $gg \rightarrow HH$ in the HEFT

- HEFT singles out the anomalous Higgs couplings as the leading effects of new physics in the EW sector
  - At LO it provides a field-theory basis for the empirical  $\kappa$ -framework
- For  $gg \rightarrow HH$  the relevant terms in the HEFT Lagrangian are

$$\Delta \mathscr{L}_{\chi} = -m_t \left( c_t \frac{h}{v} + c_{tt} \frac{h^2}{v^2} \right) \bar{t} t - c_{hhh} \frac{m_h^2}{2v} h^3 + \frac{\alpha_s}{8\pi} \left( c_{ggh} \frac{h}{v} + c_{gghh} \frac{h^2}{v^2} \right) G^a_{\mu\nu} G^{a,\mu\nu}$$

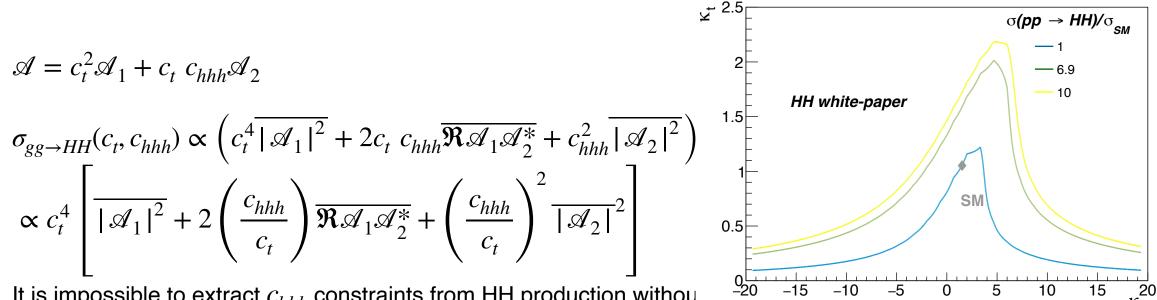
- . The SMEFT relations between  $c_{t}$  and  $c_{tt}$  ,  $\ c_{gghh}$  and  $c_{ggh}$  are not present here
- $\cdot c_{hhh}$  can be extracted only through a global analysis
  - +  $c_{tt}$  and  $c_{gghh}\!$  , appear only in processes involving HH and we cannot use single Higgs data to fix them as in SMEFT
  - +  $t\bar{t}HH$  production can allow a determination of  $c_{tt}$  independently from  $c_{hhh}$





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If  $c_{tt} = c_{gghh} = c_{ggh} = 0$ ,  $gg \to HH$  depends only from  $c_t$  and  $c_{hh}$ 



It is impossible to extract  $c_{hhh}$  constraints from HH production withou

- $c_t$  and  $c_{hhh}$  can be constrained also using single Higgs measurements.
- A combination of single H and HH measurements would provide a more model independent determination of  $c_{hhh}$

#### Perspectives for HH at HL-LHC

Estimates of the sensitivity to HH at HL-LHC are based on:

extrapolations from Run-2 analyses

dedicated studies with smeared/parametric detector response, corresponding to pile-up of 200 A combined significance to the **SM HH process of 4** $\sigma$  can be achieved with all systematic

uncertainties

	Statistical-only		Statistical + Systematic	
	ATLAS	CMS	ATLAS	CMS
$HH \rightarrow b\bar{b}b\bar{b}$	1.4	1.2	0.61	0.95
$HH  ightarrow b \bar{b}  au  au$	2.5	1.6	2.1	1.4
$HH  ightarrow b ar b \gamma \gamma$	2.1	1.8	2.0	1.8
$HH \rightarrow b\bar{b}VV(ll\nu\nu)$	-	0.59	-	0.56
$HH \rightarrow b\bar{b}ZZ(4l)$	-	0.37	-	0.37
combined	3.5	2.8	3.0	2.6
	Combined 4.5		Combined	
			4.0	

