

Constraining Higgs-Higgs-Z couplings in the 3HDM

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2HDM vs 3HDM gauge couplings

focus on

vector-vector-scalar, VVS

vector-scalar-scalar, VSS

(also quartic couplings $VVSS$)

Notation - Higgs basis

Consider first 3HDM, later “simplify” to 2HDM

Field notation:

$$\phi_i = e^{i\theta_i} \begin{pmatrix} \phi_i^+ \\ (v_i + \eta_i + i\chi_i)/\sqrt{2} \end{pmatrix}, \quad i = 1, 2, 3$$

Higgs basis:

$$H_1 = \begin{pmatrix} G^+ \\ (v + \eta_1^{\text{HB}} + iG_0)/\sqrt{2} \end{pmatrix}, \quad H_i = \begin{pmatrix} \varphi_i^{\text{HB}+} \\ (\eta_i^{\text{HB}} + i\chi_i^{\text{HB}})/\sqrt{2} \end{pmatrix}, \quad i = 2, 3$$

Mass eigenstates:

$$h_i = R_{ij} \varphi_j^{\text{HB}}$$

$$\varphi_i^{\text{HB}} = \{ \eta_1^{\text{HB}}, \eta_2^{\text{HB}}, \eta_3^{\text{HB}}, \chi_2^{\text{HB}}, \chi_3^{\text{HB}} \}, \quad i = 1, \dots, 5.$$

Kinetic terms:

$$\mathcal{L}_{\text{kin}} = \sum_{i=1,2,3} (D_\mu \phi_i)^\dagger (D^\mu \phi_i)$$

3HDM Gauge couplings

Cubic gauge-gauge-scalar part:

$$\mathcal{L}_{VVh} = \left(gm_W W_\mu^+ W^{\mu-} + \frac{gm_Z}{2 \cos \theta_W} Z_\mu Z^\mu \right) \sum_{i=1}^5 \frac{e_i}{v} h_i$$

$m_W = \frac{1}{2} g v$

\swarrow VVS

For the cubic gauge-scalar-scalar terms, we find

$$\mathcal{L}_{Vhh} = -\frac{g}{2 \cos \theta_W} \sum_{i=1}^5 \sum_{j=1}^5 \frac{\lambda_{ij}}{v} (h_i \overset{\leftrightarrow}{\partial}_\mu h_j) Z^\mu + \frac{g}{2} \sum_{i=1}^5 \sum_{k=1}^2 \left[\frac{f_{ki}}{v} (h_k^+ \overset{\leftrightarrow}{\partial}_\mu h_i) W^{\mu-} + \text{h.c.} \right]$$

$$+ \left(ieA^\mu + \frac{ig \cos 2\theta_W}{2 \cos \theta_W} Z^\mu \right) \sum_{k=1}^2 (h_k^+ \overset{\leftrightarrow}{\partial}_\mu h_k^-)$$

\swarrow VSS

\swarrow VSS+

2HDM Gauge couplings

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$$\begin{aligned} \mathcal{L}_{Vhh} = & -\frac{g}{2 \cos \theta_W} \sum_{i=1}^3 \sum_{j=1}^3 \overset{\text{VSS}}{\lambda_{ij}} \frac{1}{v} (h_i \overset{\leftrightarrow}{\partial}_\mu h_j) Z^\mu + \frac{g}{2} \sum_{i=1}^3 \left[\overset{\text{VSS}^+}{\frac{f_i}{v}} (h^+ \overset{\leftrightarrow}{\partial}_\mu h_i) W^{\mu-} + \text{h.c.} \right] \\ & + \left(ieA^\mu + \frac{ig \cos 2\theta_W}{2 \cos \theta_W} Z^\mu \right) (h^+ \overset{\leftrightarrow}{\partial}_\mu h^-) \end{aligned}$$

VVS & VSS Gauge couplings

Higgs basis:

$$\text{3HDM, VVS: } e_i = vR_{i1}$$

$$\text{3HDM, VSS: } \lambda_{ij} = v(R_{i2}R_{j4} + R_{i3}R_{j5}) - (i \leftrightarrow j)$$

$$\text{2HDM, VVS: } e_i = vR_{i1}$$

$$\text{2HDM, VSS: } \lambda_{ij} = v(R_{i2}R_{j3}) - (i \leftrightarrow j) = v\epsilon_{ijk}R_{k1} = \epsilon_{ijk}e_k$$

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Note that this is given in terms of the VVS coupling!

Not true for the 3HDM!

2HDM Gauge couplings

For a 3×3 orthogonal matrix;

$$R_{i2}R_{j3} - R_{j2}R_{i3} = \epsilon_{ijk}R_{k1} = \epsilon_{ijk}e_k/v$$

Example decomposition:

$$R = \begin{pmatrix} R_{11} & R_{12} & R_{13} \\ R_{21} & R_{22} & R_{23} \\ R_{31} & R_{32} & R_{33} \end{pmatrix} = \begin{pmatrix} c_1 c_2 & s_1 c_2 & s_2 \\ -(c_1 s_2 s_3 + s_1 c_3) & c_1 c_3 - s_1 s_2 s_3 & c_2 s_3 \\ -c_1 s_2 c_3 + s_1 s_3 & -(c_1 s_3 + s_1 s_2 c_3) & c_2 c_3 \end{pmatrix}$$

Pick $i = 1, j = 2$:

$$\begin{aligned} R_{12}R_{23} - R_{22}R_{13} &= (s_1 c_2)(c_2 s_3) - (c_1 c_3 - s_1 s_2 s_3)s_2 \\ &= s_1 c_2^2 s_3 - c_1 s_2 c_3 + s_1 s_2^2 s_3 = s_1 s_3 - c_1 s_2 c_3 = R_{31} \end{aligned}$$

$$i.e., \quad \lambda_{ij} = \epsilon_{ijk}e_k$$

3HDM Rotation matrix

$$h_i V V \quad \text{vertex} \quad e_i = v R_{i1}$$

R can be constructed from **10 independent angles**, two constructions:

$$1 : \quad R = \mathcal{O}^{(12)} \mathcal{O}^{(13)} \mathcal{O}^{(14)} \mathcal{O}^{(15)} \mathcal{O}^{(23)} \mathcal{O}^{(24)} \mathcal{O}^{(25)} \mathcal{O}^{(34)} \mathcal{O}^{(35)} \mathcal{O}^{(45)}$$

$$2 : \quad R = \mathcal{O}^{(45)} \mathcal{O}^{(35)} \mathcal{O}^{(34)} \mathcal{O}^{(25)} \mathcal{O}^{(24)} \mathcal{O}^{(23)} \mathcal{O}^{(15)} \mathcal{O}^{(14)} \mathcal{O}^{(13)} \mathcal{O}^{(12)}$$

with non-zero elements:

$$\begin{aligned} \mathcal{O}_{ii}^{(ij)} &= \cos \theta_{ij}, & \mathcal{O}_{ij}^{(ij)} &= \sin \theta_{ij} \\ \mathcal{O}_{ji}^{(ij)} &= -\sin \theta_{ij}, & \mathcal{O}_{ii}^{(ij)} &= \cos \theta_{ij} \\ \mathcal{O}_{kk}^{(ij)} &= 1 \text{ for } k \neq i, j \end{aligned}$$

By construction “1” above:

$$\begin{aligned} R_{11} &= \cos \theta_{12} \cos \theta_{13} \cos \theta_{14} \cos \theta_{15} \\ R_{21} &= -\sin \theta_{12} \cos \theta_{13} \cos \theta_{14} \cos \theta_{15} \\ R_{31} &= -\sin \theta_{13} \cos \theta_{14} \cos \theta_{15} \\ R_{41} &= -\sin \theta_{14} \cos \theta_{15} \\ R_{51} &= -\sin \theta_{15} \end{aligned}$$

Only four angles are needed for e_1, e_2, e_3, e_4, e_5 , independent of other six angles

3HDM Rotation matrix

Modulo the factor v , the e_i are “embedded” in R as follows:

$$R = \begin{pmatrix} e_1 & x & x & x & x \\ e_2 & x & x & x & x \\ e_3 & x & x & x & x \\ e_4 & x & x & x & x \\ e_5 & x & x & x & x \end{pmatrix}$$

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whereas λ_{ij} involves the remaining elements of the rows i and j :

$$\lambda_{ij} = (R_{i2}R_{j4} + R_{i3}R_{j5}) - (i \leftrightarrow j)$$
$$R = \begin{pmatrix} e_1 & \boxed{x} & \boxed{x} & \boxed{x} & \boxed{x} \\ e_2 & \boxed{x} & \boxed{x} & \boxed{x} & \boxed{x} \\ e_3 & \boxed{x} & \boxed{x} & \boxed{x} & \boxed{x} \\ e_4 & \boxed{x} & \boxed{x} & \boxed{x} & \boxed{x} \\ e_5 & \boxed{x} & \boxed{x} & \boxed{x} & \boxed{x} \end{pmatrix}$$

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illustrate for $i = 1, j = 2$:

$$\lambda_{12} = (R_{12}R_{24} - R_{22}R_{14})$$

“cross product”
blue \times red

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$\eta_1^{\text{HB}}, \eta_2^{\text{HB}}, \eta_3^{\text{HB}}, \chi_2^{\text{HB}}, \chi_3^{\text{HB}}$

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$\eta_1^{\text{HB}}, \eta_2^{\text{HB}}, \eta_3^{\text{HB}}, \chi_2^{\text{HB}}, \chi_3^{\text{HB}}$

numerology

- 10 angles
- 5 e_i [four angles + v]
- 10 non-zero λ_{ij} $[(5 \times 4)/2 - 5]$

The e_i and λ_{ij} are not all independent!

Constraints:

- $\sum_i e_i^2 = v^2$
- $\sum_j \lambda_{ij}^2 = v^2 - e_i^2$
- $\sum_{i,j} \lambda_{ij}^2 = v^2(5 - 1) = 4v^2$

non-trivial relations

Relation 1

$$e_i e_j = - \sum_k \lambda_{ik} \lambda_{jk}$$

Relation 2

$$e_i = -\frac{1}{8v} \sum_{j,k,l,m} \epsilon_{ijklm} \lambda_{jk} \lambda_{lm}$$

Relation 3

$$\lambda_{ij} = -\frac{1}{2v} \sum_{k,l,m} \epsilon_{ijklm} e_k \lambda_{lm}$$

May write seven λ_{ij} in terms of e 's and **three** other λ 's. Choose

$$\lambda_{12} = a\lambda_{34} + b\lambda_{35} + c\lambda_{45},$$

and similarly for $\lambda_{13}, \lambda_{14}, \lambda_{15}, \lambda_{23}, \lambda_{24}, \lambda_{35}$

non-trivial relations

Adding the squares (subject to the normalization), we arrive at

$$(v^2 - e_5^2)\lambda_{34}^2 + (v^2 - e_4^2)\lambda_{35}^2 + (v^2 - e_3^2)\lambda_{45}^2 \\ + 2e_4e_5\lambda_{34}\lambda_{35} - 2e_3e_5\lambda_{34}\lambda_{45} + 2e_3e_4\lambda_{35}\lambda_{45} - (e_1^2 + e_2^2)v^2 = 0$$

λ_{34} , λ_{35} , λ_{45} located on ellipsoid

May solve quadratic equation for one of them

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Bottom line: In addition to **five** e_i need only **TWO** λ_{ij} to describe all λ_{ij}

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Bottom line: In addition to **five** e_i need only **TWO** λ_{ij} to describe all λ_{ij}


4 angles


2 functions of 10 angles

non-trivial relations

Consider CP conservation in the 3HDM
5 neutral states: 3 even, 2 odd.

From the “non-trivial” relations, it follows that

$$e_i = e_j = \lambda_{ij} = 0$$

is a sufficient condition for no CP violation
in couplings between gauge bosons and neutral scalars

Alignment

Alignment:

$$e_1 = v \quad e_2 = e_3 = e_4 = e_5 = 0$$

Then

$$\lambda_{12} = \lambda_{13} = \lambda_{14} = \lambda_{15} = 0$$

The remaining ones are pairwise equal:

$$\lambda_{23} = -\lambda_{45}, \quad \lambda_{24} = \lambda_{35}, \quad \lambda_{25} = -\lambda_{34}$$

We also have

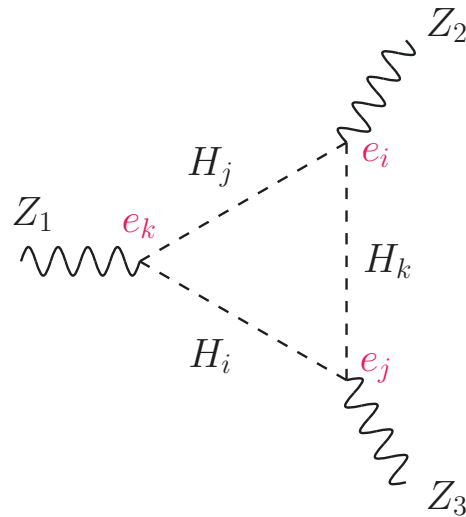
$$\lambda_{34}^2 + \lambda_{35}^2 + \lambda_{45}^2 = v^2$$

The moduli are located on a sphere

Solving for one, we are left with **two free parameters**

CP violation and alignment

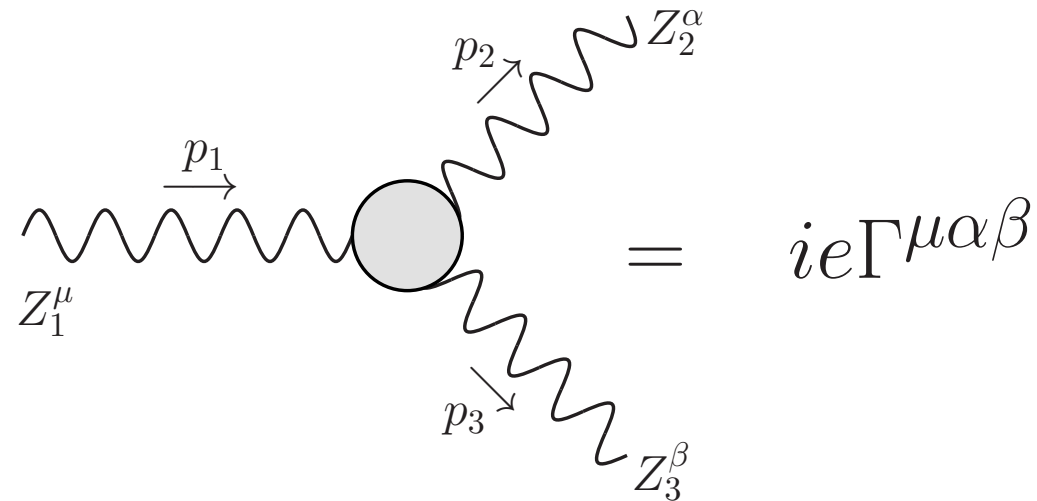
In a CP-violating 2HDM, all pairs of neutral scalars couple to the Z , allowing the triangle diagram



The existence of these couplings induces a CP-violating amplitude,

CP violation and alignment

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$$e\Gamma_{ZZZ}^{\alpha\beta\mu} = ie \frac{p_1^2 - M_Z^2}{M_Z^2} \left[f_4^Z (p_1^\alpha g^{\mu\beta} + p_1^\beta g^{\mu\alpha}) + f_5^Z \epsilon^{\mu\alpha\beta\rho} \ell_\rho \right]$$

↑
violates CP

CP violation and alignment

Also other diagrams, but importantly

f_4^Z is proportional to the invariant $\text{Im } J_2 \propto e_1 e_2 e_3$

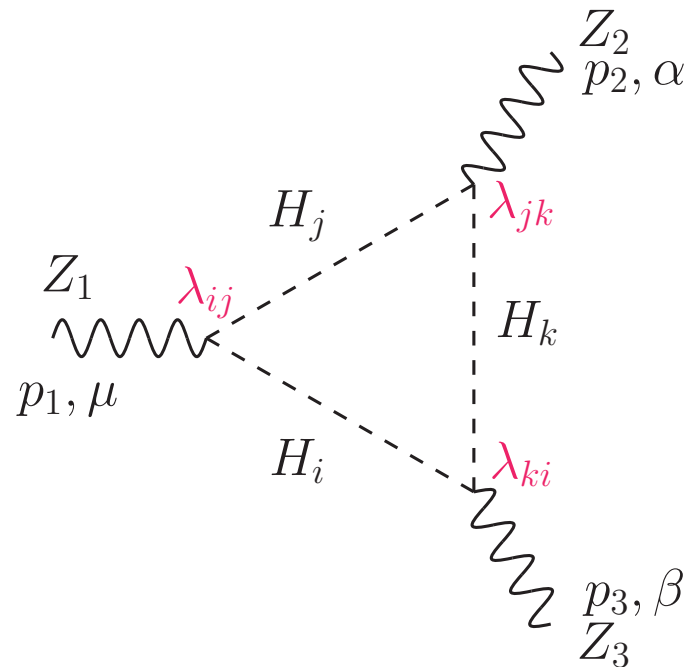
CP-violating

In the alignment limit, two of the e_i vanish, the ZZZ amplitude vanishes

$$e_i \rightarrow v \quad \rightarrow \quad e_j, e_k \rightarrow 0 \quad \text{Im } J_2 \rightarrow 0$$

CP violation and alignment

In a 3HDM



contributions proportional to $\lambda_{ij} \lambda_{jk} \lambda_{ki}$

This does not vanish in the alignment limit!

CP violation and alignment

Alignment limit

Let h_i be aligned, $e_i \rightarrow v$, $e_j \rightarrow 0$ for $j \neq i$

Consequence: $\lambda_{ij} \rightarrow 0$, but λ_{jk} unconstrained, for $j, k \neq i$

Example, Weinberg potential $Z_2 \times Z_2$ symmetric:

$$V = V_2 + V_4$$

$$V_2 = -[m_{11}(\phi_1^\dagger \phi_1) + m_{22}(\phi_2^\dagger \phi_2) + m_{33}(\phi_3^\dagger \phi_3)]$$

$$V_4 = V_0 + V_{\text{ph}},$$

$$\begin{aligned} V_0 = & \lambda_{11}(\phi_1^\dagger \phi_1)^2 + \lambda_{12}(\phi_1^\dagger \phi_1)(\phi_2^\dagger \phi_2) + \lambda_{13}(\phi_1^\dagger \phi_1)(\phi_3^\dagger \phi_3) + \lambda_{22}(\phi_2^\dagger \phi_2)^2 \\ & + \lambda_{23}(\phi_2^\dagger \phi_2)(\phi_3^\dagger \phi_3) + \lambda_{33}(\phi_3^\dagger \phi_3)^2 \\ & + \lambda'_{12}(\phi_1^\dagger \phi_2)(\phi_2^\dagger \phi_1) + \lambda'_{13}(\phi_1^\dagger \phi_3)(\phi_3^\dagger \phi_1) + \lambda'_{23}(\phi_2^\dagger \phi_3)(\phi_3^\dagger \phi_2) \end{aligned}$$

$$V_{\text{ph}} = \lambda_1(\phi_2^\dagger \phi_3)^2 + \lambda_2(\phi_3^\dagger \phi_1)^2 + \lambda_3(\phi_1^\dagger \phi_2)^2 + \text{h.c.}$$

CP violation and alignment

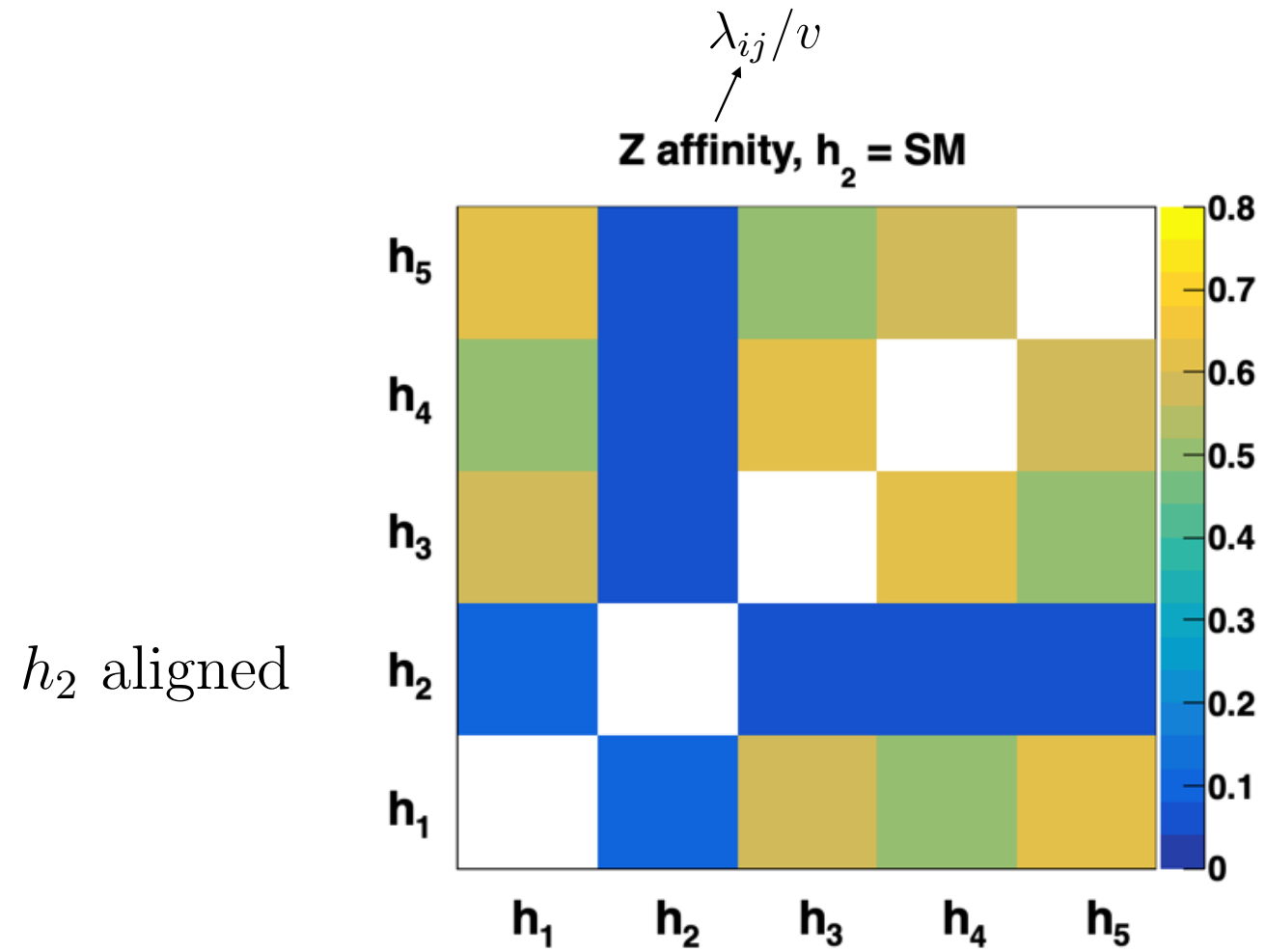
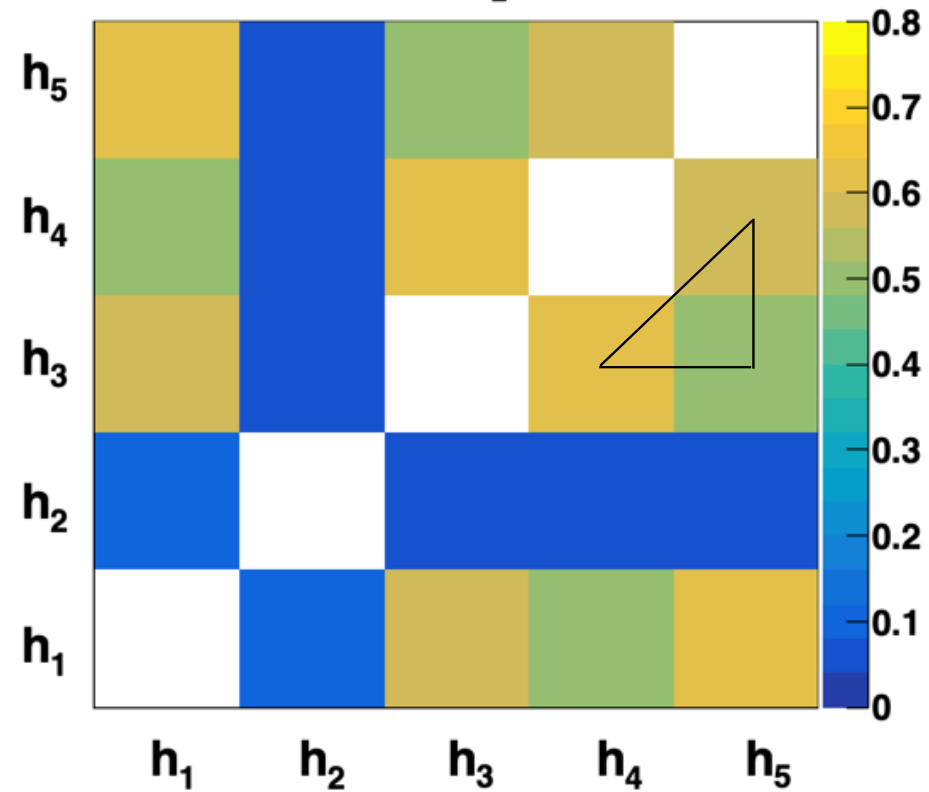


Figure from parameter scan in 2209.06499

CP violation and alignment

λ_{ij}/v
 ↑
Z affinity, $h_2 = \text{SM}$



h_2 aligned

for example:

$$\lambda_{34} \lambda_{45} \lambda_{53} \neq 0$$

Figure from parameter scan in 2209.06499

Summary

- vector (gauge)-scalar couplings:
 - VVS ($WW h_i, ZZ h_i$): e_i
 - VSS ($h_i h_j Z$): λ_{ij}
 - VSS ($h_i h_k^+ W^-$): d_{ki}
- determined by potential, relations among them
- 2HDM: $\lambda_{ij} = \epsilon_{ijk} e_k$ (*)
- 2HDM: ZZZ amplitude CP conserving in Alignment limit
- 3HDM: Relation (*) not valid
- 3HDM: ZZZ amplitude CP violating also in Alignment limit

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- work in progress