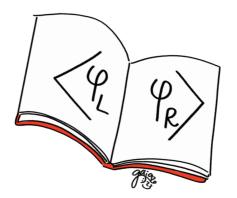
Intersection numbers via p(z)-adic expansions



Amplitudes, CERN, August 7, 2023



Based on arXiv:2304.14336 in collaboration with Gaia Fontana

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Loop integrals

- Loop integrals are the building blocks of perturbative QFT
 - essential for a deeper understanding of **amplitudes**
 - a key component of phenomenological predictions

$$I = \int_{-\infty}^{\infty} \left(\prod_{i=1}^{\ell} d^d k_i \right) \frac{1}{D_1^{\alpha_1} D_2^{\alpha_2} \cdots D_n^{\alpha_n}}, \quad \alpha_j \leq 0$$

Inverse propagators
$$D_j = l_j^2 - m_j^2$$

+ auxiliaries

Integral decomposition

Chetyrkin, Tkachov (1981), Laporta (2000)

• Feynman integrals obey linear relations, e.g. IBPs

$$\int \left(\prod_{j} d^{d} k_{j}\right) \frac{\partial}{\partial k_{j}^{\mu}} v^{\mu} \frac{1}{D_{1}^{\alpha_{1}} D_{2}^{\alpha_{2}} \cdots} = 0, \qquad v^{\mu} \in \{p_{i}^{\mu}, k_{i}^{\mu}\}$$

- Very large and sparse linear system
- Solution = reduction into a basis of linearly independent master integrals (MIs) $\{G_i\} \subset \{I_i\}$

$$I_{j} = \sum c_{jk}G_{k}$$
 master integrals

IBP reduction

- An essential ingredient of higher-order computations...
- ...but one of the main bottlenecks

Recent improvements:

Finite fields and rational reconstruction [Kant (2014), von Manteuffel, Schabinger (2014), T.P. (2016)]

- **reconstruct** results from numerical evaluations modulo a prime
- delay reconstruction to amplitude coefficients
- pushed state of art of modern amplitude calculations

Symbolic solutions

- reduction rules for symbolic exponents
- hybrid methods (e.g. syzygy eq.s + Laporta) (see e.g. Gluza, Kosower (2010))

See also recent developments in:

Fire, Kira+FireFly, FiniteFlow, NeatIBP, Blade, ...

"Direct" decomposition methods

Goal

Seek a more "direct" way of projecting out the coefficients of integral decomposition

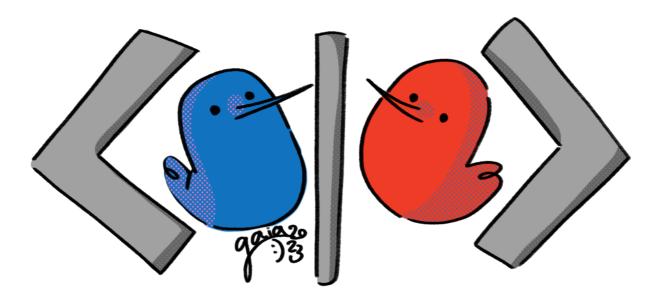
bypass the solution of large systems

potential for performance gain

but not quite there yet… ⁻∖_(ツ)_/⁻ investigate the **vector-space** structure obeyed by **loop** integrals in a family

connections with new areas of mathematics (intersection theory)

Intersection theory



The main idea

- Reinterpret Feynman integrals as elements of a vector space
- Master integrals (MIs) form a basis w.r.t. IBP relations

$$I = \sum_{j} c_{j} G_{j}$$

- Define scalar products (intersection numbers)
 - they must be consistent with IBPs!
- Project integrals into their components c_i w.r.t. the basis

$$c_j = \sum_k (G^{-1})_{jk} (G_k \cdot I)$$
 with $G_{jk} \equiv G_j \cdot G_k$

The vector space

and

Mizera (2018), Mastrolia, Mizera (2019), Frellesvig, Gasparotto, Mandal, Mastrolia, Mattiazzi, Mizera (2019)

• We consider integrals (or right integrals)

$$|\varphi_{R}\rangle = \int dz_{1} \cdots dz_{n} \frac{1}{u(\mathbf{z})} \varphi_{R}(\mathbf{z})$$

and **dual integrals (or left integrals)**
$$\langle \varphi_{L}| = \int dz_{1} \cdots dz_{n} u(\mathbf{z}) \varphi_{L}(\mathbf{z})$$

$$u(\mathbf{z}) = \prod_{j} B_{j}(\mathbf{z})^{\gamma_{j}}$$

• multivalued function
• regulates the singularities of $\varphi_{R,L}$
• $B_{j} = polynomials, \gamma_{j} = generic exponents$

• consider a set a integrals with same u(z) and integration domain, but different $\varphi(z)$ (integral family)

IBPs

• We assume regulated integrands to vanish at integration boundary

$$\sum_{j=1}^{n} \int dz_1 \cdots dz_n \,\partial_{z_j} \left(\frac{1}{u} \,\xi_j^{(R)} \right) = 0, \qquad \sum_{j=1}^{n} \int dz_1 \cdots dz_n \,\partial_{z_j} \left(u \,\xi_j^{(L)} \right) = 0$$
$$\sum_{j=1}^{n} \left| \left(\partial_{z_j} - (\partial_{z_j} u)/u \right) \xi_j^{(R)} \right\rangle = 0, \qquad \sum_{j=1}^{n} \left\langle \left(\partial_{z_j} + (\partial_{z_j} u)/u \right) \xi_j^{(L)} \right| = 0$$

- we can formally define the vector space via these equations*
- reduction to **bases** $\{ |e_j^{(R)}\rangle \}_{j=1}^{\nu}$ and $\{ \langle e_j^{(L)} | \}_{j=1}^{\nu} \text{ of MIs} independent modulo IBPs (<math>\nu = \text{dimension of vector space})$

* Additional identities, such as some symmetry relations, may exist but are formally not taken into account at this stage. They can be easily identified and implemented after the decomposition via IBPs.

Intersection numbers

Intersection numbers are rational scalar products btw integrals and their duals

$$\begin{split} |\varphi_{R}\rangle &= \sum_{i=1}^{\nu} c_{i}^{(R)} |e_{i}^{(R)}\rangle \\ \hline \textbf{rational coefficients} \\ c_{i}^{(R)} &= \sum_{j=1}^{\nu} (\mathbf{C}^{-1})_{ij} \langle e_{j}^{(L)} | \varphi_{R}\rangle \\ \hline \textbf{metric} \\ \mathbf{C}_{ij} &\equiv \langle e_{i}^{(L)} | e_{j}^{(R)}\rangle \end{split}$$

 $\langle \varphi_L | \varphi_R \rangle$

Computing intersection numbers Univariate case

- **One-fold** integrals $|\varphi_R\rangle = \int dz \frac{1}{u(z)} \varphi_R(z), \quad \langle \varphi_L| = \int dz u(z) \varphi_L(z)$
- Intersection numbers: $\langle \varphi_L | \varphi_R \rangle = \sum_{p \in \mathscr{P}_{\omega}} \operatorname{Res}_{z=p}(\psi \varphi_R)$
 - ψ local solution of the DE $(\partial_z + \omega)\psi = \varphi_L$, $\omega \equiv (\partial_z u)/u$
 - $\mathscr{P}_{\omega} = \left\{ z \mid z \text{ is a pole of } \omega \right\} \bigcup \{\infty\}$
 - solution for ψ as Laurent series around each pole

$$\psi = \sum_{i=min}^{max} c_i (z-p)^i + O\left((z-p)^{max+1}\right)$$

find c_i via linear algebra

Computing intersection numbers Multivariate case

Recursive algorithm

$$\varphi_R \rangle = \int dz_n \, |\varphi_R\rangle_{n-1}, \qquad \langle \varphi_L | = \int dz_n \, \langle \varphi_R |_{n-1}$$

- Similar procedure:
 - local solutions of DEs around poles of rational functions (system of DEs)
 - sums of residues at poles
 - depends on (n 1)-variate intersection numbers $\langle \varphi_L | \varphi_R \rangle_{n-1}$ and decompositions of (n - 1)-fold integrals

$$\varphi_R \rangle_{n-1} = \sum_{j=1}^{\nu_{(n-1)}} \varphi_{R,j} | e_j^{(R)} \rangle_{n-1}$$
(*n*-1)-fold basis of MIs

Application to loop integrals

• Use e.g. the Baikov representation

$$I = \int \left(\prod_{i=1}^{\ell} d^d k_i \right) \frac{1}{D_1^{\nu_1} \cdots D_n^{\nu_n}} = K \int dz_1 \cdots dz_n B(\mathbf{z})^{\gamma} \frac{1}{z_1^{\alpha_1} \cdots z_n^{\alpha_n}}$$

 $B = Baikov polynomial, \quad \gamma = (d - \ell - e - 1)/2$

Identifications:

$$I \sim |\varphi\rangle = \int d^n z \, \frac{1}{u} \, \varphi$$

 $\varphi = \frac{1}{z_1^{\alpha_1} \cdots z_n^{\alpha_n}},$

$$u = B^{-\gamma} \prod_{\substack{j=1 \\ j=1 \\ \text{ analytic regulators} \\ \text{ of } z_j \to 0 \text{ singularities}}}^n$$

take $\rho_j \rightarrow 0$ limit after the decomposition

Pros/Cons

:) Pros

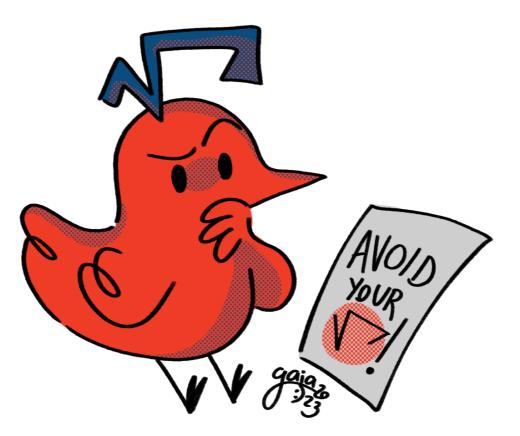
- makes vector-space structure of families of loop integrals manifest
- "direct" decomposition (not a byproduct of solving a huge system of identities)

:(Cons

- irrational contributions in intermediate stages of calculation*
 - algebraic bottleneck
 - no easy implementation over finite fields
- need of analytic regulators ρ_i

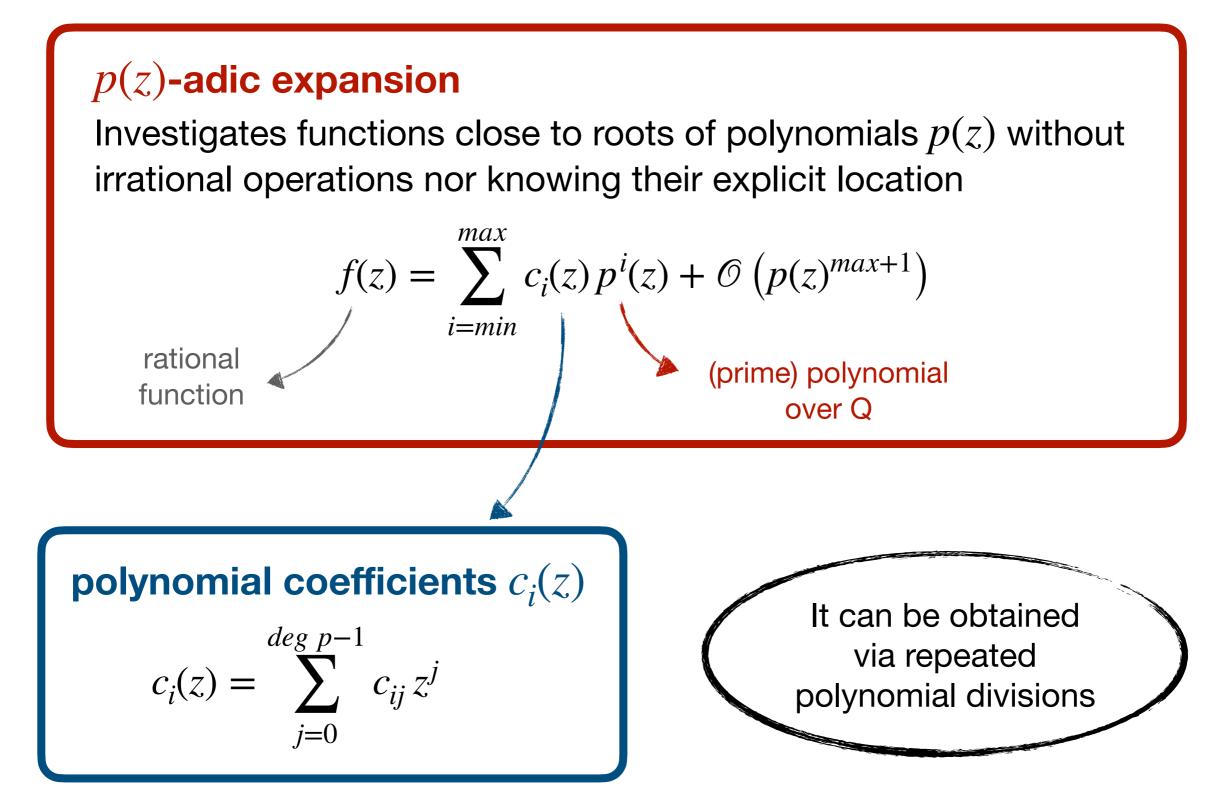
* A way out is a reduction to simple poles [Weinzierl (2020)] but requires non-trivial sequences of changes of bases and integral transformations

p(z)-adic expansions

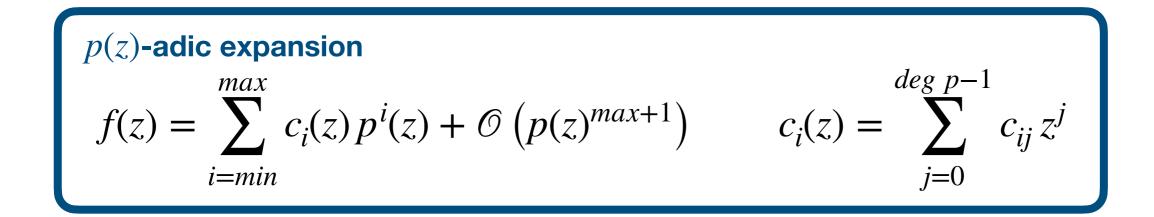


Polynomial expansions

G. Fontana, T.P. (2023)



p(z)-adic expansions and residues



Univariate global residue theorem (generalization)

Taking a sum of residues of f(z) at the roots of p(z) from their p(z)-adic expansion is trivial and does not require knowing their location

$$\operatorname{Res}_{p(z)}(f(z)) \equiv \sum_{\substack{y \mid p(y)=0}} \operatorname{Res}_{z=y}(f(z)) = \frac{c_{-1,deg \ p-1}}{l_c}$$
$$l_c \equiv \text{leading coefficient of } p(z)$$

Back to intersection numbers

Sum over "denominator factors" rather than "poles"

$$\langle \varphi_L | \varphi_R \rangle = \sum_{p(z) \in \mathscr{P}_{\omega}[z]} \operatorname{Res}_{p(z)} \left(\psi \varphi_R \right)$$

•
$$\mathscr{P}_{\omega}[z] = \{\text{factors of the denominator of } \omega\} \bigcup \{\infty\}$$

- ψ local solution of the DE $(\partial_z + \omega)\psi = \varphi_L$, $\omega \equiv (\partial_z u)/u$
- solution for ψ as p(z)-adic series expasion around each factor

$$\psi = \sum_{i=min}^{max} \sum_{j=0}^{deg} \sum_{j=0}^{p(z)-1} c_{ij} z^j p^i(z) + \mathcal{O}\left(p(z)^{max+1}\right)$$

find c_{ij} via linear algebra

... and similar for multivariate case

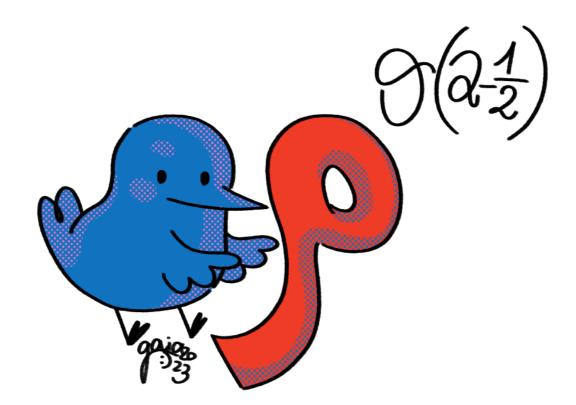
Int. numbers via p(z)-adic expansion

- The p(z)-adic expansion method
 - yields a fully rational algorithm for computing int. numbers
 - no integral transformation or change of basis needed

- Proof-of-concept implementation over finite fields
 - using FiniteFlow [T.P. (2019)] (in Mathematica)
 - delay full kinematic reconstruction
 - most operations recast as linear algebra problems

More details on Gaia's poster!

Dual integrals and analytic regulators



Analytic regulators

Analytic regulators ρ_j , regulate integrands $\varphi_L \sim \frac{1}{z_i}$

$$u = B^{-\gamma} \prod_{j=1}^{n} z_{j}^{\rho_{j}}$$

- DE for ψ has otherwise no solution - if $\varphi_L \sim 1/z^{v_j}$ ($v_j > 0$) then $\psi \sim 1/\rho_j$

• limit $\rho_i \rightarrow 0$ after the decomposition

Drawbacks

- additional variables in intermediate stages
- obscures block-triangular structure of decompositions
- more master integrals in intermediate steps of recursion

Dual integrals

Recall the decomposition:

$$|\varphi_{R}\rangle = \sum_{i=1}^{\nu} c_{i}^{(R)} |e_{i}^{(R)}\rangle, \qquad c_{i}^{(R)} = \sum_{j=1}^{\nu} (\mathbf{C}^{-1})_{ij} \langle e_{j}^{(L)} |\varphi_{R}\rangle$$

Observation

Coefficients $c_i^{(R)}$ are **independent** of the **choice** of the **dual basis** of "left" integrals $\{\langle e_j^{(L)} | \}_{j=1}^{\nu}\}$

Idea Exploit the freedom of choice of the dual basis to simplify the calculation

Choice of dual integrals

Two interesting approaches (different formalisms, similar outcomes):

- Alternative formalism for defining dual space of loop integrals
 [Caron-Huot, Pokraka (2021)]
- Simple choice of dual integrals

Choose dual integrals of the form

[G. Fontana, T.P. (2023)]

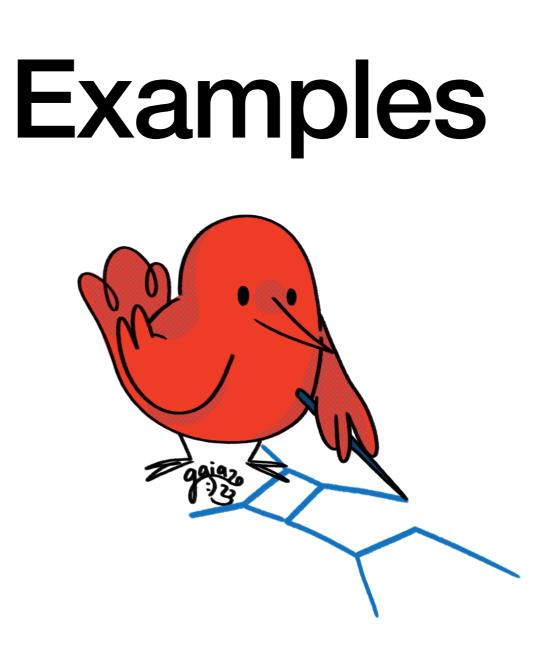
$$\varphi_L(\mathbf{z}) = \rho_1^{\Theta(\alpha_1 - \frac{1}{2})} \cdots \rho_n^{\Theta(\alpha_n - \frac{1}{2})} \frac{1}{z_1^{\alpha_1} \cdots z_n^{\alpha_n}}$$

- if there's a **denominator** factor $z_i^{\alpha_j}$ (with $\alpha_j > 0$), **multiply** by ρ_j
- systematically work in the limit $\rho_j \rightarrow 0$ (i.e. only keep leading terms in a $\rho_i \rightarrow 0$ expansion in each step)

Advantages

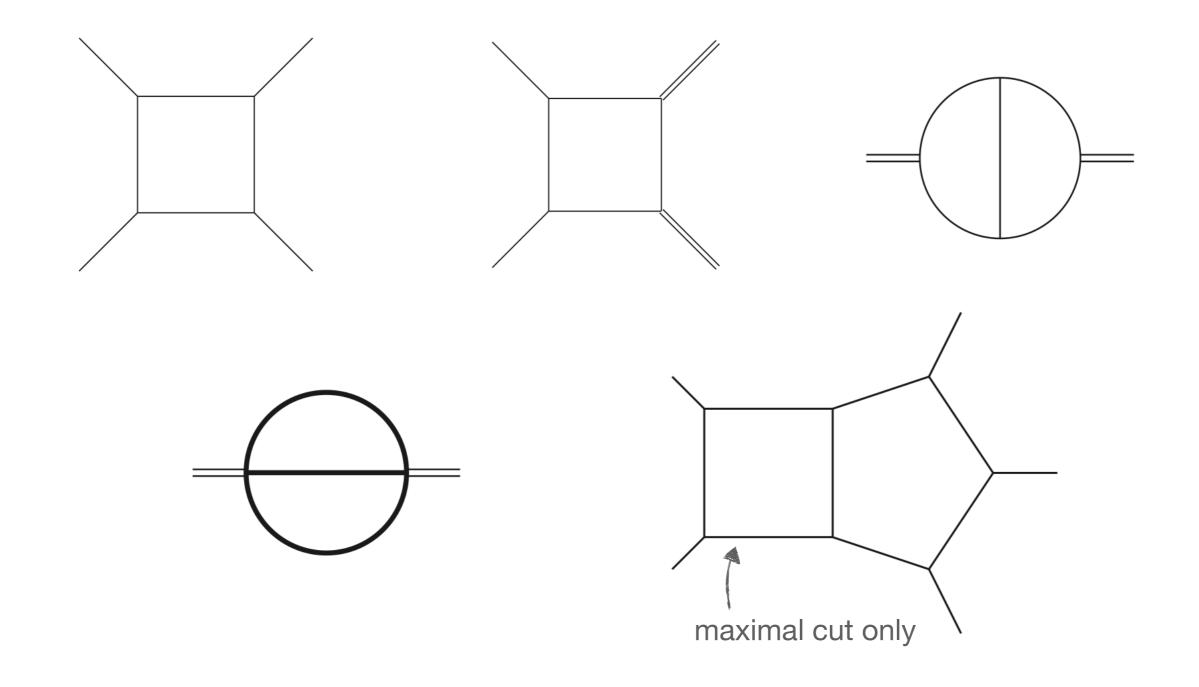
- Calculation "effectively" independent of ρ_i
 - working on leading coefficients in $\rho_j \rightarrow 0$ expansion, often just one
 - over finite fields, never sample or reconstruct ρ_i dependence
- Drastically **simpler** intermediate expressions
- Metric and reduction tables are **block triangular** (blocks ~ sectors)

- Many intersection numbers and contributions of poles to them vanish (ρ_j prefactors must cancel a $1/\rho_j$ singularities, only possible at $z_j \sim 0,\infty$)
- Fewer master integrals in intermediate steps of recursion



Simple examples and checks

...just checking that things work as expected!



Conclusions & Outlook

- Intersection theory unveils new mathematical structures in loops
- p(z)-adic expansions simplify study of functions close to singular points
 - avoid algebraic extensions
 - no need to know explicit location of irrational poles
 - avoid bottlenecks and enable finite field technologies

- Future directions:
 - simplifications/optimizations and application to different integral representations (loop-by-loop Baikov, Lee-Pomeransky)
 - Non-recursive multivariate generalization (based on Chestnov, Frellesvig, Gasparotto, Mandal, Mastrolia (2022))