## **Intersection numbers via** *p*(*z*)**-adic expansions**



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### **Loop integrals**

- **Loop integrals** are the **building blocks** of **perturbative QFT**
	- essential for a deeper understanding of **amplitudes**
	- a key component of **phenomenological predictions**

$$
I = \int_{-\infty}^{\infty} \left( \prod_{i=1}^{e} d^d k_i \right) \frac{1}{D_1^{\alpha_1} D_2^{\alpha_2} \cdots D_n^{\alpha_n}}, \qquad \alpha_j \le 0
$$
  
Inverse propagators  

$$
D_j = l_j^2 - m_j^2
$$

$$
+ \text{auxiliaries}
$$

### **Integral decomposition**

**Chetyrkin, Tkachov (1981), Laporta (2000)**

• Feynman integrals obey linear relations, e.g. **IBPs**

$$
\int \left(\prod_j d^d k_j\right) \frac{\partial}{\partial k_j^{\mu}} \nu^{\mu} \frac{1}{D_1^{\alpha_1} D_2^{\alpha_2} \cdots} = 0, \qquad \nu^{\mu} \in \{p_i^{\mu}, k_i^{\mu}\}
$$

- Very **large** and sparse linear system
- Solution = **reduction** into a **basis** of linearly independent  $\textsf{master}\ \textsf{integrals}\ (\textsf{MIs}\ \{G_j\}\subset \{I_j\}$

$$
I_j = \sum c_{jk} G_k
$$

### **IBP reduction**

- An **essential** ingredient of higher-order computations…
- …but one of the **main bottlenecks**

Recent improvements:

**Finite fields and rational reconstruction [Kant (2014), von Manteuffel, Schabinger (2014), T.P. (2016)]**

- **reconstruct** results from numerical evaluations modulo a prime
- delay reconstruction to amplitude coefficients
- pushed state of art of modern amplitude calculations

#### **Symbolic solutions**

- **reduction rules** for **symbolic exponents**
- hybrid methods (e.g. syzygy eq.s + Laporta) (see e.g. **Gluza, Kosower (2010)**)

See also recent developments in:

**Fire, Kira+FireFly, FiniteFlow, NeatIBP, Blade, …**

### **"Direct" decomposition methods**

#### **Goal**

Seek a more "direct" way of projecting out the coefficients of integral decomposition

**bypass** the **solution** of large **systems**

> potential for performance gain

> > but not quite there yet… ¯\\_(ツ)\_/¯

investigate the **vector-space** structure obeyed by **loop integrals** in a family

**connections** with new areas of **mathematics**  (**intersection theory**)

# Intersection theory



### **The main idea**

- Reinterpret Feynman integrals as elements of a **vector space**
- **Master integrals (MIs)** form a **basis** w.r.t. IBP relations

$$
I = \sum_j c_j G_j
$$

- Define **scalar products** (**intersection numbers**)
	- they must be consistent with IBPs!
- **Project** integrals into their components  $c_j$  w.r.t. the **basis**

$$
c_j = \sum_k (G^{-1})_{jk} (G_k \cdot I) \qquad \text{with } G_{jk} \equiv G_j \cdot G_k
$$

#### **The vector space**

**Mizera (2018), Mastrolia, Mizera (2019), Frellesvig, Gasparotto, Mandal, Mastrolia, Mattiazzi, Mizera (2019)**

• We consider **integrals** (or **right integrals**)

$$
|\varphi_R\rangle = \int dz_1 \cdots dz_n \frac{1}{u(\mathbf{z})} \varphi_R(\mathbf{z})
$$
  
\n• and **dual integrals** (or **left integrals**)  
\n
$$
\langle \varphi_L | = \int dz_1 \cdots dz_n u(\mathbf{z}) \varphi_L(\mathbf{z})
$$
 function  
\n
$$
u(\mathbf{z}) = \prod_j B_j(\mathbf{z})^{\gamma_j}
$$
  
\n• *multivalued function*  
\n• *regularities of*  $\varphi_{R,L}$   
\n•  $B_j = \text{polynomials}, \gamma_j = \text{generic exponents}$ 

• consider a set a integrals with same  $u(\mathbf{z})$  and integration domain, but different  $\varphi(\mathbf{z})$  (integral family)

#### **IBPs**

• We assume regulated integrands to vanish at integration boundary

$$
\sum_{j=1}^{n} \int d z_1 \cdots d z_n \partial_{z_j} \left( \frac{1}{u} \xi_j^{(R)} \right) = 0, \qquad \sum_{j=1}^{n} \int d z_1 \cdots d z_n \partial_{z_j} \left( u \xi_j^{(L)} \right) = 0
$$
  

$$
\sum_{j=1}^{n} \left| \left( \partial_{z_j} - (\partial_{z_j} u) / u \right) \xi_j^{(R)} \right| = 0, \qquad \sum_{j=1}^{n} \left\langle \left( \partial_{z_j} + (\partial_{z_j} u) / u \right) \xi_j^{(L)} \right| = 0
$$

- we can formally define the vector space via these equations\*
- reduction to bases  $\{~|~e^{(R)}_j\rangle\}_{j=1}^\nu$  and  $\{~\langle e^{(L)}_j|~\}_{j=1}^\nu$  of MIs independent modulo IBPs  $(\nu=$  dimension of vector space)

\* Additional identities, such as some symmetry relations, may exist but are formally not taken into account at this stage. They can be easily identified and implemented after the decomposition via IBPs.

### **Intersection numbers**

• **Intersection numbers** are **rational scalar products** btw integrals and their duals

• They project out integrals into their IBP decomposition

$$
|\varphi_R\rangle = \sum_{i=1}^{\nu} c_i^{(R)} |e_i^{(R)}\rangle
$$
\n
$$
c_i^{(R)} = \sum_{j=1}^{\nu} (\mathbf{C}^{-1})_{ij} \langle e_j^{(L)} | \varphi_R \rangle
$$
\n
$$
\mathbf{C}_{ij} \equiv \langle e_i^{(L)} | e_j^{(R)} \rangle
$$
\n
$$
\mathbf{C}_{ij} \equiv \langle e_i^{(L)} | e_j^{(R)} \rangle
$$
\n
$$
\mathbf{C}_{ij} \equiv \langle e_i^{(L)} | e_j^{(R)} \rangle
$$
\n
$$
\mathbf{C}_{ij} \equiv \langle e_i^{(L)} | e_j^{(R)} \rangle
$$

 $\langle \varphi_L | \varphi_R \rangle$ 

#### **Computing intersection numbers Univariate case**

- **One-fold** integrals  $|\varphi_R\rangle = \int dz$ 1  $u(z)$   $\varphi_R(z)$ ,  $\langle \varphi_L | = \int dz u(z) \varphi_L(z)$
- Intersection numbers:  $\langle \varphi_L | \varphi_R \rangle = \sum' \text{Res}_{z=p}(\psi \varphi_R)$  $p \in \mathscr{P}_\omega$ 
	- $\psi$  local solution of the DE  $(\partial_z + \omega)\psi = \varphi_L$ ,  $\omega \equiv (\partial_z u)/u$
	- $\mathscr{P}_{\omega} = \{ z | z \text{ is a pole of } \omega \} \bigcup {\{\infty\}}$
	- solution for  $\psi$  as Laurent series around each pole

$$
\psi = \sum_{i=\text{min}}^{\text{max}} c_i (z - p)^i + O((z - p)^{\text{max}+1})
$$
find  $c_i$  via linear algebra

#### **Computing intersection numbers Multivariate case**

• **Recursive** algorithm

$$
|\varphi_R\rangle = \int dz_n |\varphi_R\rangle_{n-1}, \qquad \langle \varphi_L| = \int dz_n \langle \varphi_R|_{n-1}
$$
\n
$$
\langle n-1\rangle \text{-fold integrals}
$$

- Similar procedure:
	- **local solutions** of **DEs** around **poles** of rational functions (system of DEs)
	- sums of **residues** at poles
	- depends on  $(n 1)$ -variate intersection numbers  $\langle \varphi_L | \varphi_R \rangle_{n-1}$ and decompositions of  $(n-1)$ -fold integrals

$$
|\varphi_R\rangle_{n-1} = \sum_{j=1}^{\nu_{(n-1)}} \varphi_{R,j} |e_j^{(R)}\rangle_{n-1}
$$
  
\n
$$
\longrightarrow
$$

# **Application to loop integrals**

• Use e.g. the **Baikov representation**

$$
I = \int \left( \prod_{i=1}^{\ell} d^d k_i \right) \frac{1}{D_1^{\nu_1} \cdots D_n^{\nu_n}} = K \int dz_1 \cdots dz_n B(\mathbf{z})^{\nu} \frac{1}{z_1^{\alpha_1} \cdots z_n^{\alpha_n}}
$$

*B* = Baikov polynomial,  $\gamma = (d - \ell - e - 1)/2$ 

• Identifications:

$$
I \sim |\varphi\rangle = \int d^n z \frac{1}{u} \varphi
$$

 $\varphi =$ 1  $z_1^{\alpha_1} \cdots z_n^{\alpha_n}$ 

$$
u = B^{-\gamma} \prod_{j=1}^{n} z_j^{\rho_j}
$$
  
analytic regulators  
of  $z_i \rightarrow 0$  singularities

take  $\rho_i \rightarrow 0$  limit after the decomposition  $\rho_j \rightarrow 0$ 

### **Pros/Cons**

#### **:) Pros**

- makes vector-space structure of families of loop integrals manifest
- "**direct**" decomposition (not a byproduct of solving a huge system of identities)

#### **:( Cons**

- **irrational** contributions in **intermediate** stages of calculation\*
	- algebraic bottleneck
	- no easy implementation over **finite fields**
- need of **analytic regulators** *ρ<sup>j</sup>*

\* A way out is a reduction to simple poles **[Weinzierl (2020)]** but requires non-trivial sequences of changes of bases and integral transformations

# *p*(*z*)-adic expansions



## **Polynomial expansions**

**G. Fontana, T.P. (2023)**



#### *p*(*z*)**-adic expansions and residues**



#### **Univariate global residue theorem (generalization)**

Taking a sum of residues of  $f(z)$  at the roots of  $p(z)$  from their  $p(z)$ -adic expansion is trivial and does not require knowing their location

$$
\operatorname{Res}_{p(z)} (f(z)) \equiv \sum_{y \mid p(y)=0} \operatorname{Res}_{z=y} (f(z)) = \frac{c_{-1,deg p-1}}{l_c}
$$
  
\n
$$
l_c \equiv \text{leading coefficient of } p(z)
$$

### **Back to intersection numbers**

Sum over "denominator factors" rather than "poles"

$$
\langle \varphi_L | \varphi_R \rangle = \sum_{p(z) \in \mathcal{P}_\omega[z]} \text{Res}_{p(z)} \left( \psi \varphi_R \right)
$$

- $\mathscr{P}_{\omega}[z] = \{\text{factors of the denominator of } \omega\} \bigcup {\infty}$
- $\psi$  local solution of the DE  $(\partial_z + \omega)\psi = \varphi_L$ ,  $\omega \equiv (\partial_z u)/u$
- solution for  $\psi$  as  $p(z)$ -adic series expasion around each factor

$$
\psi = \sum_{i=\text{min}}^{\text{max}} \sum_{j=0}^{\text{deg } p(z)-1} c_{ij} z^j p^i(z) + \mathcal{O}\left(p(z)^{\text{max}+1}\right)
$$
\nfind  $c_{ij}$  via linear algebra

… and similar for multivariate case

#### **Int. numbers via** *p*(*z*)**-adic expansion**

- The  $p(z)$ -adic expansion method
	- yields a fully **rational algorithm** for computing int. numbers
	- no integral transformation or change of basis needed

- Proof-of-concept **implementation** over **finite fields** 
	- using **FiniteFlow [T.P. (2019)]** (in Mathematica)
	- delay full kinematic reconstruction
	- most operations recast as linear algebra problems

#### ➡ **More details on Gaia's poster!**

# Dual integrals and analytic regulators



# **Analytic regulators**

Analytic regulators  $\rho_j$ , regulate integrands  $\varphi_L \sim$ 1  $z_j$ 

$$
u = B^{-\gamma} \prod_{j=1}^{n} \widehat{\mathcal{L}^{(j)}_{j}}
$$

- DE for  $\psi$  has otherwise no solution - if  $\varphi_L \sim 1/z^{\nu_j}$   $(\nu_j > 0)$  then *ψ*  $\varphi_L \sim 1/z^{\nu_j}$   $(\nu_j > 0)$  then  $\psi \sim 1/\rho_j$ 

• limit  $\rho_j \rightarrow 0$  after the decomposition

#### **Drawbacks**

- **additional variables** in intermediate stages
- **obscures block-triangular structure** of decompositions
- **more master integrals** in intermediate steps of recursion

### **Dual integrals**

Recall the decomposition:

$$
|\varphi_R\rangle = \sum_{i=1}^{\nu} c_i^{(R)} |e_i^{(R)}\rangle, \qquad c_i^{(R)} = \sum_{j=1}^{\nu} (\mathbf{C}^{-1})_{ij} \langle e_j^{(L)} | \varphi_R \rangle
$$

#### **Observation**

Coefficients  $c_i^{(R)}$  are independent of the choice of the dual **basis** of "left" integrals  $\{ \langle e_j^{(L)} | \ \}_{j=1}^\nu$ *i*

#### **Idea**  M Exploit the **freedom of choice** of the **dual basis** to simplify the calculation

# **Choice of dual integrals**

Two interesting approaches (different formalisms, similar outcomes):

- Alternative formalism for defining dual space of loop integrals **[Caron-Huot, Pokraka (2021)]**
- Simple **choice of dual integrals**

Choose **dual integrals** of the form **[G. Fontana, T.P. (2023)]** 

$$
\varphi_L(\mathbf{z}) = \rho_1^{\Theta(\alpha_1 - \frac{1}{2})} \cdots \rho_n^{\Theta(\alpha_n - \frac{1}{2})} \frac{1}{z_1^{\alpha_1} \cdots z_n^{\alpha_n}}
$$

- if there's a **denominator** factor  $z_j^{\alpha_j}$  (with  $\alpha_j > 0$ ), **multiply** by *αj*  $\alpha_j^{a_j}$  (with  $\alpha_j > 0$ ), **multiply** by  $\rho_j$
- systematically work in the limit  $\rho_j \to 0$ (i.e. only keep leading terms in a  $\rho_j \rightarrow 0$  expansion in each step)

## **Advantages**

- Calculation "effectively" independent of *ρj*
	- working on leading coefficients in  $\rho_j \rightarrow 0$  expansion, often just one
	- over finite fields, never sample or reconstruct  $\rho_j$  dependence
- Drastically **simpler** intermediate expressions
- Metric and reduction tables are **block triangular** (blocks ~ sectors)

$$
C = \begin{pmatrix} * & * & 0 & 0 & 0 & 0 \\ * & * & 0 & 0 & 0 & 0 \\ * & * & * & * & 0 & 0 \\ * & * & * & * & * & * \\ * & * & 0 & 0 & * & * \\ * & * & 0 & 0 & * & * \end{pmatrix}
$$

- Many intersection numbers and contributions of poles to them **vanish**  $(\rho_j$  prefactors must cancel a  $1/\rho_j$  singularities, only possible at  $z_j \thicksim 0, \infty)$
- **Fewer master integrals** in intermediate steps of recursion



### **Simple examples and checks**

#### **…just checking that things work as expected!**



# **Conclusions & Outlook**

- **Intersection theory** unveils new mathematical structures in loops
- $p(z)$ -adic expansions simplify study of functions close to singular points
	- avoid algebraic extensions
	- no need to know explicit location of irrational poles
	- avoid bottlenecks and enable finite field technologies

- Future directions:
	- simplifications/optimizations and application to different integral representations (loop-by-loop Baikov, Lee-Pomeransky)
	- Non-recursive multivariate generalization (based on **Chestnov, Frellesvig, Gasparotto, Mandal, Mastrolia (2022)**)