

Knowns and Unknowns in $B \rightarrow D^*$ FFs

Martin Jung

Flavour@TH 2023

CERN, 9th of May 2023



**UNIVERSITÀ
DI TORINO**



Istituto Nazionale di Fisica Nucleare
SEZIONE DI TORINO



Lepton-non-Universality in $b \rightarrow c\tau\nu$

$$R(X) \equiv \frac{\text{Br}(B \rightarrow X\tau\nu)}{\text{Br}(B \rightarrow X\ell\nu)}$$

- Partial cancellation of uncertainties
- ➡ Precise predictions (and measurements)

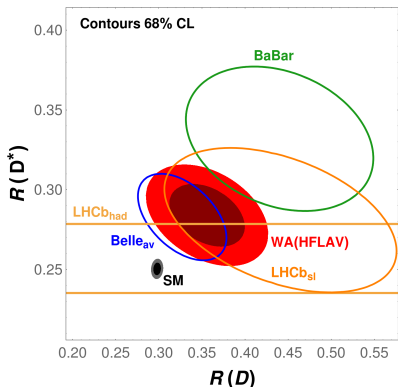
- $R(D^{(*)})$: BaBar, Belle, LHCb
- ➡ average $\sim 3 - 4\sigma$ from SM

More flavour $b \rightarrow c\tau\nu$ observables:

- τ -polarization ($\tau \rightarrow \text{had}$) [1608.06391]
- $B_c \rightarrow J/\psi\tau\nu$ [1711.05623] : huge
- Differential rates from Belle, BaBar
- Total width of B_c
- $b \rightarrow X_c\tau\nu$ by LEP
- D^* polarization (Belle)
- $R(\Lambda_c) \rightarrow$ below SM

Note: only 1 result $\geq 3\sigma$ from SM

In the following: discuss SM + NP predictions



Lepton-non-Universality in $b \rightarrow c\tau\nu$

$$R(X) \equiv \frac{\text{Br}(B \rightarrow X\tau\nu)}{\text{Br}(B \rightarrow X\ell\nu)}$$

- Partial cancellation of uncertainties
- ➡ Precise predictions (and measurements)

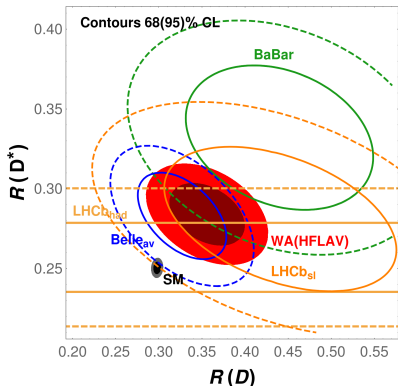
- $R(D^{(*)})$: BaBar, Belle, LHCb
- ➡ average $\sim 3 - 4\sigma$ from SM

More flavour $b \rightarrow c\tau\nu$ observables:

- τ -polarization ($\tau \rightarrow \text{had}$) [1608.06391]
- $B_c \rightarrow J/\psi\tau\nu$ [1711.05623]: huge
- Differential rates from Belle, BaBar
- Total width of B_c
- $b \rightarrow X_c\tau\nu$ by LEP
- D^* polarization (Belle)
- $R(\Lambda_c) \rightarrow$ below SM

Note: only 1 result $\geq 3\sigma$ from SM

In the following: discuss SM + NP predictions



q^2 dependence

- q^2 range can be large, e.g. $q^2 \in [0, 12]$ GeV² in $B \rightarrow D$
- Calculations give usually one or few points
- Knowledge of **functional dependence** on q^2 crucial
- This is where discussions start. . .
- Most $B \rightarrow D^*$ data not usable due to model dependence!

Give as much information as possible **independently of this choice!**

q^2 dependence

- q^2 range can be large, e.g. $q^2 \in [0, 12]$ GeV² in $B \rightarrow D$
- Calculations give usually one or few points
- ➡ Knowledge of **functional dependence** on q^2 crucial
- This is where discussions start. . .
- ➡ Most $B \rightarrow D^*$ data not usable due to model dependence!

Give as much information as possible **independently of this choice!**

Even with FF-model-dependent data:

Consistent HFLAV $B \rightarrow D^*$ fit in CLN

➡ Experimental w -dependence well established!

In the following: mostly **BGL** and **HQE** (\rightarrow CLN) parametrizations

HQE parametrization

Heavy-Quark Expansion (HQE) employs additional information:

- $m_{b,c} \rightarrow \infty$: **all** $B \rightarrow D^{(*)}$ FFs given by **1 Isgur-Wise function**
- Systematic expansion in $1/m_{b,c}$ and α_s
- Higher orders in $1/m_{b,c}$: FFs remain related
 - ↳ Parameter reduction, necessary for NP analyses!

HQE parametrization

Heavy-Quark Expansion (HQE) employs additional information:

- $m_{b,c} \rightarrow \infty$: all $B \rightarrow D^{(*)}$ FFs given by **1 Isgur-Wise function**
- Systematic expansion in $1/m_{b,c}$ and α_s
- Higher orders in $1/m_{b,c}$: FFs remain related
 - ➡ Parameter reduction, necessary for NP analyses!

CLN parametrization [Caprini+'97] :

HQE to order $1/m_{b,c}$, α_s plus (approx.) constraints from unitarity [Bernlochner/Ligeti/Papucci/Robinson'17] : identical approach, updated and consistent treatment of correlations

Problem: Contradicts Lattice QCD (both in $B \rightarrow D$ and $B \rightarrow D^*$)

Dealt with by varying calculable ($\mathcal{O}(1/m_{b,c})$) parameters, e.g. $h_{A_1}(1)$

- ➡ **Not** a systematic expansion in $1/m_{b,c}$ anymore!
- ➡ Related uncertainty remains $\mathcal{O}[\Lambda^2/(2m_c)^2] \sim 5\%$, insufficient

HQE parametrization

Heavy-Quark Expansion (HQE) employs additional information:

- $m_{b,c} \rightarrow \infty$: all $B \rightarrow D^{(*)}$ FFs given by **1 Isgur-Wise function**
- Systematic expansion in $1/m_{b,c}$ and α_s
- Higher orders in $1/m_{b,c}$: FFs remain related
 - ➔ Parameter reduction, necessary for NP analyses!

CLN parametrization [Caprini+'97] :

HQE to order $1/m_{b,c}$, α_s plus (approx.) constraints from unitarity [Bernlochner/Ligeti/Papucci/Robinson'17] : identical approach, updated and consistent treatment of correlations

Problem: Contradicts Lattice QCD (both in $B \rightarrow D$ and $B \rightarrow D^*$)
Dealt with by varying calculable ($\mathcal{O}(1/m_{b,c})$) parameters, e.g. $h_{A_1}(1)$

- ➔ **Not** a systematic expansion in $1/m_{b,c}$ anymore!
- ➔ Related uncertainty remains $\mathcal{O}[\Lambda^2/(2m_c)^2] \sim 5\%$, insufficient

Solution: Include systematically $1/m_c^2$ corrections

[Bordone/MJ/vDyk'19, Bordone/Gubernari/MJ/vDyk'20] , using [Falk/Neubert'92]

[Bernlochner+'22] : model for $1/m_c^2$ corrections \rightarrow fewer parameters

Theory determination of $b \rightarrow c$ Form Factors

[Bordone/MJ/vanDyk'19,Bordone/Gubernari/MJ/vanDyk'20]

For general NP analysis, FF shapes needed from theory!

Fit to all $B \rightarrow D^{(*)}$ FFs, using lattice, LCSR, QCDSR and unitarity

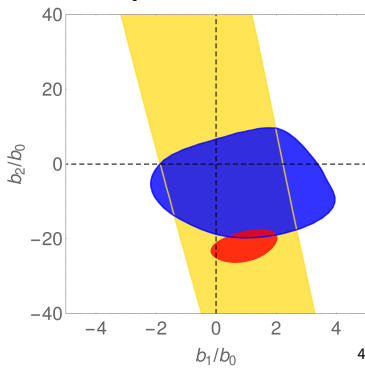
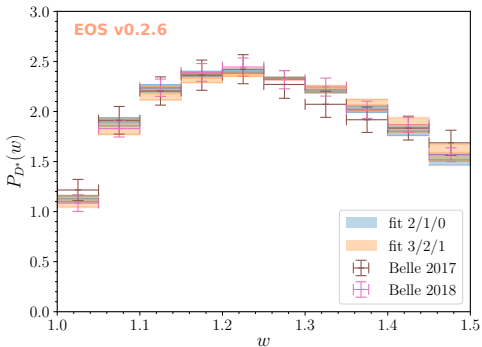
[CLN,BGL,HPQCD'15'17,FNAL/MILC'14'15,Gubernari+'18,Ligeti+'92'93]

k/l/m: order in z for leading/subleading/subsubleading IW functions

➡ 2/1/0 works, but only 3/2/1 captures uncertainties

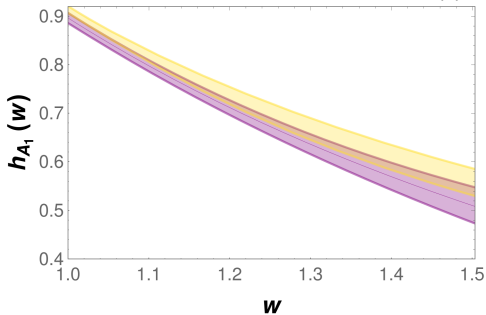
➡ Consistent V_{cb} value from Belle'17+'18

➡ Predictions for diff. rates, perfectly confirmed by data



Comparison with new lattice calculations

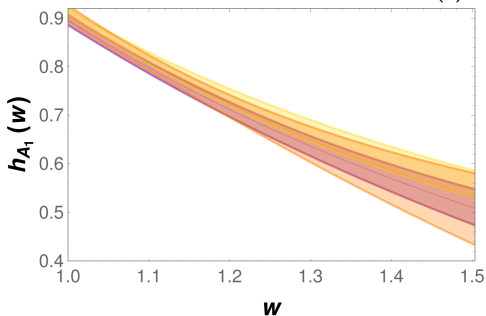
Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_s : [Harrison+'22])



- FNAL/MILC'21
- HQE@ $1/m_c^2$
- Exp (BGL)
- JLQCD prel
- HPQCD'23

Comparison with new lattice calculations

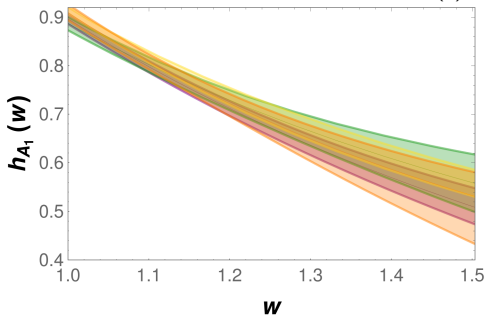
Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_S : [Harrison+'22])



- FNAL/MILC'21
- HQE@ $1/m_c^2$
- Exp (BGL)
- JLQCD prel
- HPQCD'23

Comparison with new lattice calculations

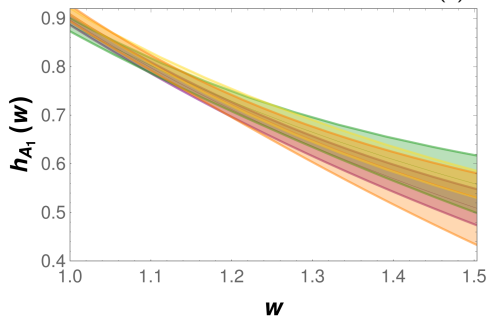
Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_s : [Harrison+'22])



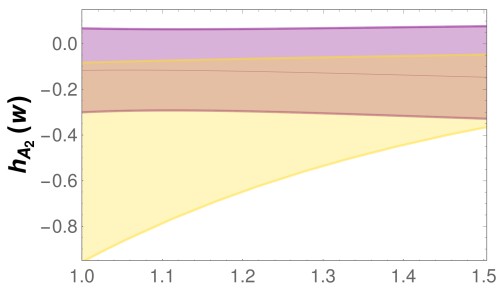
- FNAL/MILC'21
- HQE@1/ m_c^2
- Exp (BGL)
- JLQCD prel
- HPQCD'23
- Compatible

Comparison with new lattice calculations

Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_s : [Harrison+'22])



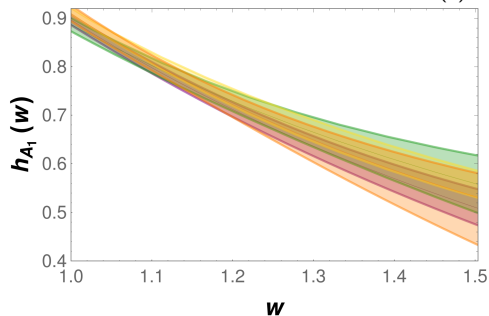
- FNAL/MILC'21
- HQE@1/ m_c^2
- Exp (BGL)
- JLQCD prel
- HPQCD'23
- Compatible



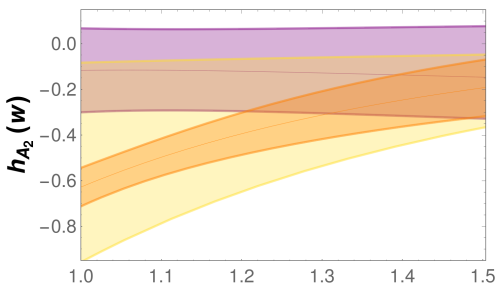
- HPQCD and BGJvD compatible

Comparison with new lattice calculations

Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_s : [Harrison+'22])



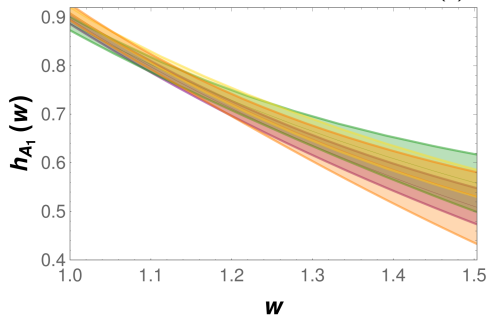
- FNAL/MILC'21
- HQE@ $1/m_c^2$
- Exp (BGL)
- JLQCD prel
- HPQCD'23
- ➡ Compatible



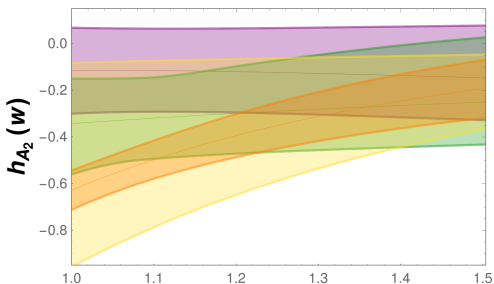
- HPQCD and BGJvD compatible
- Slope HPQCD-FNAL/MILC?

Comparison with new lattice calculations

Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_s : [Harrison+'22])



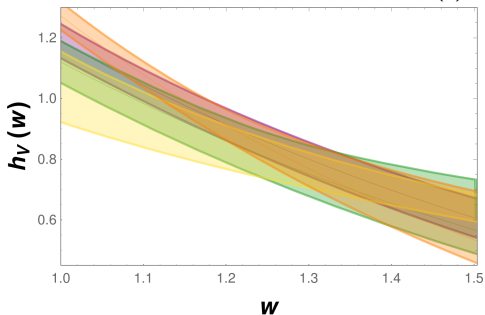
- FNAL/MILC'21
- HQE@1/ m_c^2
- Exp (BGL)
- JLQCD prel
- HPQCD'23
- ➔ Compatible



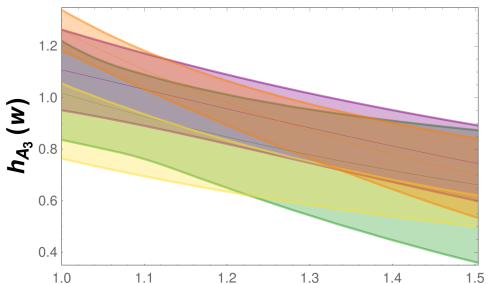
- HPQCD and BGJvD compatible
- Slope HPQCD-FNAL/MILC?
- JLQCD "diplomatic"

Comparison with new lattice calculations

Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_S : [Harrison+'22])



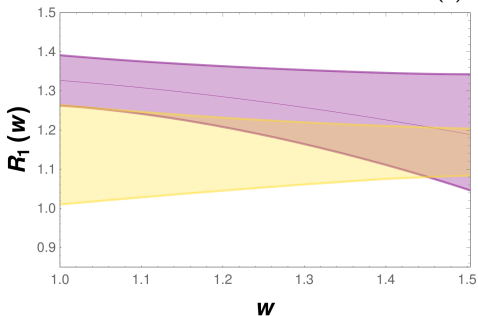
- FNAL/MILC'21
- HQE@ $1/m_c^2$
- Exp (BGL)
- JLQCD prel
- HPQC'D'23
- ➔ Compatible



- Similar pattern in h_V and h_{A_3}
- Tension between BGJvD and FNAL/MILC in h_{A_2}

Comparison with new lattice calculations

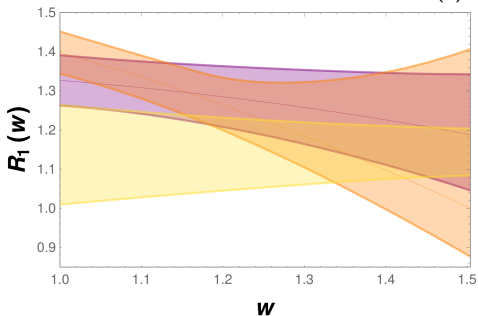
Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_s : [Harrison+'22])



- FNAL/MILC'21
- HQE@ $1/m_c^2$
- Exp (BGL)
- JLQCD prel
- HPQCD'23

Comparison with new lattice calculations

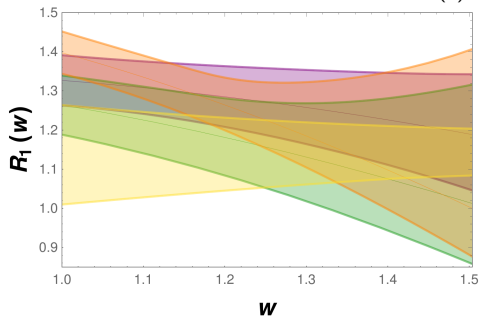
Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_S : [Harrison+'22])



- FNAL/MILC'21
- HQE@ $1/m_c^2$
- Exp (BGL)
- JLQCD prel
- HPQC'D'23

Comparison with new lattice calculations

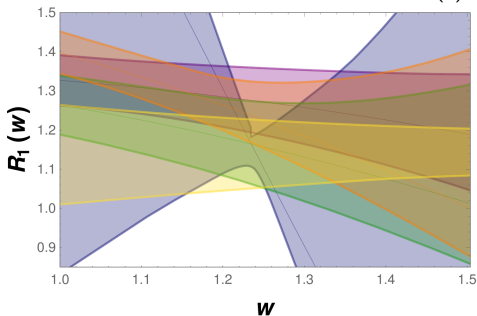
Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_S : [Harrison+'22])



- FNAL/MILC'21
- HQE@ $1/m_c^2$
- Exp (BGL)
- JLQCD prel
- HPQCD'23

Comparison with new lattice calculations

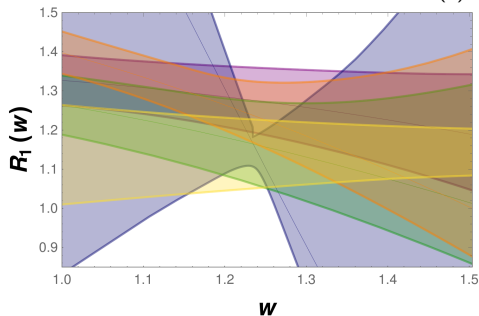
Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_s : [Harrison+'22])



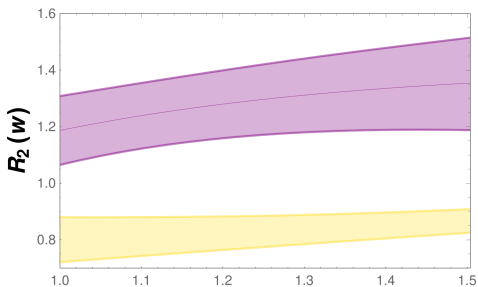
- FNAL/MILC'21
 - HQE@1/ m_c^2
 - Exp (BGL)
 - JLQCD prel
 - HPQCD'23
- ➡ Compatible. Slope?

Comparison with new lattice calculations

Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_S : [Harrison+'22])



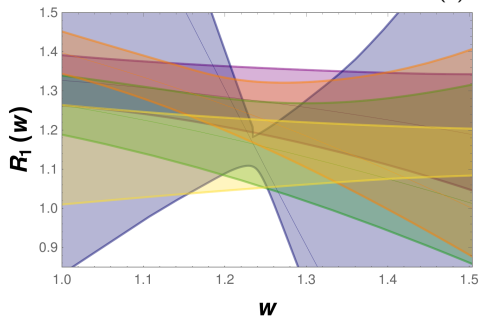
- FNAL/MILC'21
- HQE@ $1/m_c^2$
- Exp (BGL)
- JLQCD prel
- HPQCD'23
- ➡ Compatible. Slope?



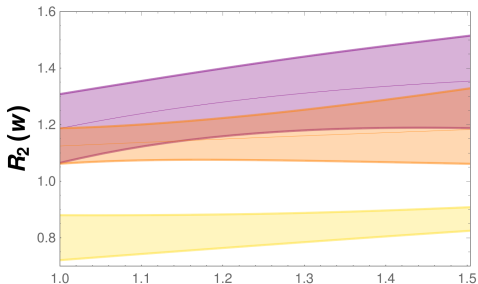
- Deviation HPQCD-BGJvD

Comparison with new lattice calculations

Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_S : [Harrison+'22])



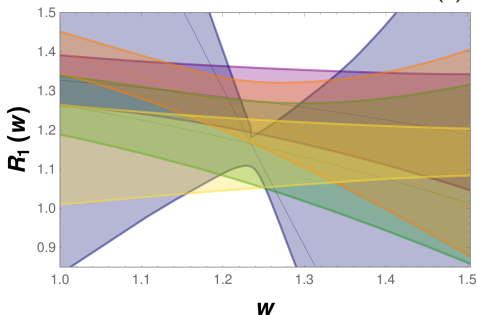
- FNAL/MILC'21
 - HQE@ $1/m_c^2$
 - Exp (BGL)
 - JLQCD prel
 - HPQCD'23
- ➡ Compatible. Slope?



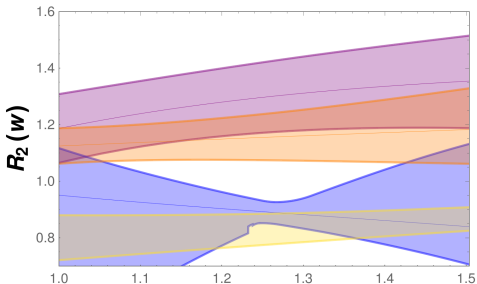
- Deviation HPQCD-BGJvD
- FNAL/MILC close to HPQCD

Comparison with new lattice calculations

Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_S : [Harrison+'22])



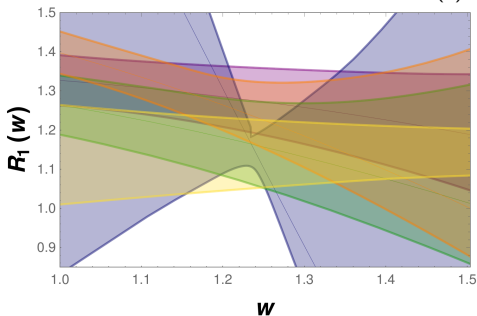
- FNAL/MILC'21
 - HQE@ $1/m_c^2$
 - Exp (BGL)
 - JLQCD prel
 - HPQCD'23
- ➡ Compatible. Slope?



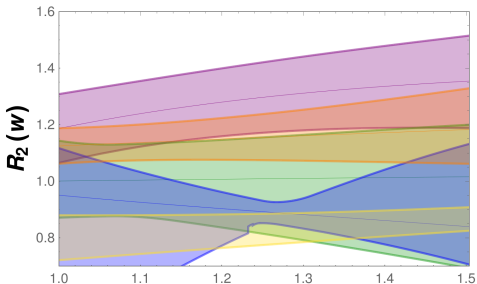
- Deviation HPQCD-BGJvD
 - FNAL/MILC close to HPQCD
 - Deviation wrt experiment
($R_2^{\text{HF}LAV}(1) = 0.853(17)$)
- ➡ Requires further investigation!

Comparison with new lattice calculations

Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_S : [Harrison+'22])



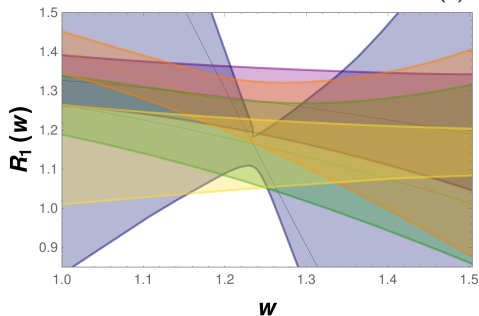
- FNAL/MILC'21
 - HQE@ $1/m_c^2$
 - Exp (BGL)
 - JLQCD prel
 - HPQCD'23
- ➡ Compatible. Slope?



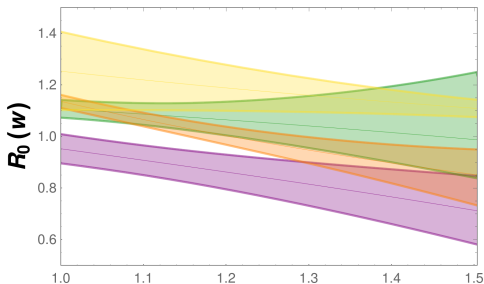
- Deviation HPQCD-BGJvD
 - FNAL/MILC close to HPQCD
 - Deviation wrt experiment ($R_2^{\text{HF}LAV}(1) = 0.853(17)$)
- ➡ Requires further investigation!
- JLQCD “diplomatic”

Comparison with new lattice calculations

Major improvement: $B_{(s)} \rightarrow D_{(s)}^*$ FFs@ $w > 1$! (B_s : [Harrison+'22])

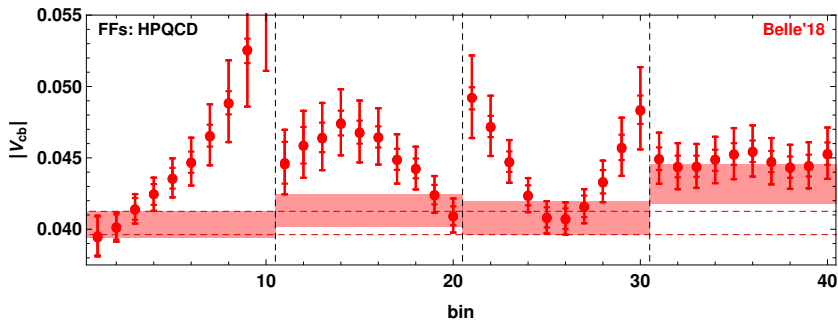
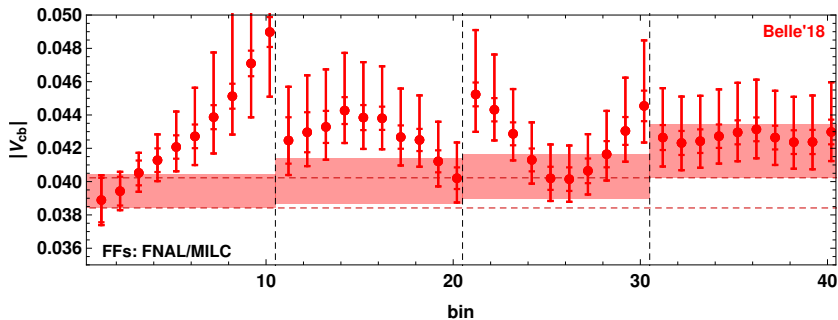


- FNAL/MILC'21
- HQE@ $1/m_c^2$
- Exp (BGL)
- JLQCD prel
- HPQCD'23
- ➡ Compatible. Slope?

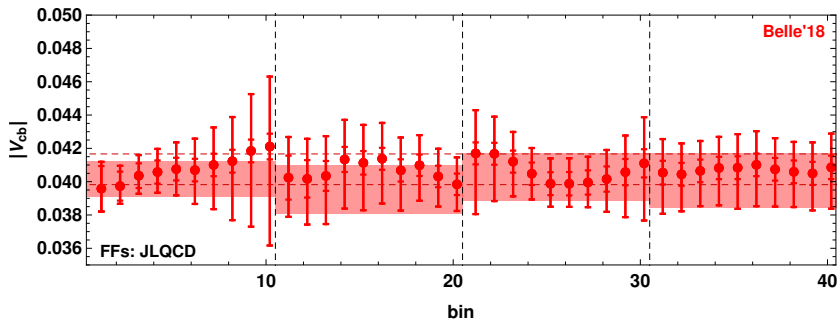
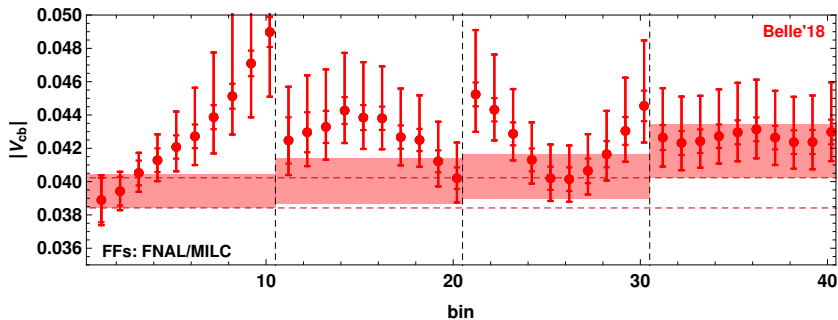


- Also in R_0 deviation wrt BGJvD
- JLQCD again “diplomatic”
- ➡ Requires further investigation!
- ➡ Correlations?

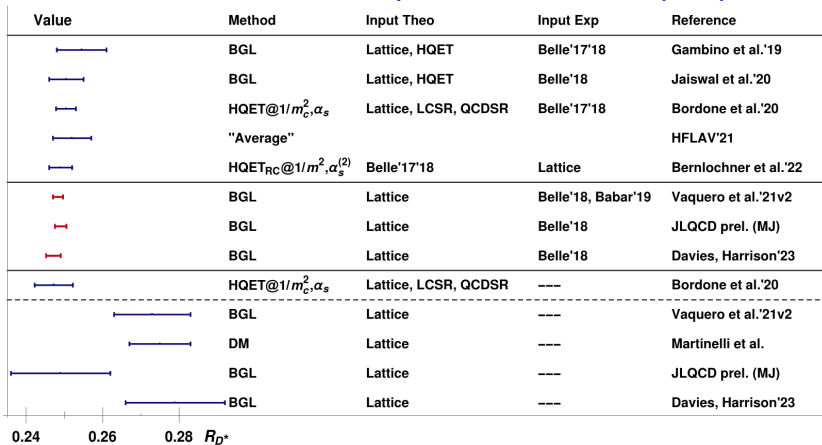
Binned V_{cb} from Belle'18 data: FNAL/MILC vs HPQCD



Binned V_{cb} from Belle'18 data: FNAL/MILC vs JLQCD



Overview over predictions for $R(D^*)$



Lattice $B \rightarrow D^*$: $h_{A_1}(w=1)$ [FNAL/MILC'14, HPQCD'17], [FNAL/MILC'21]

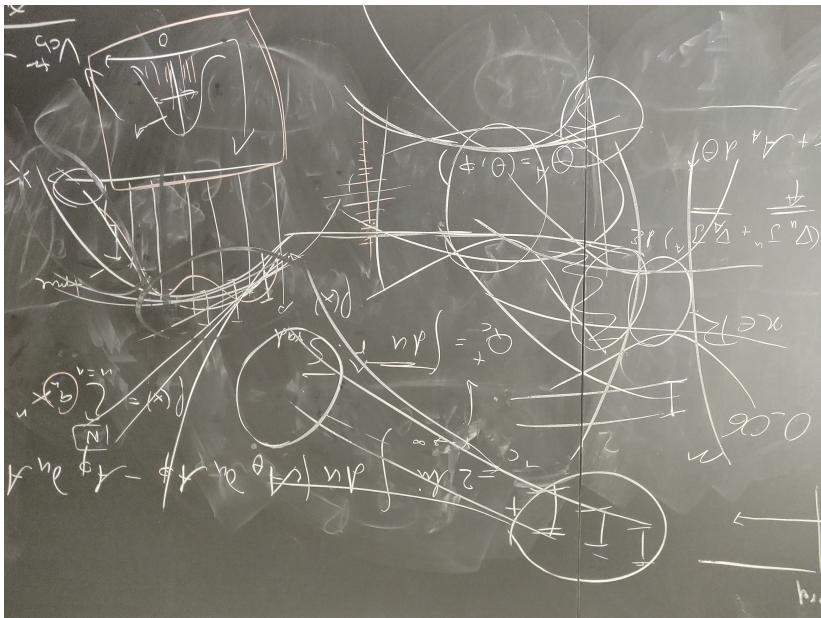
Other lattice: $f_{+,0}^{B \rightarrow D}(q^2)$ [FNAL/MILC, HPQCD'15]

QCDSR: [Ligeti/Neubert/Nir'93, '94], LCSR: [Gubernari/Kokulu/vDyk'18]

Overall consistent SM predictions!

"Explaining" $R(D^*)$ by FM/HPQCD \rightarrow NP in $B \rightarrow D^*(e, \mu)\nu$!

Form-factor truncation



Form-factor truncation

Key question: Where do we truncate our expansions?

- ➡ A [Bernlochner+'19]: include parameter only if χ^2 decreases significantly
- ➡ B (GJS, BGJvD): include one “unnecessary” order

Comments:

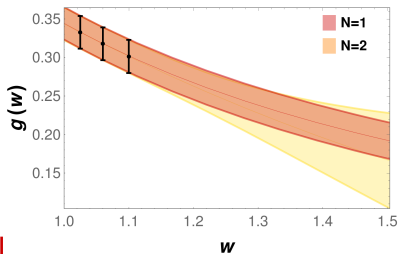
- Large difference, $\sim 50\%$ difference in uncertainty
- Motivation for A: convergence, avoid overfitting
- Motivation for B: avoid underestimating uncertainties
- ➡ Different perspectives: only describing data, A is ok.

However: we **extrapolate** to regions where we lack sensitivity

Example: $g(w)$ from FNAL/MILC

- perfect description at $\mathcal{O}(z)$
- large impact from $\mathcal{O}(z^2)$
- Nevertheless: $\mathcal{O}(z^2) \leq 6\% \times \mathcal{O}(z)$
 - ➡ overfitting limited

Just because you're not sensitive,
doesn't mean it's not there!



Priors and potential biases I

Common error estimate:

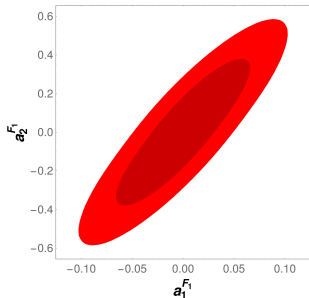
$$\delta X \sim \mathcal{O}(1) \times \text{known factor}$$

➡ What's " $\mathcal{O}(1)$ "?

- Answer seems to be community-dependent
- Often in lattice analyses: gaussian around 0, width 1
- BGJvD HQET FFs: flat range $[-20, 20]$
- JLQCD: $d_w(h_{A_2}) = 10.3(7.1)$
- ➡ potentially large differences
- ➡ needs to be checked and **communicated**

Similarly: treatment of BGL coefficients

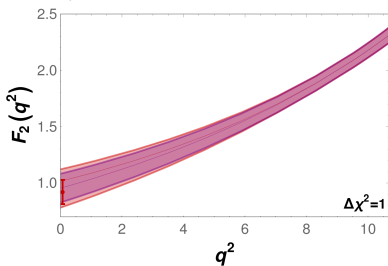
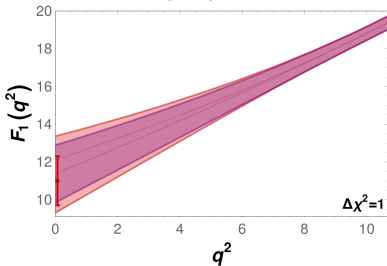
- FNAL/MILC and HPQCD:
series in $(w - 1)^n$
- Priors can be strong in BGL space
- Plot: prior information, only (HPQCD)
- ➡ Comparable to final result



Priors and potential biases II

Different conclusions starting from identical information

Example: $R(D^*)$ extraction from FNAL/MILC data



$R(D^*)$ including kinematical identities and weak unitarity

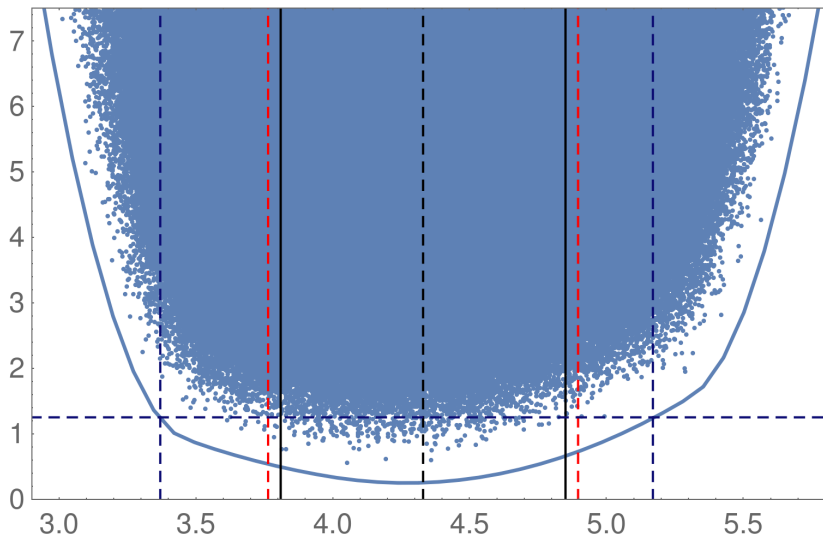
$$R(D^*) \stackrel{\text{WU}}{=} 0.269_{-0.008}^{+0.020} \quad \stackrel{\text{FM}}{=} 0.274 \pm 0.010 \quad \stackrel{\text{Rome}}{=} 0.275 \pm 0.008.$$

Difference WU-FM: FM apply prior on BGL coefficients

Difference WU-Rome (educated guess): iterated “unitarity filter”
+ different error estimate

Applying data: $R(D^*) = 0.249 \pm 0.001(!)$ **universally.**

Uncertainty determination



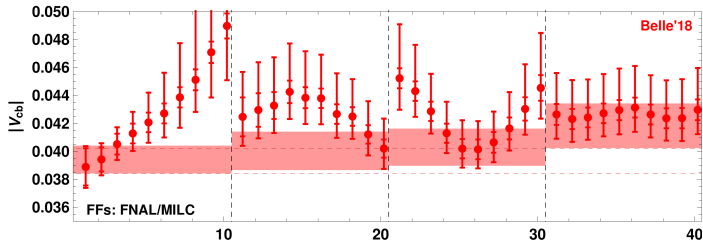
MC points together with χ^2 profile (minimizing for each FF value)

Vertical: CV MC, " 1σ " MC, symmetric 68.3% interval MC, $\Delta\chi^2 = 1$

The Dispersive Matrix (DM) Method

Alternative implementation of unitarity [Bourrely+'81,Lellouch'95] :

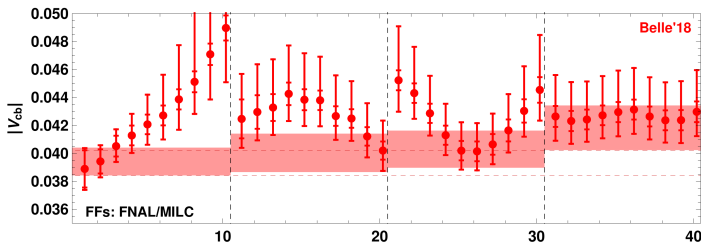
- Identical starting point as BGL: dispersion relation
- Known information in a matrix with positive determinant
 - ➡ Form-factor bounds, equivalent to GJS [Caprini'19]
- Enables parametrization-free analysis



Implemented recently for $B \rightarrow D^* \ell \nu$ [DiCarlo+'21,Martinelli+'21,22] :

- Use DM w/ new FNAL/MILC data to obtain FF bands
- Calculate V_{cb} bin-wise, combine $d\Gamma/dx$ bins ($x = q^2, \cos\theta, \dots$) (including experimental and theoretical correlations)
 - ➡ 2×4 V_{cb} values. Claim: 0.5σ to V_{cb}^{incl} , 1.3σ to $R(D^*)$

The Dispersive Matrix (DM) Method



Differences between DM and GJS [Gambino/MJ/Schacht'19] :

- GJS: **Combined fit** of lattice and experiment, imposing unitarity
- DM: **Unweighted, uncorrelated** average of the 4 V_{cb} values:

$$\mu = \frac{1}{N} \sum_{k=1}^N x_k, \quad \sigma_x^2 = \frac{1}{N} \sum_{k=1}^N \sigma_k^2 + \frac{1}{N} \sum_{k=1}^N (x_k - \mu_x)^2$$

$$\rightarrow V_{cb}^{\text{GJS}} = (39.2_{-1.2}^{+1.4}) \times 10^{-3}, \quad V_{cb}^{\text{DM}} = (40.8 \pm 1.7) \times 10^{-3}$$

Conclusions

We have work ahead of us!

1. q^2 dependence of FFs critical
 - ➡ Need **parametrization-independent data**
2. Inclusion of higher-order (theory) uncertainties essential
3. HQE: systematic expansion in $1/m, \alpha_s$, relates FFs
 - ➡ $\mathcal{O}(1/m_c)$ (\rightarrow CLN) **not sufficient anymore**
4. Important first LQCD analyses in $B_{(s)} \rightarrow D_{(s)}^*$ @ finite recoil
 - ➡ HPQCD: First 2+1+1 results, full q^2 range!
 - ➡ Tensions in ratios – correlations?
5. Despite complications: $R(D^{(*)})$ SM prediction robust!

Central lesson:

Experiment and theory (lattice + pheno) need to work closely together!

Exclusive decays: Form factors

In exclusive decays, hadronic information encoded in **Form Factors**
They parametrize fundamental mismatch:

Theory (e.g. SM) for **partons** (quarks)

vs.

Experiment with **hadrons**

$$\langle D_q(p') | \bar{c} \gamma^\mu b | \bar{B}_q(p) \rangle = (p + p')^\mu f_+^q(q^2) + (p - p')^\mu f_-^q(q^2), \quad q^2 = (p - p')^2$$

Most general matrix element parametrization, given **symmetries**:
Lorentz symmetry plus P- and T-symmetry of QCD
 $f_\pm(q^2)$: real, scalar functions of **one** kinematic variable

How to obtain these functions?

- ➡ **Calculable** w/ **non-perturbative** methods (Lattice, LCSR, . . .)
Precision?
- ➡ **Measurable** e.g. in semileptonic transitions
Normalization? Suppressed FFs? NP?

The BGL parametrization [Boyd/Grinstein/Lebed, 90's]

FFs are parametrized by a few coefficients the following way:

1. Consider **analytical structure**, make poles and cuts explicit
2. Without poles or cuts, the rest can be **Taylor-expanded** in z
3. Apply QCD symmetries (unitarity, crossing)
↳ **dispersion relation**
4. Calculate **partonic part** (mostly) perturbatively

The BGL parametrization [Boyd/Grinstein/Lebed, 90's]

FFs are parametrized by a few coefficients the following way:

1. Consider **analytical structure**, make poles and cuts explicit
2. Without poles or cuts, the rest can be **Taylor-expanded** in z
3. Apply QCD symmetries (unitarity, crossing)
↳ **dispersion relation**
4. Calculate **partonic part** (mostly) perturbatively

Result: Model-independent parametrization

$$F(t) = \frac{1}{P(t)\phi(t)} \sum_{n=0}^{\infty} a_n [z(t, t_0)]^n.$$

- a_n : **real** coefficients, the only unknowns
- $P(t)$: **Blaschke factor(s)**, information on poles below t_+
- $\phi(t)$: **Outer function**, chosen such that $\sum_{n=0}^{\infty} a_n^2 \leq 1$

Series in z with **bounded coefficients** (each $|a_n| \leq 1$)!

↳ Uncertainty related to truncation is **calculable**!

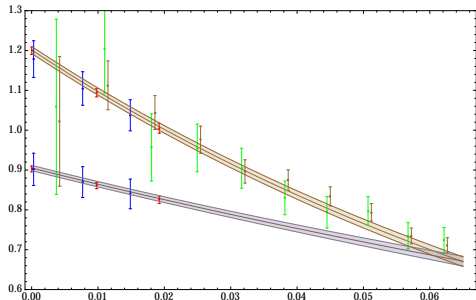
$B \rightarrow D\ell\nu$

$B \rightarrow D\ell\nu$, aka “The teacher’s pet”:

- Excellent agreement between experiments [BaBar'09,Belle'16]
- Excellent agreement between two lattice determinations [FNAL/MILC'15,HPQCD'16]
- ➡ Lattice data inconsistent with CLN parametrization! (but consistent w/ HQE@1/m, discussed later)
- BGL fit [Bigi/Gambino'16] :

$$R(D) = 0.299(3).$$

See also [Jaiswal+,Berlochner+'17,MJ/Straub'18,Bordone/MJ/vanDyk'19]



$f_{+,0}(z)$, inputs:

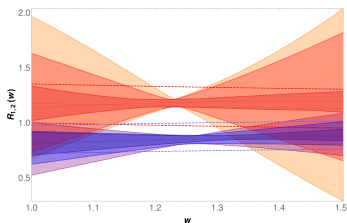
- FNAL/MILC'15
- HPQCD'16
- BaBar'09
- Belle'16

$V_{cb} + R(D^*)$ w/ data + lattice + unitarity [Gambino/MJ/Schacht'19]

Belle'18(+ '17) provide FF-independent data for 4 single-differential rates

BGL analysis:

- Datasets compatible
- d'Agostini bias + syst. important
- **Expand FFs to z^2**
 - ↳ **50% increased uncertainties**
- Belle'18: no parametrization dependence
- Belle'17 never published → replace w/ Belle'23, not available yet
- Tension w/ inclusive reduced, but not removed



$$R(D^*) = 0.253^{+0.007}_{-0.006} \quad (\text{including LCSR point})$$

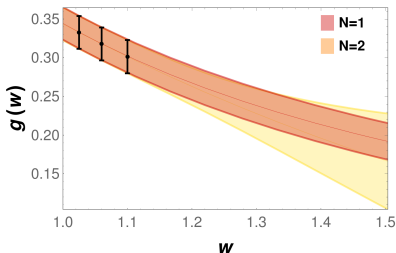
Comparison to Bernlochner+'22

Bernlochner et al. also perform HQE analysis @ $1/m_c^2$. Differences:

- Postulate different counting within HQET
 - ➔ Highly constraining model for higher-order corrections
- Avoid use of LCSR (and mostly QCDSR) results
- Include partial α_s^2 corrections
- Include FNAL/MILC results partially
- Expansion in z : 2/1/0 (justified in [Bernlochner+'19])

Observations:

- $1/m_c^2$ corrections necessary
- Overall small uncertainties
- $V_{cb} = (38.7 \pm 0.6) \times 10^{-3}$
 - ➔ smaller due to larger $\mathcal{F}(1)$
- $R(D^*)$: agreement w/ BGJvD
- $R(D) \sim 3\sigma$ from GJS + BGJvD
 - ➔ In my opinion due to model

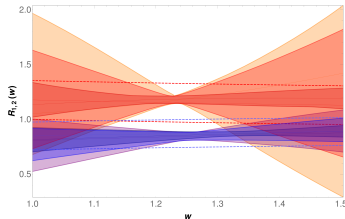


$V_{cb} + R(D^*)$ w/ data + lattice + unitarity [Gambino/MJ/Schacht'19]

Belle'17+'18 provide FF-independent data for 4 single-differential rates

Analysis of these data with **BGL form factors**:

- Datasets roughly compatible
- d'Agostini bias + syst. important
- All FFs to z^2 to include uncertainties
➡ 50% increased uncertainties
- 2018: no parametrization dependence



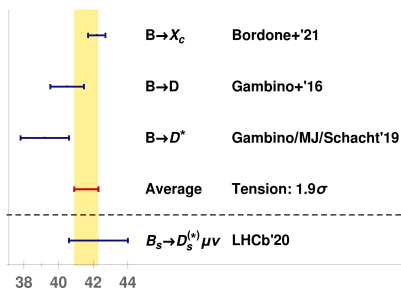
$$|V_{cb}^{D^*}| = 39.6_{-1.0}^{+1.1} [39.2_{-1.2}^{+1.4}] \times 10^{-3}$$

$$R(D^*) = 0.254_{-0.006}^{+0.007} [0.253_{-0.006}^{+0.007}]$$

In brackets: 2018 only ($\Delta V_{cb}^{\text{Belle}} = 0.9$)

Updating the $|V_{cb}|$ puzzle:

- Tension 1.9σ (larger $\delta V_{cb}^{B \rightarrow D^*}$)
- $B_s \rightarrow D_s^{(*)}$ reduces tension further
- $V_{cb}^{B \rightarrow D^*}$ vs. V_{cb}^{incl} still problematic



See also [Bigi+, Bernlocher+, Grinstein+'17, Jaiswal+'17'19, MJ/Straub'18, Bordone+'19/20]

Theory determination of $b \rightarrow c$ Form Factors

SM: BGL fit to data + FF normalization $\rightarrow |V_{cb}|$

NP: can affect the q^2 -dependence, introduces additional FFs

➡ To determine general NP, FF shapes needed from theory

[MJ/Straub'18, Bordone/MJ/vDyk'19] used all available theory input:

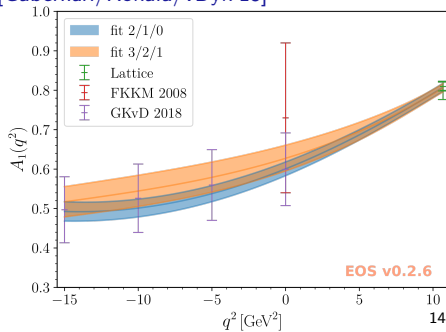
- Unitarity bounds (using results from [CLN, BGL])
 - ➡ non-trivial $1/m$ vs. z expansions
- LQCD for $f_{+,0}(q^2)$ ($B \rightarrow D$), $h_{A_1}(q_{\max}^2)$ ($B \rightarrow D^*$)

[HPQCD'15,'17, Fermilab/MILC'14,'15]

- LCSR for all FFs (mod f_T) [Gubernari/Kokulu/vDyk'18]
- QCDSR results for $1/m$ IW functions [Ligeti+'92'93]
- HQET expansion to $\mathcal{O}(\alpha_s, 1/m_b, 1/m_c^2)$

FFs under control;
 $R(D^*) = 0.247(6)$

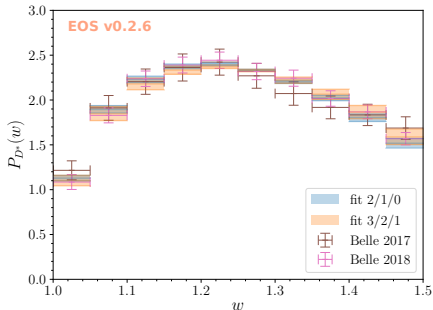
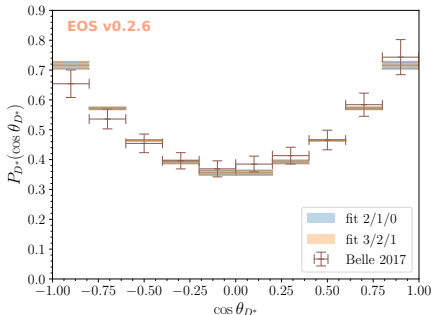
[Bordone/MJ/vDyk'19]



Robustness of the HQE expansion up to $1/m_c^2$

[Bordone/MJ/vDyk'19]

Testing FFs by comparing to data and fits in BGL parametrization:

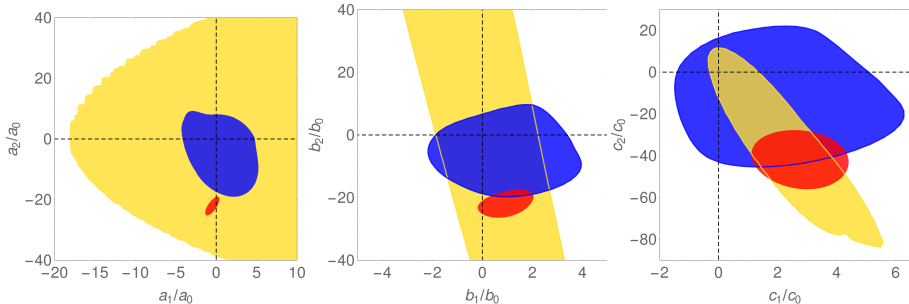


- Fits 3/2/1 and 2/1/0 are **theory-only fits(!)**
- $k/l/m$ denotes orders in z at $\mathcal{O}(1, 1/m_c, 1/m_c^2)$
- w -distribution yields information on FF shape $\rightarrow V_{cb}$
- Angular distributions more strongly constrained by theory, only
- Predicted shapes perfectly confirmed by $B \rightarrow D^{(*)} \ell \nu$ data
- V_{cb} from Belle'17 compatible between HQE and BGL!

Robustness of the HQE expansion up to $1/m_c^2$

[Bordone/MJ/vDyk'19]

Testing FFs by comparing to data and fits in BGL parametrization:



- $B \rightarrow D^*$ BGL coefficient ratios from:
 1. Data (Belle'17+'18) + weak unitarity (yellow)
 2. HQE theory fit 2/1/0 (red)
 3. HQE theory fit 3/2/1 (blue)
- ➡ Again compatibility of theory with data
- ➡ 2/1/0 underestimates the uncertainties massively
- ➡ For b_i, c_i ($\rightarrow f, \mathcal{F}_1$) data and theory complementary

Including $\bar{B}_s \rightarrow D_s^{(*)}$ Form Factors [Bordone/Gubernari/MJ/vDyk'20]

Dispersion relation *sums* over hadronic intermediate states

- ➡ Includes $B_s D_s^{(*)}$, included via SU(3) + conservative breaking
- ➡ Explicit treatment can improve also $\bar{B} \rightarrow D^{(*)} \ell \nu$

Experimental progress in $\bar{B}_s \rightarrow D_s^{(*)} \ell \nu$:

2 new LHCb measurements [2001.03225, 2003.08453]

Improved theory determinations required, especially for NP

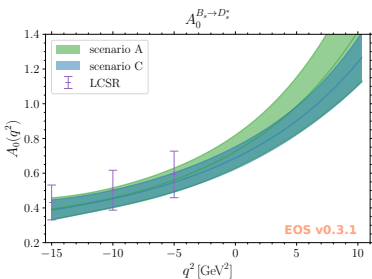
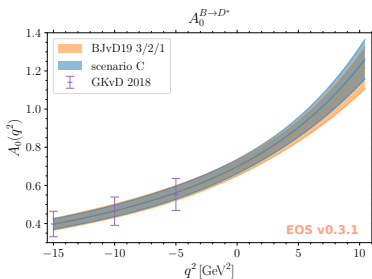
We extend our $1/m_c^2$ analysis by including:

- Available lattice data:
(2 $\bar{B}_s \rightarrow D_s$ FFs (q^2 dependent), 1 $\bar{B}_s \rightarrow D^*$ FF (only q_{\max}^2))
 - Adaptation of existing QCDSR results [Ligeti/Neubert/Nir'93'94], including SU(3) breaking
 - New LCSR results extending [Gubernari+'18] to B_s , including SU(3) breaking
- ➡ Fully correlated fit to $\bar{B} \rightarrow D^{(*)}$, $\bar{B}_s \rightarrow D_s^{(*)}$ FFs

Including $\bar{B}_s \rightarrow D_s^{(*)}$ Form Factors, Results

We observe the following:

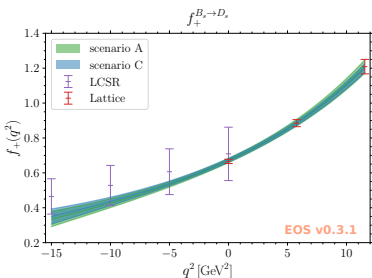
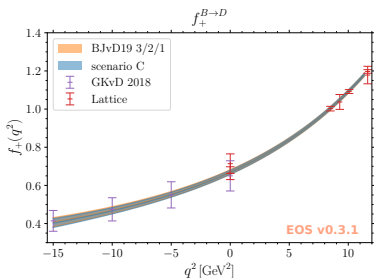
- Theory constraints fitted consistently in an HQE framework
- $\mathcal{O}(1/m_c^2)$ power corrections have $\mathcal{O}(1)$ coefficients
- No indication of sizable SU(3) breaking
- Slight influence of strengthened unitarity bounds
- Improved determination of $\bar{B}_s \rightarrow D_s^{(*)}$ FFs



Including $\bar{B}_s \rightarrow D_s^{(*)}$ Form Factors, Results

We observe the following:

- Theory constraints fitted consistently in an HQE framework
- $\mathcal{O}(1/m_c^2)$ power corrections have $\mathcal{O}(1)$ coefficients
- No indication of sizable SU(3) breaking
- Slight influence of strengthened unitarity bounds
- Improved determination of $\bar{B}_s \rightarrow D_s^{(*)}$ FFs



Including $\bar{B}_s \rightarrow D_s^{(*)}$ Form Factors, Results

We observe the following:

- Theory constraints fitted consistently in an HQE framework
- $\mathcal{O}(1/m_c^2)$ power corrections have $\mathcal{O}(1)$ coefficients
- No indication of sizable SU(3) breaking
- Slight influence of strengthened unitarity bounds
- Improved determination of $\bar{B}_s \rightarrow D_s^{(*)}$ FFs

Theory-only predictions:

$$R(D) = 0.299(3)$$

$$R(D^*) = 0.247(5)$$

$$R(D_s) = 0.297(3)$$

$$R(D_s^*) = 0.245(8)$$

Theory+Experiment (Belle'17) predictions:

$$R(D) = 0.298(3)$$

$$R(D^*) = 0.250(3)$$

$$R(D_s) = 0.297(3)$$

$$R(D_s^*) = 0.247(8)$$

Application: Flavour universality in $B \rightarrow D^*(e, \mu)\nu$

[Bobeth/Bordone/Gubernari/MJ/vDyk'21]

So far: Belle'18 data used in SM fits, **flavour-averaged**

However: Bins 40×40 covariances given **separately** for $\ell = e, \mu$

➡ Belle'18: $R_{e/\mu}(D^*) = 1.01 \pm 0.01 \pm 0.03$

➡ What can we learn about flavour-non-universality? \rightarrow 2 issues:

1. $e - \mu$ correlations not given, but constructible from Belle'18
2. 3 bins linearly dependent, but covariances not singular

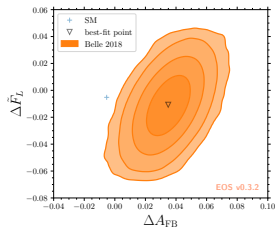
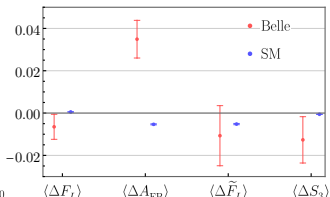
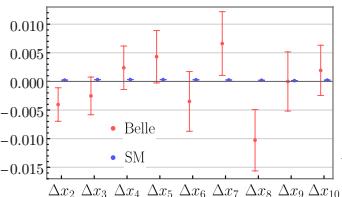
Two-step analysis:

1. 2×4 angular observables suffice for 2×30 angular bins

➡ Model-independent description **including** NP!

2. Compare with SM predictions, using FFs@ $1/m_c^2$ [Bordone+'19]

➡ $\sim 4\sigma$ discrepancy in $\Delta A_{\text{FB}} = A_{\text{FB}}^\mu - A_{\text{FB}}^e$



Generalities regarding this anomaly

- ~ 15% of a SM tree decay $\sim V_{cb}$: This is a huge effect!
- ➡ Need contribution of $\sim 5 - 10\%$ (w/ interference)
or $\gtrsim 40\%$ (w/o interference) of SM

What do we do about this?

- Check the SM prediction!
[→ Bigi+, Bordone+, Gambino+, Grinstein+, Bernlochner+]
- ➡ $\delta R(D^*)$ larger, anomaly remains
- Combined analysis of all $b \rightarrow c\tau\nu$ observables [100+ papers]
➡ First model discrimination
- Related indirect bounds (partly model-dependent)
➡ High p_T searches, lepton decays, LFV, EDMs, ...
- Analyze flavour structure of potential NP contributions
➡ quark flavour structure, e.g. $b \rightarrow u$
➡ lepton flavour structure, e.g. $b \rightarrow c\ell(= e, \mu)\nu$



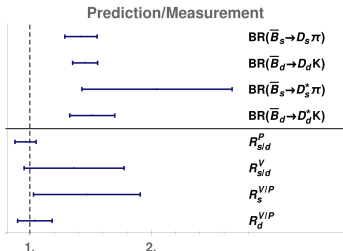
A puzzle in non-leptonic $b \rightarrow c$ transitions

[Bordone/Gubernari/Huber/MJ/vDyk'20]

FFs also of central importance in non-leptonic decays:

- Complicated in general, $B \rightarrow M_1 M_2$ dynamics
- Simplest cases: $\bar{B}_d \rightarrow D_d^{(*)} \bar{K}$ and $\bar{B}_s \rightarrow D_s^{(*)} \pi$ (5 diff. quarks)
 - ➡ Colour-allowed tree, $1/m_b^0 @ \mathcal{O}(\alpha_s^2)$ [Huber+'16], **factorizes at $1/m_b$**
 - ➡ Amplitudes dominantly $\sim \bar{B}_q \rightarrow D_q^{(*)}$ FFs
 - ➡ Used to determine f_s/f_d at hadron colliders [Fleischer+'11]

Updated and extended calculation: tension of 4.4σ w.r.t. exp.!



- Large effect, $\sim -30\%$ for BRs
- Ratios of BRs ok
- QCdf uncertainty $\mathcal{O}(1/m_b^2, \alpha_s^3)$
- Data consistent (**too few abs. BRs**)
- NP? $\Delta_P \sim \Delta_V \sim -20\%$ **possible**
 - ➡ We will learn something important!