



# Beautiful and Charming Energy Correlators

Loopfest 2023 - SLAC

**Bianka Meçaj - Yale University**

Based on [2210.09311](#) and ongoing work



Evan Craft



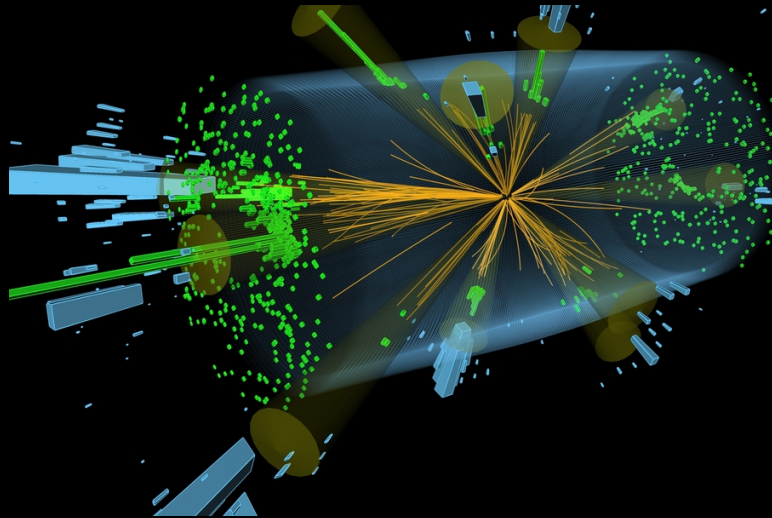
Kyle Lee



Ian Moulton

# QCD at Hadron Colliders

Almost every LHC event contains jets



Jets are reconstructed using jet algorithms (anti- $k_T$ )

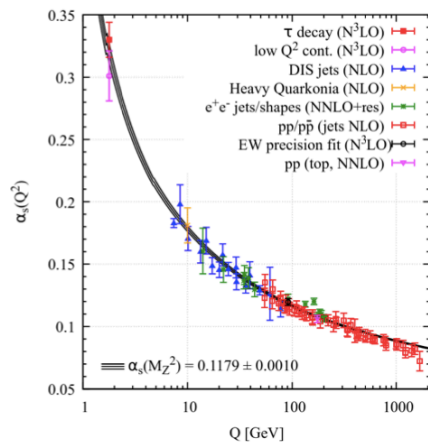
Cacciari, Salam 2006  
Salam, Soyez 2007

**How can we learn the most about underlying physics from the reconstructed jets?**

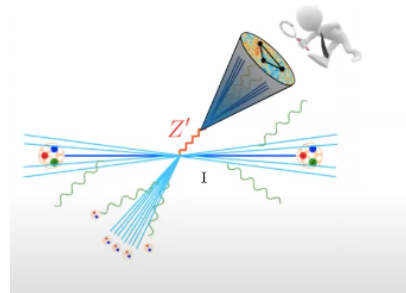
# Jets at the LHC

## Jet substructure

### QCD precision tests



### New Physics

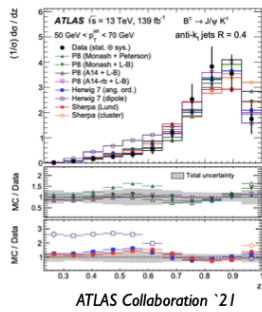


**Both precision measurements and New Physics searches require precise description of jet cross sections.**

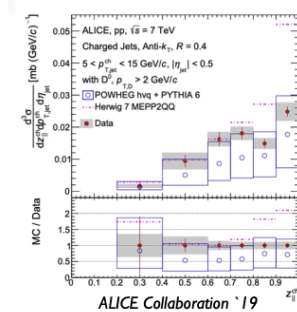
# Jets at the LHC

## Jet substructure and heavy quarks

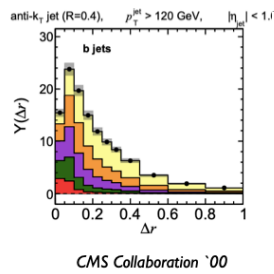
Many interesting processes include heavy quark effects:  $h \rightarrow b\bar{b}, h \rightarrow c\bar{c}$



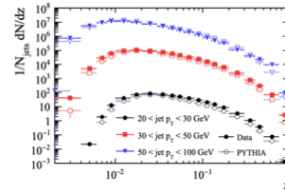
ATLAS



ALICE



CMS

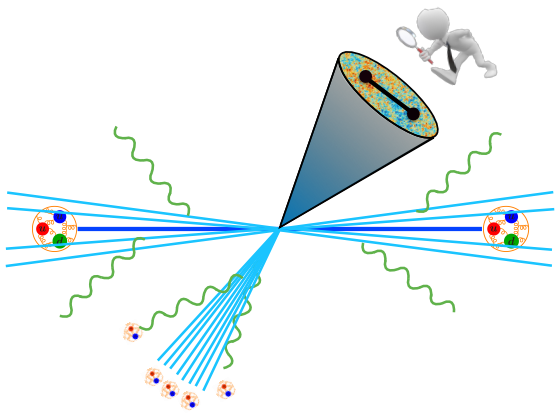


LHCb

Can we probe intrinsic mass effects with jet substructure?

# Jet substructure

Study the internal structure of a jet



Any physics dynamics will be imprinted in the energy distributions inside the jet.

## Well-defined in QFT!

- Distribution of energy inside the jet is described by correlation functions of the energy flow operators  $\Rightarrow$  energy correlators.

$$\langle \Psi | \varepsilon(\vec{n}_1) \varepsilon(\vec{n}_2) \dots \varepsilon(\vec{n}_n) | \Psi \rangle$$

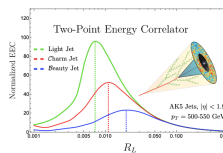
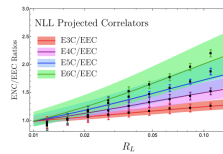
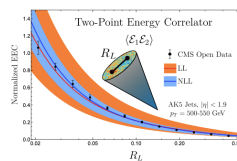
[Basham, Brown, Ellis, Love]

$$\mathcal{E}(\vec{n}) = \lim_{r \rightarrow \infty} \int_0^{\infty} dt r^2 n^i T_{0i}(t, r\vec{n})$$

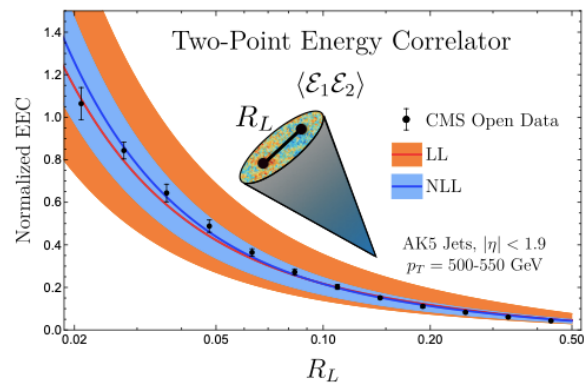
# Energy correlators for jet substructure at LHC

## Outline

- Scaling behavior
- Spectrum of the jet
- Heavy quark jets



# Scaling behavior



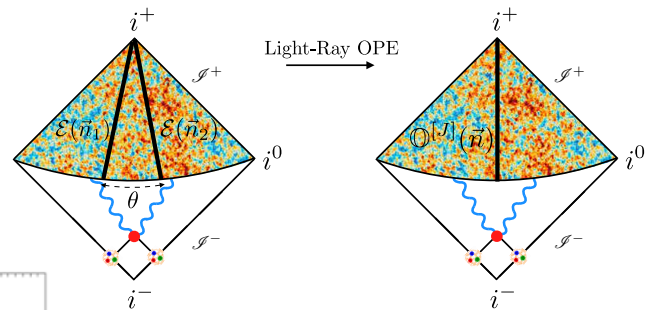
# Scaling behavior

We will study energy correlators inside high energy jets at the LHC: small angle behavior

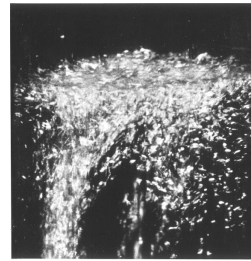
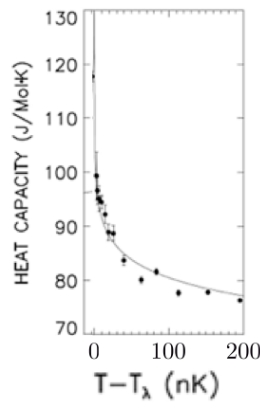
- Energy correlators admit an OPE

$$\langle \Psi | \varepsilon(\vec{n}_1) \varepsilon(\vec{n}_2) | \Psi \rangle \sim \sum \theta^{\gamma_i} \mathcal{O}_i(\vec{n}_1)$$

- Universal scaling behavior in QFT as operators are brought together!



[Hofman, Maldacena]  
[Chang, Kologlu, Kravchuk, Simmons Duffin, Zhiboedov]

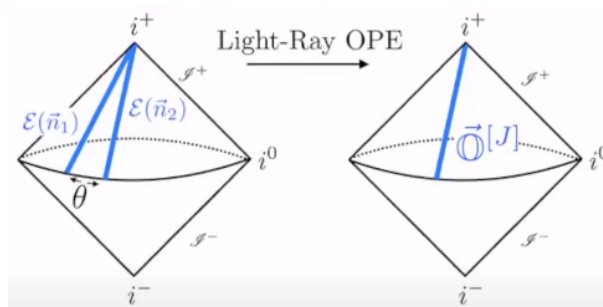




# The light-ray OPE

- The leading scaling behavior at the LHC is described by the leading terms in the OPE: **twist two light-ray operators**.
- Light-ray OPE is a rigorous and convergent expansion in CFT.

$$\varepsilon(\vec{n}_1) \varepsilon(\vec{n}_2) = \sum c_i \theta^{r_i-4} \mathbb{O}_i(\vec{n}_1)$$



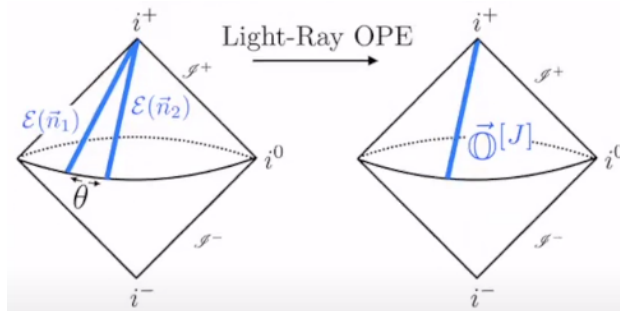
[Hofman, Maldacena]  
[Chang, Kologlu, Kravchuk, Simmons Duffin, Zhiboedov]

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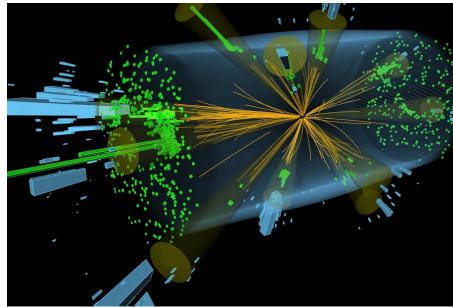
Have to understand the structure of these operators in QCD!



[Hofman, Maldacena]  
[Chang, Kologlu, Kravchuk, Simmons Duffin, Zhiboedov]

# Energy Correlators at the LHC

## Factorization Formula



$$\frac{d\Sigma}{dp_T d\eta dz} = \sum_i \mathcal{H}_i(p_T z, \eta, \mu) \otimes \int_0^1 dx x^N \mathcal{F}_{ij}(z, x, p_T R, \mu) J_j^{[N]}(z, x, \mu)$$

Hard function: includes pdfs

Matching coefficient, jet algorithm

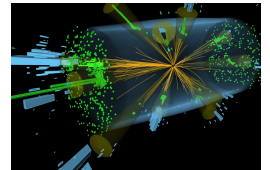
Can calculate any higher point correlator at the LHC

[Lee, Mecaj, Moult]

Energy correlator jet function

# Two-point energy correlator

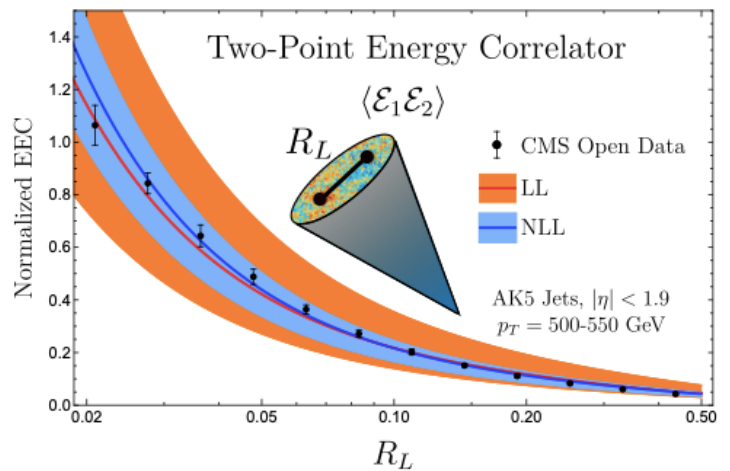
## The simplest jet substructure observable



- The complicated LHC environment is described by a simple observable!
- Probe the OPE structure of  $\langle \varepsilon(\vec{n}_1)\varepsilon(\vec{n}_2) \rangle$

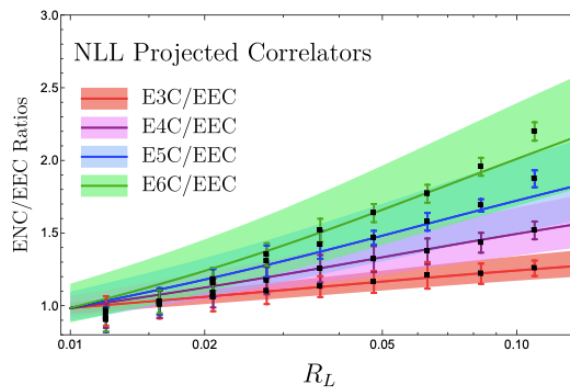
$$\langle \Psi | \varepsilon(\vec{n}_1)\varepsilon(\vec{n}_2) | \Psi \rangle \sim \sum \theta^i \mathcal{O}_i(\vec{n}_1)$$

- A jet substructure observable that can test quantum scaling behavior of operators.



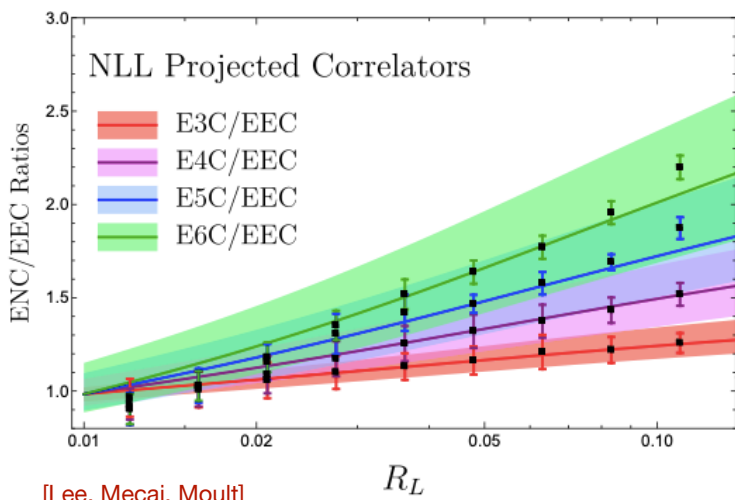
[Lee, Mecaj, Moult]

# Jet Spectrum



# The jet spectrum

## Higher-point correlators



[Lee, Mecej, Mout]

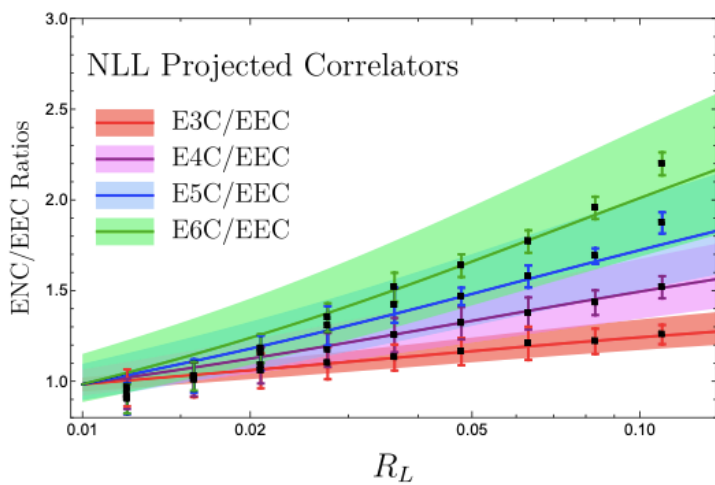
[Chen, Mout, Zhang, Zhu]

Asymptotic energy flux directly probes the spectrum of (twist-2) lightray operators at the quantum level!

$$\frac{\langle \mathcal{E}_1 \mathcal{E}_2 \dots \mathcal{E}_{J-1} \rangle}{\langle \mathcal{E}_1 \mathcal{E}_2 \rangle} \sim \frac{\langle \mathbb{O}^{[J]} \rangle}{\langle \mathbb{O}^{[3]} \rangle}$$

# The jet spectrum

## Higher-point correlators



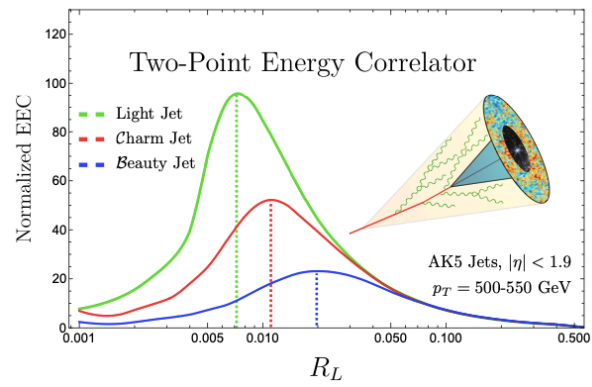
[Lee, Mecaj, Moul]t

[Chen, Moul, Zhang, Zhu]

$$\theta^r \rightarrow \exp\left(\frac{\hat{\gamma}}{2\beta_0} \ln \frac{\alpha_s(\theta Q)}{\alpha_s(Q)}\right)$$

- Can be observed at the high energies at the LHC at high precision
- Ratio of the higher-point correlators with the two-point isolates anomalous scaling!
- The anomalous scaling behavior depends on N (slope increases with N)
  - ↓
- First hand probe of the anomalous dimensions of QCD operators.
- Non-perturbative effects cancel in the ratio

# Heavy quark jets

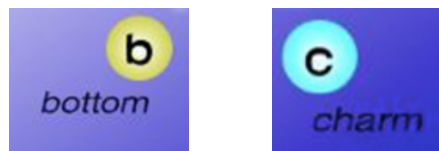




# Energy Correlators on heavy jets

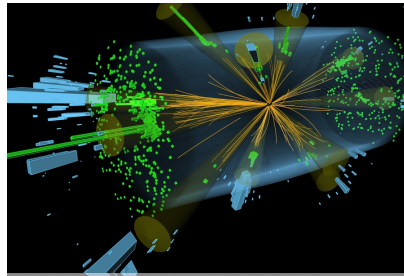
Introduce an additional scale

- At the LHC energies there is access to the transition phase from massless to massive behaviour  $\Rightarrow$  more complexity
- **Also very interesting!**
  - Can probe intrinsic mass effects of quarks before confinement into hadrons



# Factorization theorem

Can compute any higher point correlators on massive quarks at LHC at NLL



Hard function (NNLO)

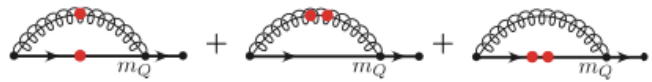
$$\Sigma^{[N]}(R_L, p_T^2, m_Q, \mu) = \int_0^1 dx x^N \underbrace{\vec{J}^{[N]}(R_L, x, m_Q, \mu)}_{\text{Massive Energy Correlator Jet Function (NLO)}} \cdot \underbrace{\vec{H}(x, p_T^2, \mu)}_{\text{Hard function (NNLO)}}$$

[Czakon, Generet, Mitov, Poncelet; 2021]

$$\begin{aligned} \mu_H &\sim p_T \\ \mu_J &\sim p_T R \end{aligned}$$

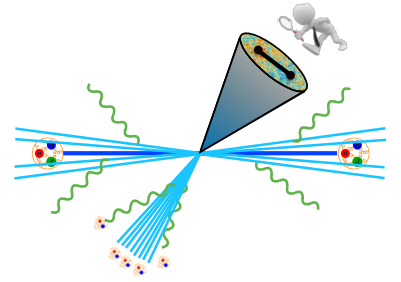
Massive Energy Correlator Jet Function (NLO)

[Craft, Lee, BM, Moulit]



# Heavy quark jet function

## Result



$$J_Q^{\text{bare}}(z, M, \mu) = \delta(z) \left( 1 + \frac{\alpha_s C_F}{4\pi} \left[ - \left( \gamma_{qq}^{(0)}(3) + \gamma_{gq}^{(0)}(3) \right) \left( \frac{1}{\epsilon_{\text{UV}}} + \ln \frac{\mu^2}{M^2} \right) - \frac{19}{6} \right] \right) \\ + \frac{\alpha_s C_F}{\pi} \frac{1}{z} \left[ \frac{3}{4} - \frac{5}{2} \delta^2 - \frac{\delta^4}{1 + \delta^2} + 3\delta^3 \arctan \left( \frac{1}{\delta} \right) + \frac{1}{2} \delta^2 (1 - \delta^2) \ln \frac{\delta^2}{1 + \delta^2} \right]$$

The mass should not affect the UV behavior of the jet function.

This can be seen from comparing the UV poles with the light quark jet function.

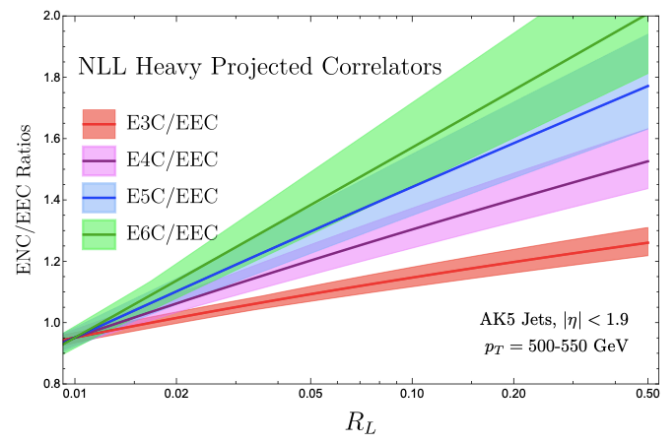
$$J_q^{\text{bare}}(z, \mu) = \delta(z) + \frac{\alpha_s C_F}{4\pi} \left[ \delta(z) \left( -\frac{3}{\epsilon_{\text{UV}}} - \frac{37}{3} \right) + 3 \frac{Q^2}{\mu^2} \mathcal{L}_0 \left( \frac{Q^2}{\mu^2} z \right) \right] \\ = \delta(z) + \frac{\alpha_s C_F}{4\pi} \left[ \delta(z) \left( - \left( \gamma_{qq}^{(0)}(3) + \gamma_{gq}^{(0)}(3) \right) \frac{1}{\epsilon_{\text{UV}}} - \frac{37}{3} \right) + 3 \frac{Q^2}{\mu^2} \mathcal{L}_0 \left( \frac{Q^2}{\mu^2} z \right) \right]$$

[Craft, Lee, BM, Moul]t

# Projected energy correlators

Resolve the UV scaling behaviour

- Ratios of higher point correlators with the two point EEC are independent of IR effects, including quark mass.
- The exact behavior as the massless case.
- Non-trivial cross check of the factorization theorem!
- Anomalous dimensions should not be affected by the IR physics.

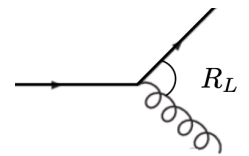


[Craft, Lee, BM, Moulton]

# Massive jets

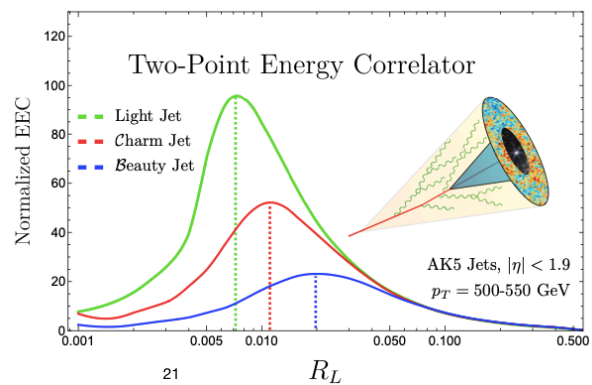
## Massive Energy Correlator Jet Function

$$\Sigma^{[N]}(R_L, p_T^2, m_Q, \mu) = \int_0^1 dx x^N \overbrace{\vec{J}^{[N]}(R_L, x, m_Q, \mu)} \cdot \underbrace{\vec{H}(x, p_T^2, \mu)}_{\text{Hard function}}$$



Virtuality  $\sim p_T R_L + m_Q^2$

- Formation time changes with the mass of the quark.
- Can clearly see this from the two-point EEC.

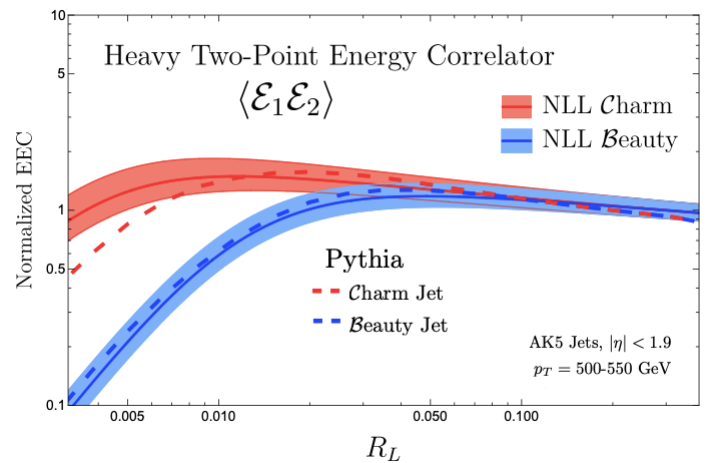


[Craft, Lee, BM, Moutl]

# Massive two point correlator

## First massive jet substructure observable at NLL

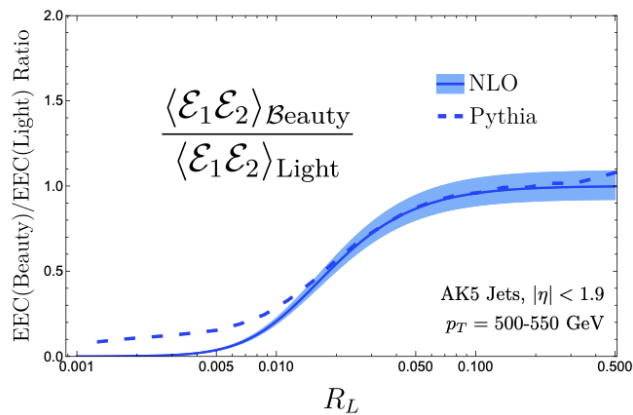
- Scaling behaviour identical to massless case for larger scales.
- A turn-over for  $R_L \rightarrow m_Q/p_T$
- The change in the slope is perturbative effect contrary to massless jets:  $R_L \rightarrow \Lambda_{QCD}/p_T$
- The turn-over region is of interest for improving heavy quark description in parton shower.



[Craft, Lee, BM, Moutl]

# Intrinsic mass effects

## Dead-cone effect

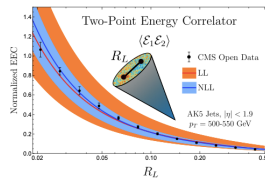


[Craft, Lee, BM, Moutl]

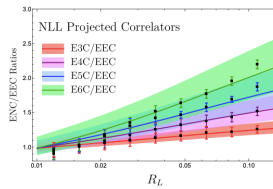
- Ratios of the massive and massless EEC isolate mass (IR) effects.
- A transition region related to the quark mass, which is perturbatively calculable.
- Excellent agreement with MC.
- Small angle suppression can be interpreted as a dead-cone effect.

# Conclusions

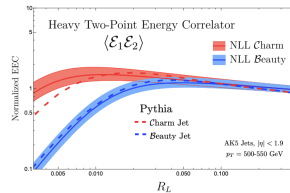
- Can probe a universal scaling behavior of QFT in the complicated LHC environment.



- Higher-point correlators are calculated for LHC and are a direct probe of anomalous scaling dimension of QCD operators.



- Energy Correlators for heavy quark jets probe intrinsic mass effects of elementary particles.





# What is next?

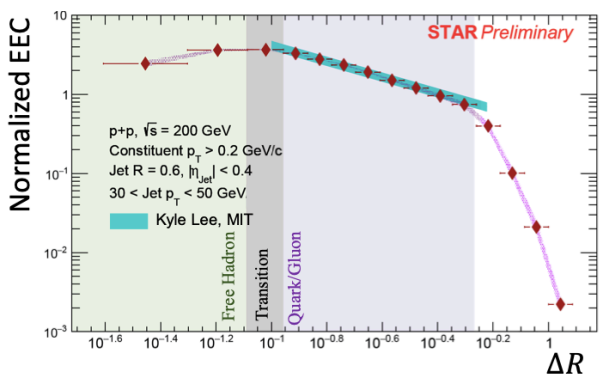
**Experimental Measurements for both light and heavy quark energy correlators.**

# Exciting experimental results!



Talk by N.Sahoo and A.Tamis at  
HARD PROBES-March 2023

- STAR collaboration  $\sqrt{s} = 200\text{GeV}$



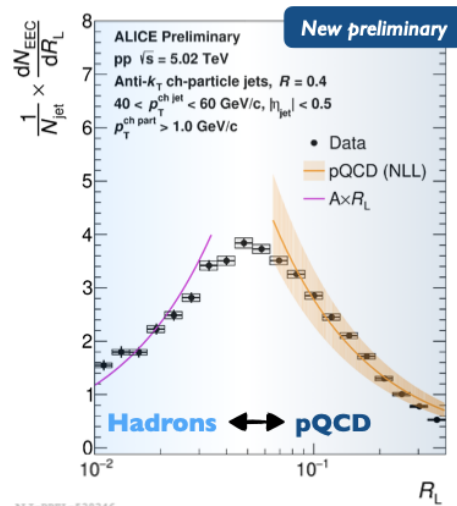
$$\text{Normalized EEC} = \frac{1}{\sum_{\text{jets}} \sum_{i \neq j} \frac{E_i E_j}{p_{T,i}^2 p_{T,j}^2}} \frac{d(\sum_{\text{jets}} \sum_{i \neq j} \frac{E_i E_j}{p_{T,i}^2 p_{T,j}^2})}{d(\Delta R)}$$

Direct observation of the transition from free hadrons to quarks/gluons at a universal scaling!



Talk by J.Mulligan and R.Cruz-Torres  
at HARD PROBES-March 2023

- ALICE collaboration  $\sqrt{s} = 5\text{TeV}$



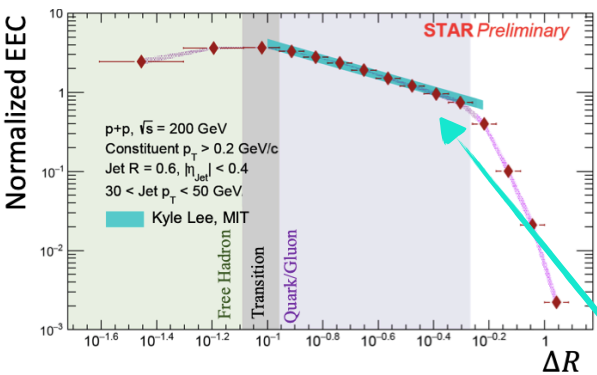
Universal behavior of the transition region.

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$$\text{Normalized EEC} = \frac{1}{\sum_{\text{jets}} \sum_{i \neq j} \frac{E_i E_j}{p_{T,i}^2}} \frac{d(\sum_{\text{jets}} \sum_{i \neq j} \frac{E_i E_j}{p_{T,i}^2})}{d(\Delta R)}$$

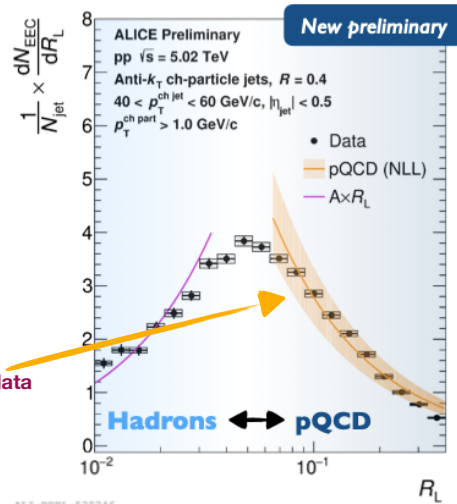
Direct observation of the transition from free hadrons to quarks/gluons at a universal scaling!

Excellent agreement of data with the theory NLL calculation



Talk by J.Mulligan and R.Cruz-Torres at HARD PROBES-March 2023

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Universal behavior of the transition region.

**Thank you!**