

Higgs Boson Production at the Next-to-Next-to-Leading Power

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Topics discussed

- High-energy or small-mass limit of $gg \rightarrow H$

- *leading logs at $\mathcal{O}(m_q^2/s)$*

T. Liu, A.A. Penin, Phys.Rev.Lett. **119**, 262001 (2017)

- *next-to-leading logs at $\mathcal{O}(m_q^2/s)$*

C. Anastasiou, A.A. Penin, JHEP **07**, 195 (2020)

- *leading logs at $\mathcal{O}(m_q^4/s^2)$*

T. Liu, S. Modi, A.A. Penin, JHEP **02**, 170 (2022)

The advent of power corrections

● Power vs perturbative corrections

- Λ_X^2/Q^2 vs α_X^n

- e.g. $\Lambda_X = m_b$, $Q = m_H$, $\alpha_X = \alpha_s \Rightarrow n = 3$

● Logarithmically enhanced power corrections

- *phenomenologically relevant*

- *intriguing from QFT point of view*

- ➡ *weird RG structure, magic relations, etc.*

● Recent stream of the NLP results for

- *mass, angle, soft momentum, threshold, jetiness, ...*

Higgs production at the LHC

- Total cross section at 13 GeV

$$\sigma_{pp \rightarrow H+X} = 48.68 \text{ pb}$$

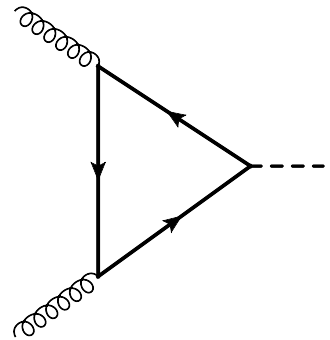
- Dominant theory uncertainties

- *scale choice* $+0.10$
 -1.15 pb
- *PDF N^3LO* $\pm 0.56 \text{ pb}$
- $m_b > 0$ *NNLO+* $\pm 0.40 \text{ pb}$

Anastasiou et al. JHEP 1605, 058 (2016)

Bottom quark mass effect

● Leading contribution



The diagram shows a triangle loop of bottom quarks. Two external lines, represented by curly lines for gluons, enter the loop from the left. The top and bottom lines of the loop have arrows pointing to the right. The right side of the loop is a dashed line, representing a quark propagator. To the right of the diagram is the equation: $\propto \alpha_s \ln^2(m_H^2/m_b^2) \frac{m_b^2}{m_H^2}$

➡ *large logs at subleading power*

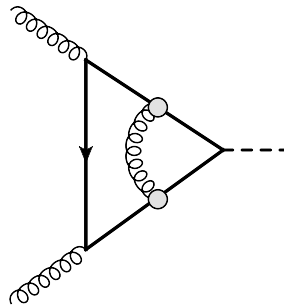
● *effective expansion parameter* $\alpha_s \ln^2(m_H^2/m_b^2) \sim 40\alpha_s$

➡ *resummation is mandatory*

$gg \rightarrow H$ amplitude at NLP LL

● Non-Sudakov logs

T. Liu, A.A. Penin, Phys.Rev.Lett. 119, 262001 (2017)



● Factorization formula

$$\mathcal{M}_{gg \rightarrow H}^b = Z_g^{2LL} g(z) \mathcal{M}_{gg \rightarrow H}^{b(0)}$$

● *gluon Sudakov factor* $Z_g^{2LL} = \exp \left[-\frac{C_A}{\epsilon^2} \frac{\alpha_s}{2\pi} \frac{\mu^{2\epsilon}}{Q^{2\epsilon}} \right]$

● *non-Sudakov double logarithms*

$$g(z) = 2 \int_0^1 d\xi \int_0^{1-\xi} d\eta e^{2z\eta\xi} = {}_2F_2(1, 1; 3/2, 2; z/2)$$

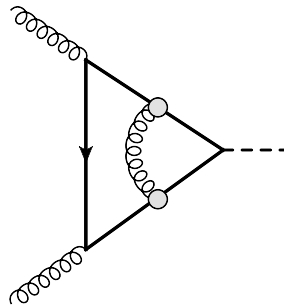
● *double-log variable* $z = (C_A - C_F) x$, $x = \frac{\alpha_s}{4\pi} L^2$, $L = \ln(m_H^2/m_q^2)$

eikonal color nonconservation

$gg \rightarrow H$ amplitude at NLP LL

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T. Liu, A.A. Penin, Phys.Rev.Lett. 119, 262001 (2017)



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➡ *Magic #1: same function for NLP QED scattering in Regge limit color (Sudakov, soft) \Leftrightarrow kinematics (Regge, Glauber)*

$gg \rightarrow H$ amplitude at NLP NLL

C. Anastasiou, A.A. Penin, JHEP **07**, 195 (2020)

$$\mathcal{M}_{gg \rightarrow H}^{bNLL} = C_b \left(\frac{\alpha_s(m_H)}{\alpha_s(m_b)} \right)^{\gamma_m^{(1)}/\beta_0} Z_g^{2NLL} \left[-\frac{3}{2} \frac{m_b^2}{m_H^2} L^2 \mathcal{M}_{gg \rightarrow H}^{t(0)} \right]$$

Yukawa RG factor
gluon Sudakov form factor
LO amplitude

$$C_b = \left[g(z) + \frac{\alpha_s L}{4\pi} (2\gamma_q^{(1)} g_\gamma(z) - \beta_0 g_\beta(z)) \right] = 1 + \sum_{n=1}^{\infty} c_n$$

$$c_1 = \frac{z}{6} + C_F \frac{\alpha_s L}{4\pi}, \quad c_2 = \frac{z^2}{45} + \frac{z}{5} \frac{\alpha_s L}{4\pi} \left[\frac{3}{2} C_F - \beta_0 \left(\frac{5}{6} \frac{L_\mu}{L} - \frac{1}{3} \right) \right],$$

$$c_3 = \frac{z^3}{420} + \frac{z^2}{5} \frac{\alpha_s L}{4\pi} \left[\frac{5}{21} C_F - \beta_0 \left(\frac{2}{9} \frac{L_\mu}{L} - \frac{2}{21} \right) \right], \quad \dots$$

$$L = \ln(s/m_q^2), \quad L_\mu = \ln(s/\mu^2)$$

c_2 agrees with 3-loop num. calc. Czakon, Niggetiedt, JHEP **05**, 149 (2020)

Top-bottom interference in threshold cross section

C. Anastasiou, A.A. Penin, JHEP **07**, 195 (2020)

	LO	NLO	NNLO	N ³ LO
$\delta\sigma_{pp\rightarrow H+X}^{\text{LL}}$	-1.420	-1.640	-1.667	-1.670
$\delta\sigma_{pp\rightarrow H+X}^{\text{NLL}}$	-1.420	-2.048	-2.170	-2.189
$\delta\sigma_{pp\rightarrow H+X}$	-1.023	-2.000		

• **NLL K-factors with full threshold** $\delta\sigma_{pp\rightarrow H+X}^{\text{NLO}}$

$$\delta\sigma_{gg\rightarrow H+X}^{\text{NNLO}} \approx -0.12 \text{ pb}$$

$$\delta\sigma_{gg\rightarrow H+X}^{\text{N}^3\text{LO}} \approx -0.02 \text{ pb}$$

• **New uncertainty interval** $-0.32 \text{ to } 0.08 \text{ pb}$ (factor 2 reduction)

Next-to-next-to leading power: why?

● Theory

- *proof of principle*
- *structure of RG: terra incognita*

● Phenomenology

- *validate the quark mass expansion*
- *NNLO Higgs pair production at high p_{\perp} , etc.*

Next-to-next-to leading power: how?

● Brut-force approach

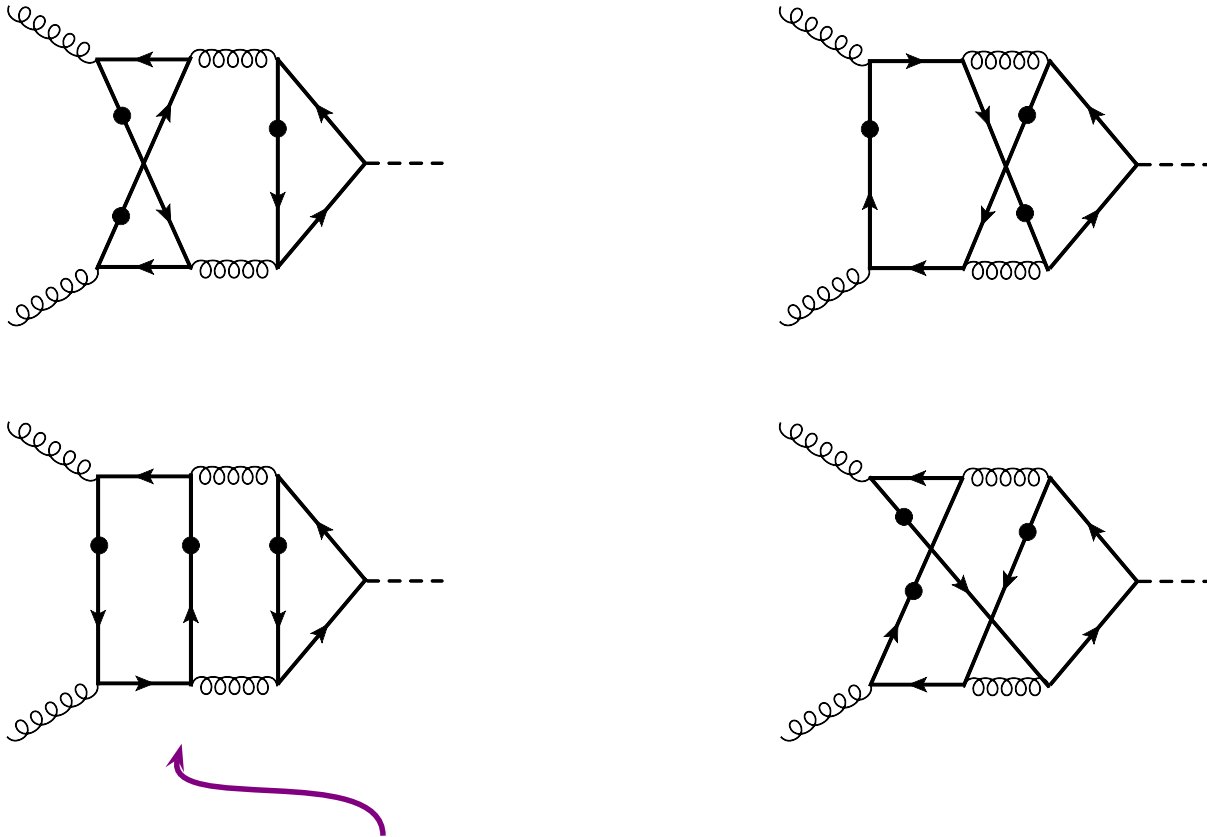
A.A. Penin, Phys.Lett. B **119** (2015) 262001

- *expansion by regions* \Leftrightarrow *homogeneous integrals* \Leftrightarrow *log corrections*
- *Ward identities + momentum shifts + eikonal factorization* \Leftrightarrow *resummation*

● Classification of NNLP terms

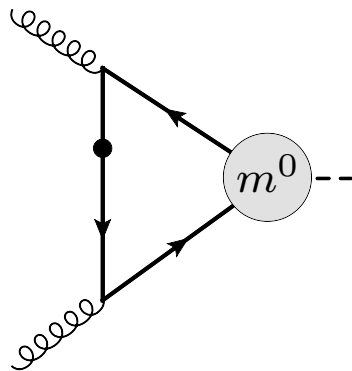
- *soft quarks exchange: single/triple*
- *topology: factorizable/nonfactorizable*

Triple soft quark exchange

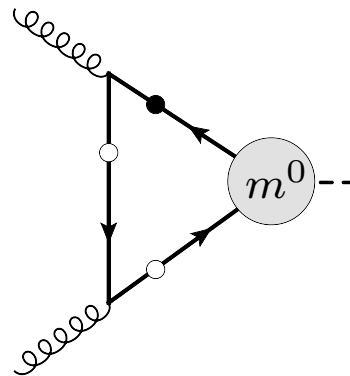


only planar graph contributes \Leftrightarrow *factorizable topology*

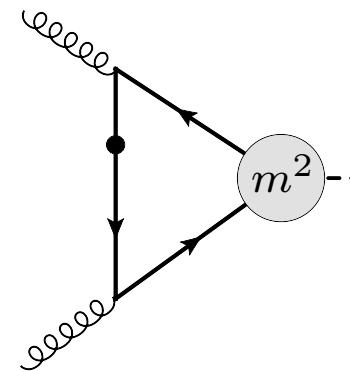
Factorizable contribution



NLP



NNLP (a)



NNLP (b)

- (a) power corrections to $gg \rightarrow q\bar{q}$ amplitude
 - non helicity-flip contribution for the first time
- (b) power corrections to the off-shell FF
 - can be inferred from on-shell result

Scalar form factor NLP LL

● On-shell FF

T. Liu, A.A. Penin, Phys.Rev.Lett. **119**, 262001 (2017)



$$F_S = Z_q^2 \sum_{n=0}^{\infty} \frac{m_q^{2n}}{Q^2} F_S^{(n)}, \quad F_S^{(1)} = -\frac{C_F T_F}{3} x^2 f(-z)$$

● Magic #2: universal function for all NLP form factors

$$f(z) = 1 + \frac{z}{5} + \frac{11}{420} z^2 + \frac{z^3}{378} + \dots$$

confirmed to 3 loops in M. Fael, et al. Phys.Rev.D **106**, 034029 (2022)

Scalar form factor NLP LL

On-shell FF

T. Liu, A.A. Penin, Phys.Rev.Lett. **119**, 262001 (2017)



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Magic #2: universal function for all NLP form factors

$$f(z) \sim 6 \left[\ln \left(\frac{z}{2} \right) + \gamma_E \right] \left(\frac{2\pi e^z}{z^5} \right)^{1/2}$$

$$f(-z) \sim \left[(\ln(2z) + \gamma_E)^2 - \frac{\pi^2}{2} \right] \frac{3}{z^2},$$

asymptotic behaviour at $z \rightarrow \infty$

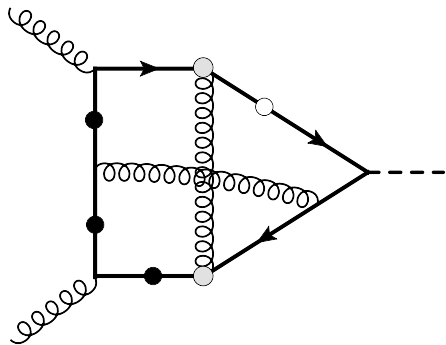
Nonfactorizable contribution

- *Abelian diagram*

- *eikonal gluon coupled to soft quark*



- *Non-Abelian case* $C_F \rightarrow C_F - C_A$



Results

T. Liu, S. Modi, A.A. Penin, JHEP **02**, 170 (2022)

● Scalar amplitude

$$M_{ggH}^q = Z_g^2 \ln^2 \left(\frac{m_q^2}{m_H^2} \right) \sum_{n=0}^{\infty} \left(\frac{m_q^2}{m_H^2} \right)^n M_{ggH}^{(n)}, \quad M_{ggH}^{(0)} = g(z)$$

● NNLP

$$M_{ggH}^{(1)} = \left[-4g(z) + \left(\frac{T_F C_F}{45} h(z) - \frac{(C_A - C_F)(C_A - 2C_F)}{9} j(z) \right) x^2 \right]$$

● three loops

large- N_c

$$M_{ggH}^{(1)} = \left[-4 - \frac{2}{3}(C_A - C_F)x + \left(\frac{T_F C_F}{45} - \frac{14}{45}C_F^2 + \frac{23}{45}C_F C_A - \frac{9}{45}C_A^2 \right) x^2 \right]$$

● agrees with num. calc. M.Czakon, M.Niggetiedt, JHEP **2005**, 149 (2020)

Results

● All-order

● factorizable single soft quark: $g(z) = {}_2F_2(1, 1; 3/2, 2; x/2)$

● factorizable triple soft quark

$$h(z) = 6! \int_0^1 d\eta \int_0^{1-\eta} d\xi \int_0^\eta d\eta_2 \int_0^\xi d\xi_2 \int_0^{\eta_2} d\eta_1 \int_0^{\xi_2} d\xi_1 e^{2z(\eta\xi - \eta_2\xi_2 + \eta_1\xi_1)}$$

● nonfactorizable

$$j(z) = 72 \int_0^1 d\eta \int_0^{1-\eta} d\xi \int_0^{1-\xi} d\eta_1 \int_0^{1-\eta_1-\xi} d\xi_1 \eta\xi_1 e^{2z\eta(\xi+\xi_1)} \\ \times \left[1 + \frac{e^{-2z\eta\xi} - 1}{2} + \frac{e^{-2z\eta\xi} - 1 + 2z\eta\xi}{4z\eta\xi_1} \right],$$

● $C_F \rightarrow C_F - C_A$ rule is still to be proven beyond three loops

Summary

- Bottom effect in Higgs boson production in gluon fusion (*NLP-NLL-threshold*)

➔ *uncertainty interval reduced to 0.40 pb (factor 2)*

- Power corrections in mass
 - *first ever NNLP LL result*
 - *small quark mass expansion parameter $4m_q^2/s$*